## **METHODS OF THEORETICAL PHYSICS**

# **Mirror Pairs of Quintic Orbifolds**

**A. Belavin***a***,** *<sup>c</sup>***,** *<sup>d</sup>***, \* and B. Eremin***a***,** *<sup>b</sup>***,** *<sup>c</sup>***,** *<sup>d</sup>*

*a Landau Institute for Theoretical Physics, Russian Academy of Sciences, Chernogolovka, Moscow region, 142432 Russia b Skolkovo Institute of Science and Technology (Skoltech), Skolkovo, Moscow region, 143025 Russia c Moscow Institute of Physics and Technology (National Research University), Dolgoprudnyi, Moscow region, 141700 Russia d Institute for Information Transmission Problems (Kharkevich Institute), Russian Academy of Sciences, Moscow, 127994 Russia \*e-mail: sashabelavin@gmail.com* Received September 3, 2020; revised September 3, 2020; accepted September 3, 2020

Two constructions of mirror pairs of Calabi–Yau manifolds are compared by example of quintic orbifolds 2. The first, Berglund–Hubsch–Krawitz, construction is as follows. If *X* is the factor of the hypersurface 2 by a certain subgroup  $H'$  of the maximum allowed group  $SL$ , the mirror manifold Y is defined as the factor by

the dual subgroup  $H^{T}$ . In the second, Batyrev, construction, the toric manifold *T* containing the mirror *Y* as a hypersurface specified by zeros of the polynomial  $W_Y$  is determined from the properties of the polynomial  $W_X$  specifying the Calabi–Yau manifold X. The polynomial  $W_Y$  is determined in an explicit form. The group of symmetry of the polynomial  $W_Y$  is found from its form and it is tested whether it coincides with that predicted by the Berglund–Hubsch–Krawitz construction.

**DOI:** 10.1134/S002136402018006X

Calabi–Yau manifolds appear when six of ten dimensions are compactified in the superstring theory [1], which is a way to unify the Standard Model and Quantum Gravity. The effective theory is specified in terms of the so-called special geometry arising in the moduli space [2] of Calabi–Yau manifolds and their mirrors. We study mirror pairs by example of three cases of quintic orbifolds. We consider Calabi–Yau manifolds defined as hypersurfaces in the weighted projective space

$$
\text{jective space} \qquad \qquad \mathbb{P}^4_{(k_1, k_2, k_3, k_4, k_5)} = \{ (x_1, \dots, x_5) \in \mathbb{C}^5 \setminus \{0\} | x_i \sim \lambda^{k_i} x_i \}. \qquad (1)
$$

Let *M* be the matrix of degrees of the quasihomogeneous polynomial  $W_M(x) := \sum_{i=1}^5 \prod_{j=1}^5 x_j^{M_{ij}}$ , i.e.,

$$
W_M(\lambda^{k_i} x) = \lambda^d W_M(x), \qquad (2)
$$

which specifies a hypersurface by the equation  $W_M(x) = 0$ . The matrix *M* is invertible and has a chain or loop Fermat type [3], which, together with the equality

$$
d = \sum_{i}^{5} k_i,
$$
 (3)

is the condition that the hypersurface is a Calabi–Yau manifold.

The whole family of Calabi–Yau manifolds  $X_M$  is obtained by the deformation of the initial polynomial  $W_M$  by allowed monomials:

$$
W(x, \varphi) = W_M(x) + \sum_{l=1}^{h} \varphi_l \prod_{i=1}^{5} x_i^{S_{ll}}.
$$
 (4)

Here,  $\varphi_l$  are moduli of the complex structure of the manifold  $X_M$ .

We note that it follows from Eqs. (1) and (2) that, if

$$
\lambda = \omega, \quad \omega^d = 1,\tag{5}
$$

the polynomial  $W_M(x)$  is invariant under the action of the group of transformations  $Q_M$  , which acts as

$$
x_i \mapsto \omega^{k_i} x_i. \tag{6}
$$

The group  $Q_M$  is called the quantum symmetry group. The polynomial  $W_M$  can have the maximal symmetry group Aut(*M*),  $Q_M \subset \text{Aut}(M)$ , and its order is [3]

$$
|\text{Aut}(M)| = \det M. \tag{7}
$$

The group  $Aut(M)$  is generated by the generators  $q(M)_{i}$ ,  $i = 1, ..., 5$ , which act at each coordinate  $x_{j}$  as

$$
q_i(M): x_j \mapsto e^{2\pi i B_{ji}} x_j, \tag{8}
$$

where the matrix *B* is inverse to the matrix *M*.

A generator of the quantum group  $Q_M$  is the product  $\prod_{i=1}^{5} q_i(M)$ , which acts on  $x_i$  as in Eq. (6). The group  $Aut(M)$  has subgroups that conserve the holomorphic 3-form Ω nonvanishing on  $X_M$  or, equivalently, conserve  $\prod_{i=1}^{5} x_i$ . Such subgroups are called allowed. Let  $SL(M)$  be the maximal allowed subgroup.

There is the chain of embeddings  $Q_M \subseteq$  $SL(M) \subseteq Aut(M)$ .

Now, we can define an orbifold in  $X_M$  as follows. We choose in  $SL(M)$  a certain subgroup  $H(M)$  that includes (or equals)  $Q_M$ . Further, the factor-group by  $Q_M$  is obtained from it as

$$
H'(M) := H(M)/Q_M. \tag{9}
$$

Then, the orbifold is defined as

$$
Z(M, H) := X_M / H^*(M). \tag{10}
$$

A similar construction can be made for the matrix  $M^T$  and four groups are obtained:  $Q_{M^T} \subseteq H(M^T) \subseteq$  $SL(M^T) \subseteq Aut(M^T)$ , where  $H(M^T)$  is a certain subgroup  $SL(M^T),$  and the orbifold  $Z(M^T, H^T).$ 

Now, there is the question: Can these groups be chosen, and, if so, can they be chosen such that  $Z(M, H)$  and  $Z(M^T, H^T)$  are mirrors of each other?

The Berglund–Hubsch–Krawitz construction [4, 5] is as follows. Let the group  $H(M)$  leave invariant the monomials

$$
e_{l}(x) = \prod_{i=1}^{5} x_{i}^{S_{li}} \tag{11}
$$

for a certain choice of  $S_{li}$ . These monomials, together with  $W_M(x)$ , specify the orbifold  $Z(M, H)$ . Then, by definition, the generators of the group  $H(M^T)$  are

$$
q(l) := \prod_{i=1}^{5} q_i(M^T)^{S_{li}}, \qquad (12)
$$

where  $S_{li}$  are the same as in Eq. (11) and the generators of the group  $SL(M^T)$  are specified by their action on  $x_j$ 

$$
q_i(H^T): x_j \mapsto e^{2\pi i B_{ij}} x_j. \tag{13}
$$

JETP LETTERS Vol. 112 No. 6 2020

According to Eq. (13), the generators of the group  $H(M<sup>T</sup>)$  act on the coordinates as

$$
q(l): x_j \mapsto e^{2\pi i B_{ij} S_{il}} x_j. \tag{14}
$$

We now define monomials as  $e_m^T := \prod_{i=1}^5 z_i^{R_{mi}}$  from the requirement of their invariance under the action of the group  $H(M<sup>T</sup>)$ .  $e_m^T := \prod_{i=1}^5 z_i^{R_{mi}}$ 

It is convenient to reproduce this requirement in the form

$$
\langle l,m\rangle := B_{ij} S_{li} R_{mj} \in \mathbb{Z}.
$$
 (15)

Below, we consider quintic orbifolds for which  $B_{ij} = \delta_{ij}/5$ . In this case, (15) is transformed to the condition

$$
\sum_{i} S_{ii} R_{mi} = 0 \pmod{5}.
$$
 (16)

The Berglund–Hubsch–Krawitz hypothesis implies that the polynomial  $W_{M}$ , together with the superposition of  $e_m^T$ , specifies the mirror orbifold  $Z(M^T, H^T)$ .

The Chiodo–Ruan theorem [6] states the mirror symmetry of the orbifolds  $Z(M, H)$  and  $Z(M^T, H^T)$ in the cohomology level.

In the simplest case, we can set  $H(M) = Q_M$ . Then, the orbifold is the initial Calabi–Yau manifold  $X_M$ , i.e.,  $Z(M, Q_M) = X_M$ . Then, the mirror Calabi–Yau manifold denoted as  $Y_M$  is the orbifold by the group  $SL^{r}(M^{T}) := SL(M^{T})/Q_{M^{T}}$ :

$$
Y_M \equiv Z(M^T, S L^T) = X_{M^T} / S L^{\prime} (M^T). \tag{17}
$$

According to the second, Batyrev, construction of mirror pairs of Calabi–Yau manifolds [7, 8], the mirror *Y* of the initial orbifold *X* is specified as a hypersurface in the toric manifold. First, we describe the construction of the toric manifold in *X*. The polynomial specifying the family  $X_M$  can be represented in the form

$$
W_X(x, \varphi) = \sum_{i=1}^{5} \prod_{j=1}^{5} x_j^{M_{ij}} + \sum_{l=1}^{h_X} \varphi_l \prod_{j=1}^{5} x_j^{S_{jl}}
$$
  
= 
$$
\sum_{a=1}^{5+h_X} C_a(\varphi) \prod_{j=1}^{5} x_j^{V_{aj}}.
$$
 (18)

The degrees of the monomials of the initial polynomial (18) specifying the orbifold *X* are vectors in the lattice  $\mathbb{Z}^5$ , namely,  $(\mathbf{V}_a)_j = V_{aj}$ . They can also be represented in the matrix form  $V^T = (M^T | S^T)$ . In view of quasihomogeneity,  $\sum_{j=1}^{5} V_{aj} k_j = d$ . Consequently, the vectors  $(\mathbf V_a)_j$  are located in four-dimensional lattice  $\mathbb Z^4$ 

and constitute a Batyrev polygon [8]. At the same time, the same vectors  $V_a$  are the edges of the fan [7], which determines the toric manifold *T*. Since  $h_X + 5$ vectors  $V_a$  lie in the four-dimensional space, they should satisfy  $h_X$  linear relations

$$
\sum_{a=1}^{h_X+5} Q_{la} \mathbf{V}_a = 0, \quad l = 1, ..., h_X.
$$
 (19)

These relations determine the weights of the action of the Abelian group  $(\mathbb{C}^*)^{h_X}$  in  $\mathbb{C}^{h_X+4}$ 

$$
y_a \to \lambda^{Q_{la}} y_a
$$
,  $a = 1,..., h_X + 4$ ,  $l = 1,..., h_X$ , (20)  
which specifies the toric manifold

$$
T := \frac{\left(\mathbb{C}^{h_X + 4} - Z\right)}{\left(\mathbb{C}^*\right)^{h_X}},\tag{21}
$$

where the manifold  $Z$  is invariant under this action [7].

The next step is the construction of the hypersurface in *T*.

The mirror Calabi–Yau manifold *Y* is specified by the zeros of a certain quasihomogeneous polynomial  $[9-11]$   $W_Y(y_1, \ldots, y_{h_X+4})$ , i.e.,

$$
Y := \{(y_1, \dots, y_{h_X + 4}) \in T | W_Y(y) = 0\}.
$$
 (22)

We find the explicit form of the polynomial  $W_Y$ from the homogeneity conditions:

$$
W_Y(\lambda^{Q_{i_1}}y_1,\ldots,\lambda^{Q_{i,h_X+4}}y_{h_X+4})
$$
  
=  $\lambda^d W_Y(y_1,\ldots,y_{h_X+4}), \quad i=1,\ldots,h_X.$  (23)

We represent the polynomial  $W_Y$  in the form of some monomials  $L_i(y)$ , which will be found below:  $W_Y(y_1, ..., y_{h_X+4}),$   $i = 1, ..., h_X.$ <br>
nt the polynomial  $W_Y$  in the form<br>  $L_i(y)$ , which will be found below:<br>  $= \sum_{i=1}^{h_Y+5} \tilde{C}_i(\psi) L_i(y),$   $L_i(y) = \prod_{i=1}^{h_X+4} y_a^{n_a^i}$ .

$$
W_Y(y) = \sum_{i=1}^{h_Y+5} \tilde{C}_i(\psi) L_i(y), \quad L_i(y) = \prod_{a=1}^{h_X+4} y_a^{n_a^i}.
$$
 (24)

From Eqs. (23) and (24), we obtain the systems of equations for the exponents  $n_a$  (for brevity, we omit the index *i*):

$$
\sum_{a=1}^{h_X+4} Q_{la} n_a = \delta_{l,h_X} d, \quad l = 1,...,h_X.
$$
 (25)

We note that six solutions of the system (25) are always known. Since

$$
\sum_{i=1}^{5} V_{ai} k_i = d, \quad \sum_{a=1}^{h_X + 5} Q_{la} V_{ai} = 0, \quad Q_{h_X, h_X + 5} = -d, \quad (26)
$$

these  $5 + 1$  solutions of the system (25) have the form

$$
n_a^i = V_{ai}, \quad i = 1, ..., 5,
$$
 (27)

$$
n_a^6 = 1,\t\t(28)
$$

where  $a = 1, ..., h_X + 4$ .

The polynomial  $W_Y$  consists of only  $h_Y + 5$  invariant monomials whose degrees are found from the solutions (25) with the imposition of additional conditions of nonnegativity of degrees  $n_a^i \ge 0$ . Here, is the Hodge number of the Calabi–Yau manifold *Y*. Thus, it is necessary to solve the system  $h_Y := h_{21}^Y$ 

$$
\sum_{a=1}^{h_X+4} Q_{la} n_a = \delta_{l,h_X} d, \quad n_a \ge 0,
$$
  
\n
$$
l = 1,...,h_X, \quad a = 1,...,h_X + 4.
$$
 (29)

Using the symmetry of the action of the torus, one can make a change to the following (projective) coordinates invariant under the action of  $(\mathbb{C}^*)^{h_X}$ :

$$
z_j = \prod_{a=1}^{h_X + 4} y_a^{V_{aj}/5}.
$$
 (30)

Thus, we reduce the toric manifold to the projective space  $\mathbb{P}^4$ . In coordinates  $z_j$ , the polynomial  $W_\gamma$ has the form

$$
W_Y(z) = \sum_{i=1}^5 \prod_{j=1}^5 x_j^{M_{ji}} + \sum_{m=1}^{h_Y} \Psi_m e_m^T(z). \tag{31}
$$

It is seen that the mirror *Y* is determined by the transposed matrix  $M^T$ .

From the form of deformations  $e_m^T(z)$  of the complex moduli space of the manifold *Y*, we find the group of symmetry of the polynomial  $W_Y$  and can verify whether it coincides with the group  $H<sup>T</sup>$  obtained in the Krawitz procedure.

Below, we apply these two constructions to quintic orbifolds. The whole family in this case is specified by the following equation in the projective space:

$$
\mathcal{Q} = \left\{ (x_1, ..., x_5) \right\}
$$

$$
\in \mathbb{P}^4 \Big| W_{\mathcal{Q}}(x, \varphi) := W_0(x) + \sum_{s=1}^{101} \varphi_i e_i(x) = 0 \right\},
$$
(32)

where

$$
W_0(x) := x_1^5 + x_2^5 + x_3^5 + x_4^5 + x_5^5,
$$
  
\n
$$
e_l = \prod_{i=1}^5 x_i^{S_{li}}, \quad 0 \le S_{li} \le 3,
$$
\n(33)

and  $\varphi$ <sub>*l*</sub> are the moduli of the complex structure of the  $manifold 2.$ 

In the case of quintic, the matrix of degrees is diagonal and is  $M = 5I$ , where *I* is the identity matrix, and  $W_M = W_{M^T} \equiv W_0$ . The initial polynomial  $W_0(x)$  has the quantum symmetry group  $Q_{5I} = \mathbb{Z}_5[1, 1, 1, 1, 1]$ . The full symmetry group of the monomial is  $Aut(5I) = \mathbb{Z}_5^5$ 

JETP LETTERS Vol. 112 No. 6 2020

and its generators act as  $q_i(5I): x_i \mapsto e^{2\pi i/d}x_i$ . The group conserving the product  $\prod_{i=1}^5 x_i$  has the form

$$
SL(5I) = \mathbb{Z}_5^4 = \mathbb{Z}_5[1, 0, 0, 0, 4] \oplus \mathbb{Z}_5[0, 1, 0, 0, 4]
$$
  

$$
\oplus \mathbb{Z}_5[0, 0, 1, 0, 4] \oplus \mathbb{Z}_5[0, 0, 0, 1, 4].
$$
 (34)

The group  $SL(5I)$  is sometimes called the maximal allowed group of the polynomial  $W_0$ . It contains the diagonal  $Q_{5I} = \mathbb{Z}_5[1, 1, 1, 1, 1]$ , and the factor group is

$$
SL'(5I) := SL(5I)/Q_{5I} = \mathbb{Z}_5[0, 1, 0, 0, 4]
$$
  
\n
$$
\oplus \mathbb{Z}_5[0, 0, 1, 0, 4] \oplus \mathbb{Z}_5[0, 0, 0, 1, 4].
$$
\n(35)

We deal with quintic orbifolds  $X = \mathcal{Q}/H'$ , which are specified by the following equation in  $\mathbb{P}^4$ :

$$
W_X(x, \varphi)
$$
  
 :=  $x_1^5 + x_2^5 + x_3^5 + x_4^5 + x_5^5 + \sum_{l=1}^{h_X} \varphi_l e_l(x) = 0.$  (36)

Here,  $h_X := h_{21}^X$  is the Hodge number. The monomials  $e_l$  are invariant under the action of the group *H*. According to the procedure described above, we construct the mirror Calabi–Yau orbifold

$$
Y := \mathcal{Q}/H^{T},\tag{37}
$$

or, in terms of the initial polynomial,

$$
W_Y(z, \varphi)
$$
  
 :=  $z_1^5 + z_2^5 + z_3^5 + z_4^5 + z_5^5 + \sum_{m=1}^{h_Y} \Psi_m e_m^T(z) = 0.$  (38)

The monomials  $e_m^T(z)$  are in turn invariant under the action of the group  $H^T$ .

We now perform calculations for three orbifolds.

**Example 1.** As the first example, we consider the orbifold [12]

$$
X = \frac{2}{\mathbb{Z}_5[0, 1, 1, 4, 4]}, \quad h_X = 17, \tag{39}
$$

which is specified by the zeros of the polynomial

$$
W_X(x) = x_1^5 + x_2^5 + x_3^5 + x_4^5 + x_5^5 + \sum_{l=1}^{17} \varphi_l e_l(x)
$$
  
= 
$$
\sum_{a=1}^{22} C_a \prod_{j=1}^5 x_j^{V_{aj}}.
$$
 (40)

JETP LETTERS Vol. 112 No. 6 2020

Here,  $e_i(x) := \prod_{i=1}^5 x_i^{S_{ii}}$  are invariant under the action of  $H = \mathbb{Z}_5[1, 1, 1, 1, 1] \oplus \mathbb{Z}_5[0, 1, 1, 4, 4]$  and have the form

$$
e_1 = x_2^3 x_3^2, \t e_2 = x_2^2 x_3^3, \t e_3 = x_4^2 x_5^3,
$$
  
\n
$$
e_4 = x_4^3 x_5^2, \t e_5 = x_1^3 x_2 x_4, \t e_6 = x_1^3 x_3 x_4,
$$
  
\n
$$
e_7 = x_1^3 x_3 x_5, \t e_8 = x_1^3 x_2 x_5, \t e_9 = x_1 x_2^2 x_4^2,
$$
  
\n
$$
e_{10} = x_1 x_3^2 x_5^2, \t e_{11} = x_1 x_3^2 x_4^2, \t e_{12} = x_1 x_2^2 x_5^2,
$$
  
\n
$$
e_{13} = x_1 x_2 x_3 x_4^2, \t e_{14} = x_1 x_2 x_3 x_5^2,
$$
  
\n
$$
e_{15} = x_1 x_2^2 x_4 x_5, \t e_{16} = x_1 x_3^2 x_4 x_5,
$$
  
\n
$$
e_{17} = x_1 x_2 x_3 x_4 x_5.
$$

According to the Berglund–Hubsch–Krawitz construction, the mirror Calabi–Yau manifold is constructed as the orbifold  $Y = \mathcal{Q}/H^{T}$ . The group  $H^{T}$  is specified by the generators  $q(l) = \prod_{i=1}^{5} q_i (5I)^{S_{li}}$ , where  $S_{li}$  are the degrees of monomials  $e_l(x)$  in Eq. (41), and  $q_i(5I)$  generate the groups  $\mathbb{Z}_5[1,0,0,0,0]$ ...  $\mathbb{Z}_5[0, 0, 0, 0, 1]$ . In other words, the group  $H^{T}$  is generated by additive groups  $\mathbb{Z}_5[S_{11}, S_{12}, S_{13}, S_{14}, S_{15}]$  in the modulus of the action  $\mathbb{Z}_5[1,1,1,1,1]$ . In this case, the group  $H^{\prime T}$  has two generators and it has the form

$$
H^{T} = \mathbb{Z}_{5}[0, 1, 4, 0, 0] \oplus \mathbb{Z}_{5}[0, 0, 0, 1, 4].
$$
 (42)

The explicit form of monomials  $e_m(z) = \prod_{i=1}^5 z_i^{R_m}$ in Eqs. (44) can be alternatively obtained by solving Eq. (16), i.e.,

$$
\sum_{i} S_{ii} R_{mi} = 0 \pmod{5}.
$$
 (43)

The solutions of Eqs. (43) give the degrees of monomials and have the form

$$
e_1^T(z) = z_1 z_2^2 z_3^2, \t e_2^T(z) = z_1 z_4^2 z_5^2,
$$
  
\n
$$
e_3^T(z) = z_1^3 z_4 z_5, \t e_4^T(z) = z_1^3 z_2 z_3,
$$
  
\n
$$
e_5^T(z) = z_1 z_2 z_3 z_4 z_5.
$$
\n(44)

Then, the family of the mirror orbifold

$$
Y = \frac{2}{\mathbb{Z}_5[0, 1, 4, 0, 0] \oplus \mathbb{Z}_5[0, 0, 0, 1, 4]}, \quad h_Y = 5 \quad (45)
$$

is specified as the solutions of the equation

$$
W_Y(z) = z_1^5 + z_2^5 + z_3^5 + z_4^5 + z_5^5 + \sum_{m=1}^5 \Psi_m e_m^T(z) = 0.
$$
 (46)

We now apply the Batyrev approach and represent the polynomial  $W_X$  in the form

$$
W_X(x, \varphi) = \sum_{a=1}^{22} C_a(\varphi) \prod_{j=1}^5 x_j^{V_{aj}}.
$$
 (47)

The weights of the action of the torus are determined from the vectors  $V_a$ . We write them in the matrix form

$$
(V)_{aj} := V = \begin{pmatrix} 5I_{5\times 5} \\ S \end{pmatrix},
$$
  
\n
$$
(Q)_{la} := Q = (S \mid -5I_{17\times 17}),
$$
\n(48)

where  $S \in Mat(\mathbb{Z})_{17\times 5}$  is determined from Eq. (40). Using these weights  $Q_{la}$ , we obtain the toric manifold

$$
T = \frac{\mathbb{C}^{21} - Z}{\mathbb{C}^{*^{17}}}
$$
 (49)

and determine the polynomial specifying the mirror orbifold *Y*

$$
I = \frac{1}{\binom{x^{17}}{a}}
$$
\nthe polynomial specifying the mirror

\n
$$
W_Y(y) = \sum_{i=1}^{h_Y + 5} \tilde{C}_i \prod_{a=1}^{21} y_a^{n_a^i}.
$$
\n(50)

Imposing the necessary conditions (23), we arrive at the system of equations

$$
\sum_{a=1}^{21} Q_{la} n_a = 5, \quad n_b \ge 0, \quad b = 1, ..., 21.
$$
 (51)

Solving Eqs.  $(51)$ , we obtain the sets  $n_a$  and change the coordinates in *T*, reducing it to  $\mathbb{P}^4$ . The polynomial specifying the orbifold *Y* has the form

$$
W_Y(z) = z_1^5 + z_2^5 + z_3^5 + z_4^5 + z_5^5 + \sum_{m=1}^5 \psi_m e_m^T(z), \quad (52)
$$

where the monomials  $e_m^T(z)$  are specified by Eqs. (44).

Thus, the two constructions give the same results. **Example 2.** We now consider the orbifold

4

$$
X = \frac{2}{\mathbb{Z}_5[0, 1, 2, 3, 4]}, \quad h_X = 21. \tag{53}
$$

The Calabi–Yau manifolds of *X* are specified by the zeros of the polynomial

$$
W_X(x) = x_1^5 + x_2^5 + x_3^5 + x_4^5 + x_5^5 + \sum_{l=1}^{21} \varphi_l e_l(x)
$$
  
= 
$$
\sum_{a=1}^{26} C_a \prod_{j=1}^5 x_j^{V_{aj}},
$$
 (54)

and the deformations of the polynomial  $e_l(x)$  are invariant under the action of the group  $H =$  $\mathbb{Z}_5[1,1,1,1,1] \oplus \mathbb{Z}_5[0,1,2,3,4]$ :

$$
e_1 = x_1^3 x_3 x_4, \t e_2 = x_1^3 x_2 x_5, \t e_3 = x_1^2 x_2 x_3^2,
$$
  
\n
$$
e_4 = x_1^2 x_2^2 x_4, \t e_5 = x_1^2 x_3 x_5^2, \t e_6 = x_1^2 x_4^2 x_5,
$$
  
\n
$$
e_7 = x_1 x_2 x_4^3, \t e_8 = x_1 x_2^2 x_5^2, \t e_9 = x_1 x_4 x_5^3,
$$
  
\n
$$
e_{10} = x_1 x_3^3 x_5, \t e_{11} = x_1 x_2^3 x_3, \t e_{12} = x_1 x_3^2 x_4^2,
$$
  
\n
$$
e_{13} = x_3^2 x_4 x_5^2, \t e_{14} = x_2 x_3^3 x_4, \t e_{15} = x_2 x_3 x_5^3,
$$
\n(55)

$$
e_{16} = x_2 x_4^2 x_5^2
$$
,  $e_{17} = x_2^2 x_3 x_4^2$ ,  $e_{18} = x_2^2 x_3^2 x_5$ ,  
\n $e_{19} = x_2^3 x_4 x_5$ ,  $e_{20} = x_3 x_4^3 x_5$ ,  
\n $e_{21} = x_1 x_2 x_3 x_4 x_5$ .

Similar to the preceding example, the Berglund– Hubsch–Krawitz construction gives the mirror orbifold  $Y = \mathcal{Q}/H^{T}$ , where

$$
H^{T} = \mathbb{Z}_{5}[0, 1, 3, 1, 0] \oplus \mathbb{Z}_{5}[0, 1, 1, 0, 3].
$$
 (56)

The only monomial invariant under the action of  $H^T$ is  $\prod_{i} z_i$ . Then, the mirror Calabi–Yau orbifold

$$
Y = \frac{9}{\mathbb{Z}_5[0, 1, 3, 1, 0] \oplus \mathbb{Z}_5[0, 1, 1, 0, 3]}, \quad h_Y = 1 \quad (57)
$$

is specified by the zeros of the polynomial

$$
W_Y(z) = z_1^5 + z_2^5 + z_3^5 + z_4^5 + z_5^5 + \psi z_1 z_2 z_3 z_4 z_5. \tag{58}
$$

We now apply the Batyrev construction. The set of vectors  $V_a$ ,  $a = 1,..., 26$ , and integers  $Q_{la}$ ,  $l = 1,..., 21$ , are determined similarly. The mirror of *Y* is specified by the zeros of the polynomial  $W_Y$  in the toric manifold

$$
T = \frac{\mathbb{C}^{25} - Z}{\mathbb{C}^{*21}}.
$$
 (59)

Imposing similar constraints on the degrees of monomials of the polynomial  $W_Y$  and making a change to the projective coordinates, we arrive at its explicit form that coincides with Eq. (58).

**Example 3.** We consider now the orbifold

$$
X = \frac{2}{\mathbb{Z}_5[0, 0, 0, 1, 4]}, \quad h_X = 25. \tag{60}
$$

The whole family of the orbifold of the Calabi–Yau manifold  $X$  is specified by the equation

$$
W_X = x_1^5 + x_2^5 + x_3^5 + x_4^5 + x_5^5 + \sum_{l=1}^{25} \varphi_l e_l(x)
$$
  
= 
$$
\sum_{a=1}^{30} C_a \prod_{j=1}^5 x_j^{V_{aj}},
$$
 (61)

where the monomials  $e_l$  are invariant under the action of  $H = \mathbb{Z}_5[1, 1, 1, 1, 1] \oplus \mathbb{Z}_5[0, 0, 0, 1, 4]$ :

$$
e_1 = x_1^3 x_2^2, \t e_2 = x_1^3 x_3^2, \t e_3 = x_2^3 x_3^2,
$$
  
\n
$$
e_4 = x_1^2 x_2^3, \t e_5 = x_1^2 x_3^3, \t e_6 = x_2^2 x_3^3,
$$
  
\n
$$
e_7 = x_1^3 x_2 x_3, \t e_8 = x_1 x_2^3 x_3, \t e_9 = x_1 x_2 x_3^3,
$$
  
\n
$$
e_{10} = x_1^2 x_2^2 x_3, \t e_{11} = x_1^2 x_2 x_3^3, \t e_{12} = x_1 x_2^2 x_3^2,
$$
  
\n
$$
e_{13} = x_1^3 x_4 x_5, \t e_{14} = x_2^3 x_4 x_5, \t e_{15} = x_3^3 x_4 x_5,
$$
  
\n
$$
e_{16} = x_1^2 x_2 x_4 x_5, \t e_{17} = x_1^2 x_3 x_4 x_5,
$$
  
\n(62)

JETP LETTERS Vol. 112 No. 6 2020

 $21$ 

$$
e_{18} = x_2^2 x_3 x_4 x_5, \t e_{19} = x_1 x_2^2 x_4 x_5, \n e_{20} = x_1 x_3^2 x_4 x_5, \t e_{21} = x_2 x_3^2 x_4 x_5, \n e_{22} = x_1 x_4^2 x_5^2, \t e_{23} = x_2 x_4^2 x_5^2, \n e_{24} = x_3 x_4^2 x_5^2, \t e_{25} = x_1 x_2 x_3 x_4 x_5.
$$

In [12], the total Hodge number is

$$
h_X^{\text{Full}} = h_X + h_X^{\text{Laurent}} = 49. \tag{63}
$$

Thus, there are additional 24 Laurent monomials among the deformations of the initial polynomial  $W_0$ :

$$
g_1 = x_1^{-1} x_2 x_3^3 x_4 x_5, \t g_2 = x_1 x_2^{-1} x_3^3 x_4 x_5, \ng_3 = x_1^3 x_2 x_3^{-1} x_4 x_5, \t g_4 = x_1^{-1} x_2^2 x_3^2 x_4 x_5, \ng_5 = x_1^2 x_2^{-1} x_3^2 x_4 x_5, \t g_6 = x_1^2 x_2^2 x_3^{-1} x_4 x_5, \ng_7 = x_1^{-1} x_2^3 x_3 x_4 x_5, \t g_8 = x_1 x_2^3 x_3^{-1} x_4 x_5, \ng_9 = x_1^3 x_2^{-1} x_3 x_4 x_5, \t g_{10} = x_1^{-1} x_3^2 x_4^2 x_5^2, \ng_{11} = x_2^{-1} x_3^2 x_4^2 x_5^2, \t g_{12} = x_1^2 x_3^{-1} x_4^2 x_5^2, \ng_{13} = x_1^{-1} x_2 x_3 x_4^2 x_5^2, \t g_{14} = x_1 x_2^{-1} x_3 x_4^2 x_5^2, \ng_{15} = x_1 x_2 x_3^{-1} x_4^2 x_5^2, \t g_{16} = x_1^{-1} x_2^2 x_4^2 x_5^2, \ng_{17} = x_2^2 x_3^{-1} x_4^2 x_5^2, \t g_{18} = x_1^2 x_2^{-1} x_4^2 x_5^2, \ng_{19} = x_1^{-1} x_4^3 x_5^3, \t g_{20} = x_2^{-1} x_4^3 x_5^3, \ng_{21} = x_3^{-1} x_4^3 x_5^3, \t g_{22} = x_1^{-1} x_2^3 x_3^3, \ng_{23} = x_1^3 x_2^{-1} x_3^3, \t g_{24} = x_1^3 x_2^3 x_3^3.
$$

According to the Berglund–Hubsch–Krawitz construction, the mirror Calabi–Yau manifold is

$$
Y = \frac{\mathcal{D}}{\mathbb{Z}_5[0, 1, 4, 0, 0] \oplus \mathbb{Z}_5[0, 3, 0, 1, 1]}, \quad h_Y = 5, \quad (65)
$$

and is specified by zeros of the polynomial

$$
W_Y = \sum_{j=1}^5 z_j^5 + \psi_1 z_4^2 z_5^3 + \psi_2 z_4^3 z_5^2 + \psi_3 z_1 z_2 z_3 z_5^2
$$
  
+  $\psi_4 z_1 z_2 z_3 z_4^2 + \psi_5 z_1 z_2 z_3 z_4 z_5.$  (66)

In the Batyrev construction, the mirror orbifold is realized as a hypersurface in the manifold

$$
T = \frac{\mathbb{C}^{29} - Z}{\mathbb{C}^{*25}}.
$$
 (67)

Calculations similar to the above give the explicit form of the polynomial in the projective coordinates, which coincides with Eq. (66).

Calculations also give two additional monomials  $z_4 z_5^4$  and  $z_4^4 z_5$ , but they lie in the kernel of the Milnor

JETP LETTERS Vol. 112 No. 6 2020

ring 
$$
\mathbb{C}[z_1,\ldots,z_5]/\frac{\partial W_0}{\partial z_i}
$$
. The two constructions give the same results

same results.

To summarize, we have shown that the two constructions of mirror pairs give the same results for the case of quintic orbifolds. It is interesting to understand the relation between these constructions in the general case.

#### ACKNOWLEDGMENTS

We are grateful to A. Litvinov for stimulating discussions.

### FUNDING

This work was performed at the Landau Institute for Theoretical Physics, Russian Academy of Sciences. A. Belavin acknowledges the support of the Ministry of Science and Higher Education of the Russian Federation (state assignment no. 0033-2019-0004). B. Eremin acknowledges the support of the Russian Science Foundation (project no. 18-12-00439).

#### REFERENCES

- 1. P. Candelas, G. T. Horowitz, A. Strominger, and E. Witten, Nucl. Phys. B **258**, 46 (1985). https://doi.org/10.1016/0550-3213(85)90602-9
- 2. A. Strominger, Commun. Math. Phys. **133**, 163 (1990). https://doi.org/10.1007/BF02096559
- 3. M. Kreuzer, Phys. Lett. B **328**, 312 (1994); arXiv: hepth/9402114 [hep-th]. https://doi.org/10.1016/0370-2693 (94)91485-0
- 4. P. Berglund and T. Hubsch, AMS/IP Stud. Adv. Math. **9**, 327 (1998); arXiv:hep-th/9201014 [hep-th]. https://doi.org/10.1016/0550-3213(93)90250-S
- 5. M. Krawitz, arXiv: 0906.0796.
- 6. A. Chiodo and Y. Ruan, arXiv: 0908.0908.
- 7. K. Hori, S. Katz, A. Klemm, R. Pandharipande, R. Thomas, C. Vafa, R. Vakil, and E. Zaslow, *Mirrorsymmetry* (Am. Math. Soc., Providence, RI, 2003), p. 101.
- 8. V. V. Batyrev, J. Alg. Geom. **3**, 493 (1994); arXiv: alggeom/9310003 [math.AG].
- 9. K. Aleshkin, A. Belavin, and A. Litvinov, arXiv: 1812.00478 [hep-th]. https://doi.org/10.1088/1742-5468/ab081a
- 10. K. Aleshkin and A. Belavin, JETP Lett. **110**, 705 (2019); arXiv: 1911.11678 [hep-th].
- 11. A. Belavin and B. Eremin, Theor. Math. Phys. **201**, 1606 (2019); arXiv:1907.11102 [hep-th]. https://doi.org/10.1134/S0040577919110060
- 12. B. R. Greene and M. R. Plesser, Nucl. Phys. B **338**, 15 (1990). https://doi.org/10.1016/0550-3213(90)90622-K

*Translated by R. Tyapaev*