RESEARCH

Open Access

Two unconditionally stable difference schemes for time distributed-order differential equation based on Caputo–Fabrizio fractional derivative



Haili Qiao¹, Zhengguang Liu² and Aijie Cheng^{1*}

*Correspondence: aijie@sdu.edu.cn ¹ School of Mathematics, Shandong University, Jinan, China Full list of author information is available at the end of the article

Abstract

We consider distributed-order partial differential equations with time fractional derivative proposed by Caputo and Fabrizio in a one-dimensional space. Two finite difference schemes are established via Grünwald formula. We show that these two schemes are unconditionally stable with convergence rates $O(\tau^2 + h^2 + \Delta \alpha^2)$ and $O(\tau^2 + h^4 + \Delta \alpha^4)$ in discrete L^2 , respectively, where $\Delta \alpha$, h, and τ are step sizes for distributed-order, space, and time variables, respectively. Finally, the performance of difference schemes is illustrated via numerical examples.

Keywords: Distributed-order; Caputo–Fabrizio derivatives; Compact finite difference; Stability and convergence; Numerical experiments

1 Introduction

The concept of fractional derivative can be traced back to the letter of Leibniz to L'Hospital in 1695. Mathematician Euler discovered the Γ function in 1729, which established the foundation of fractional derivative. Compared with the classical integer-order models, fractional derivatives provide a more profound and comprehensive explanation about memory, heredity, non-locality of complex phenomena and processes. At present, the most commonly used definitions of fractional derivatives are the Grünwald-Letnikov derivative, the Riemann-Liouville derivative, the Riesz derivative, and the Caputo derivative [1]. Some new fractional derivatives, such as the Caputo-Fabrizio derivative and the Atangana–Baleanu fractional derivative, have tremendously promoted the capability of modeling complex physical phenomena and processes. The connatural behavior has been analyzed for evolution equations generated by three fractional derivatives namely the Riemann-Liouville, Caputo-Fabrizio, and Atangana-Baleanu fractional derivatives [2–4]. The Caputo–Fabrizio derivative is defined via an integral operator without singular kernel [5, 6], and it has supplementary motivating properties compared to the others. For example, it can describe substance heterogeneities and configurations with different scales, which noticeably cannot be characterized by the local theories. The properties

© The Author(s) 2020. This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons licence, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons licence, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons licence and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this licence, visit http://creativecommons.org/licenses/by/4.0/.



of numerical approximation are discussed for the initial value problems with Caputo– Fabrizio fractional derivative operator [7, 8]. Several valuable applications to the real world data, such as blood ethanol concentration system, chickenpox disease model, and dengue fever outbreak model, have been discussed; the discussions were based on the fractional derivatives in the Caputo sense, Caputo–Fabrizio sense, and Atangana–Baleanu–Caputo sense, respectively [9–15].

In the past decades, anomalous transport has been discussed in a wide range of applications, for example, turbulence [16–19], geoscience [20], bioscience [21, 22], porous media [23–27], viscoelastic material [28–30], and so on. The underlying anomalous features, manifesting in memory-effects, sharp peak, power-law distribution, self-similar structure, and nonlocal interactions, can be well described by fractional partial differential equations [31–33]. However, in some applications, a single power law cannot characterize more complex physical processes, such as decelerating sub-diffusion [34, 35], accelerating super-diffusion, and random processes subordinated to Wiener processes [31, 36]. These processes can be described by the distributed-order derivatives first introduced by Caputo, see [35] and the references therein. Recently, it is believed that distributedorder differential equations are more suitable for describing complex dynamic systems than fractional-order models.

Usually, it is difficult to obtain analytical solutions of these models, so different numerical schemes for fractional partial differential equations(FPDE), which can exhibit history dependence and nonlocal features, have been developed. The common numerical methods for fractional partial differential equations include finite-difference [37-41], finiteelement methods [42, 43], finite volume schemes [44-46], mixed finite element schemes [47], spectral/spectral-element schemes [48, 49], the decomposition method [50, 51], and others [52]. In [53–56], numerical analysis for distributed-order FPDEs was extensively investigated. More recently, Tomovski and Sandev [57] expressed the solutions of generalized distributed-order diffusion equations equipped with fractional time derivatives by use of the Fourier-Laplace transform. The authors [58] obtained the existence about Cauchy problems for the diffusion equations with time distributed-order derivative, they also computed it by Laplace transform and Fourier transform. Luchko [56] discussed the dependence of the uniqueness and continuity of the generalized time fractional diffusion equations on initial conditions. However, there are few numerical studies about the distributed-order equations, especially the study about high-order schemes is almost blank. It has motivated us to find efficient numerical schemes for distributed-order FPDEs.

Consider the following distributed-order fractional diffusion equation equipped with the fractional derivative developed by Caputo and Fabrizio [5]:

$$\begin{cases} \mathcal{D}_t^w u(x,t) = \frac{\partial^2 u(x,t)}{\partial x^2} + f(x,t), & (x,t) \in \Omega \times (0,T], \\ u(x,0) = 0, & x \in \Omega, \\ u(x,t) = 0, & x \in \partial\Omega, 0 \le t \le T, \end{cases}$$
(1)

where the distributed-order derivative in the Caputo-Fabrizio sense is defined as follows:

$$\mathcal{D}_{t}^{w}u(x,t) = \int_{v_{1}}^{v_{2}} w(\alpha)_{0}^{CF} D_{t}^{\alpha}u(x,t) \, d\alpha, \quad w(\alpha) \geq 0, \int_{v_{1}}^{v_{2}} w(\alpha) \, d\alpha = c_{0} > 0$$

$${}_{0}^{CF}D_{t}^{\alpha}u(x,t)=\frac{1}{1-\alpha}\int_{0}^{t}u'(x,s)\exp\left[-\alpha\frac{t-s}{1-\alpha}\right]ds,$$

and $0 < v_1 < v_2 < 1$, $\Omega = (a, b)$, and f(x, t) represents the source term.

This paper has the following organization. Section 2 provides some preliminary work, including the space partition, the inner products, and norms. Section 3 is devoted to the finite difference scheme, which is proved to be second-order with respect to distributed-order, space, and time respectively, and the stability and convergence rates are also studied. In Sect. 4, a compact finite difference method is proposed, the stability and convergence are rigorously proved. Section 5 provides numerical tests that illustrate the reliability of theoretical analysis. In Sect. 6 we summarize this paper and indicate the possible work in the future.

2 Preliminary

In order to establish finite difference methods, the space interval [a, b] is partitioned by $x_i = a + ih$ $(0 \le i \le M)$, and the time interval [0, T] is partitioned by $t_k = k\tau$ $(0 \le k \le N)$, with M, N positive integers and h = (b - a)/M the spatial grid size and $\tau = T/N$ the temporal step size. Divide the integral interval $[v_1, v_2]$ to 2J subintervals, $\Delta \alpha = \frac{v_2 - v_1}{2J}$. Denote $\alpha_l = v_1 + l\Delta\alpha$, l = 0, 1, 2, ..., 2J. Let $\Omega_h = \{x_i : 0 \le i \le M\}$, $\Omega_\tau = \{t_k : 0 \le k \le N\}$, thus the computational range $[a, b] \times [0, T]$ has a discretization $\Omega_h \times \Omega_\tau$. A grid function u is represented as $u_i^k = u(x_j, t_k)$. U_i^k denotes the finite difference solution at the grid point (x_j, t_k) .

Denote $V_h = \{v | v = (v_0, v_1, ..., v_M)\}$. For any grid function $v \in V_h$, we define the difference quotient operators

$$\delta_x v_j = \frac{v_j - v_{j-1}}{h}, \quad 1 \le j \le M; \qquad \delta_x^2 v_j = \frac{v_{j+1} - 2v_j + v_{j-1}}{h^2}, \quad 1 \le j \le M - 1,$$

and the average operator

$$\mathcal{A}\nu_{j} = \begin{cases} \frac{1}{12}(\nu_{j+1} + 10\nu_{j} + \nu_{j-1}), & 1 \le j \le M - 1, \\ \nu_{j}, & j = 0, M. \end{cases}$$

It is easy to see

$$\mathcal{A} v_j = \left(I + \frac{h^2}{12} \delta_x^2 \right) v_j,$$

where *I* denotes the identical operator. We also denote $Av = (Av_1, Av_2, ..., Av_M)$ for $v = (v_1, v_2, ..., v_M)$, and A(u, v) = (Au, v) with (\cdot, \cdot) the discrete L^2 inner product defined below.

For two grid functions $u, v \in V_h^0 = \{v | v \in V_h, v_0 = v_M = 0\}$, the discrete inner products and norms are defined as follows:

$$(u,v)=h\sum_{j=1}^{M-1}u_jv_j,\qquad \langle u,v\rangle=h\sum_{j=1}^M\delta_xu_j\delta_xv_j,$$

and

$$\|u\|_{2}^{2} = (u, u), \qquad |u|_{1}^{2} = \langle u, u \rangle, \qquad \|u\|_{\infty} = \max_{1 \le j \le M-1} |u_{j}|.$$

By summation by parts and the boundary conditions, it is easy to get that

$$(\delta_x^2 u, v) = -\langle u, v \rangle = (u, \delta_x^2 v).$$

For the average operator A, we also define

$$\|\nu\|_A^2 \triangleq \mathcal{A}(\nu,\nu) = (\mathcal{A}\nu,\nu).$$

Additionally, denote by $V_{\tau} = \{\nu | \nu = (\nu^0, \nu^1, \dots, \nu^N)\}$ the space of all grid functions with respect to Ω_{τ} . Given a grid function $\nu \in V_{\tau}$, we denote the difference quotient in time direction as

$$\delta_t \nu^k = \frac{\nu^k - \nu^{k-1}}{\tau}.$$

Lemma 1 ([59]) The following inequalities hold for any grid function $v \in V_h^0$:

$$\|\nu\|_{2} \leq \frac{b-a}{\sqrt{6}} |\nu|_{1}, \qquad \frac{1}{3} \|\nu\|_{2}^{2} \leq \|\nu\|_{A}^{2} \leq \|\nu\|_{2}^{2}.$$

Throughout this paper, the symbol C will indicate a genetic constant that depends on the function u and may assume different values at different occurrences, and the symbol O(S) indicates a quantity bounded above by S with a constant.

3 A second-order scheme for space and distributed order

Lemma 2 ([60]) *If* $s(\alpha) \in C^2[v_1, v_2]$, then we have

$$\int_{\nu_1}^{\nu_2} s(\alpha) d\alpha = \Delta \alpha \sum_{l=0}^{2J} c_l s(\alpha_l) - \frac{\Delta \alpha^2}{12} s^{(2)}(\xi), \quad \xi \in (\nu_1, \nu_2),$$

where

$$c_l = \begin{cases} \frac{1}{2}, & l = 0, 2J, \\ 1, & l = 1, 2, \dots, 2J - 1. \end{cases}$$

Differential equation (1) has the form at the node (x_i, t_n)

$$\mathcal{D}_{t}^{w}u(x_{i},t_{n}) = \frac{\partial^{2}u(x_{i},t_{n})}{\partial x^{2}} + f(x_{i},t_{n}), \quad 1 \le i \le M-1, 1 \le n \le N-1.$$
(2)

Let $s(\alpha, x_i, t_n) = w(\alpha)_0^{CF} D_t^{\alpha} u(x_i, t_n)$, and we set $w(\alpha) \in C^2[v_1, v_2]$, ${}_0^{CF} D_t^{\alpha} u(x_i, t)|_{t=t_n} \in C^2[v_1, v_2]$. By Lemma 2 we obtain

$$\begin{split} \mathcal{D}_t^w u(x_i, t_n) &= \int_{\nu_1}^{\nu_2} s(\alpha, x_i, t_n) \, d\alpha \\ &= \Delta \alpha \sum_{l=0}^{2J} c_l s(\alpha_l, x_i, t_n) - \frac{\Delta \alpha^2}{12} \frac{\partial^2 s(\alpha, x_i, t_n)}{\partial \alpha^2} \bigg|_{\alpha = \xi_i^n} \\ &= \Delta \alpha \sum_{l=0}^{2J} c_l w(\alpha_l)_0^{CF} D_t^\alpha u(x_i, t_n) + O(\Delta \alpha^2), \end{split}$$

where $\xi_i^n \in (v_1, v_2)$.

We have [61]

where

$$G_k^{\alpha_l} = \exp\left[-\alpha_l \tau \frac{k}{1-\alpha_l}\right] - \exp\left[-\alpha_l \tau \frac{k+1}{1-\alpha_l}\right].$$

Noting that u(x, 0) = 0, we have

$$\mathcal{D}_{t}^{w}u(x_{i},t_{n}) = \Delta\alpha \sum_{l=0}^{2J} c_{l}w(\alpha_{l}) \frac{1}{\alpha_{l}\tau} \left[G_{0}^{\alpha_{l}}u(x_{i},t_{n}) - \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_{l}} - G_{n-k}^{\alpha_{l}} \right) u(x_{i},t_{k}) \right] + O(\tau^{2} + \Delta\alpha^{2})$$
(4)

and

$$\frac{\partial^2 u}{\partial x^2}(x_i, t_n) = \delta_x^2 u_i^n + O(h^2).$$
(5)

Substituting (4) and (5) into (2), we get

$$\Delta \alpha \sum_{l=0}^{2J} c_l w(\alpha_l) \frac{1}{\alpha_l \tau} \left[G_0^{\alpha_l} u_i^n - \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) u_i^k \right]$$

= $\delta_x^2 u_i^n + f_i^n + (r_1)_i^n, \quad 1 \le i \le M - 1, 1 \le n \le N,$ (6)

and there exists a constant C satisfying

$$(r_1)_i^n \leq C(\tau^2 + h^2 + \Delta \alpha^2).$$

The boundary and initial conditions show that

$$u_{i}^{0} = 0, \quad 1 \le i \le M - 1, u_{0}^{k} = u_{M}^{k} = 0, \quad 0 \le k \le N.$$
(7)

Canceling the infinitesimal part and replacing the true solution u_i^k with approximate solution U_i^k , we get the finite difference method as follows:

$$\Delta \alpha \sum_{l=0}^{2J} c_l w(\alpha_l) \frac{1}{\alpha_l \tau} \left[G_0^{\alpha_l} U_i^n - \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) U_i^k \right] \\ = \delta_x^2 U_i^n + f_i^n, \quad 1 \le i \le M - 1, 1 \le n \le N,$$

$$U_i^0 = 0, \quad 1 \le i \le M - 1,$$

$$U_0^n = U_M^n = 0, \quad 0 \le n \le N.$$
(8)

3.1 Stability analysis

We give a lemma about $G_k^{\alpha_l}$ which is useful in the stability analysis.

Lemma 3 ([61]) *From the definition of* $G_k^{\alpha_l}$ *, we have* $G_k^{\alpha_l} > 0$ *and* $G_{k+1}^{\alpha_l} < G_k^{\alpha_l}$ *,* $\forall k \leq N$.

Rearranging equation(8), we get

$$\Delta \alpha \sum_{l=0}^{2J} c_l w(\alpha_l) \frac{1}{\alpha_l \tau} \sum_{k=1}^n G_{n-k}^{\alpha_l} U_i^k - \delta_x^2 U_i^n$$

= $\Delta \alpha \sum_{l=0}^{2J} c_l w(\alpha_l) \frac{1}{\alpha_l \tau} \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} U_i^k + f_i^n, \quad 1 \le i \le M-1, 1 \le n \le N.$ (9)

Denote $\eta_1 = \Delta \alpha \sum_{l=0}^{2J} c_l w(\alpha_l) \frac{1}{\alpha_l \tau}$. Multiplying both sides of equation (9) by $h U_i^n$ and making summation with respect to *i* from 1 to M - 1, we obtain

$$\eta_1 \sum_{k=1}^n G_{n-k}^{\alpha_l} (U^k, U^n) + |U^n|_1^2 = \eta_1 \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} (U^k, U^n) + (f^n, U^n).$$

Noting Lemma 1 and Young's inequality, we have

$$\frac{6}{L^2} \| U^n \|_2^2 + \eta_1 G_0^{\alpha_l} \| U^n \|_2^2 \le \eta_1 \sum_{k=1}^{n-1} (G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l}) (U^k, U^n) + \frac{L^2}{24} \| f^n \|_2^2 + \frac{6}{L^2} \| U^n \|_2^2.$$

Using the triangle inequality, we obtain

$$\eta_1 G_0^{\alpha_l} \left\| U^n \right\|_2^2 \le \frac{\eta_1}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left(\left\| U^k \right\|_2^2 + \left\| U^n \right\|_2^2 \right) + \frac{L^2}{24} \left\| f^n \right\|_2^2$$

Consequently, we get

$$\eta_{1}G_{0}^{\alpha_{l}} \left\| U^{n} \right\|_{2}^{2} + \frac{\eta_{1}}{2} \sum_{k=1}^{n-1} \left(G_{n-k}^{\alpha_{l}} - G_{n-k-1}^{\alpha_{l}} \right) \left\| U^{n} \right\|_{2}^{2}$$

$$\leq \frac{\eta_{1}}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_{l}} - G_{n-k}^{\alpha_{l}} \right) \left\| U^{k} \right\|_{2}^{2} + \frac{L^{2}}{24} \left\| f^{n} \right\|_{2}^{2}.$$

Namely,

$$\frac{\eta_1}{2} \left(G_0^{\alpha_l} + G_{n-1}^{\alpha_l} \right) \left\| U^n \right\|_2^2 \le \frac{\eta_1}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left\| U^k \right\|_2^2 + \frac{L^2}{24} \left\| f^n \right\|_2^2$$

Then it follows that

$$\frac{\eta_1}{2} \sum_{k=1}^n G_{n-k}^{\alpha_l} \| U^k \|_2^2 \le \frac{\eta_1}{2} \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} \| U^k \|_2^2 + \frac{L^2}{24} \| f^n \|_2^2.$$

Denote

$$Q(U^n) = \frac{\eta_1}{2} \sum_{k=1}^n G_{n-k}^{\alpha_l} \| U^k \|_2^2.$$

Making summation with respect to i from 1 to N shows that

$$Q(U^N) \le Q(U^0) + \frac{L^2}{24} \sum_{n=1}^N ||f^n||_2^2.$$

Observing that $U^0 = 0$ implies $Q(U^0) = 0$, we have

$$Q(U^N) \le \frac{L^2}{24} \sum_{n=1}^N \|f^n\|_2^2.$$

Theorem 1 For scheme (8), stability inequality is as follows:

$$Q(\mathcal{U}^m) \le \frac{L^2}{24} \sum_{n=1}^m \|f^n\|_2^2, \quad \forall 0 \le m \le N.$$

3.2 Optimal error estimates

Combining equations (6), (7) with (8), we have the following error equation:

$$\eta_1 \sum_{k=1}^n G_{n-k}^{\alpha_l} e_i^k - \delta_x^2 e_i^n = \eta_1 \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} e_i^k + (r_1)_i^n.$$
(10)

Here $(r_1)_i^n = O(\tau^2 + h^2 + \Delta \alpha^2)$, $e_i^k = u_i^k - U_i^k$, $\forall 0 \le k \le N$.

Multiply two sides of equation (10) with he_i^n and make summation with respect to *i* from 1 to M - 1 to get

$$\eta_1 \sum_{k=1}^n G_{n-k}^{\alpha_l}(e^k, e^n) + |e^n|_1^2 = \eta_1 \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l}(e^k, e^n) + ((r_1)^n, e^n).$$

Noting Lemma 1 and Young's inequality, we have

$$\frac{6}{L^2} \left\| e^n \right\|_2^2 + \eta_1 G_0^{\alpha_l} \left\| e^n \right\|_2^2 \le \eta_1 \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left(e^k, e^n \right) + \frac{L^2}{24} \left\| (r_1)^n \right\|_2^2 + \frac{6}{L^2} \left\| e^n \right\|_2^2.$$

By the triangle inequality, we obtain

$$\eta_1 G_0^{\alpha_l} \|e^n\|_2^2 \leq \frac{\eta_1}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left(\|e^k\|_2^2 + \|e^n\|_2^2 \right) + \frac{L^2}{24} \|(r_1)^n\|_2^2.$$

Transposition gives rise to

$$\begin{split} &\eta_1 G_0^{\alpha_l} \left\| e^n \right\|_2^2 + \frac{\eta_1}{2} \sum_{k=1}^{n-1} \left(G_{n-k}^{\alpha_l} - G_{n-k-1}^{\alpha_l} \right) \left\| e^n \right\|_2^2 \\ &\leq \frac{\eta_1}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left\| e^k \right\|_2^2 + \frac{L^2}{24} \left\| (r_1)^n \right\|_2^2, \end{split}$$

namely,

$$\frac{\eta_1}{2} \left(G_0^{\alpha_l} + G_{n-1}^{\alpha_l} \right) \left\| e^n \right\|_2^2 \le \frac{\eta_1}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left\| e^k \right\|_2^2 + \frac{L^2}{24} \left\| (r_1)^n \right\|_2^2.$$

Then we have

$$\frac{\eta_1}{2} \sum_{k=1}^n G_{n-k}^{\alpha_l} \| e^k \|_2^2 \le \frac{\eta_1}{2} \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} \| e^k \|_2^2 + \frac{L^2}{24} \| (r_1)^n \|_2^2.$$

Denote

$$Q(e^{n}) = \frac{\eta_{1}}{2} \sum_{k=1}^{n} G_{n-k}^{\alpha_{l}} \|e^{k}\|_{2}^{2}.$$

Making summation with respect to n from 1 to N shows that

$$Q(e^N) \le Q(e^0) + \frac{L^2}{24} \sum_{n=1}^N ||(r_1)^n||_2^2.$$

Observing that $e^0 = 0$ implies $Q(e^0) = \frac{\eta_1}{2} G_0^{\alpha_l} \|e^0\|_2^2 = 0$, we have

$$Q(e^N) \leq \frac{L^2}{24} \sum_{n=1}^N ||(r_1)^n||_2^2.$$

Theorem 2 For scheme (8), the following stability inequality holds:

$$\|e^m\|_2 \le C(\tau^2 + h^2 + \Delta \alpha^2), \quad \forall 0 \le m \le N.$$

4 A fourth-order method for space and distributed order

The following lemma is necessary for establishing a scheme with fourth-order accuracy in spatial variable,

Lemma 4 ([62]) Let $\theta(s) = (1-s)^3 [5-3(1-s)^2]$. Suppose that $f(x) \in C^6[a, b]$, h = (b-a)/M, $x_i = a + ih \ (0 \le i \le M)$, then we have

$$\frac{1}{12} \Big[f''(x_{i-1}) + 10f''(x_i) + f''(x_{i+1}) \Big] \\ = \frac{1}{h^2} \Big[f(x_{i-1}) - 2f(x_i) + f(x_{i+1}) \Big] \\ + \frac{h^4}{360} \int_0^1 \Big[f^{(6)}(x_i - sh) + f^{(6)}(x_i + sh) \Big] \theta(s) \, ds, \quad 1 \le i \le M - 1.$$

Lemma 5 ([60]; The composite Simpson formula) Let $s(\alpha) \in C^4[\nu_1, \nu_2]$, then we have

$$\int_{\nu_1}^{\nu_2} s(\alpha) d\alpha = \Delta \alpha \sum_{l=0}^{2J} d_l s(\alpha_l) - \frac{\Delta \alpha^4}{180} s^{(4)}(\xi), \quad \xi \in (\nu_1, \nu_2),$$

where

$$d_{l} = \begin{cases} \frac{1}{3}, & l = 0, 2J, \\ \frac{2}{3}, & l = 2, 4, \dots, 2J - 4, 2J - 2, \\ \frac{4}{3}, & l = 1, 3, \dots, 2J - 3, 2J - 1. \end{cases}$$

The differential equation at the node (x_i, t_n) has the form

$$\mathcal{D}_t^w u(x_i, t_n) = \frac{\partial^2 u(x_i, t_n)}{\partial x^2} + f(x_i, t_n), \quad 1 \le i \le M - 1, 1 \le n \le N - 1.$$

$$(11)$$

Denote $s(\alpha, x_i, t_n) = w(\alpha)_0^{CF} D_t^{\alpha} u(x_i, t_n)$ and suppose that $w(\alpha) \in C^4[v_1, v_2]$, ${}_0^{CF} D_t^{\alpha} u(x_i, t_n)|_{t=t_n} \in C^4[v_1, v_2]$. By Lemma 5, we get

$$\begin{split} \mathcal{D}_t^w u(x_i, t_n) &= \int_{\nu_1}^{\nu_2} s(\alpha, x_i, t_n) \, d\alpha \\ &= \Delta \alpha \sum_{l=0}^{2J} d_l s(\alpha_l, x_i, t_n) - \frac{\Delta \alpha^4}{180} \frac{\partial^4 s(\alpha, x_i, t_n)}{\partial \alpha^4} \Big|_{\alpha = \xi_i^n} \\ &= \Delta \alpha \sum_{l=0}^{2J} d_l w(\alpha_l)_0^{CF} D_t^{\alpha_l} u(x_i, t_n) + O(\Delta \alpha^4), \end{split}$$

where $\xi_i^n \in (\nu_1, \nu_2)$.

Noting u(x, 0) = 0 and taking use of equation (3), we have

$$\mathcal{D}_{t}^{w}u(x_{i},t_{n}) = \Delta\alpha \sum_{l=0}^{2J} d_{l}w(\alpha_{l})\frac{1}{\alpha_{l}\tau} \left[G_{0}^{\alpha_{l}}u(x_{i},t_{n}) - \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_{l}} - G_{n-k}^{\alpha_{l}} \right)u(x_{i},t_{k}) \right] + O(\tau^{2} + \Delta\alpha^{4}).$$
(12)

By Lemma 4, we obtain

$$\mathcal{A}\frac{\partial^2 u}{\partial x^2}(x_i, t_n) = \delta_x^2 u_i^n + O(h^4).$$
(13)

Substituting (12)-(13) into (11), we get

$$\mathcal{A}\Delta\alpha \sum_{l=0}^{2j} d_l w(\alpha_l) \frac{1}{\alpha_l \tau} \left[G_0^{\alpha_l} u_i^n - \sum_{k=1}^{n-1} (G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l}) u_i^k \right]$$

= $\delta_x^2 u_i^n + \mathcal{A} f_i^n + (r_2)_i^n, \quad 1 \le i \le M - 1, 1 \le n \le N,$ (14)

and there exists a constant C satisfying

$$(r_2)_i^n \le C(\tau^2 + h^4 + \Delta \alpha^4).$$

Noting the boundary and initial conditions, we have

$$u_i^0 = 0, \quad 1 \le i \le M - 1,$$

 $u_0^n = u_M^n = 0, \quad 0 \le n \le N.$
(15)

Canceling the infinitesimal part $(r_2)_i^n$ and replacing the exact solution u_i^k with the approximate solution U_i^k , a compact difference method is obtained as follows:

$$\mathcal{A}\Delta\alpha \sum_{l=0}^{2J} d_{l}w(\alpha_{l}) \frac{1}{\alpha_{l}\tau} \left[G_{0}^{\alpha_{l}} U_{i}^{n} - \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_{l}} - G_{n-k}^{\alpha_{l}} \right) U_{i}^{k} \right]$$

= $\delta_{x}^{2} U_{i}^{n} + \mathcal{A}f_{i}^{n}, \quad 1 \le i \le M - 1, 1 \le n \le N,$ (16)
 $U_{i}^{0} = 0, \quad 1 \le i \le M - 1,$
 $U_{0}^{n} = U_{M}^{n} = 0, \quad 0 \le n \le N.$

4.1 Stability analysis

Sorting equation (16), we get the following equation:

$$\mathcal{A}\Delta\alpha \sum_{l=0}^{2J} d_l w(\alpha_l) \frac{1}{\alpha_l \tau} \sum_{k=1}^n G_{n-k}^{\alpha_l} U_i^k - \delta_x^2 U_i^n$$
$$= \mathcal{A}\Delta\alpha \sum_{l=0}^{2J} d_l w(\alpha_l) \frac{1}{\alpha_l \tau} \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} U_i^k + \mathcal{A}f_i^n, \quad 1 \le i \le M-1, 1 \le n \le N.$$
(17)

We set $\eta_2 = \Delta \alpha \sum_{l=0}^{2J} d_l w(\alpha_l) \frac{1}{\alpha_l \tau}$. Multiplying both sides of equation (17) by $h U_i^n$ and making summation with respect to *i* from 1 to M - 1 leads to

$$\eta_2 \sum_{k=1}^{n} G_{n-k}^{\alpha_l} (U^k, U^n)_A + |U^n|_1^2 = \eta_2 \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} (U^k, U^n)_A + (f^n, U^n)_A.$$

Noting Lemma 1 and Young's inequality, we have

$$\frac{6}{L^2} \left\| U^n \right\|_2^2 + \eta_2 G_0^{\alpha_l} \left\| U^n \right\|_A^2 \le \eta_2 \sum_{k=0}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left(U^k, U^n \right)_A + \frac{L^2}{24} \left\| f^n \right\|_A^2 + \frac{6}{L^2} \left\| U^n \right\|_2^2.$$

Using the triangle inequality, we obtain

$$\eta_2 G_0^{\alpha_l} \| U^n \|_A^2 \le \frac{\eta_2}{2} \sum_{k=1}^{n-1} (G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l}) (\| U^k \|_A^2 + \| U^n \|_A^2) + \frac{L^2}{24} \| f^n \|_A^2.$$

It follows that

$$\eta_{2}G_{0}^{\alpha_{l}}\left\|U^{n}\right\|_{A}^{2} + \frac{\eta_{2}}{2}\sum_{k=1}^{n-1}\left(G_{n-k}^{\alpha_{l}} - G_{n-k-1}^{\alpha_{l}}\right)\left\|U^{n}\right\|_{A}^{2}$$
$$\leq \frac{\eta_{2}}{2}\sum_{k=1}^{n-1}\left(G_{n-k-1}^{\alpha_{l}} - G_{n-k}^{\alpha_{l}}\right)\left\|U^{k}\right\|_{A}^{2} + \frac{L^{2}}{24}\left\|f^{n}\right\|_{A}^{2},$$

namely,

$$\frac{\eta_2}{2} \left(G_0^{\alpha_l} + G_{n-1}^{\alpha_l} \right) \left\| U^n \right\|_A^2 \le \frac{\eta_2}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left\| U^k \right\|_A^2 + \frac{L^2}{24} \left\| f^n \right\|_A^2$$

Transposition leads to

$$\frac{\eta_2}{2} \sum_{k=1}^n G_{n-k}^{\alpha_l} \| U^k \|_A^2 \le \frac{\eta_2}{2} \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} \| U^k \|_A^2 + \frac{L^2}{24} \| f^n \|_A^2.$$

Denote

$$Q(U^n) = \frac{\eta_2}{2} \sum_{k=1}^n G_{n-k}^{\alpha_l} \| U^k \|_A^2.$$

Making summation from n = 1 to n = N shows that

$$Q(U^N) \le Q(U^0) + \frac{L^2}{24} \sum_{n=1}^N ||f^n||_A^2$$

Observing that $u^0 = 0$ implies $Q(u^0) = 0$, then we have

$$Q(U^N) \le \frac{L^2}{24} \sum_{n=1}^N ||f^n||_A^2.$$

Theorem 3 For scheme (16), the following stability inequality holds:

$$Q(U^m) \le \frac{L^2}{24} \sum_{n=1}^m \|f^n\|_A^2, \quad \forall 0 \le m \le N.$$

4.2 Convergence analysis

Combining equations (14), (15) with (16), an error equation can be obtained as follows:

$$\mathcal{A}\eta_2 \sum_{k=1}^n G_{n-k}^{\alpha_l} e_i^k - \delta_x^2 e_i^n = \mathcal{A}\eta_2 \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} e_i^k + (r_2)_i^n,$$
(18)

where $(r_2)_i^n = O(\tau^2 + h^4 + \Delta \alpha^4)$, $e_i^k = u_i^k - U_i^k$, $\forall 0 \le k \le N$.

Multiplying two sides of equation (18) by he_i^n and making summation with respect to *i* from 1 to M - 1, we get

$$\mathcal{A}\eta_{2}\sum_{k=1}^{n}G_{n-k}^{\alpha_{l}}(e^{k},e^{n})_{A}+\left|e^{n}\right|_{1}^{2}=\mathcal{A}\eta_{2}\sum_{k=1}^{n-1}G_{n-k-1}^{\alpha_{l}}(e^{k},e^{n})_{A}+((r_{2})^{n},e^{n})_{A}.$$

Noting Lemma 1 and Young's inequality, we have

$$\frac{6}{L^2} \left\| e^n \right\|_2^2 + \eta_2 G_0^{\alpha_l} \left\| e^n \right\|_A^2 \le \eta_2 \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left(e^k, e^n \right)_A + \frac{L^2}{24} \left\| (r_2)^n \right\|_A^2 + \frac{6}{L^2} \left\| e^n \right\|_2^2.$$

Using the triangle inequality, we obtain

$$\eta_2 G_0^{\alpha_l} \|e^n\|_A^2 \le \frac{\eta_2}{2} \sum_{k=1}^{n-1} (G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l}) (\|e^k\|_A^2 + \|e^n\|_A^2) + \frac{L^2}{24} \|(r_2)^n\|_A^2$$

Rearranging the above equation leads to

$$\begin{split} \eta_2 G_0^{\alpha_l} \left\| e^n \right\|_A^2 &+ \frac{\eta_2}{2} \sum_{k=1}^{n-1} \left(G_{n-k}^{\alpha_l} - G_{n-k-1}^{\alpha_l} \right) \left\| e^n \right\|_A^2 \\ &\leq \frac{\eta_2}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left\| e^k \right\|_A^2 + \frac{L^2}{24} \left\| (r_2)^n \right\|_A^2, \end{split}$$

namely

$$\frac{\eta_2}{2} \left(G_0^{\alpha_l} + G_{n-1}^{\alpha_l} \right) \left\| e^n \right\|_A^2 \le \frac{\eta_2}{2} \sum_{k=1}^{n-1} \left(G_{n-k-1}^{\alpha_l} - G_{n-k}^{\alpha_l} \right) \left\| e^k \right\|_A^2 + \frac{L^2}{24} \left\| (r_2)^n \right\|_A^2$$

Finally we get

$$\frac{\eta_2}{2} \sum_{k=1}^n G_{n-k}^{\alpha_l} \|e^k\|_A^2 \le \frac{\eta_2}{2} \sum_{k=1}^{n-1} G_{n-k-1}^{\alpha_l} \|e^k\|_A^2 + \frac{L^2}{24} \|(r_2)^n\|_A^2.$$

Denote

$$Q(e^n) = \frac{\eta_2}{2} \sum_{k=1}^n G_{n-k}^{\alpha_l} \|e^k\|_A^2.$$

Make summation with respect to n from 1 to N to get

$$Q(e^N) \le Q(e^0) + \frac{L^2}{24} \sum_{n=1}^N ||(r_2)^n||_A^2.$$

Observing that $e^0 = 0$ implies $Q(e^0) = 0$, it follows that

$$Q(e^N) \leq \frac{L^2}{24} \sum_{n=1}^N ||(r_2)^n||_A^2.$$

Theorem 4 For scheme (16), the following stability inequality holds:

$$\left\|e^{m}\right\|_{2} \leq C\left(\tau^{2} + h^{4} + \Delta\alpha^{4}\right), \quad \forall 0 \leq m \leq N.$$

5 Numerical tests

We perform numerical experiments to verify the proposed numerical format. Here, we consider the convergence speed of the numerical solution format. For the problem to be solved, we take the domain $\Omega = [0, \pi]$, T = 0.5. The numerical simulations are performed with MATLAB2015 on a 8GB memory computer.

Consider time distributed-order model as follows.

Example 1 For equation (1), if we give the true solution $u(x, t) = 4t^2 \sin(x)$ with $w(\alpha) = \alpha^2$, accordingly, f(x, t) can be derived as follows:

$$f(x,t) = 8\sin(x) \left[\int_{\nu_1}^{\nu_2} (\alpha t + \alpha - 1) \, d\alpha - e^t \int_{\nu_1}^{\nu_2} (\alpha - 1) e^{\frac{t}{\alpha - 1}} \, d\alpha \right] + 4t^2 \sin(x).$$

The error norms for difference schemes (8) and (16) are denoted as $\|\cdot\|_l^{\#}$ and $\|\cdot\|_l^{*}$, where l = 2 or ∞ , respectively.

Firstly, the numerical convergence orders of the two difference formats are tested by Example 1. Let the spatial step size *h* and the distributed-order subinterval size $\Delta \alpha$ be fixed small enough. We take $h = \pi/900$, $\Delta \alpha = 1/300$, and $\tau = 1/4$, 1/8, 1/16, 1/32 to observe the convergence rates about time step τ for both of the two difference formats. Table 1 shows the simulation results for three different distribution intervals [ν_1 , ν_2] = [0.1, 0.4], [0.5, 0.9], and [0.2, 0.8], including the errors and convergence orders in discrete L^2 norm. From these results we can see that the time convergence rates for both formats are second-order.

Next, we fix $v_1 = 0.1$, $v_2 = 0.9$, and take several different step sizes satisfying N = J and $M = \text{floor}(\pi J)$, i.e., $\tau = \frac{1}{2J}$, $h \approx \frac{1}{J}$, and $\Delta \alpha = \frac{(v_2 - v_1)}{2J}$. The simulation results of the numerical format (8) for Example 1 are shown in Table 2. From the table we can see that, when the spatial step size h, the discrete step of the distributed order $\Delta \alpha$, and the time step τ are halved simultaneously, the errors between the numerical solution and the true solution will be reduced to quarter in both of maximum norm and the discrete L^2 norm. These results show that the numerical convergence orders for difference format (8) are approximately 2 about time variable, space variable, and the distributed order.

V ₁ , V ₂	τ	$\ e^{N}\ _{2}^{\#}$	Order	$\ e^{N}\ _{2}^{*}$	Order
0.1, 0.4	1/2 ²	3.1795e-04		3.1810e-04	
	1/2 ³	7.9364e-05	2.0022	7.9543e-05	1.9997
	1/24	1.9701e-05	2.0102	1.9887e-05	1.9999
	1/25	4.7846e-06	2.0418	4.9718e-06	2.0000
0.5, 0.9	1/2 ²	0.0288		0.0288	
	1/2 ³	0.0074	1.9605	0.0074	1.9605
	1/24	0.0019	1.9615	0.0019	1.9615
	1/25	4.6265e-04	2.0380	4.6265e-04	2.0380
0.2, 0.8	1/2 ²	0.0124		0.0124	
	1/2 ³	0.0031	2.0000	0.0031	2.0000
	1/24	7.7660e-04	1.9970	7.8234e-04	1.9864
	1/2 ⁵	1.8990e-04	2.0319	1.9569e-04	1.9992

Table 1 The discrete L^2 error and convergence rates in time for schemes (8) and (16) for Example 1 with $h = \pi/900$ and $\Delta \alpha = 1/300$

Table 2 The discrete L^2 errors, maximum errors, and convergence rates for scheme (8) for Example 1 with $v_1 = 0.1$ and $v_2 = 0.9$

h	τ	$\Delta \alpha$	$\ e^N\ _{\infty}^{\#}$	Order	$\ e^{N}\ _{2}^{\#}$	Order
1/8	1/16	1/20	7.1433e-04		8.9705e-04	
1/16	1/32	1/40	1.7223e-04	2.0523	2.1586e-04	2.0551
1/32	1/64	1/80	4.2569e-05	2.0165	5.3352e-05	2.0165
1/64	1/128	1/160	1.0500e-05	2.0194	1.3160e-05	2.0194
1/128	1/256	1/320	2.6233e-06	2.0009	3.2879e-06	2.0009

Table 3	The discrete L ²	errors, maximum e	errors, and	convergence	rates for schem	ie (<mark>16</mark>) for
Example	1 with $v_1 = 0.1$	and $v_2 = 0.9$				

h	τ	$\Delta \alpha$	$\ e^N\ _{\infty}^*$	Order	$\ e^{N}\ _{2}^{*}$	Order
1/4	1/4	1/10	0.0060		0.0075	
1/8	1/16	1/20	3.7641e-04	3.9946	4.7269e-04	3.9879
1/16	1/64	1/40	2.3576e-05	3.9969	2.9548e-05	3.9998
1/32	1/256	1/80	1.4731e-06	4.0004	1.8462e-06	4.0004

Table 4 The discrete L^2 errors, maximum errors, and their convergence rates for scheme (8) for Example 2 when $v_1 = 0.1$ and $v_2 = 0.9$

h	τ	$\Delta \alpha$	$\ e^N\ _{\infty}^{\#}$	Order	$\ e^{N}\ _{2}^{\#}$	Order
$\pi/4$	1/8	1/10	3.2542e-2		4.0785e-2	
$\pi/8$	1/16	1/20	8.0780e-3	2.0102	1.0124e-2	2.0102
$\pi/16$	1/32	1/40	2.0157e-3	2.0027	2.5263e-3	2.0027
$\pi/32$	1/64	1/80	5.0368e-4	2.0007	6.3126e-4	2.0007
$\pi/64$	1/128	1/160	1.2590e-4	2.0002	1.5780e-4	2.0002
$\pi/128$	1/256	1/320	3.1475e-5	2.0000	3.9448e-5	2.0000
$\pi/256$	1/512	1/640	7.8687e-6	2.0000	9.8619e-6	2.0000
$\pi/512$	1/1024	1/1280	1.9672e-6	2.0000	2.4655e-6	2.0000

Table 5 The discrete L^2 errors, maximum errors, and convergence rates for scheme (16) for Example 2 with $v_1 = 0.1$ and $v_2 = 0.9$

h	τ	$\Delta \alpha$	$\ e^N\ _{\infty}^*$	Order	$\ e^{N}\ _{2}^{*}$	Order
$\pi/4$	1/4	1/10	2.3994e-2		3.0072e-2	
$\pi/8$	1/16	1/20	1.5429e-3	3.9590	1.9337e-3	3.9590
$\pi/16$	1/64	1/40	9.6510e-5	3.9988	1.2096e-4	3.9988
$\pi/32$	1/256	1/80	6.0314e-6	4.0001	7.5593e-6	4.0001
$\pi/64$	1/1024	1/160	3.7695e-7	4.0000	4.7244e-7	4.0000
$\pi/128$	1/4096	1/320	2.3559e-8	4.0000	2.9527e-8	4.0000

In addition, we choose a group of mesh steps satisfying $8N = J^2$ and $M = \text{floor}(\pi J)$, which means that τ varies synchronously with $\Delta \alpha^2$ and h^2 . The simulation results of numerical format (16) for Example 1 are given in Table 3. We observe that when the spatial step h and the distributed-order step size $\Delta \alpha$ are halved, the corresponding numerical errors in discrete norm are reduced to 1/16 times. This means both of the convergence rates for spatial variable and distributed order are of fourth order and thus the time convergence rate is of second order.

Example 2 For equation (1), if we give the true solution $u(x, t) = 2t \sin(x)$ with $w(\alpha) = \alpha^2$, accordingly, f(x, t) can be derived as follows:

$$f(x,t) = 2\sin(x) \left[\int_{\nu_1}^{\nu_2} \alpha \, d\alpha - e^t \int_{\nu_1}^{\nu_2} \alpha e^{\frac{t}{\alpha-1}} \, d\alpha \right] + 2t \sin(x).$$

For Example 2, we set the distributed-order interval $[\nu_1, \nu_2] = [0.1, 0.9]$ and the numbers of discrete points are required to satisfy M = N = J for scheme (8). The simulation results in Table 4 show that scheme (8) has convergence rates of second order for all of the spatial step *h*, the temporal step τ , and the distributed-order step $\Delta \alpha$.

In addition, we choose a group of mesh steps satisfying $8N = J^2$ and M = J, which means that τ varies synchronously with $\Delta \alpha^2$ and h^2 . The simulation results of numerical format (16) for Example 2 are given in Table 5. We observe that when the spatial step h and the

distributed-order step size $\Delta \alpha$ are halved, the corresponding numerical errors in discrete norm are reduced to 1/16 times. This means both of the convergence rates for spatial variable and distributed order are fourth order, and accordingly the convergence rate for time variable is second order.

6 Conclusion

In this paper, two effective finite difference schemes are developed for solving time distributed-order partial differential equations equipped with the new Caputo–Fabrizio fractional derivative in a one-dimensional space. The first scheme (8) is based on second-order central difference quotient in space direction and the other one (16) is based on compact difference. We have proved that both of the two finite difference schemes are unconditionally stable. Theoretical analysis shows that the first scheme has convergence rates of second order for all of the spatial step h, the temporal step τ , and the distributed-order step $\Delta \alpha$; and the second scheme (compact difference) has convergence rates of fourth order for both of the spatial step h and the distributed-order step $\Delta \alpha$, while second order for the temporal step τ . Finally, some numerical experiments, which demonstrate the performance of the schemes and verify the correctness of the theoretical results, are carried out.

The basic model is in a one-dimensional space, an extension to multi-dimensional models is possible and feasible. With the development of the theory about fractional calculus, fractional partial differential equations with variable order become the potentially powerful tool for describing complex phenomena [63]. We hope to extend the idea of this paper to this kind of fractional equations in the future. Moreover, because of the nonlocal property of fractional derivatives, the requirement for memory and computational complexity will increase rapidly along with the size of discrete systems. So we plan to develop some techniques of fast solution in the future work.

Acknowledgements

The authors are grateful to the reviewers. Their comments helped to improve the quality of the paper significantly.

Funding

This work is funded by the National Natural Science Foundation of China under Grants 91630207 and 11971272, which will be used for the payment of possible publication charge.

Availability of data and materials

All data generated or analysed during this study are included in this published article.

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

HQ contributed to the numerical schemes, error analysis, and numerical experiments. ZL contributed to the mathematical model and numerical schemes. AC contributed to the analysis of numerical results. All authors read and approved the final manuscript.

Author details

¹School of Mathematics, Shandong University, Jinan, China. ²School of Mathematics and Statistics, Shandong Normal University, Jinan, China.

Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Received: 23 October 2019 Accepted: 13 January 2020 Published online: 17 January 2020

References

- 1. Podlubny, I.: Fractional Differential Equations (1999)
- 2. Atangana, A.: Non validity of index law in fractional calculus: a fractional differential operator with Markovian and non-Markovian properties. Phys. A, Stat. Mech. Appl. **505**, 688–706 (2018)
- Atangana, A., Gómez-Aguilar, J.F.: Fractional derivatives with no-index law property: application to chaos and statistics. Chaos Solitons Fractals 114, 516–535 (2018)
- Atangana, A., Gómez-Aguilar, J.F.: Decolonisation of fractional calculus rules: breaking commutativity and associativity to capture more natural phenomena. Eur. Phys. J. Plus 133(4), 166 (2018)
- Caputo, M., Fabrizio, M.: A new definition of fractional derivative without singular kernel. Prog. Fract. Differ. Appl. 1(2), 73–85 (2015)
- Atangana, A.: On the new fractional derivative and application to nonlinear Fisher's reaction–diffusion equation. Appl. Math. Comput. 273, 948–956 (2016)
- 7. Atangana, A., Araz, S.I.: New numerical method for ordinary differential equations: Newton polynomial. J. Comput. Appl. Math. (2019). https://doi.org/10.1016/j.cam.2019.112622
- Qureshi, S., Rangaig, N.A., Baleanu, D.: New numerical aspects of Caputo–Fabrizio fractional derivative operator. Mathematics 7(4), 374 (2019)

 Atangana, A., Nieto, J.J.: Numerical solution for the model of RLC circuit via the fractional derivative without singular kernel. Adv. Mech. Eng. 7(10), 1–7 (2015). https://doi.org/10.1177/1687814015613758

- Qureshi, S., Yusuf, A.: Fractional derivatives applied to MSEIR problems: comparative study with real world data. Eur. Phys. J. Plus 134(4), 171. https://doi.org/10.1140/epjp/i2019-12661-7 (2019)
- 11. Qureshi, S., Yusuf, A., Shaikh, A.A., Inc, M., Baleanu, D.: Fractional modeling of blood ethanol concentration system with real data application. Chaos, Interdiscip. J. Nonlinear Sci. **29**(1), 013143 (2019). https://doi.org/10.1063/1.5082907
- 12. Qureshi, S., Atangana, A.: Mathematical analysis of dengue fever outbreak by novel fractional operators with field data. Phys. A, Stat. Mech. Appl. 526, 121127 (2019). https://doi.org/10.1016/j.physa.2019.121127
- 13. Qureshi, S., Yusuf, A.: Modeling chickenpox disease with fractional derivatives: from Caputo to Atangana–Baleanu. Chaos Solitons Fractals **122**, 111–118 (2019)
- Qureshi, S., Kumar, P.: Using Shehu integral transform to solve fractional order Caputo type initial value problems. J. Appl. Math. Comput. Mech. 18(2), 75–83 (2019)
- Atangana, A., Qureshi, S.: Modeling attractors of chaotic dynamical systems with fractal–fractional operators. Chaos Solitons Fractals 123, 320–337 (2019)
- Cheng, A., Wang, H., Wang, K.: A Eulerian–Lagrangian control volume method for solute transport with anomalous diffusion. Numer. Methods Partial Differ. Equ. 31(1), 253–267 (2015)
- Iwayama, T., Murakami, S., Watanabe, T.: Anomalous eddy viscosity for two-dimensional turbulence. Phys. Fluids 27(4), 045104 (2015)
- Metzler, R., Jeon, J.H., Cherstvy, A.G., Barkai, E.: Anomalous diffusion models and their properties: non-stationarity, non-ergodicity, and ageing at the centenary of single particle tracking. PCCP, Phys. Chem. Chem. Phys. 16(44), 24128–24164 (2014)
- 19. Shraiman, B.I., Siggia, E.D.: Scalar turbulence. Nature 405(6787), 639–646 (2000)
- Armour, K.C., Marshall, J., Scott, J.R., Donohoe, A., Newsom, E.R.: Southern Ocean warming delayed by circumpolar upwelling and equatorward transport. Nat. Geosci. 9(7), 549–554 (2016)
- Naghibolhosseini, M., Long, G.R.: Fractional-order modeling and simulation of human ear. Int. J. Comput. Math. 95, 1257–1273 (2018)
- Perdikaris, P., Karniadakis, G.Em.: Fractional-order viscoelasticity in one-dimensional blood flow models. Ann. Biomed. Eng. 42(5), 1012–1023 (2014)
- Tyukhova, A., Dentz, M., Kinzelbach, W., Willmann, M.: Mechanisms of anomalous dispersion in flow through heterogeneous porous media. Phys. Rev. Fluids 1(7), 074002 (2016)
- Ardakani, A.G.: Investigation of Brewster anomalies in one-dimensional disordered media having Lévy-type distribution. Eur. Phys. J. B 89(3), 1–6 (2016)
- Zhang, Y., Meerschaert, M.M., Neupauer, R.M.: Backward fractional advection dispersion model for contaminant source prediction. Water Resour. Res. 52(4), 2462–2473 (2016)
- Edery, Y., Dror, I., Scher, H., Berkowitz, B.: Anomalous reactive transport in porous media: experiments and modeling. Phys. Rev. E, Stat. Nonlinear Soft Matter Phys. 91(5), 052130 (2015)
- Zhang, Y., Meerschaert, M.M., Baeumer, B., Labolle, E.M.: Modeling mixed retention and early arrivals in multidimensional heterogeneous media using an explicit Lagrangian scheme. Water Resour. Res. 51(8), 6311–6337 (2015)
- Suzuki, J.L., Zayernouri, M., Bittencourt, M.L., Karniadakis, G.E.: Fractional-order uniaxial visco-elasto-plastic models for structural analysis. Comput. Methods Appl. Mech. Eng. 308, 443–467 (2016)
- 29. Goychuk, I.: Anomalous transport of subdiffusing cargos by single kinesin motors: the role of mechano-chemical coupling and anharmonicity of tether. Phys. Biol. **12**(1), 016013 (2015)
- Mashelkar, R.A., Marrucci, G.: Anomalous transport phenomena in rapid external flows of viscoelastic fluids. Rheol. Acta 19(4), 426–431 (1980)
- 31. Meerschaert, M.M., Sikorskii, A.: Stochastic Models for Fractional Calculus, vol. 43 (2012)
- 32. Klages, R., Radons, G., Sokolov, I.M.: Anomalous Transport: Foundations and Applications. Wiley-VCH, Weinheim (2008)
- Metzler, R., Klafter, J.: The random walk's guide to anomalous diffusion: a fractional dynamics approach. Phys. Rep. 339(1), 1–77 (2000)
- Gorenflo, R., Luchko, Y., Yamamoto, M.: Time-fractional diffusion equation in the fractional Sobolev spaces. Fract. Calc. Appl. Anal. 18(3), 799–820 (2015)
- 35. Sokolov, I.M., Chechkin, A.V., Klafter, J.: Distributed-order fractional kinetics. Acta Phys. Pol. B 35(4), 1323–1341 (2004)
- Konjik, S., Oparnica, L., Zorica, D.: Distributed order fractional constitutive stress-strain relation in wave propagation modeling. arXiv preprint arXiv:1709.01339 (2017)
- 37. Zhang, Y., Sun, Z., Wu, H.: Error estimates of Crank–Nicolson-type difference schemes for the subdiffusion equation. SIAM J. Numer. Anal. 49(6), 2302–2322 (2011)

- Gao, G., Sun, Z.: A compact finite difference scheme for the fractional sub-diffusion equations. J. Comput. Phys. 230(3), 586–595 (2011)
- Chen, C., Liu, F., Turner, I., Anh, V.: Numerical approximation for a variable-order nonlinear reaction-subdiffusion equation. Numer. Algorithms 63(2), 265–290 (2013)
- Fu, H., Ng, M.K., Wang, H.: A divide-and-conquer fast finite difference method for spacetime fractional partial differential equation. Comput. Math. Appl. 73(6), 1233–1242 (2017)
- Fu, H., Wang, H.: A preconditioned fast finite difference method for space-time fractional partial differential equations. Fract. Calc. Appl. Anal. 20(1), 88–116 (2017)
- 42. Jin, B., Lazarov, R., Thomée, V., Zhou, Z.: On nonnegativity preservation in finite element methods for subdiffusion equations. Math. Comput. 86(2), 37–45 (2017)
- Ainsworth, M., Glusa, C.: Aspects of an adaptive finite element method for the fractional Laplacian: a priori and a
 posteriori error estimates, efficient implementation and multigrid solver. Comput. Methods Appl. Mech. Eng. 327,
 4–35 (2017)
- Coronel-Escamilla, A., Gómez-Aguilar, J.F., Torres, L., Escobar-Jiménez, R.F.: A numerical solution for a variable-order reaction-diffusion model by using fractional derivatives with non-local and non-singular kernel. Phys. A, Stat. Mech. Appl. 491, 406–424 (2018)
- Macías-Díaz, J.E.: An explicit dissipation-preserving method for Riesz space-fractional nonlinear wave equations in multiple dimensions. Commun. Nonlinear Sci. Numer. Simul. 59, 67–87 (2018)
- Ammi, M.R.S., Jamiai, I.: Finite difference and Legendre spectral method for a time-fractional diffusion-convection equation for image restoration. Discrete Contin. Dyn. Syst., Ser. S 11(1), 103–117 (2018)
- Liu, Y., Du, Y., Li, H., Li, J., He, S.: A two-grid mixed finite element method for a nonlinear fourth-order reaction-diffusion problem with time-fractional derivative. Comput. Math. Appl. 70(10), 2474–2492 (2015)
- Yamamoto, M.: Weak solutions to non-homogeneous boundary value problems for time-fractional diffusion equations. J. Math. Anal. Appl. 460(1), 365–381 (2018)
- 49. Mao, Z., Shen, J.: Efficient spectral–Galerkin methods for fractional partial differential equations with variable coefficients. J. Comput. Phys. **307**, 243–261 (2016)
- Al-Khaled, K., Momani, S.: An approximate solution for a fractional diffusion-wave equation using the decomposition method. Appl. Math. Comput. 165(2), 473–483 (2005)
- 51. Fu, H., Wang, H., Wang, Z.: POD/DEIM reduced-order modeling of time-fractional partial differential equations with applications in parameter identification. J. Sci. Comput. **74**(1), 220–243 (2018)
- 52. Chen, W., Ye, L., Sun, H.: Fractional diffusion equations by the Kansa method. Comput. Math. Appl. 59(5), 1614–1620 (2010)
- 53. Zaky, M.A.: A Legendre collocation method for distributed-order fractional optimal control problems. Nonlinear Dyn. 91(4), 2667–2681 (2018)
- Li, X., Rui, H.: Two temporal second-order H¹-Galerkin mixed finite element schemes for distributed-order fractional sub-diffusion equations. Numer. Algorithms 79, 1107–1130 (2018)
- Fan, W., Liu, F.: A numerical method for solving the two-dimensional distributed order space-fractional diffusion equation on an irregular convex domain. Appl. Math. Lett. 77, 114–121 (2018)
- Luchko, Y.: Boundary value problems for the generalized time-fractional diffusion equation of distributed order. Fract. Calc. Appl. Anal. 12(4), 409–422 (2009)
- Tomovski, Ž., Sandev, T.: Distributed-order wave equations with composite time fractional derivative. Int. J. Comput. Math. 95(6–7), 1100–1113 (2018)
- Atanackovic, T.M., Pilipovic, S., Zorica, D.: Existence and calculation of the solution to the time distributed order diffusion equation. Phys. Scr. **T136**, 014012 (2009)
- 59. Sun, Z.: Numerical Methods for Partial Differential Equations, 2nd ed. (2012) (in Chinese)
- 60. Gautschi, W.: Numerical Analysis, 2nd ed. (2012)
- 61. Liu, Z., Cheng, A., Li, X.: A fast-high order compact difference method for the fractional cable equation. Numer. Methods Partial Differ. Equ. **34**(6), 2237–2266 (2018)
- 62. Liao, H., Sun, Z.: Maximum norm error bounds of ADI and compact ADI methods for solving parabolic equations. Numer. Methods Partial Differ. Equ. 26(1), 37–60 (2010)
- Wang, H., Zheng, X.: Wellposedness and regularity of the variable-order time-fractional difusion equations. J. Math. Anal. Appl. 475, 1778–1802 (2019)

Submit your manuscript to a SpringerOpen[®] journal and benefit from:

- ► Convenient online submission
- ► Rigorous peer review
- Open access: articles freely available online
- ► High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at > springeropen.com