

Jeong Young Park and Gabor A. Somorjai

Understanding fundamental mechanisms for surface electronic excitation is of great importance in surface chemistry. Charge transport through metal–oxide interfaces plays a significant role in heterogeneous catalysis. Over the last several decades, a number of experimental and theoretical results suggest that this charge flow through metal-support interfaces leads to catalytic enhancement often observed in mixed catalysts. Direct measurement of charge flow on actual catalysts is a rather challenging task because it requires the use of an electronic circuit. This approach has been enabled by a catalytic nanodiode that is mainly composed of a catalytic metal and semiconducting oxides that form a Schottky contact. In this article, we describe the advances in this approach. We show that there is close connection between the phenomena of hot-electron creation and chemical reaction that occur at both gas–solid and liquid-solid interfaces. The intensity of hot-electron flow is well correlated with the turnover rates of corresponding reactions, which opens the possibility for developing new operando methodologies to monitor catalytic reactions as well as a novel scheme for the electronic control of chemical reactions.

Introduction

 In heterogeneous catalysis, one of the key technologies for increasing the activity and selectivity of a chemical reaction is the use of catalytically active metal nanoparticles (NPs) in combination with a suitable support.¹⁻⁶ This approach is based on earlier work that revealed significant changes in the activity of fine metal particles when placed on the support of a metal oxide, even though the support itself is inactive for this reaction. Specifically, in the 1960s, Schwab et al. discovered significant changes in the activity of fine metal particles when placed on the support of a metal oxide, $7/8$ and attributed this effect to the formation of a Schottky barrier at the metal– oxide interface with the subsequent transfer of charge carriers through the barrier, which affects the course of the surface reaction.^{3,7,9} This effect was also investigated by Boffa et al.¹⁰ using rhodium deposited on various reducible oxides. They observed a remarkable 14-fold increase in catalytic activity for $CO₂$ hydrogenation on three different oxides—TiOx, NbOx, and TaOx. The activity was highest at a half monolayer of oxide coverage where the oxide–metal interface area was at a maximum. Later, these phenomena were referred to as

the strong metal–support interaction (SMSI) effect,¹¹ which indicates the enhancement of catalytic activity when Group VIII metal catalysts (including Pt, Pd, Rh, Fe, Ni, and Ir) are supported on reducible oxides such as $CeO₂$, $Nb₂O₅$, and $TiO₂$.

 To understand the electronic origin of the SMSI effect, it is desirable to directly measure the flow of charge between the metal and the oxide. To achieve this goal, metal–oxide catalysts need to be combined with a Schottky diode, and thus, the catalytic nanodiode was developed.^{1,2}

 This article considers the main aspects of research aimed at developing methods for studying the transfer of charge carriers through metal–support interfaces under catalytic reaction conditions. First, we review the mechanisms for hot-electron excitation and transport through metal–oxide interfaces. We then show various schemes for detecting hot electrons that are generated during catalytic processes. We provide an overview of the latest results from detecting hot electrons in supported catalysts during chemical reactions at both gas–solid and liquid–solid interfaces. Hot electrons can be generated by photon absorption on the surface, therefore, the photocatalytic

Jeong Young Park, Center for Nanomaterials and Chemical Reactions, Institute for Basic Science; and Department of Chemistry, Korea Advanced Institute of Science and Technology, Republic of Korea; jeongypark@kaist.ac.kr

Gabor A. Somorjai , University of California, Berkeley, USA ; somorjai@berkeley.edu doi:10.1557/mrs.2019.295

process amplified by surface plasmons is quite relevant to hot electron flow. We discuss this aspect as well.

Principle of catalytic nanodiode

One explanation to justify the generation of hot electrons is related to the difference in the heat capacity of electrons and phonons. The electronic heat capacity (C_{electron}) of most metals in thermal equilibrium is about one hundred times smaller than the lattice heat capacity (*Clattice*) at 300 K. For example, for copper,

$$
\frac{C_{electron}}{C_{lattice}} \approx \frac{\gamma \cdot T}{3R} = \frac{0.70 \times 10^{-3} \times 300}{3 \times 8.3} = 0.008,
$$
 (1)

where R is the gas constant, *T* is the temperature, and γ is the Sommerfeld constant $(0.70 \text{ mJ/mol} \cdot \text{K}^2 \text{ for Cu})$. When heat is deposited from exothermic surface reactions or photon flux, the electrons heat up much faster (femtoseconds) than the lattice (picoseconds) because their heat capacity is much lower. Another mechanism to explain hot-electron generation is described in the framework of a nonadiabatic mechanism. The transition of the system to a new state occurs with one of the electronically excited states leading to the creation of an electron–hole (*e–h*) pair.

Relaxation of hot electrons happens on the femtosecond to picosecond time scale, and their mean free path is on the order of tens of nanometers.12 This implies two detection strategies. The first is to obtain sufficient time resolution for observing this excitation. The second is to use a nanometer-scale energy

barrier for irreversible transfer of the hot-electron flux. For the first approach, two-photon timeresolved photoemission spectroscopy has been employed to directly study the dynamics of optically excited electrons at metal and semiconductor surfaces. This technique has been applied to the direct measurement of hot electron relaxation in noble and transition metals¹³ and surface-state dynamics on clean and adsorbate-covered metal surfaces,¹⁴ as well as charge-carrier dynamics in semiconductors, where much work has been performed. For the second approach, Nienhaus et al. used Schottky diodes in high-vacuum experiments aimed at detecting *e–h* pairs excited during the chemisorption of atomic hydrogen on the surface of Ag thin films supported on *n*-Si and *p*-Si substrates, as shown in **Figure 1**a.12,15,16 The hot-electron flow generated by hydrogen adsorption was detected as the current, as shown in Figure 1b.

It is important to consider the question of how to detect hot charge carriers that are excited in metal catalysts by an exothermic chemical reaction. As previously mentioned, in the case of highly exothermic reactions taking place on the surface of metals with a low work function, nonadiabatic reactions can be detected by observing chemiluminescence or the emission of excited electrons into a vacuum. However, for conventional catalytic reactions, this approach does not work because the work function of most catalytic metals (e.g., Pt, Pd, and Rh) is much larger than the excess energy released by the surface reaction. Therefore, for a long time, it was unclear if low-energy chemical reactions create hot charge carriers in metal catalysts.

A solution to this problem was found by Somorjai and Park, who proposed the use of thin-film Schottky diodes, called catalytic nanodiodes, for detecting hot electrons during catalytic reactions.17–19 An example of this scheme for the planar Pt/Si catalytic nanodiode under H_2O_2 decomposition reaction is shown in Figure 1c.20 A thin metal film acts as both the catalyst and electrode, thus allowing the flux of the chemically excited hot electrons to be measured. To ensure a reliable supply of hot carriers from the surface of the film to the Schottky contact, the film thickness should not exceed the ballistic mean free path of the electrons in the metal.20

An exothermic chemical reaction catalyzed on a metal surface creates distributions of hot charge carriers (electrons and holes) located between levels corresponding to the Fermi energy in the metal and the maximum chemical energy liberated by the reaction. To detect hot carriers via a catalytic nanodiode, excited electrons are detected because of internal emission through the Schottky barrier formed at the metal– semiconductor interface, as shown in the energy-band diagram

Figure 1. (a) Ag/Si Schottky diode used to measure hot-electron flow generated by chemisorption of atomic hydrogen. (b) Chemicurrent as a function of time measured with a Ag/Si diode. Reprinted with permission from Reference 12. © 2002 Elsevier. (c) Scheme of the Pt/Si catalytic nanodiodes during the H₂O₂ decomposition process, and (d) the corresponding energy-band diagram. From Reference 20.

(Figure 1d) of the Pt/Si catalytic nanodiode. Since the thickness of the metal film is on the order of the mean free path of the hot electrons (i.e., typically ≤ 15 nm), these excited charge carriers cross the film without significant attenuation and reach the metal–semiconductor interface while still energetic enough to overcome the Schottky barrier.

Various schemes of catalytic nanodiodes

In earlier experiments on the detection of hot electrons during catalytic chemical reactions, planar Schottky nanodiodes were used that consist of an ultrathin Pt film deposited on an n -TiO₂ or n -GaN substrate.^{19,21–23} Since the detection of hot electrons occurs during their transfer through the Schottky contact, a necessary condition for the operation of a catalytic nanodiode is a properly sized Schottky barrier.^{12,24} The height of the barrier must be large enough to limit the transfer of thermal carriers through the metal–semiconductor interface, which allows the detection of only the hot charge carriers excited in nonadiabatic reactions. However, the Schottky barrier must not exceed the exothermicity of the surface reaction. Thus, considerations for the transparency of the barrier for charge carriers with a certain excess energy should be taken into account when choosing a semiconductor support for the nanodiodes.24,25

For an ideal Schottky diode, the height of the barrier, φ_{b} , can be found using $\varphi_b = \Phi_m - \chi_s$, where Φ_m is the metal work function and χ _s is the electron affinity of the semiconductor.²⁶ However, for real nanodiodes, significant deviations from this equation can be observed. The height of the barrier is determined by fitting an experimentally measured current–voltage curve to the diode equation:

$$
I(V) = I_{\rm s} \left[\exp\left(\frac{e_0 (V - IR_{\rm ser})}{nk_{\rm B} T}\right) - 1 \right],\tag{2}
$$

where $I_{\rm s} = AA^*T^2 \exp\left(-\frac{\epsilon_0 G}{k_{\rm B}}\right)$ $I_s = AA^* T^2 \exp\left(-\frac{e_0 \varphi_b}{k_B T}\right)$ is the saturation current, *A* is

the area of a Schottky contact, A^* is the Richardson constant, *T* is the temperature, e_0 is the elementary charge, φ_h is the Schottky barrier height, k_B is the Boltzmann constant, *V* is the voltage, R_{ser} is the series resistance, and *n* is the ideality factor, which is a measure of how closely the diode follows the ideal diode equation.

The use of planar Schottky nanodiodes (**Figure 2**a) has made it possible to carry out proof-of-concept experiments demonstrating the feasibility of studying nonadiabatic processes in catalytic reactions using solid-state devices. However, the structure of such thin-film nanodiodes differs significantly from the supported metal–semiconductor catalysts used in industrial processes. To fill this gap, several research groups have experimented with more complex nanodiodes whose structures simulate real-world catalysts.

For instance, Park et al. demonstrated the possibility of detecting nonadiabatic electronic excitation in colloid NPs of platinum (Pt) using the catalytic H_2 and CO oxidation reactions.27,28 The idea behind this experiment is seen in Figure 2b—Pt NPs of a defined size were deposited as two-dimensional arrays on an Au thin film supported on the TiO₂ surface of the nanodiode. The Au film creates an electrical connection between the Pt NPs and the external circuit, thus allowing continuous hot-electron flow. Another idea for making nanodiodes that are close to real three-dimensional

> (3D) catalysts was proposed by Schierbaum et al., who studied the effects of charge carrier creation during catalytic reactions on $Pt/TiO₂$ nanodiodes based on porous titanium oxide layers fabricated using plasma electrolytic oxidation of a Ti metal foil,^{29,30} as shown in Figure 2c. Karpov et al. and Jeon et al. also conducted chemicurrent studies using similar nanodiodes consisting of a Pt mesh supported on mesoporous TiO₂ and ZrO₂ layers.^{31–34} This approach is related to building a 3D catalytic nanodiode that allows a higher flux of hot electrons and catalytic activity. Recently, Goddeti et al. demonstrated a semi-3D catalytic diode using $Pt/TiO₂$ nanotubes that resulted in much higher hot-electron generation.^{35,36} In addition, an inverse catalyst with cobalt oxide NPs on Pt/TiO₂ catalytic nanodiodes was demonstrated by Lee et al. (Figure 2d).³⁷ The interface between Pt and CoO revealed a higher chemicurrent as well as higher catalytic activity during hydrogen oxidation, emphasizing the importance of the metal–oxide interface in heterogeneous catalysis and hot-electron generation.38

Hot-electron detection during gas-phase and liquid-phase reactions

Using a catalytic nanodiode, hot-electron fluxes were detected during various gas-phase reactions (**Figure 3**a), including CO oxidation,¹⁹ NO/CO reaction,³⁹ methanol oxidation,⁴⁰ and hydrogen oxidation.23,28,41 A typical chemicurrent signal that corresponds to the transfer of hot electrons excited by the catalytic reaction in a H₂ (15 Torr) + O_2 (745 Torr) gas mixture on 1.7-nm Pt NPs on $Au/TiO₂$ nanodiodes is shown in Figure 3b.28 The chemicurrent demonstrates a noticeable temperature dependence that is caused by an increase in the rate of H2 oxidation at the elevated temperature of the Pt catalyst. The chemicurrent value is also dependent on the metal thickness. As can be seen, an increase in both the Pt NPs size and Au film thickness leads to a sharp decrease in the chemicurrent, which is associated with a decrease in hot charge carrier flux because of electron–electron and electron–phonon scattering.28

The Au layer that is used for the ohmic contact can be replaced by other conductive thin layers such as graphene. To study the chemicurrent effect on Pt NPs supported on a graphene layer, a graphene-based catalytic nanodiode was developed.42 Graphene has several extraordinary properties that make it a promising candidate for use as a catalytic support. Graphene is a particularly excellent conductor of ballistic electrons and is thus an interesting material for detecting

nonadiabatic electronic excitation in catalytic reactions using Schottky nanodiodes.

As mentioned earlier, the detection of nonadiabatic reactions using catalytic nanodiodes occurs because of internal emission of hot electrons through the Schottky barrier and therefore does not depend on the density of the reacting medium. This provides an opportunity to study catalytic reactions occurring at liquid–solid interfaces that are of great practical importance for heterogeneous catalysis.43–47 The scheme for nanodiode-based hot-electron detection during solid–liquid reactions is shown in Figure 3c. To date, most data from the experimental detection of chemically excited hot electrons during liquid-phase reactions were reported for the catalytic decomposition reaction of hydrogen peroxide

 $(2H_2O_2 \rightarrow 2H_2O+O_2)$, which is commonly used in industry as an oxidizing agent for pulp and paper bleaching, wastewater treatment, textile production, chemical synthesis, and many other applications.^{48,49} In addition, the decomposition of H_2O_2 is highly exothermic $(\Delta H = -25.3 \text{ kJ/mol})^{47}$ and can therefore serve as a model system for mechanistic studies of energy and charge transfer in catalysts at liquid–solid interfaces.

Typical chemicurrent signals measured on planar 10-nm Pt/*n*-Si, Au/*n*-Si, and Ag/*n*-Si nanodiodes in an aqueous solution containing H_2O_2 are shown in Figure 3d. The chem-

> icurrent is measured when the nanodiode is immersed in the H_2O_2 solution.^{43,47} The shape of the chemicurrent signal reflects the surface reaction rate. Initially, when the surface reaction rate is the greatest, the chemicurrent reaches a peak value. Then, as the steady-state reaction is established, the chemicurrent decreases and acquires a value that varies only little with time. The value of this steady-state chemicurrent is a clear function of the H_2O_2 concentration in the solution.46 Thus, chemicurrent measurements provide an easy-to-implement method for studying catalytic processes at the liquid–solid interface in real time.

Surface plasmon driven hotelectron flux and its influence on photocatalytic reactions

The application of plasmonics to the field of heterogeneous catalysis has received a lot of attention in recent years because of the potential to efficiently harvest and utilize solar energy for chemical processes. Plasmonic metal nanostructures or hybrid materials (composed of catalytically active semiconductor and plasmonic metal NPs) under light irradiation are efficient sources for generating hot electrons. This concept has been demonstrated by several groups who observed the amplification of hot-electron flux on plasmonic nanodiodes.⁵⁰⁻⁵²

Figure 4. (a) Schematic of the MAPbI₃/plasmonic-Au/TiO₂ nanodiode, and cross-sectional elemental energy-dispersive spectroscopy mapping image of the MAPbI₃-covered plasmonic nanodiode. (b) Incident photon to current efficiency (IPCE) as a function of incident photon energy measured on the MAPbI₃/plasmonic-Au/TiO₂ nanodiodes according to the number of MAPbI₃ layers.⁵³ (c) Schematic of photoanode composed of Au nanoparticles (NPs) on $TiO₂$ nanotube arrays (TNAs) and the working principle of the photoelectrochemical cell. (d) Incident-photon-to-electron conversion plots of electrodes composed of Au NPs on TNA with different sized Au NPs and bare TNA collected at 1.23 V versus reversible hydrogen electrode (RHE).⁵⁴ Note: CB, conduction band; VB, valence band.

The efficiency of hot-electron generation is also influenced by the lifetime of hot electrons. This aspect was investigated by Park et al. who modified a plasmonic nanodiode with perovskite materials. In this study, an MAPbI₃-modified plasmonic nanodiode combining $MAPbI₃$ (perovskite) with the local surface plasmon resonance effect of nanostructured Au was utilized, as shown in **Figure 4**a.⁵³ A higher hot-electron flux was observed in this device (as shown in Figure 4b) due to the presence of the perovskite material; this was attributed to the extended lifetime of the hot electrons, which was measured using femtosecond transient absorption experiments.

A higher hot-electron flux can influence the rates of catalytic chemical reactions through excitation by the photoinduced surface plasmon resonances of the metal nanostructures. Moon et al.⁵⁴ investigated this aspect by studying photocatalytic water splitting mediated by the injection of hot carriers using TiO₂ nanotube arrays decorated with various sized plasmonic Au NPs, as shown in the experimental scheme in Figure 4c. The internal quantum efficiency of the Au NPs on TiO₂ nanotube arrays photoanodes for the oxygen evolution reaction in visible light increased from 1.96 to 7.68%, when the size of the Au NPs decreased from 30 to 5 nm (Figure 4d). Moreover, the smaller NP-decorated $TiO₂$ nanotube arrays showed about a ninefold lower Au loading compared with the larger NPs with the same coverage.⁵⁴

Enhancement of the local electric field that serves as "hot spots" in plasmonic photocatalysis boosts the generation of electrons and holes. An increased photoinduced charge is generated locally in the $TiO₂$ via local field enhancement of

the plasmonic NPs.^{55,56} Considering that strong electric enhancement occurs at the edges of the NPs, smaller NPs will increase the total perimeter length compared with larger sizes with the same coverage; thus, more hot carriers can travel over the interface. Interesting future studies involve the direct probing of hot-electron transfer using ultrafast studies or metal–semiconductor diodes during photocatalytic processes.

Conclusion

In this article, we highlighted recent studies aimed at understanding the nature of charge transfer arising during catalytic reactions on supported metal catalysts. First, we discussed the basic mechanisms of catalytic nanodiodes where nonadiabatic effects could be potentially responsible for creating hot charge carriers in metal catalysts during exothermic reactions and photon irradiation. We also provided an overview of the various schemes and applications for catalytic nanodiodes, including hot-electron detection during solid–gas and solid–liquid reactions. The detection of hot electrons in the form of a chemicurrent allows

for *operando* studies of nonadiabatic effects during gas–solid and liquid–solid reactions. Because of the intrinsic relation between surface plasmons and hot-electron flux, the direct measurement of hot electrons can be a useful tool to monitor the local electric field and "hot spot" during photocatalytic reactions.

A more fundamental understanding of hot-electron phenomena is required at both theoretical and experimental levels. Hot electrons on a metallic surface can be created by external energy deposition in the form of photons, ions, electrons, and chemical reactions. Therefore, we can consider hot electrons to be a major mediator for general energy conversion. The scheme of energy conversion from chemical (i.e., catalytic reactions) and photon-to-electrical (i.e., hot-electron current) energy may give insight into other fields, including solar cells and electrochemical and photocatalytic devices.

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Jeong Young Park is a professor in the Department of Chemistry at the Korea Advanced Institute of Science and Technology, Republic of Korea, and an associate director of the Center for Nanomaterials and Chemical Reactions at the Institute for Basic Science. He received his BS and PhD degrees from the Department of Physics, Seoul National University, Republic of Korea, in 1993 and 1999. After completing postdoctoral research at the University of Maryland, he joined the Lawrence Berkeley National Laboratory where he worked as a postdoctoral fellow from 2002 to 2005, and then as a staff scientist from 2006 to 2009. His research focuses on

surface phenomena covering energy dissipation and conversion at surfaces, nanotribology, catalysis, and hot-electron transport through metal–oxide interfaces. He has published 230 peer-reviewed articles and has edited one book in the fields of surface science and catalysis. Park can be reached by email at jeongypark@kaist.ac.kr.

Gabor A. Somorjai is a professor of chemistry at the University of California, Berkeley, and a university professor of the University of California System. He received his PhD degree in chemistry from the University of California, Berkeley, in l960. He is the recipient of numerous awards, including the American Chemical Society's William H. Nichols Medal (2015), the National Academy of Sciences NAS Award in Chemical Sciences (2013), Banco Bilbao Vizcaya Argentaria Foundation Frontiers of Knowledge Award in Basic Sciences (2011), Ente Nazionale Idrocarburi New Frontiers of Hydrocarbons Prize (2011), the Priestley Medal (2008), the

Irving Langmuir Prize (2007), the National Medal of Science (2002), and the Wolf Prize (1998). Somorjai was awarded the Materials Research Society's highest honor, the Von Hippel Award in 1997. He has educated more than 130 doctoral candidates and 250 postdoctoral fellows. He has written three textbooks, and more than 1200 scientific papers in the fields of surface chemistry, heterogeneous catalysis, and solid-state chemistry. Somorjai can be reached by email at somorjai@berkeley.edu .

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