

Instability of Thyatron Operation in Power Supply Units of Metal-Vapor Lasers

N. A. Yudin

Tomsk State University, pr. Lenina 36, Tomsk, 634050 Russia

e-mail: yudin@tic.tsu.ru

Received March 31, 2014; in final form, April 23, 2014

Abstract—The causes of an increase in the ТГН1-1000/25 thyatron jitter from 5 to ~30 ns in pumping sources of metal-vapor lasers (MVLs) upon changes in the excitation-pulse repetition frequency from 700 Hz to ~10 kHz are analyzed. It is shown that to provide the thyatron operation with an instability of the grid–cathode gap breakdown of ~1 ns, a hydrogen generator, and a device that forms an igniting pulse at the thyatron grid must be connected to stabilized dc power-supply units. Engineering solutions are proposed that allow control of the anode-current delay time at a thyatron jitter of ~1 ns.

DOI: 10.1134/S0020441215010224

INTRODUCTION

The unstable operation of pulse hydrogen thyatrons with respect to a triggering pulse (jitter) is one of the main causes that hinder the practical use of thyatrons connected in parallel for increasing the switched energy in lasers on self-contained transitions of metal atoms (self-contained lasers, metal-vapor lasers (MVLs)) and in systems for summing laser radiation of MVLs: “master generator–amplifier” [1, 2]. The jitter of a ТГН1-1000/25 thyatron, which is most widely used in power supply units of MVLs, is ~5 ns according to the certificate characteristics at a pulse repetition frequency of ~700 Hz [3]. A typical value of the jitter of a ТГН1-1000/25 thyatron that operates in the discharge circuit of an MVL with a repetition rate of excitation pulses of ~10 kHz is ~30 ns, which is comparable to the duration of MVL-generated pulses. The true causes of the appearance of this jitter are unknown.

It was supposed in [4] that the jitter may be associated to a certain degree with the instability of the thyatron filament current, which is in turn determined by voltage fluctuations in the mains. However, the thyatron jitter can be hardly related to the filament current of the thyatron cathode because the time of changes in the temperature of a massive thermionic cathode substantially exceeds ~0.01 s. In fact, supplying the filament voltage of the ТГН1-1000/25 thyatron cathode from the source of a dc voltage of 6.3 V in a copper-vapor laser with a ГЛ-201 commercial gas-discharge tube [5] did not change the observed value of the thyatron jitter.

Let us consider the discharge development in a thyatron in more detail in order to assess the possible causes of appearance of a ТГН1-1000/25 thyatron jitter of ~30 ns for a pulse repetition frequency of ~10 kHz. As is known [3, 6], pulse thyatrons have a

high electric strength. The thyatron grid has a low penetrability, so that the anode field does not actually act in the near-cathode region. To initiate a discharge between the anode and cathode, it is first necessary to apply a positive voltage pulse (ignition pulse) between the grid and cathode, which produces an auxiliary discharge in this gap. The latter initiates a discharge in the main anode–grid gap.

At the beginning of the switching period when the concentration of charged particles between the cathode and grid is still low, the resistance of the thyatron is high. As the charged-particle concentration increases, the ion current to the grid in the anode–grid gap also increases and at a certain moment, the total current of the grid $I_g = I_e + I_i$ (where I_e and I_i are the electron and ion currents, respectively) changes its sign.

The moment when the grid current changes its direction coincides with the thyatron enabling onset (a noticeable anode current appears). The instantaneous value of the electron current before the change in its direction is the grid current that enables the thyatron in the pulsed mode. A high ion current flows to the grid for a very short time, as long as there is a considerable potential difference between the grid and anode and the charged-particle concentration in the cathode–grid discharge gap is insufficient.

An increase in the plasma concentration in the thyatron reduces the grid–anode voltage, and the electron current of the cathode again becomes the main fraction of the grid current. An ion-current surge is accompanied by a short rise of the grid voltage relative to the cathode, which is called the “grid peak.” A grid peak increases a discharge current between the cathode and grid and accelerates the termination of the switching process. The grid-peak amplitude

reaches 20–30% of the anode voltage, and the peak duration is $\sim 10^{-7}$ s.

To maintain the gas-pressure constancy in hydrogen thyratrons, there is a special generator–accumulator of hydrogen, which is an element heated with a special heater. Hydrogen is liberated under heating, and a required pressure is maintained during long-term operation of the device. To compensate for the effect of the ambient temperature and stabilize the filament current of the hydrogen generator, a wire resistor with a high temperature coefficient of resistance, which is placed beyond the housing of the instrument, is connected in series to the hydrogen generator. The heater of the hydrogen generator with the compensating resistor is connected in parallel to the thyatron filament circuit or has a separate lead, thus allowing the hydrogen generator to be powered from a separate source.

Along with the jitter, another important characteristic of the thyatron is the anode-current delay time (τ_{del}), which is the interval between the moment corresponding to the onset of the voltage-front rise of an igniting pulse in the grid circuit and the moment of the discharge onset in the thyatron-anode circuit. The anode-current delay time for the ТГП1-1000/25 thyatron is $\sim 0.35 \pm 0.15$ μs [3] and is determined by the sum of two components: $\tau_{\text{del}} = \tau_{\text{brd}} + \tau_{\text{dis}}$, where τ_{brd} is the time from the onset of the voltage-front rise of the igniting pulse to the breakdown voltage of the grid–cathode gap, and τ_{dis} is the time from the breakdown moment to the moment of the discharge initiation in the thyatron anode circuit. An increase in the thyatron jitter with an increase in the pulse repetition rate may be correspondingly related to both instabilities of the parameters of the igniting pulse and breakdown voltage of the grid–cathode gap and an instability in the discharge development in the thyatron anode circuit after the breakdown.

INFLUENCE OF THE PARAMETERS OF THE IGNITING PULSE ON THE THYRATRON JITTER

From the above description of the discharge-development process in the thyatron, it can be presumed that an increase in the thyatron jitter with an increase in the pulse repetition rate is determined by the power supplies of the hydrogen generator and the igniting-pulse forming device. The hydrogen generator in the MVL is connected (as a rule) in parallel to the thyatron filament circuit and is powered from an ac power supply of ~ 6.3 V/50 Hz; this may determine a hydrogen-pressure modulation in the thyatron. According to the Paschen law, the breakdown voltage is a function of the gas pressure. Correspondingly, a hydrogen-pressure modulation in the thyatron must modulate the breakdown voltage across the thyatron grid–cathode gap and determine the thyatron jitter.

Connecting the hydrogen generator of the ТГП1-1000/25 thyatron to a stabilized source of a 6.3-V dc voltage allowed the observed thyatron jitter to be reduced by 18 ns at an igniting-pulse repetition frequency of ~ 10 kHz. This directly indicates that the hydrogen pressure in the thyatron is modulated by the electric-mains frequency of ~ 50 Hz.

When the hydrogen generator was supplied by the dc voltage, a change in the ТГП1-1000/25 thyatron jitter from 5 to 12 ns was observed under changes in the pulse repetition rate from 700 Hz to ~ 10 kHz, and the instability in the breakdown of the grid–cathode gap changed from ~ 1 to 7 ns in the absence of a voltage at the thyatron anode. This allows one to conclude that the minimum thyatron jitter (~ 5 ns) is not associated with the discharge-development instability in the thyatron–anode circuit after a breakdown of the grid–cathode gap.

Our studies showed that a breakdown of the grid–cathode gap in the ТГП1-1000/25 thyatron occurs when the amplitude of the igniting-pulse voltage reaches ~ 500 V, if the hydrogen generator is powered by a stabilized voltage. The observed increase in the thyatron jitter with an increase in the repetition frequency of igniting pulses is caused by an amplitude instability of the igniting-pulse voltage and determined by the pulsation value at the output of the rectifier, which provides powering of the igniting-pulse-forming unit.

Measurements have shown that 10% voltage pulsations at the rectifier output lead to a variation in the leading edge of the igniting-pulse voltage (Fig. 1) and, correspondingly, in the time within which the breakdown voltage is reached. This leads to an instability of the grid–cathode gap breakdown of ~ 20 ns. To reduce the instability of the grid–cathode gap breakdown, it is necessary to reduce the pulsations at the rectifier output by increasing the filter capacitance. To provide the thyatron operation with a grid–cathode gap breakdown instability of ~ 1 ns, it is necessary to connect the hydrogen generator and the igniting-pulse-forming unit to stabilized dc voltage sources with voltage pulsations at the source output of $\leq 1\%$. In this case, the thyatron jitter is determined only by the discharge-development instability in the thyatron-anode circuit after the grid–cathode gap breakdown.

A specific feature of thyratrons [3, 6] is that the higher the slope of the grid current dI_g/dt , the higher the instantaneous grid-current value at which the thyatron is. This regularity can be explained in the following way. The thyatron is enabled at a certain value of the anode current, which is extracted from plasma of the auxiliary grid discharge. This current at a given anode voltage corresponds to a definite plasma concentration in the grid region into which the anode field penetrates.

The critical concentration is formed as a result of two processes: the accumulation of charged particles and their diffusion to the walls that surround the dis-

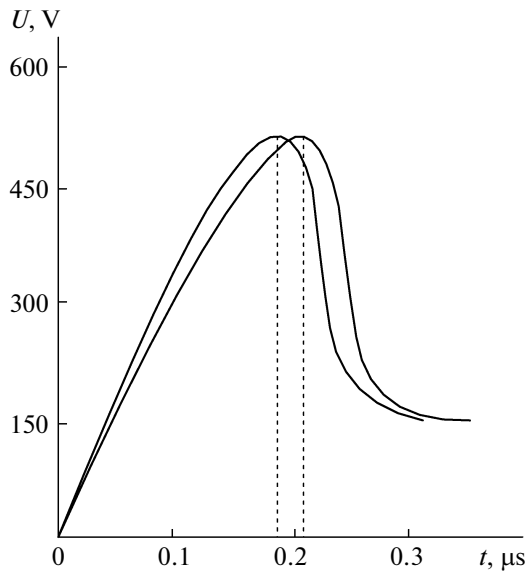


Fig. 1. Voltage of the discharge initiation in the grid–cathode gap and the ТГН1-1000/25 thyatron as a function of 10% voltage pulsations of the power source in the device for forming a positive voltage pulse at the thyatron grid. Measurements were performed in the absence of a voltage at the thyatron anode.

charge space. The thyatron is enabled upon pulsed ignition when the grid-plasma concentration has not yet been settled and charged particles are accumulated more rapidly than their diffusion to the walls occurs. Therefore, the plasma concentration is proportional not to the instantaneous current value but to the amount of the electric charges that passed through the discharge gap.

The aforementioned features can be used for reducing the thyatron unblanking time and its jitter time (<5 ns).

FORMATION OF A THYRATRON TRIGGERING PULSE IN THE FORM OF A GRID PEAK

The possibility of forming a thyatron-triggering pulse in the form of a grid peak is determined by the fact that a grid peak arises at the thyatron grid in any case and, consequently, its artificial formation, at least, must not impair the operating characteristics of the thyatron and in the best case, must result in reducing both the thyatron-enabling and jitter times. Investigations showed that the use of a thyatron-triggering pulse in the form of a grid peak allows the enabling time of the ТГН1-1000/25 thyatron to be reduced to ~30 ns, and the thyatron jitter is virtually absent (<1 ns) in this case.

The use of a thyatron-triggering pulse in the form of a grid peak should be preferred in systems intended for synchronizing the operation of several lasers, e.g., in master oscillator–amplifier laser systems, because,

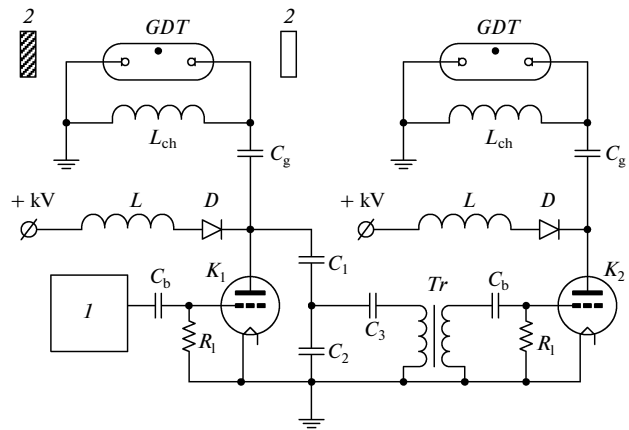


Fig. 2. Schematic diagram of the oscillator–amplifier laser system with synchronization of the operation of two thyatrons: (I) master oscillator, (2) resonator, (GDT) gas-discharge tube, (C_g) storage capacitors, (C_b) blocking capacitors, (C_1 , C_2) capacitive voltage divider, (C_3) storage capacitor, (R_1) leakage resistor, (L , D) charging choke and diode, (L_{ch}) shunting inductance, (Tr) pulse transformer, and (K_1 , K_2) thyatrons.

in this case, a device for forming a grid peak can be manufactured most easily (Fig. 2). Figure 2 shows a diagram of the oscillator–amplifier laser complex with synchronization of the operation of thyatrons in pumping generators.

The device for forming a grid peak in this circuit was based on the elements C_1 – C_3 and a pulse transformer Tr . The operation of the oscillator–amplifier complex with synchronization of the operation of two thyatrons is accomplished as follows. The storage capacitors C_g and capacitors C_1 – C_3 are charged from high-voltage rectifiers through charging chokes L and diodes D .

After the storage capacitors are charged, a triggering pulse arrives at the thyatron K_1 from master oscillator I . The thyatron K_1 is unblanked and a discharge of the corresponding storage capacitor, which forms an excitation pulse across the gas-discharge tube (GDT), occurs. In this case, the capacitors C_1 – C_3 are also discharged, and the discharge of the capacitor C_3 through the transformer Tr (the transformation ratio is 1 : 1) forms a grid peak at the grid of the thyatron K_2 , which unblanks the latter.

The capacitors C_1 and C_2 form a voltage divider and determine the charging voltage of the capacitor C_3 . The capacitance of the capacitor C_3 is chosen on the basis of the certificate characteristics of the thyatron with allowance for the provision of the required energy for unblanking the thyatron. For the ТГН1-1000/25 thyatron, it is ~100–150 pF at a voltage of ~5–7 kV. In this case, the grid-peak duration is ~30 ns. The thyatron K_2 is enabled 30 ns after the thyatron K_1 . As the capacitors C_1 – C_3 , КВИ-3-type capacitors are used.

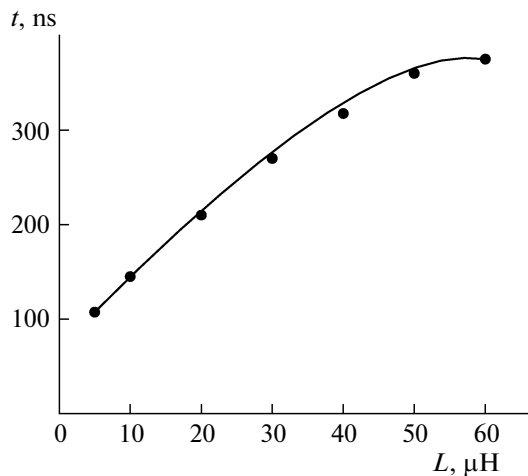


Fig. 3. Time of the discharge initiation in the ТГН1-1000/25 thyatron grid–cathode gap as a function of the choke L inductance. Measurements were performed in the absence of a voltage at the thyatron anode.

CONTROLLING THE ANODE-CURRENT DELAY TIME

The performed studies showed that a breakdown of the grid–cathode gap in the ТГН1-1000/25 thyatron occurs when the amplitude of the positive voltage pulse at the grid is ~ 500 V. In this case, the anode-current delay time depends on the slope of the grid-voltage pulse front, and the thyatron jitter decreases with an increase in the current in the grid–cathode gap. A simple technical solution that allows control of the anode-current delay time at a thyatron jitter of ~ 1 ns follows from the above.

In fact, if the voltage rise time across the thyatron grid–cathode gap is varied, this must lead to a change in the onset time of the discharge initiation in this gap; thus, the anode-current delay time can be controlled. The possibility of changing the moment of the discharge initiation in the thyatron grid–cathode gap in a wide range was experimentally tested using different circuits for forming thyatron-triggering pulses. Direct changes in the discharge-initiation moment for the thyatron grid–cathode gap were performed via connection of an additional choke L between the thyatron grid and the device for forming a positive voltage pulse at the grid.

Figure 3 shows the dependence of the time of the discharge initiation in the ТГН1-1000/25 thyatron grid–cathode gap on the choke L inductance. By changing the inductance of the choke L from 5 to 60 μH , one can control the time of the discharge initiation in the thyatron grid–cathode gap within a range of 200–300 ns. In this case, the inductance of L limits the discharge current in the grid–cathode gap. In this case, the grid-current development rate decreases; this increases the thyatron enabling time and must inevitably result in the appearance of a thyatron jitter.

However, if a capacitor is connected in parallel to the thyatron grid, in this case, the voltage rise rate at the thyatron grid determines the time of the capacitor charging to a breakdown voltage of ~ 500 V of the grid–cathode gap. After the thyatron breakdown, the current development in the grid–cathode gap is determined by the discharge time of this capacitor and is independent of the choke- L value. This simple technical solution allows elimination of the thyatron jitter in the mode of controlling the discharge initiation time in the grid–cathode gap. The optimal value of the capacitance of this capacitor (КВИ-3) is ~ 470 – 1000 pF for the ТГН1-1000/25 thyatron.

On the basis of the performed investigations, a unit for forming a positive voltage pulse at the ТГН1-1000/25 thyatron grid was developed. A circuit diagram of this unit is shown in Fig. 4. The unit for forming a positive voltage pulse is triggered by a generator of rectangular pulses with an amplitude of 15 V and a duration of ~ 0.5 – 1 μs . The generator specifies the repetition frequency of pulses that initiate a discharge in the grid–cathode gap.

The resistor R_2 limits the current in the gate circuit of the field-effect transistor Q_1 . The resistor R_3 protects the transistor in the case of a break in the resistor R_2 resistor. The elements R_1 , C_1 , D_1 , C_2 , R_4 , and D_2 form the required shape of pulses, eliminate the parasitic generation at the fronts, and protect the transistor against high-power voltage pulses. The inclusion of the capacitor C_3 into the circuit is determined by the necessity of connecting the source of a negative bias voltage to the thyatron grid through the resistor R_5 . The time-setting elements L and C_4 determine the anode-current delay time, and the delay time is controlled via a change in the choke L inductance.

It is expedient to use this control method for providing the simultaneous triggering of thyatrons connected in parallel, because each of them is triggered from its individual positive-pulse forming unit. The necessity of adjusting the L and C_4 circuits in the pulse-formation units arises because of the fact that circuit elements have a certain spread of parameters, and the breakdown voltage of the grid–cathode gap of the ТГН1-1000/25 thyatrons is $\sim 500 \pm 10$ V.

The thyatron-triggering delays in the master oscillator–amplifier systems in the absence of a jitter can be also performed via introduction of an appropriate delay between pulses that arrive at the gate of the transistor Q_1 from the rectangular-pulse generator. In this case, a choke of the ДПМ-0.6-16 or ДПМ-1.2-30 type with a constant inductance of ~ 10 – 30 μH can be used as the choke L . As was mentioned above, the development of a current in the grid–cathode gap after the breakdown is determined by a discharge of the capacitor C_4 and does not depend on the choke L inductance. This eliminates the necessity of minimizing the leakage inductance of the pulse transformer Tr , thus simplifying its production. The transformer Tr is wound on a ferrite core with an ETD39 frame (Sie-

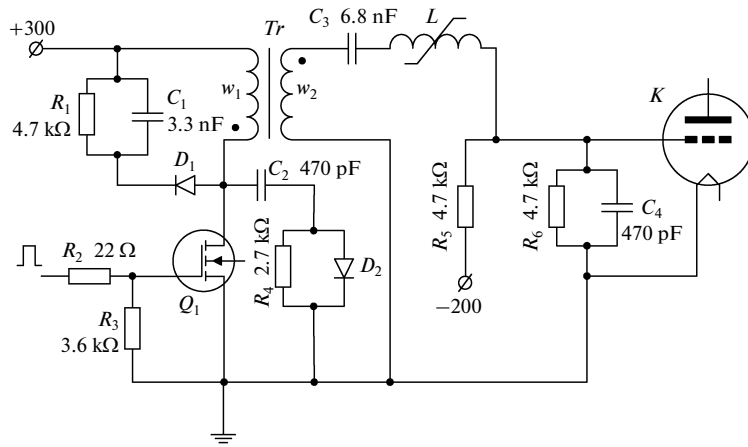


Fig. 4. Circuit diagram of the device for forming a positive voltage pulse at the thyatron grid: (K) ТГП1-1000/25, (Q_1) 2SK1357, (D_1, D_2) FR207, and (Tr) ETD39.

mens + Matsushita firm, $w_1 = 20$ turns and $w_2 = 40$ turns).

CONCLUSIONS

The performed investigations allowed us to establish the causes of the appearance of a jitter in pulse hydrogen thyratrons and to develop devices for their triggering, which allow control of the anode-current delay time at a thyatron jitter of ~ 1 ns. These studies were initiated by the necessity of practically realizing the master oscillator–amplifier system in a single gas-discharge tube at a two-pulse excitation of the active medium [7], when the first excitation pulse forms a light field in the resonator and the second pulse provides its amplification. Under these pumping conditions, the thyatron jitter determined a substantial instability in the copper-vapor-laser pulse energy.

REFERENCES

1. Coutts, D.W. and Brown, D.J.W., *CLEO'93. Tech. Dig. Opt. Soc. Amer.*, Washington, DC, 1993, p. 460.

2. Amit, M., Lavi, S., Erez, G., and Miron, E., *Opt. Commun.*, 1987, vol. 62, no. 2, p. 110.
 3. Katsnel'son, B.V., Kalugin, A.M., and Larionov, A.S., *Elektrovakuumnye elektronnye i gazorazryadnye pribory: Spravochnik* (Electrovacuum Electron and Gas-Discharge Devices: A Handbook), Moscow: Radio i Svyaz', 1985.
 4. Batenin, V.M., Buchanov, V.V., Kazaryan, M.A., Klimovskii, I.I., and Molodykh, E.I., *Lazery na samoorganizatsionnykh perekhodakh atomov metallov* (Lasers on Self-Contained Transitions of Metal Atoms), Moscow: Nauchnaya Kniga, 1998.
 5. Grigor'yants, A.G., Kazaryan, M.A., and Lyabin, N.A., *Lazery na parakh medi: konstruktziya, kharakteristiki i primeneniya* (Copper Vapor Lasers: Construction, Characteristics, and Applications), Moscow: Fizmatlit, 2005, p. 312.
 6. Fogel'son, T.B., Breusova, L.N., and Vagin, L.I., *Impul'snye vodorodnye tiratrony* (Pulse Hydrogen Thyratrons), Moscow: Sovetskoe Radio, 1974.
 7. Polunin, Yu.P. and Yudin, N.A., *Quantum Electron.*, 2003, vol. 33, no. 9, p. 833.

Translated by A. Seferov