## **QUASIDIFFERENTIAL OPERATOR AND QUANTUM ARGUMENT SHIFT METHOD**

#### **Y. Ikeda**∗

*We describe an explicit formula for the first-order quasiderivation of an arbitrary central element of the universal enveloping algebra of a general linear Lie algebra. We apply it to show that derivations of any two central elements of the universal enveloping algebra commute. This contributes to the Vinberg problem of finding commutative subalgebras in universal enveloping algebras with the underlying Poisson algebras determined by the argument shift method.*

**Keywords:** universal enveloping algebra, Lie algebra, quantum argument shift method, deformation quantization

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#### **1. Introduction**

The argument shift method introduced by Mischenko and Fomenko [\[1\]](#page-6-0) is a well-known method devised to construct commutative families of elements in Poisson algebras. Loosely speaking, it consists in shifting the central elements of this algebra (central in the sense of a Poisson bracket) in the direction of a vector field  $\xi$  that satisfies an algebraic relation involving the Poisson tensor

$$
\mathcal{L}_{\xi}^2 \pi = 0,\tag{1}
$$

where  $\mathcal{L}_{\xi}\pi$  is the Lie derivative of the Poisson tensor  $\pi$  with respect to  $\xi$ . This relation is often called the "Nijenhuis equation," or "Nijenhuis property" and can be regarded as a sort of linearity condition. If this relation holds, then the shifted (in the direction of  $\xi$ ) central elements, although no more central, continue to commute with each other.

The most common nontrivial example of this Nijenhuis property and the one that was originally discovered by Manakov, Mischenko, and Fomenko [\[1\]](#page-6-0), [\[2\]](#page-6-1) is the case where the Poisson structure is given by the usual Kirillov–Kostant brackets on the dual space of a Lie algebra and the field is any constant vector field on this space: the Kirillov–Kostant brackets have linear coefficients and hence the Nijenhuis equation holds automatically. This method is quite productive and gives rise to maximal commutative Poisson subalgebras in polynomials on the dual space.

The purpose of this note is to suggest a construction that generalizes the argument shift method to the domain of universal enveloping algebras. Namely, Vinberg asked whether the commutative subalgebras

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generated by the shift can be lifted to the universal enveloping algebra of the Lie algebra. In about 30 years several constructions have appeared that determine such a lift, see, e.g.,  $[3]-[5]$  $[3]-[5]$  $[3]-[5]$ . But all these constructions are based on the considerations of generating elements in the center of the Poisson algebras and do not address the question of whether the shift itself can be transfered into the noncommutative situation. In this note, we address this question and suggest a method that hypothetically leads to its solution, at least for the Lie algebra  $g = gl_d$ .

Our approach is based on the considerations of the so-called "quasiderivations" of the algebra  $Ugl<sub>d</sub>$ , introduced in [\[6\]](#page-6-4). Namely, we prove an explicit formula for the first quasiderivation of such an element (Theorem [1\)](#page-2-0) and thus show that the first quasiderivations in any direction of *any* central element in the universal enveloping algebra commute (Corollary [1\)](#page-2-1). These observations support the conjecture that *quasiderivations induce a shift in the universal enveloping algebras analogous to the shift in Poisson algebras*.

#### **2. Preliminaries and statement of the results**

We let  $\mathbb F$  denote the real number field  $\mathbb R$  or the complex number field  $\mathbb C$  and let d be a nonnegative integer. The  $d \times d$  matrix units are denoted by

$$
E_i^j = \begin{pmatrix} 0 \\ \vdots \\ 0 \\ 1 \\ 0 \\ \vdots \\ 0 \end{pmatrix} \quad (0 \dots 010 \dots 0)
$$
 (2)

with 1 in the *i*th place of the column vector and in the *j*th place of the row vector. We then have the commutation relations  $[E_p^q, E_r^s] = \delta_r^q E_p^s - E_r^q \delta_p^s$ . The universal enveloping algebra of the general linear Lie algebra  $gl_d = gl_d(\mathbb{F})$  is denoted by  $U = Ugl_d$  and its center is denoted by  $Z = Z(Ugl_d)$ . We let  $\pi$  denote the canonical Lie-algebra homomorphism of the general linear Lie algebra  $gl<sub>d</sub>$  into its universal enveloping algebra  $Ugl_d$  and set  $e_i^j = \pi(E_i^j)$  for each i and j. The generators  $e_i^j$  satisfy the same commutation relations  $[e_p^q, e_r^s] = \delta_r^q e_p^s - e_r^q \delta_p^s$  as the matrix units  $E_i^j$ . We define

$$
e = \sum_{i,j=1}^{d} e_j^i \otimes E_i^j \in U \otimes M(d, \mathbb{F}) = \begin{pmatrix} e_1^1 & e_2^1 & \dots & e_d^1 \\ e_1^2 & e_2^2 & \dots & e_d^2 \\ \vdots & \vdots & \ddots & \vdots \\ e_1^d & e_2^d & \dots & e_d^d \end{pmatrix} \in M(d, U).
$$
 (3)

We remark that the  $(i, j)$  entry of the matrix e is the generator  $e_i^i = \pi(E_i^i)$ , whereas the  $(j, i)$  entry of the corresponding matrix unit  $E_i^i$  is equal to 1. The role of the lower and upper indices is opposite here. In other words, we let e denote the *transposed* matrix of generators of the algebra  $Ugl_d$ ; this choice of notation simplifies our computations below. We have

$$
(e^2)^i_j = \sum_{k=1}^d e^i_k e^k_j, \qquad [e^q_p, (e^n)^s_r] = [(e^n)^q_p, e^s_r] = \delta^q_r (e^n)^s_p - (e^n)^q_r \delta^s_p. \tag{4}
$$

<span id="page-1-0"></span>It can be shown that the center of  $Ugl_d$  is generated by the set  $\left\{\text{tr} \, e^n\right\}_{n=1}^d$ .

**Definition 1** (see [\[6\]](#page-6-4)). Quasiderivations of  $Ugl_d$  is a unique family of linear operators  $\partial_i^i$  on U satisfying the equations

- $\partial_i^i 1 = 0$  for any *i* and *j*;
- $\partial_i^i e_p^q = \delta_p^i \delta_i^q$  for any *i*, *j*, *p*, and *q*;
- the twisted Leibniz rule

$$
\partial_j^i(fg) = (\partial_j^i f)g + f(\partial_j^i g) + \sum_{k=1}^d (\partial_k^i f)(\partial_j^k g) \tag{5}
$$

for any  $i$  and  $j$  and for any elements  $f$  and  $g$  of  $U$ .

<span id="page-2-0"></span>We define a map  $\partial$  of the universal enveloping algebra U into the matrix algebra  $M(d, U)$  by  $(\partial f)^i_j = \partial^i_j f$ for any  $f\in U.$  The following theorem is the main result in this paper.

**Theorem 1.** *We have*

$$
\partial \operatorname{tr} e^n = \frac{(e+\delta)^n - (e-\delta)^n}{2} + \sum_{k=0}^{n-1} (\operatorname{tr} e^k) \frac{(e+\delta)^{n-k-1} - (e-\delta)^{n-k-1}}{2} \tag{6}
$$

<span id="page-2-1"></span>*for any n*, where  $\delta$  *is the*  $d \times d$  *identity matrix.* 

**Corollary 1.** We suppose that  $\xi$  is an arbitrary element of the Lie algebra  $gl_d$  and let

$$
\partial_{\xi} = \sum_{i,j=1}^{d} \xi_i^j \partial_j^i. \tag{7}
$$

*Then we have*  $[\partial_{\xi} f, \partial_{\xi} g] = 0$  *for any central elements* f *and* g.

**Proof.** It is sufficient to show that  $[tr(\xi e^m), tr(\xi e^n)] = 0$  because we have

$$
\partial \prod_{k=1}^{m} \text{tr} \, e^{n_k} = (\text{tr} \, e^{n_1} + \partial \, \text{tr} \, e^{n_1}) \dots (\text{tr} \, e^{n_m} + \partial \, \text{tr} \, e^{n_m}) - \prod_{k=1}^{m} \text{tr} \, e^{n_k}.
$$
 (8)

From

$$
[\text{tr}(\xi e^m), e_j^i] = \sum_{p,q=1}^d \xi_p^q [(e^m)_q^p, e_j^i] = \sum_{p,q=1}^d \xi_p^q (\delta_j^p (e^m)_q^i - (e^m)_j^p \delta_q^i) = [e^m, \xi]_j^i,
$$

we deduce

$$
[\text{tr}(\xi e^m), \text{tr}(\xi e^n)] = \sum_{k=1}^n \text{tr}(\xi e^{k-1} [e^m, \xi] e^{n-k}) = \sum_{k=1}^n \text{tr}[\xi e^{m+n-k}, \xi e^{k-1}].\tag{9}
$$

We can assume that  $m < n$ . Then

$$
[\text{tr}(\xi e^m), \text{tr}(\xi e^n)] = \frac{[\text{tr}(\xi e^m), \text{tr}(\xi e^n)] - [\text{tr}(\xi e^m), \text{tr}(\xi e^m)]}{2} =
$$
  

$$
= \frac{1}{2} \sum_{k=m+1}^n \text{tr}[\xi e^{m+n-k}, \xi e^{k-1}] = 0.
$$
 (10)

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## **3. The proof of the main theorem**

We begin with the following observation: due to the twisted Leibniz rule in Definition [1,](#page-1-0) we have

$$
\partial_j^i (e^n)_q^p = \sum_{r=1}^d \partial_j^i (e_r^p (e^{n-1})_q^r) =
$$
\n
$$
= \sum_{r=1}^d (\partial_j^i e_r^p)(e^{n-1})_q^r + \sum_{r=1}^d e_r^p \partial_j^i (e^{n-1})_q^r + \sum_{k,r=1}^d (\partial_k^i e_r^p)(\partial_j^k (e^{n-1})_q^r) =
$$
\n
$$
= (e^{n-1})_q^i \delta_j^p + \sum_{r=1}^d e_r^p \partial_j^i (e^{n-1})_q^r + \partial_j^p (e^{n-1})_q^i.
$$
\n(11)

for  $n > 0$ .

**Lemma 1.** *The quasiderivation*  $\partial_i^i(e^n)_q^p$  *has the form* 

$$
\partial_j^i (e^n)_q^p = \sum_{k+\ell < n} \left( a_{k,\ell}^{(n)} (e^k)_q^i (e^\ell)_j^p + b_{k,\ell}^{(n)} (e^k)_q^p (e^\ell)_j^i \right),\tag{12}
$$

where  $a_{k,\ell}^{(n)}$  and  $b_{k,\ell}^{(n)}$  are nonnegative integers.

**Proof.** The proof is by induction on  $n$ . We have

$$
\partial_j^i (e^n)_q^p = (e^{n-1})_q^i \delta_j^p + \sum_{r=1}^d e_r^p \partial_j^i (e^{n-1})_q^r + \partial_j^p (e^{n-1})_q^i =
$$
\n
$$
= (e^{n-1})_q^i \delta_j^p + \sum_{k,\ell} a_{k,\ell}^{(n-1)} \sum_{r=1}^d e_r^p (e^k)_q^i (e^\ell)_j^r + \sum_{k,\ell} b_{k,\ell}^{(n-1)} (e^{k+1})_q^p (e^\ell)_j^i + \sum_{k,\ell} (a_{k,\ell}^{(n-1)} (e^k)_q^p (e^\ell)_j^i + b_{k,\ell}^{(n-1)} (e^k)_q^i (e^\ell)_j^p).
$$
\n(13)

by the induction hypothesis. Because

$$
\sum_{r=1}^d e_r^p (e^k)_q^i (e^\ell)_j^r = (e^k)_q^i (e^{\ell+1})_j^p + \sum_{r=1}^d [e_r^p, (e^k)_q^i] (e^\ell)_j^r = (e^k)_q^i (e^{\ell+1})_j^p + \delta_q^p (e^{k+\ell})_j^i - (e^k)_q^p (e^\ell)_j^i,
$$

we have

$$
\partial_j^i (e^n)_q^p = (e^{n-1})_q^i \delta_j^p + \sum_{k,\ell} a_{k,\ell}^{(n-1)} \left( (e^k)_q^i (e^{\ell+1})_j^p + \delta_q^p (e^{k+\ell})_j^i \right) + \sum_{k,\ell} b_{k,\ell}^{(n-1)} \left( (e^{k+1})_q^p (e^\ell)_j^i + (e^k)_q^i (e^\ell)_j^p \right). \tag{14}
$$

This equation gives rise to a recursive process

$$
\cdots \to (a_{k,\ell}^{(n-1)}, b_{k,\ell}^{(n-1)})_{k,\ell} \to (a_{k,\ell}^{(n)}, b_{k,\ell}^{(n)})_{k,\ell} \to \cdots
$$
\n(15)

We have the polynomial identities

$$
\sum_{k+\ell < n} a_{k,\ell}^{(n)} x^k y^\ell = x^{n-1} + \sum_{k+\ell < n-1} (a_{k,\ell}^{(n-1)} x^k y^{\ell+1} + b_{k,\ell}^{(n-1)} x^k y^\ell),\tag{16}
$$

$$
\sum_{k+\ell < n} b_{k,\ell}^{(n)} x^k y^\ell = \sum_{k+\ell < n-1} \left( a_{k,\ell}^{(n-1)} y^{k+\ell} + b_{k,\ell}^{(n-1)} x^{k+1} y^\ell \right). \tag{17}
$$

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Using

$$
\partial_j^i \operatorname{tr} e^n = \sum_{k+\ell < n} \left( a_{k,\ell}^{(n)} (e^{k+\ell})_j^i + (\operatorname{tr} e^k) b_{k,\ell}^{(n)} (e^\ell)_j^i \right),
$$

we can define polynomials  $f^{(n)}(x)$  and  $g_k^{(n)}(x)$  with integer coefficients by

$$
f^{(n)}(x) = \sum_{k,\ell} a_{k,\ell}^{(n)} x^{k+\ell}, \qquad g_k^{(n)}(x) = \sum_{\ell} b_{k,\ell}^{(n)} x^{\ell}.
$$
 (18)

Then we have

$$
f^{(n)}(x) = x^{n-1} + \sum_{k,\ell} a_{k,\ell}^{(n-1)} x^{k+\ell+1} + \sum_{k,\ell} b_{k,\ell}^{(n-1)} x^{k+\ell} =
$$
  
=  $x^{n-1} + xf^{(n-1)}(x) + \sum_{k < n-1} x^k g_k^{(n-1)}(x),$  (19)

by identity [\(16\)](#page-3-0), and also

$$
g_0^{(n)}(x) = \sum_{k,\ell} a_{k,\ell}^{(n-1)} x^{k+\ell} = f^{(n-1)}(x),
$$
  

$$
g_{k+1}^{(n)}(x) = \sum_{\ell} b_{k,\ell}^{(n-1)} x^{\ell} = g_k^{(n-1)}(x) = \dots = g_0^{(n-k-1)}(x) = f^{(n-k-2)}(x),
$$

by identity [\(17\)](#page-3-1). In this way, we obtain the formulas

$$
\partial_j^i \operatorname{tr} e^n = \left( f^{(n)}(e) + \sum_{k < n} (\operatorname{tr} e^k) g_k^{(n)}(e) \right)_j^i = \left( f^{(n)}(e) + \sum_{k < n} (\operatorname{tr} e^k) f^{(n-k-1)}(e) \right)_j^i \tag{20}
$$

and

<span id="page-4-0"></span>
$$
f^{(n)}(x) = x^{n-1} + xf^{(n-1)}(x) + \sum_{k < n-1} x^k f^{(n-k-2)}(x). \tag{21}
$$

The solution of this recursive equation is

$$
f^{(n)}(x) = \frac{(x+1)^n - (x-1)^n}{2} \tag{22}
$$

and we have

$$
\partial_j^i \operatorname{tr} e^n = \left( \frac{(e+\delta)^n - (e-\delta)^n}{2} + \sum_{k=0}^{n-1} (\operatorname{tr} e^k) \frac{(e+\delta)^{n-k-1} - (e-\delta)^{n-k-1}}{2} \right)_j^i.
$$

# $\Delta$ ppendix: Exact formula for  $\partial_j^i(e^n)_q^p$

The methods that we used above can be generalized to give formulas for quasiderivations of more general elements.

We proceed to obtain  $a_{k,\ell}^{(n)}$  and  $b_{k,\ell}^{(n)}$  and derive an exact formula for  $\partial_i^i (e^n)_q^p$ . We set

$$
f_{\pm}^{(n)}(x) = \frac{(x+1)^n \pm (x-1)^n}{2} = \sum_{k=0}^n \frac{1 \pm (-1)^{n-k}}{2} {n \choose k} x^k
$$
 (A.1)

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and

$$
h_k^{(n)}(x) = \sum_{\ell=0}^{n-k-1} a_{k,\ell}^{(n)} x^{\ell}.
$$
 (A.2)

Then we have

$$
\partial_j^i (e^n)_q^p = \sum_{k=0}^{n-1} ((e^k)_q^i h_k^{(n)}(e)_j^p + (e^k)_q^p g_k^{(n)}(e)_j^i) =
$$
  
= 
$$
\sum_{k=0}^{n-1} ((e^k)_q^i h_k^{(n)}(e)_j^p + (e^k)_q^p f_-^{(n-k-1)}(e)_j^i),
$$
 (A.3)

whence

$$
\sum_{k=0}^{n-1} y^k h_k^{(n)}(x) = y^{n-1} + \sum_{k=0}^{n-2} y^k (x h_k^{(n-1)}(x) + f_-^{(n-k-2)}(x))
$$
 (A.4)

by identity [\(16\)](#page-3-0). We have  $h_{n-1}^{(n)}(x) = 1$  and

$$
h_k^{(n)}(x) = xh_k^{(n-1)}(x) + f_-^{(n-k-2)}(x) = \dots =
$$
  

$$
= x^{n-k-1}h_k^{(k+1)}(x) + \sum_{\ell=0}^{n-k-2} x^{\ell} f_-^{(n-k-\ell-2)}(x) =
$$
  

$$
= x^{n-k-1} + \sum_{\ell=0}^{n-k-2} x^{\ell} f_-^{(n-k-\ell-2)}(x) =
$$
  

$$
= f_-^{(n-k)}(x) - x f_-^{(n-k-1)}(x) = f_+^{(n-k-1)}(x).
$$
 (A.5)

for  $k < n$  by identity [\(21\)](#page-4-0). Furthermore,

$$
a_{k,\ell}^{(n)} = \frac{1 - (-1)^{n-k-\ell}}{2} \binom{n-k-1}{\ell}, \qquad b_{k,\ell}^{(n)} = \frac{1 + (-1)^{n-k-\ell}}{2} \binom{n-k-1}{\ell} \tag{A.6}
$$

and

$$
\partial_j^i (e^n)_q^p = \sum_{k=0}^{n-1} \left( (e^k)_q^i f_+^{(n-k-1)}(e)_j^p + (e^k)_q^p f_-^{(n-k-1)}(e)_j^i \right) =
$$
\n
$$
= \sum_{k=0}^{n-1} \left( (e^k)_q^i \left( \frac{(e+\delta)^{n-k-1} + (e-\delta)^{n-k-1}}{2} \right)_j^p + (e^k)_q^p \left( \frac{(e+\delta)^{n-k-1} - (e-\delta)^{n-k-1}}{2} \right)_j^i \right) =
$$
\n
$$
= \sum_{k+\ell=0}^{n-1} {n-k-1 \choose \ell} \left( \frac{1 - (-1)^{n-k-\ell}}{2} (e^k)_q^i (e^\ell)_j^p \frac{1 + (-1)^{n-k-\ell}}{2} (e^k)_q^p (e^\ell)_j^i \right).
$$

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### <span id="page-6-0"></span>**REFERENCES**

- <span id="page-6-1"></span>1. A. S. Mishchenko and A. T. Fomenko, "Euler equations on finite-dimensional Lie groups," *Izv. Math.*, **12**, 371–389 (1978).
- <span id="page-6-2"></span>2. S. V. Manakov, "Note on the integration of Euler's equations of the dynamics of an *n*-dimensional rigid body," *Funct. Anal. Appl.*, **10**, 328–329 (1976).
- 3. A. A. Tarasov, "On some commutative subalgebras of the universal enveloping algebra of the Lie algebra gl(*n,* <sup>C</sup>)," *Sb. Math.*, **<sup>191</sup>**, 1375–1382 (2000).
- <span id="page-6-3"></span>4. L. G. Rybnikov, "The argument shift method and the Gaudin model," *Funct. Anal. Appl.*, **40**, 188–199 (2006).
- <span id="page-6-4"></span>5. A. I. Molev, "Feigin–Frenkel center in types *B*, *C* and *D*," *Invent. Math.*, **191**, 1–34 (2013).
- 6. D. Gurevich, P. Pyatov, and P. Saponov, "Braided Weyl algebras and differential calculus on *U*(*u*(2))," *J. Geom. Phys.*, **62**, 1175–1188 (2012), [arXiv: 1112.6258.](http://arxiv.org/abs/1112.6258)