

Start-to-end Simulation with Rare Isotope Beam for Post Accelerator of the RAON Accelerator

Hyunchang JIN* and Ji-Ho JANG

Rare Isotope Science Project, Institute for Basic Science, Daejeon 34047, Korea

(Received 19 November 2015, in final form 23 November 2015)

The RAON accelerator for the Rare Isotope Science Project (RISP) has been developed to create and accelerate various kinds of stable heavy ion beams and rare isotope beams for a wide range of science applications. In the RAON accelerator, the rare isotope beams generated by the Isotope Separation On-Line (ISOL) system will be transported through the post accelerator, namely, from the post Low Energy Beam Transport (LEBT) system and the post Radio Frequency Quadrupole (RFQ) to the superconducting linac (SCL3). The accelerated beams will be put to use in the low energy experimental hall or accelerated again by the superconducting linac (SCL2) in order to be used in the high energy experimental hall. In this paper, we will describe the results of the start-to-end simulations with the rare isotope beams generated by the ISOL system in the post accelerator of the RAON accelerator. In addition, the error analysis and correction at the superconducting linac SCL3 will be presented.

PACS numbers: 41.85.Ja, 52.59.Fn

Keywords: Start-to-end simulation, RAON accelerator

DOI: 10.3938/jkps.69.922

I. INTRODUCTION

The Rare Isotope Science Project (RISP) was established in December 2011 to construct the RAON (Rare isotope Accelerator Of Newness) accelerator for various science programs. The RAON accelerator project [1, 2] to generate and accelerate a variety of stable heavy ion beams and rare isotope beams for use in a wide range of basic science research and various applications is ongoing. For the production of various rare isotope beams, the RAON accelerator put to use the In-flight Fragmentation (IF) system and the Isotope Separation On-Line (ISOL) system. Initially, the IF system uses a diver linac, which consists of a 28 GHz superconducting electron cyclotron resonance ion source (ECR-IS), a main low energy beam transport (LEBT) section [3], a radio-frequency quadrupole (RFQ) accelerator, a medium energy beam transport (MEBT) section, a low energy superconducting linac (SCL1), a charge stripper section (CSS), and a high energy superconducting linac (SCL2). The beams accelerated by the SCL2 collide with the IF target, and then the rare isotope beams are created from the collisions. Next, the ISOL system uses a 70 MeV cyclotron as the driver to deliver a 70 kW beam power up to the ISOL target. The rare isotope beams created by the ISOL system are accelerated again by the post accelerator, which consists of a post LEBT [4], a post

RFQ, a post MEBT, and a low energy superconducting linac SCL3. The beams accelerated by the SCL3 will be delivered up to the low energy experimental hall or to the SCL2 after having passed the the post accelerator to the driver linac transport (P2DT) section [5]. A schematic layout of the RAON accelerator is shown in Fig. 1.

As mentioned above, the post accelerator consists of four sections: the post LEBT, the post RFQ, the post MEBT, and the SCL3. First, the lattice of the post LEBT was partly modified because that of the ISOL system was recently changed. The new lattice design and beam dynamics simulations of the post LEBT were presented in [4]. The main goal of the post LEBT is the stable transportation of the heavy ion beams and rare isotope beams up to the post RFQ. To accomplish this goal, the electro-static quadrupoles are used in the post LEBT. Second, the post RFQ of the post accelerator and the RFQ of the driver linac are the same. The post RFQ is about 5 m long, and the number of cell is 245. The maximum peak surface electric field is 17 MV/m; thus, the post RFQ can accelerate a 10 keV/u beam to 500 keV/u. Third, the post MEBT includes eight quadrupoles and three buncher cavities, which are used to satisfy the beam requirements of the SCL3. Finally, the SCL3 is divided into three sections, SCL31, SCL32 and SCL33, depending on the cavity type and the number of cavities in one cryomodule. The first section, the SCL31 is composed of 22 quarter-wave resonator (QWR) type cavities, and each cavity is surrounded by one cry-

*E-mail: hcjin@ibs.re.kr

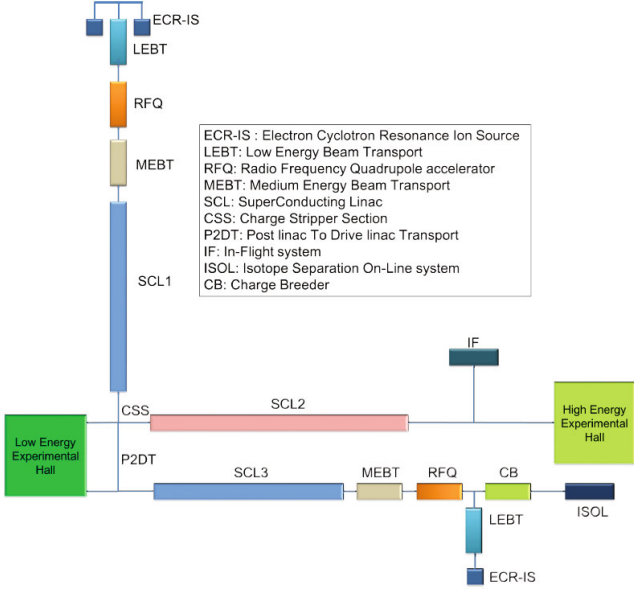


Fig. 1. (Color online) Layout of the RAON accelerator.

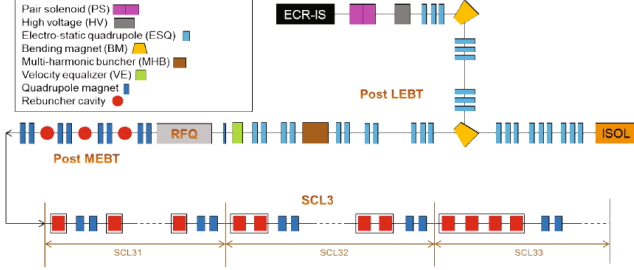


Fig. 2. (Color online) Layout of the post accelerator in the RAON accelerator.

omodule. The second section, the SCL32 has 26 half-wave resonator (HWR) type cavities, and one cryomodule surrounds two cavities. Third, the SCL33 has 76 HWR type cavities, and one cryomodule surrounds 4 cavities. A schematic layout of the post accelerator is shown in Fig. 2.

In this paper, we will present the results of the start-to-end simulations for the post accelerator with the reference beam $^{132}\text{Sn}^{33+}$. In addition, a study of the error analysis and correction at the superconducting linac SCL3 in the RAON accelerator will be described by using the graphical user interface (GUI) based on the MATLAB program and the DYNAC code [6,7].

II. START-TO-END SIMULATION

1. Beam information

The reference beam generated by the ISOL system was recently determined as a tin beam, $^{132}\text{Sn}^{33+}$. At the end

Table 1. Initial beam information.

Parameter	Value	Unit
Beam	$^{132}\text{Sn}^{33+}$	-
Energy	10.0	keV/u
Normalized transverse emittance	0.1	mm·mrad
β_0	0.2	mm/mrad
α_0	0.1	-
Number of macro-particles	100,000	-
Beam distribution	4D water-bag	-

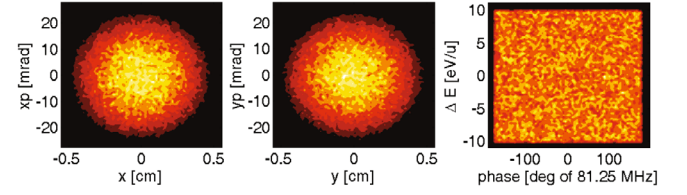


Fig. 3. (Color online) Beam distributions in the transverse and the longitudinal phase spaces.

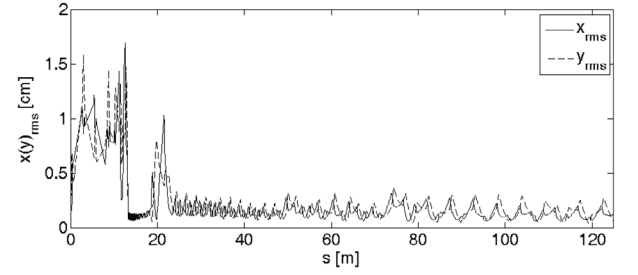


Fig. 4. Transverse rms beam size along the post accelerator.

of the ISOL system, the beam energy is 10 keV/u and the normalized transverse emittance is 0.1 mm·mrad. The basic beam information at the entrance of the post LEBT is listed in Table 1. Based on this beam information, we carry out a start-to-end simulation with 100,000 macro-particles by using the particle tracking code TRACK [8]. Figure 3 shows the beam distributions in the transverse and the longitudinal phase spaces at the entrance of the post LEBT.

2. Simulation results

Along the post accelerator, the beam pipe radius of each section varies and is 6.0 cm at the post LEBT, 2.5 cm at the post MEBT, and 2.0 cm at the SCL3. Those beam pipe radii were determined by considering the beam energy and the optimization of the cavities. Therefore, the beam size should be less than the beam pipe radius to avoid the beam loss at each section. In view of such a condition, the strength of each magnet was calculated, and the start-to-end simulation was con-

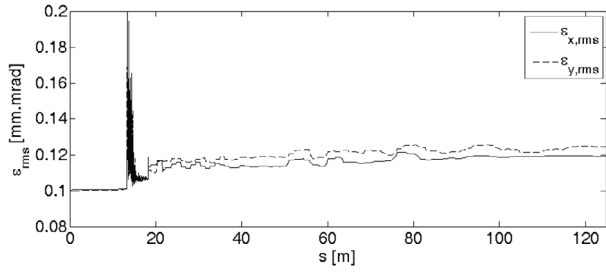


Fig. 5. Normalized transverse rms beam emittance along the post accelerator.

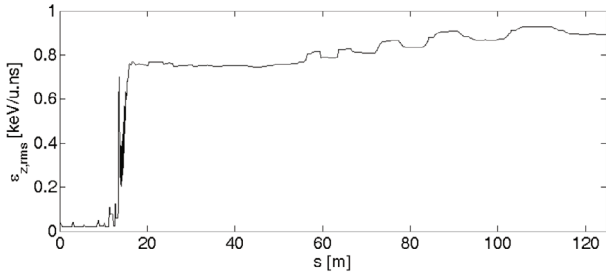


Fig. 6. Normalized longitudinal rms beam emittance along the post accelerator.

Table 2. Normalized transverse and longitudinal rms emittances at some positions along the post accelerator.

Parameter	Horizontal [mm-mrad]	Vertical [mm-mrad]	Longitudinal [keV/u.ns]
End of RFQ	0.113	0.111	0.750
End of MEBT	0.113	0.116	0.758
End of SCL31	0.114	0.119	0.742
End of SCL32	0.117	0.120	0.856
End of SCL33	0.119	0.125	0.889

ducted. Figure 4 shows the transverse root-mean-square (rms) beam size along the post accelerator. Especially, the rms beam size is kept less than 0.5 cm at the SCL3, and the beam transmission after the post RFQ is about 98%, which is one of the advantages of the RAON accelerator.

Figure 5 shows the normalized transverse rms beam emittance along the post accelerator. After the beam passes through the RFQ, the energy spread becomes larger, which affects the transverse emittance growth. Also, the vertical emittance increases at the SCL3 because of the effect of the electro-magnetic field of the QWR cavity.

The normalized longitudinal rms beam emittance along the post accelerator is shown in Fig. 6. After the beam has been longitudinally bunched by the post RFQ, the energy spread increases because of the long tail of the beam; thus, the longitudinal emittance increases. Table 2 shows the normalized transverse and longitudinal rms emittances at some positions along the post accel-

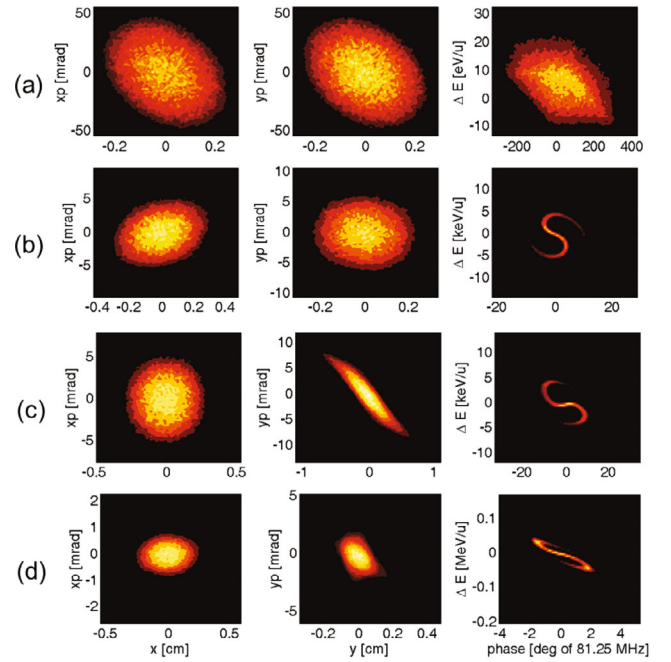


Fig. 7. (Color online) Beam distribution in the phase space along the post accelerator: (a) post LEFT exit, (b) post RFQ exit, (c) post MEBT exit, and (d) SCL3 exit.

erator. The growths of the horizontal, transverse and longitudinal emittances from the end of the RFQ to the end of the SCL33 are about 5.3%, 12.6% and 18.5%, respectively. The minimum requirement for the normalized transverse rms emittance at the low energy experimental halls is about 2 mm-mrad, which is much larger than the result of the simulation. Figure 7 shows the transverse and the longitudinal beam distributions at the end of each section. As mentioned above, the beam is longitudinally bunched by the post RFQ, and it leads to the growth of the longitudinal beam emittance.

III. ERROR SIMULATIONS

1. Error tolerance

In the SCL3, the beam orbit can be distorted by a variety of error sources, such as the magnet misalignment, the cavity field error and so on. After the beam loss caused by the orbit distortion has been checked, the effect of each error source can be verified. In this paper, the tolerance of each error source is determined when the beam loss caused by each error source becomes smaller than 0.1%. The TRACK code is used for the calculation of the tolerance and the results are listed in Table 3. In the simulations, the errors of the quadrupoles and the cavities are given by Gaussian distributions with rms values and are truncated at the 3 times the rms value. Among the error sources, the tolerance of the quadrupole

Table 3. Tolerance of error sources at the SCL3.

-	-	Value	Unit
Initial position	x	+0.68/-0.65	[cm]
Initial position	y	+0.27/-0.25	[cm]
Initial angle	xp	+8.0/-7.9	[mrad]
Initial angle	yp	+4.0/-3.7	[mrad]
Cavity	rms misalignment x,y	0.12	[cm]
Cavity	rms misalignment z	0.08	[cm]
Cavity	rms field amplitude	2.1	[%]
Cavity	rms field phase	0.9	[deg]
Quadrupole	rms misalignment x,y	0.018	[cm]
Quadrupole	rms misalignment z	1.2	[cm]
Quadrupole	rms tilt	16.2	[mrad]

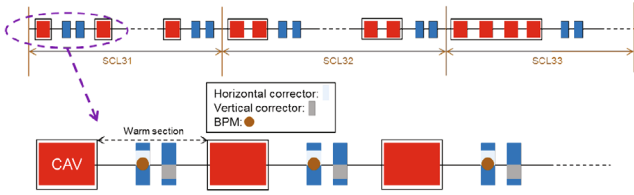


Fig. 8. (Color online) Layout of the SCL3 in the RAON accelerator.

Table 4. Errors used in the orbit correction simulations.

Parameter	Value	Unit
Initial position x,y	10	μm
Initial angle xp,yp	10	μrad
Cavity rms misalignment x,y	10	μm
Quadrupole rms misalignment x,y	100 – 500	μm

transverse rms misalignment is about 0.018 cm and is the most dominant one among all error sources.

2. Orbit correction

Figure 8 shows a schematic layout of the SCL3. For the orbit correction, a horizontal corrector and a beam position monitor (BPM) are located at the first quadrupole, and a vertical corrector is located at the second quadrupole at each warm section. With the correctors and BPMs, the orbit correction for the distorted orbit is carried out at the SCL3. The errors used in the orbit correction simulations are listed in Table 4. The beam orbit is distorted by these errors and then the orbit correction is carried out by using the correctors and the BPMs by using the singular value decomposition (SVD) method. The rms transverse misalignment of the quadrupole, which is the most dominant error source, changes from 100 μm to 500 μm .

For the orbit correction in the RAON accelerator, the

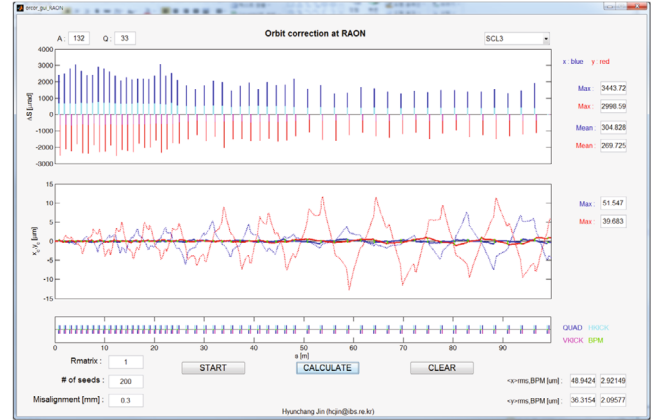

 Fig. 9. (Color online) GUI for the orbit correction in the RAON accelerator with a quadrupole misalignment of 300 μm and 200 random seeds.

 Table 5. Summary of the orbit correction for a quadrupole misalignment of 300 μm and 200 random seeds.

Parameter	Value	Unit
$\langle x \rangle_{rms,BPM}$ before correction	48.9	μm
$\langle y \rangle_{rms,BPM}$ before correction	36.3	μm
$\langle x \rangle_{rms,BPM}$ after correction	2.9	μm
$\langle y \rangle_{rms,BPM}$ after correction	2.1	μm
Average horizontal corrector strength	0.30	mrad
Average vertical corrector strength	0.27	mrad

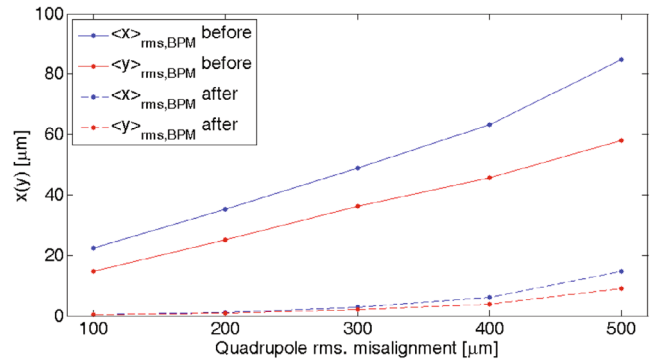


Fig. 10. (Color online) Rms orbit size at BPMs before and after the orbit correction.

GUI based on the MATLAB program and the DYNAC code has been developed. Figure 9 shows the GUI screen after the orbit correction for a quadrupole misalignment of 300 μm and 200 random seeds. As listed in Table 5, the horizontal (vertical) rms beam size decreases from 48.9 (36.3) μm to 2.9 (2.1) μm after the orbit correction. The average corrector strength is less than the mechanical maximum value.

Figure 10 shows the rms orbit size at the BPMs before and after the orbit correction for 100 – 500 μm rms quadrupole misalignment with 200 random seeds. For

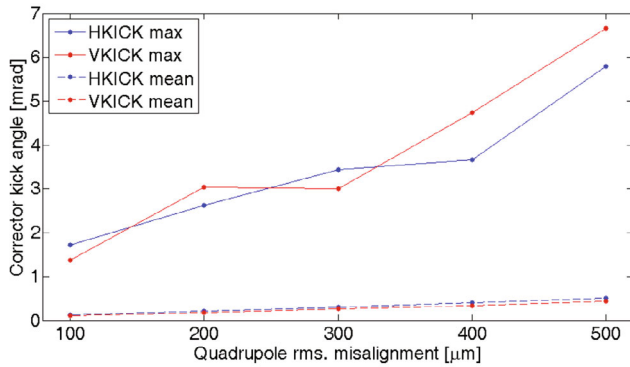


Fig. 11. (Color online) Maximum and average corrector strength for the orbit correction.

quadrupole rms misalignments from $100 \mu\text{m}$ to $500 \mu\text{m}$, the rms orbit size decreases about 90%. The maximum and the average corrector strengths for the orbit correction are shown in Fig. 11. For some random seeds, the corrector kick angle becomes larger, but the average value is less than 0.6 mrad, which is much less than the mechanical maximum kick angle of about 2.7 mrad. Research on the proper numbers and positions of the correctors and the BPMs will be continued.

IV. CONCLUSIONS

We presented the results of the start-to-end simulations with the rare isotope beam in the post accelerator of the RAON accelerator. For the post accelerator, a new reference beam, $^{132}\text{Sn}^{33+}$, from the ISOL system was tracked from the post LEBT to the superconducting linac SCL3. At the end of the SCL3, the beam energy reached about 27.7 MeV/u, and the rms beam size

along the post accelerator was kept much less than the beam pipe radii. The error analysis and correction were also performed at the SCL3. The tolerance of each error source was calculated, and the beam orbit distortion caused by the errors was corrected with the SVD method. Additionally, for various error sources, the orbit distortion was checked and corrected by using the correctors and the BPMs. As a result, the average kick angle of the correctors for the orbit correction was much less than the mechanical maximum value. Researches for the correction of the distorted beam orbit will be carried out continuously.

ACKNOWLEDGMENTS

This work was supported by the Rare Isotope Science Project of the Institute for Basic Science funded by the Ministry of Science, ICT and Future Planning and the National Research Foundation of Korea (2013M7A1A1075764).

REFERENCES

- [1] S. Kim *et al.*, Baseline Design Summary, http://risp.ibs.re.kr/orginfo/info_blds.do (2012).
- [2] D. Jeon *et al.*, J. Korean Phys. Soc. **55**, 7 (2014).
- [3] H. Jin and J. Jang, J. Korean Phys. Soc. **67**, 8 (2015).
- [4] H. Jin *et al.*, Rev. Sci. Instrum. **87**, 2 (2016).
- [5] H. Jin *et al.*, Nucl. Instr. and Meth. A **795**, 65 (2015).
- [6] E. Tanke, TH429, LINAC02 (2002).
- [7] H. Jin *et al.*, MOPJE017, IPAC15 (2015).
- [8] V. Aseev *et al.*, TPAT028, PAC05 (2005).