PHYSICAL ACOUSTICS

Low-Frequency Shear Elasticity of a Colloid Nanosuspension

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Abstract—The paper presents the experimental results of an acoustic impedance study of low-frequency $(10⁵ Hz)$ shear elasticity of a colloidal $SiO₂$ suspension of differently sized particles in polyethylene siloxane liquid PES-2. The agreement of the experimental results obtained by different variants of the acoustic resonance method confirms that the low-frequency shear elasticity of colloidal nanoparticle suspensions is a bulk property.

Keywords: nanosuspension, piezoelectric, oscillations, modulus, impedance method, nanoparticles **DOI:** 10.1134/S1063771020050103

Studies [1, 2] describe the acoustic impedance method for measuring the low-frequency (10^5 Hz) shear elasticity of liquids. At one end of the horizontal surface of a rectangular prism of piezoelectric quartz, a layer of the studied liquid is applied, which is covered with a hard cover-plate. During tangential oscillations of the piezoelectric quartz, the liquid layer undergoes shear deformations and standing shear waves are occur. As the thickness of the liquid interlayer changes, so do the resonance frequency and resonance curve width of the piezoelectric quartz. Shifts of the resonant frequency of the piezoelectric crystal were obtained from acoustic resonance method theory [1–3]:

$$
\Delta f' = \frac{S}{4\pi^2 M f_0}
$$

$$
\times \frac{(G'\beta + G''\alpha)\sin 2\beta H + (G'\alpha - G''\beta)\sin 2\alpha H}{\cosh 2\alpha H - \cos 2\beta H},
$$
 (1)

ch
$$
2\alpha H - \cos 2\beta H
$$

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$$
\Delta f'' = \frac{S}{4\pi^2 M f_0}
$$
\n
$$
\times \frac{(G''\beta - G'\alpha)\sin 2\beta H + (G''\alpha + G'\beta)\sin 2\alpha H}{\cosh 2\alpha H - \cos 2\beta H},
$$
\n(2)

where *G*' and *G*" are the real and imaginary shear moduli, *H* is the thickness of the liquid layer, *M* is the mass of the piezoelectric quartz, *S* is the contact area between the liquid and piezoelectric quartz, and β and α are the real and imaginary components of the complex wavenumber. Figure 1 shows the theoretical dependences of the real Δ*f*' and imaginary Δ*f*" frequency shifts on the thickness of the liquid layer for a liquid with $G = 3 \times 10^4$ Pa and tan $\theta = 0.3$ calculated by these formulas [1, 2].

Clearly, with an increase in the thickness of the liquid interlayer, frequency shifts yield damped oscillations. When the shear wave has completely attenuated, the fre-

quency shifts take the limiting values $\Delta f_{\infty}^{'}$ and $\Delta f_{\infty}^{''}$.

From analysis of expressions (1) and (2), three methods follow for determining the low-frequency shear elasticity of liquids [1–8]. The first method is realized for small liquid layer thicknesses, when $H \ll \lambda$. In this case, the frequency shifts exhibit a linear dependence on the inverse thickness of the liquid layer 1/*H*. The second method is based on determination of *G*' from the length of the shear wave, which is determined from the attenuation maxima. The third method, similar to Mason's well-known impedance method [9], measures the limiting values of the frequency shifts to which they tend with increasing liquid layer thickness. Since for $H \rightarrow \infty$ the shear wave is completely attenuated, the need for the cover-plate vanishes and the entire horizontal surface of the piezoelectric quartz can be loaded with a thick layer of the studied fluid.

In this case, from expression (1) for *G*' it is possible to obtain the following calculation formula $[1-3]$:

$$
G' = \frac{16\pi^2 M^2}{S^2 \rho} \left[\left(\Delta f_{\infty}^{\prime \prime} \right)^2 - \left(\Delta f_{\infty}^{\prime} \right)^2 \right],\tag{3}
$$

where *S* is the area of the entire horizontal surface of the piezoelectric quartz. From expression (3), as well as from Fig. 1, it is clear that in the presence of shear elasticity in liquids, Δ*f*" should be greater than Δ*f*′. All three methods for determining low-frequency shear elasticity were used with conventional and polymer liquids as an example, which gave quite consistent results $[1-8]$.

Fig. 1. Theoretical dependences of real (*1*) and imaginary (*2*) frequency shifts on liquid layer thickness.

Fig. 2. Profiles of resulting liquid layer: *1*, piezoelectric quartz; *2*, liquid.

Paper [10] studied the low-frequency shear elasticity of colloidal suspensions of $SiO₂$ nanoparticles in polyethylsiloxane liquid PES-2 by the acoustic resonance method for thicknesses *H* much shorter than the wavelength λ . In these experiments, the layer thickness of the studied suspensions varied within a few microns. Therefore, the presence of low-frequency shear elasticity in the studied nanosuspensions can be attributed to the special properties of the boundary layers under the action of the field of surface forces.

This paper presents the experimental results of a study of shear elasticity of suspensions of $SiO₂$ nanoparticles in PES-2 by the impedance method for

 $H \ge \lambda$. In the experiments, a piezoelectric quartz of X-18.5° cut was used, with a resonant frequency of 73.2 kHz, a mass of 6.82 g, and dimensions of $35 \times 12 \times 6$ mm. Colloidal $SiO₂$ nanosuspensions in PES-2 were obtained by prolonged dispersion using ultrasonic methods [10].

On the thoroughly cleaned horizontal surface of the piezoelectric quartz, a thick layer of the suspension was applied, in which the shear wave is completely attenuated (Fig. 2). The limiting imaginary $\Delta f_{\infty}^{\prime\prime}$ and real frequency shifts $\Delta f_{\infty}^{\prime}$ were then measured. The limiting value of the real frequency shift $\Delta f_{\infty}^{\prime}$ can be neglected, since its contribution is negligible, not exceeding 3% [1–3].

For the suspension studied, 0.5% by mass fraction of $SiO₂$ nanoparticles in polyethylsiloxane liquid PES-2 with dimensions of 100 nm, the imaginary shift limit value $\Delta f_{\infty}^{\prime\prime}$ amounted to 8 Hz. Calculation according to formula (3) for the real shear modulus *G*' gives a value of 1.06×10^5 Pa. The results obtained for other suspensions with different sizes and concentrations $c =$ 0.5 wt % are given in Table 1. The density of these suspensions is 0.94 g/cm³. " Δ*f*[∞]

Comparison of the results obtained with small liquid layer thicknesses $H \ll \lambda$ [10] and the results obtained by the impedance method for $H \ge \lambda$ demonstrates their good agreement. This confirms that the low-frequency (10⁵ Hz) shear elasticity of colloidal nanosuspensions is a bulk property. The low-frequency viscoelastic behavior of low-molecular-weight liquids is also discussed by other researchers [11–14].

Nanosuspensions are used in various nanotechnologies, in particular, the development of drugs, intensification of heat transfer, in of new functional materials, lubricants, paints, etc. The use of various rheological liquid media in many technological processes is responsible for the great interest in studying their mechanical properties [15–19].

$SiO2/PES-2$	$G \times 10^{-5}$ Pa, for $H \ll \lambda$ [10]	$tan\theta$	$G' \times 10^{-5}$ Pa, for $H \ge \lambda$
20 nm	0.09	0.73	0.08
50 nm	0.17	0.18	0.15
100 nm	1.08	0.10	1.06

Table 1. Values of shear moduli of suspension of $SiO₂/PES-2$ with different nanoparticle sizes.

REFERENCES

- 1. B. B. Badmaev, T. S. Dembelova, and B. B. Damdinov, *Viscose and Elastic Properties of Polymeric Liquids* (Buryat Scientific Centre of Siberian Department of RAS, Ulan-Ude, 2013) [in Russian].
- 2. B. B. Badmaev, T. S. Dembelova, B. B. Damdinov, and Ch. Zh. Gulgenov, Acoust. Phys. **63** (6), 642 (2017).
- 3. B. Badmaev, T. Dembelova, B. Damdinov, D. Makarova, and O. Budaev, Colloids Surf., A **383**, 90 (2011).
- 4. U. B. Bazaron, B. V. Deryagin, and O. R. Budaev, Dokl. Akad. Nauk SSSR **205** (6), 1324 (1972).
- 5. B. B. Badmaev, B. B. Damdinov, and D. S. Sanditov, Acoust. Phys. **50** (2), 121 (2004).
- 6. B. B. Badmaev, S. A. Bal'zhinov, B. B. Damdinov, and T. S. Dembelova, Acoust. Phys. **56** (5), 640 (2010).
- 7. W. P. Mason, *Piezoelectric Crystals and Their Application to Ultrasonics* (D. Van Nostrand, 1950; IL, Moscow, 1952).
- 8. B. B. Badmaev, O. R. Budaev, and T. S. Dembelova, Acoust. Phys. **45** (5), 541 (1999).
- 9. B. B. Badmaev, T. S. Dembelova, B. B. Damdinov, D. N. Makarova, and E. D. Namdakova, Uch. Zap. Fiz. Fak. Mosk. Univ., No. 5, 1751302 (2017).
- 10. T. S. Dembelova, A. B. Tsyrenzhapova, D. N. Makarova, B. B. Damdinov, and B. B. Badmaev, Uch. Zap. Fiz. Fak. Mosk. Univ., No. 5, 145301 (2014).
- 11. D. Collin and P. Martinoty, Phys. A (Amsterdam, Neth.) **320**, 235 (2003).
- 12. H. P. Kavehpoor and G. H. McKinley, Tribol. Lett. **17** (2), 327 (2004).
- 13. L. Noirez and P. Baroni, J. Mol. Struct. **972**, 16 (2010).
- 14. L. Noirez, J. Phys.: Condens. Matter **24**, 372101 (2012).
- 15. V. Ya. Rudyak, A. V. Minakov, and M. I. Pryazhnikov, Tr. Novosib. Gos. Arkhit.-Stroit. Univ. (Sibstrin) **21** $(1(67))$, 30 (2018).
- 16. J. Chevalier, O. Tillement, and F. Ayela, Appl. Phys. Lett. **91** (23), 233103 (2007).
- 17. N. B. Ur'ev, S. V. Emel'yanov, and K. A. Titov, Prot. Met. Phys. Chem. Surf. **51** (2), 226 (2015).
- 18. I. B. Esipov, O. M. Zozulya, and M. A. Mironov, Acoust. Phys. **60** (2), 169 (2014).
- 19. D. N. Makarova and I. B. Esipov, Vestn. Buryat. Gos. Univ. Khim. Fiz., Nos. **2–3**, 45 (2018).