Changes in the Electron Transfer Symmetry in the Photosystem I Reaction Centers upon Removal of Iron–Sulfur Clusters

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Abstract— In photosynthetic reaction centers of intact photosystem I (PSI) complexes from cyanobacteria, electron transfer at room temperature occurs along two symmetrical branches of redox cofactors *A* and *B* at a ratio of ~3 : 1 in favor of branch *A*. Previously, this has been indirectly demonstrated using pulsed absorption spectroscopy and more directly by measuring the decay modulation frequencies of electron spin echo signals (electron spin echo envelope modulation, ESEEM), which allows to determine the distance between the separated charges of the primary electron donor P_{700}^+ and phylloquinone acceptors A_{IA}^- and A_{IB}^- in the symmetric redox cofactors branches *A* and *B*. In the present work, these distances were determined using ESEEM in PSI complexes lacking three $4Fe-4S$ clusters, F_X , F_A , and F_B , and the PsaC protein subunit (the so-called $P_{700}-A_1$ core), in which phylloquinone molecules A_{1A} and A_{1B} serve as the terminal electron acceptors. It was shown that in the $P_{700}-A_1$ core preparations, the average distance between the centers of the $P_{700}^+\bar{A}_1^-$ ion-radical pair at a temperature of 150 K in an aqueous glycerol solution and in a dried trehalose matrix, as well as in a trehalose matrix at 280 K, is \sim 25.5 Å, which corresponds to the symmetrical electron transfer along the *A* and *B* branches of redox cofactors at a ratio of 1 : 1. Possible reasons for the change in the electron transfer symmetry in PSI upon removal of the PsaC subunit and 4Fe–4S clusters F_x , F_A , and F_B are discussed.

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INTRODUCTION

The reaction center (RC) of photosystem I (PSI) of cyanobacteria and green plants contains 11 redox cofactors, including six chlorophyll (Chl) molecules, two phylloquinone (PhQ) molecules, and three 4Fe–4S type iron–sulfur clusters – F_X , F_A , and F_B [1]. Laser flashinduced electron transfer includes the oxidation of the primary electron donor chlorophyll dimer P_{700} and subsequent transfer through the primary chlorophyll acceptor A_0 to the phylloquinone acceptors A_1 and 4Fe–4S clusters F_X , F_A , and F_B [2]. In the absence of exogenous electron acceptors, charge recombination between the reduced terminal $4Fe-4S [F_A/F_B]$ ⁻ clusters and the oxidized P_{700}^{+} donor occurs in PSI RC. Previously, it was shown that sequential removal of the F_B , F_A , and F_X clusters from PSI results in the backward electron transfer to P_{700}^{+} from the preceding acceptors F_A , F_X , and A_1 , respectively [3-5].

The PSI complex from cyanobacteria consists of 12 protein subunits, three of which (PsaA, PsaB, and PsaC)

Abbreviations: A1, phylloquinone, secondary electron acceptor; Chl, chlorophyll; Em, midpoint redox potential; ESEEM, electron spin echo envelope modulation; P₇₀₀, chlorophyll dimer, primary electron donor; PhQ, phylloquinone; PSI, photosystem I; RC, reaction center.

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bind all electron transfer cofactors. These subunits also bind 79 out of 90 Chl molecules of the light-harvesting antenna. PsaA and PsaB subunits form a heterodimer that constitutes the central part of the complex and bind six Chl molecules, two PhQ molecules, and the 4Fe–4S cluster F_x , while the PsaC subunit adjacent to the acceptor site of the PsaA/PsaB heterodimer binds the terminal iron–sulfur clusters F_A and F_B . Redox cofactors associated with PsaA/PsaB are located along the C_2 axis of symmetry that passes through the geometric center of the Chl dimer P_{700} and the F_X cluster and form two symmetrical branches (*A* and *B*) [1, 3]. The *A* branch contains the primary acceptor A0A (consisting of Chl molecules *Chl2A* and *Chl3A*) and the PhQ molecule A_{1A} ; the *B* branch contains the primary acceptor A_{0B} (consisting of Chl molecules *Chl2B* and *Chl3B*) and the PhQ molecule A_{1B} . Electron transfer occurs along both branches; it starts with the oxidation of P_{700} and ends at F_X , where the branches converge and further transfer to the terminal acceptors $[F_A/F_B]$ takes place (Fig. 1a).

The factors determining participation of both branches of redox cofactors in the electron transfer in PSI RCs have not been fully elucidated [3]. In the other two types of RCs studied, photosystem II (PSII) complexes and bacterial RCs, which also have two symmetrical branches of cofactors similar to the *A* and *B* branches, the transfer occurs only along one branch to the primary quinone acceptor Q_A , and then from Q_A^- to the secondary acceptor Q_B with the formation of the stable semiquinone Q_{B}^- . Excitation of such RCs by a second flash of light leads to complete reduction of Q_B^- with the formation of protonated secondary quinone Q_BH_2 , which, in turn, is oxidized by the cytochrome bc_1 complex ($b_6 f$ in PSII). The transfer of electrons in PSI occurs from the PhQ molecules A_{1A} and A_{1B} to the iron–sulfur cluster F_X , but not between the A_{1A} and A_{1B} . It can be assumed that the possibility of transfer along both branches of redox cofactors in PSI is an evolutionary adaptation that makes

it possible to preserve an efficient reduction of lowpotential electron acceptors, ferredoxin and NADP, in the case of random mutations near the primary acceptors in one of the branches.

The electron transfer in the RCs of cyanobacterial PSI occurs asymmetrically along the *A* and *B* branches at a ratio of \sim 3 : 1 in favor of the *A* branch [6]. In this case, the transfer from A_{0A} to A_{1A} and from A_{0B} to A_{1B} occurs with the same characteristic time $\tau \approx 25$ ps, while the transfer from A_{1A} to F_X is approximately an order of magnitude slower than from A_{1B} to F_x ($\tau \approx 200$ ns vs. $\tau \approx 20$ ns, respectively) [7-9]. The difference in the rates of the A_{1A} → F_X and A_{1B} → F_X reactions is mainly caused by the difference in the free energy (ΔG) values of these reactions, since the midpoint redox potentials (E_m) of A_{1A} and A_{1B} differ by ~150 mV [10, 11]. It is possible that the recombination kinetics of the ion-radical pairs $P_{700}^{+}A_{1B}^{-}$ and $P_{700}^+A_{IA}^-$ in the PSI complexes lacking 4Fe–4S clusters (so called $P_{700}-A_1$ core, Fig. 1c) has different characteristic times ($\tau \approx 20$ and 200 μs, respectively) due to the difference between the ΔG values of the backward electron transfer reactions $A_{IB}^- \rightarrow P_{700}^+$ and $A_{IA}^- \rightarrow P_{700}^+$ [5, 12]. However, it should be noted that the removal of the PsaC subunit and iron–sulfur clusters can lead to changes in the E_m values of A_{1A} and A_{1B} . Since the negatively charged clusters F_A and F_B are located asymmetrically with respect to the PhQ molecules A_{1A} and A_{1B} , the removal of these clusters should lead to a decrease in the difference between the redox potentials of A_{1A} and A_{1B} [10, 13].

Electron transfer asymmetry between the branches of redox cofactors in PSI continues to be a subject of discussion (see review [3]). As shown by Plato et al., the spin density center between the *Chl1A* and *Chl1B* monomers in the P₇₀₀ dimer is significantly shifted towards *Chl1B* [14, 15]; as a result, the distances between the spin density centers of the $P_{700}^+A_{IA}^-$ and $P_{700}^+A_{IB}^-$ ion-radical pairs are \sim 26 Å and \sim 25 Å, respectively (Fig. 1c). Previously, EPR spectroscopy was used to register the modulation fre-

Fig. 1. Structure of the PSI RC: a) intact $P_{700} - F_A/F_B$ complex; b) $P_{700} - A_1$ core complex lacking 4Fe–4S clusters and PsaC subunit; c) distances between the spin density centers of P_{700}^+ and A_{IA}/A_{IB}^- . The paths of backward electron transfer are indicated by arrows.

quencies of the electron spin echo signal decay (electron spin echo envelope modulation, ESEEM), demonstrating that at cryogenic temperatures, the measured distance between the centers of the spin density in the ion-radical pair $P_{700}^+A_1^-$ (~26 Å) in intact PSI complexes from cyanobacteria ($P_{700} - F_A/F_B$) corresponds to the electron transfer along the *A* branch [16, 17]. The same distance was revealed by ESEEM for PSI complexes isolated from the cyanobacterial *MenB* mutant containing a high-potential quinone derivative (2,3-dichloronaphthoquinone) in the A_{1A} and A_{1B} binding sites. In these complexes, electron transfer to the 4Fe–4S clusters was not observed for thermodynamic reasons, and charge recombination occurred between A_1^- and P_{700}^+ [16]. The measured distance between P_{700}^{+} and A_1^- was also ~26 Å, which corresponds to the predominant electron transfer along the *A* branch.

Another way to determine the degree of electron transfer asymmetry between the redox cofactor branches *A* and *B* is to measure the ratio of amplitudes of kinetic components corresponding to the forward electron transfer from A_{IA}^- and A_{IB}^- to F_X . By measuring the electrochromic shift of carotenoid absorption band, it was shown that this ratio is 2 : 1 in favor of the *A* branch [18]. However, the carotenoid molecules in PSI RCs are oriented differently with respect to the PhQ molecules in the A_{1A} and A_{1B} sites. The $P_{700} - A_1$ core complexes were used to study the contribution of kinetic components corresponding to the charge recombination of the ion-radical pairs $P_{700}^+ A_{IA}^-$ and $P_{700}^+ A_{IB}^-$ on the one hand and to the electrochromic carotenoid shift on the other hand. Comparison of the relative amplitudes of these kinetic components showed that the contribution of the transfer along the *A* branch is underestimated by a factor of 2.33 compared to the contribution of the transfer along the branch [19]. Taking into account this coefficient, the ratio of forward electron transfer along the branches of redox cofactors should be ~4 : 1 in favor of the *A* branch. Since charge recombination from the PhQ molecules is not observed in the $P_{700}-F_A/F_B$ complexes, the kinetics of backward electron transfer from A_{IA}^- and A_{IB}^- to P_{700}^+ can only be measured if the forward transfer to the iron–sulfur clusters is arrested. Partial arrest of the forward electron transfer from A_l to the 4Fe–4S clusters is observed with decreasing temperature and upon drying in a glassy trehalose matrix [20-22]. However, at room temperature in solution, charge recombination in the ion-radical pairs $P_{700}^+A_{1A}^-$ and $P_{700}^+A_{1B}^-$ can only be registered in the $P_{700}^-A_{1B}^$ core complexes.

The study of the temperature dependence of the charge recombination kinetics between A_1^- and P_{700}^+ in the P_{700} –A₁ core complexes by pulse absorption spectrometry demonstrated that at 300 K, the ratio of the kinetic components contributions of the backward electron transfer from the reduced A_{IA}^- and A_{IB}^- acceptors to P_{700}^+ is ~60 : 40 in favor of the *A* branch. As the temperature drops to 150 K, these contributions become almost equal (50 : 50) [12].

The distances between the spin density centers in the $P_{700}^{+}A_1^-$ ion-radical pairs were measured at cryogenic temperatures in water–glycerol solutions. At temperatures above 200 K, pulsed EPR spectroscopy could not be used reliably due to the fast spin relaxation. However, drying PSI complexes in a glassy trehalose matrix slows down spin relaxation and makes it possible to detect the formation of the $P_{700}^+A_1^-$ ion-radical pairs using high-frequency EPR spectroscopy in the W and Q bands, as well as to determine the distance between the spin density centers in these pairs using ESEEM [17, 23, 24]. Previously, high-frequency EPR spectroscopy and ESEEM were used to study the spin-correlated $P_{870}^+Q_A^-$ radical pair in bacterial RCs in a dried trehalose matrix at 150 K and at room temperature [25]. It was shown that drying these preparations in a trehalose matrix does not change the molecular configuration of the $P_{870}^+Q_{\rm A}^-$ radical pair. It can be assumed that measuring the distances between the spin density centers of the ion-radical pairs is also possible for other pigment–protein complexes in a dried trehalose matrix at room temperature. The use of dehydrated trehalose matrices for the long-term preservation of the activity of various protein complexes at room temperature was discussed in detail in a review by Möbius et al. [26].

Recently, we measured the distance between P_{700}^{+} and A_1^- in the $P_{700} - F_A/F_B$ complexes in a dried trehalose matrix at cryogenic and room temperatures [23]. At 150 K and 220 K, this distance was \sim 26 Å and increased to \sim 27 Å at 280 K. However, this apparent increase in the distance (based on a decrease in the modulation frequency of the ESEEM signal), was explained by the acceleration of the A_1^- spin–lattice relaxation from 3 to 1 μ s under the influence of the iron ions of the 4Fe–4S clusters. Therefore, it was concluded that the distance between the spin density centers of the $P_{700}^+A_1^-$ pair in the trehalose matrix at room temperature remains unchanged and corresponds to the preferential electron transfer along the *A* branch.

In another study [24], we used the ESEEM method to determine the distance between P_{700}^+ and A_1^- in PSI complexes lacking the PsaC subunit and terminal 4Fe–4S clusters F_A/F_B ($P_{700}-F_X$ core complexes). The obtained data demonstrated that this distance in the dried trehalose matrix $(\sim 25 \text{ Å})$ corresponded to the preferential electron transfer along the *B* branch [24].

In the present study, we measured the distance between P_{700}^+ and A_1^- in the P_{700} - A_1 core complexes of PSI in a water–glycerol solution and in a dried trehalose matrix at cryogenic temperatures, as well as in the trehalose matrix at 280 K.

MATERIALS AND METHODS

Preparation of $P_{700}-A_1$ **core complexes.** Trimeric complexes of PSI were isolated from cyanobacterium *Synechocystis* sp. PCC 6803 [27] by treating thylakoid membranes with dodecyl-β-D-maltoside, followed by centrifugation in a sucrose density gradient (5-20%) at 210,000*g* for 3 h using Beckman VTi 50 rotor (USA), obtained with the support from the Lomonosov Moscow State University Program of Development. To produce PSI complexes lacking the iron–sulfur clusters $F_{X}/F_{A}/F_{B}$ $(P_{700}-A_1 \text{ core}, \text{ see Fig. 1b})$, intact PSI complexes from *Synechocystis* sp. PCC 6803 were incubated in a medium containing 50 mM Tris-HCl (pH 8.0) and 6.8 M urea for 90 min at room temperature in the dark. The final detergent concentration did not exceed ~0.05% (*m/V*). To remove urea, the PSI suspension was washed with 50 mM Tris-HCl buffer (pH 8.0) down to the urea concentration of $~0.03\%$. After the suspension was concentrated, the Chl concentration in the PSI samples was \sim 2 mg Chl/ml; at this stage, the terminal F_A/F_B clusters had been removed from the complex.

The obtained samples were incubated in a medium containing 50 mM Tris-HCl (pH 8.0), 3.4 M urea, and 5 mM potassium ferricyanide for 2 h at room temperature in the dark. To remove urea and ferricyanide, the PSI suspension was repeatedly washed with a buffer containing 50 mM Tris-HCl (pH 8.0) to the final concentrations of urea and ferricyanide of ~0.03% and ~0.005% (*m/V*), respectively. After concentrating the $P_{700}-A_1$ core suspension, the Chl concentration in RCs was \sim 1.5 mg Chl/ml (see [28] and references therein for the method description).

To obtain trehalose matrices, a solution containing PSI complexes and trehalose at a molar ratio of 1 : 40,000 (PSI : trehalose) was dehydrated in the presence of saturated aqueous solution of LiCl [21] (under these conditions, a relative air humidity above the salt solution was maintained at 11%). Dehydration of the PSI solution for 14 days in a desiccator with a saturated LiCl solution allowed to achieve an equilibrium between the protein in the trehalose matrix and air vapor. The water–glycerol solutions of the PSI complexes contained 50% glycerol (of the total volume).

ESEEM of P₇₀₀ $-A_1$ **core complexes.** To find the distance between the separated charges in RCs, we employed the pulse primary electron spin echo method. The measurements were carried out in the Q-band (34.03 GHz) with a Bruker Elexsys E580 EPR spectrometer (USA) using an EN 5107D2 resonator with a photoexcitation window that introduced the laser beam directly to the sample and Bruker ER 4131VT nitrogen cryostat (USA). YAG:Nd³⁺ laser Quantel Brilliant B (France–USA) was used for photoexcitation (wavelength, 532 nm; energy, 1 mJ; a repetition rate, 10 Hz). Measurements were performed in quartz capillaries with an outer diameter of 1.2 mm and an inner diameter of 1 mm sealed on one side. Dried trehalose matrices were placed in the capillaries in an atmosphere of inert gas. The experimental protocol is shown in Fig. 2 and is described in more detail in [23]. The Fourier transform of the signals was performed with

the Bruker Xepr software package (USA) using the fast Fourier transform algorithm.

In ESEEM experiments, laser pulse was used to induce charge separation in the PSI RC. After some delay, a sequence of two microwave pulses generated the primary electron spin echo signal. The dependence of the echo signal amplitude on the time interval *τ* between the two microwave pulses was recorded. The duration of the microwave $\pi/2$ pulse was 40 ns. During the experiment, ESEEM was observed that was caused by the spin–spin dipole–dipole interaction between the separated charges $P_{700}^{+}A_1^-$ in the PSI RC [29]. The frequency of the observed modulation is equal to the frequency of the spin–spin dipole–dipole interaction of the separated charges, which is defined as:

$$
\omega_{d-d} = \frac{g_1 g_2 \beta^2 (1 - 3 \cos^2 \theta) r^3}{\hbar},\tag{1}
$$

where r is the distance between the separated charges; g_1 , g_2 are the *g*-factors of the separated charges; β is the Bohr magneton; \hbar is the Planck constant; θ is the angle between the radius vector *r* and the direction of the induction vector of the constant magnetic field B_0 .

Equation (1) demonstrates that the dipole frequency ω_{d-d} takes on different values with a certain distribution, depending on the angle *θ*. Orientations of the vector *r* with $\theta = \pi/2$ have the greatest statistical weight, i.e., the largest contributions to the observed modulation of the echo signal come from RCs with $\theta = \pi/2$, for which the dipole frequency is equal to:

$$
\Omega = \frac{g_1 g_2 \beta^2}{r^3}.
$$
 (2)

The EPR spectrum of the spin-correlated pair $P_{700}^+A_1^$ obtained by integrating the echo signal with a fixed value of $\tau = 300$ ns, is shown at Fig. 3. The ESEEM signals were measured for the resonant value of the magnetic field corresponding to the signal from P_{700}^+ of the ion-radical pair $P_{700}^+A_1^-$. Signal modulation was observed for only one magnetization component of the echo signal [30].

Fig. 2. Experimental protocol for measuring the modulation frequency of the primary spin echo signal caused by the dipole–dipole interaction of the ions-radicals of the charge-separated state in PSI RC.

Fig. 3. EPR spectrum of the $P_{700}^+A_1^-$ ion-radical pair in the $P_{700}-A_1$ core PSI complexes incorporated into the trehalose matrix at 280 K. The spectrum was obtained by integrating the echo signal for various magnetic field values; the time interval τ in a two-pulse sequence was 300 ns.

RESULTS

Figure 4 shows the ESEEM signals of the $P_{700}-A_1$ core PSI complexes at room and cryogenic temperatures; Fig. 5 demonstrates the corresponding Fourier spectra.

Figure 5 shows that the main contribution to the modulation of the echo signal is due to the frequency of ~3.2 MHz. Figure 5b reveals the features of the Fourier spectra in the region around 3.2 MHz at an enlarged scale, which makes it possible to discern the modulation

Fig. 4. ESEEM signals of P_{700}^{+} in the P_{700} -A₁ core PSI complex obtained in water–glycerol solution at 150 K (black curve), in trehalose matrix at 150 K (red curve), and in trehalose matrix at 280 K (blue curve); green curves represent the fitting of the experimental data according to Eq. (3). The curves are arbitrarily shifted along the y-axis.

frequency shift by $~60$ kHz upon temperature increase from 150 to 280 K for the sample in the trehalose matrix.

The most probable values of the dipole frequency for the studied systems and the distance between electron spins (separated charges) were found according to Eq. (2), based on the obtained Fourier spectra. Figure 5 shows that at 150 K, the Fourier spectra maxima in trehalose and glycerol coincide.

Fig. 5. Fourier spectrum of P_{700}^+ ESEEM signals obtained for the $P_{700}-A_1$ core PSI complex (a) and its main peak (b). The spectra are shown for the complexes in a water–glycerol solution at 150 K (black curve), in the trehalose matrix at 150 K (red curve), and in the trehalose matrix at 280 K (blue curve).

To verify the obtained estimates, the echo signal modulation curves *V* were calculated according to the following equation (see [23, 31]), based on the distance between the spins:

$$
V(\tau) = 1 - p + p \langle \cos(\omega_{d-d}\tau) \rangle, \tag{3}
$$

where $\langle \cos(\omega_{d-d} \tau) \rangle = \frac{1}{2} \int \cos(\Omega(1-3\cos^2\theta) \sin\theta \, d\theta)$ is the mean dipole–dipole interaction orientation, and *p* is the probability of spin flip by the microwave pulse at the time moment τ (see Fig. 2). 2

The obtained distances between the cation and anion in the $P_{700}^+A_1^-$ ion-radical pair for various PSI complexes and temperatures are presented in the table. The distances for intact PSI and $P_{700} - F_X$ core complex lacking terminal 4Fe–4S clusters F_A/F_B taken from our previously published works [23, 24], are also provided.

Analysis of the data for the $P_{700}-A_1$ core presented in the table leads to the two observations. Firstly, an increase in the temperature by 130 K caused an increase in the distance between the cofactors by only 0.1 Å. Secondly, the distance between the separated charges $\rm P_{700}^+$ and $\rm A_1^$ in the $P_{700}-A_1$ core complexes calculated from the modulation frequency of the observed ESEEM signal, has an intermediate value between the distances for $P_{700}^+A_{IA}^-$ and $P_{700}^{+}A_{1B}^{-}$ pairs in intact PSI.

When interpreting the results obtained for the P_{700} –A₁ core complexes, we should take into account that the observed ESEEM signal modulation frequency may be lower than the frequency of the dipole–dipole interaction according to the Eq. (1), due to the low-frequency shift caused by the spin–lattice relaxation of the cofactor A_1^- spins in intact PSI (as shown previously by us in [23, 24]). In this case, a significant contribution to the spin–lattice relaxation is made by the interaction with paramagnetic iron–sulfur clusters. In the previous studies [23, 24], we estimated the distances between the separated charges for the $P_{700}-F_A/F_B$ and $P_{700}-F_X$ core complexes as \sim 26 Å and \sim 25 Å, respectively, by taking into account the above-mentioned low-frequency ESEEM

frequency shift. Comparing these results with the X-ray data (Fig. 1), we came to the conclusion that the charge separation occurs predominantly along the *A* branch in the $P_{700}-F_A/F_B$ complexes, while in the $P_{700}-F_X$ core complexes, the separation occurs along the *B* branch of the PSI RC.

In this work, we studied the $P_{700}-A_1$ core complexes, in which the rate of spin–lattice relaxation of separated charges is much lower than the spin–lattice relaxation rate for the $P_{700}-F_A/F_B$ and $P_{700}-F_X$ core complexes of PSI. Therefore, in the case of $P_{700}-A_1$ core complexes, the low-frequency shift of the modulation frequency due to the spin–lattice relaxation is negligibly small, and the ESEEM frequency in this system is equal to the dipole frequency (1). The echo signal modulation frequency observed for the $P_{700} - A_1$ core can be represented as the average value between the signal modulation frequencies in the ESEEM experiments for the intact $P_{700}-F_A/F_B$ and $P_{700}-F_{X}$ core complexes (here, the modulation frequency for the intact RC was obtained at $T = 150$ K, when the low-frequency shift of the dipole modulation frequency is negligible):

$$
\Omega_{\text{mean}} = f^*(2.98 \text{ MHz}) + (1 - f)^*(3.35 \text{ MHz}). \quad (4).
$$

To interpret this phenomenon, we have to assume that both *A* and *B* cofactor branches of the RC are involved in the charge separation in the $P_{700}-A_1$ core complexes, with dipole frequencies of 2.98 MHz and 3.35 MHz, respectively. From empirical modeling, it can be concluded that Fourier transform of the observed signal of the separated charges in the PSI RC can produce one broadened bell-shaped curve instead of a curve with two resolved maxima provided that following two conditions are met. Firstly, the phase relaxation times of spins should be ≤ 1 μs, and, secondly, the distances between the spins in two subensembles of pairs should differ by no more than 1 Å. Both these conditions are valid for the putative signals of the $P_{700}^+A_{1A}^-$ and $P_{700}^+A_{1B}^-$ ion-radical pairs. In this case, it can be expected that the observed

Modulation frequencies of the observed echo signal and the corresponding distances between the separated charges P_{700}^{+} and A₁ for the PSI complexes in the dried trehalose matrix

PSI complex type		Modulation frequency, MHz		
		calculated from ESEEM signal according to (3)	calculated from Fourier spectrum of ESEEM signal	Distance, A
$P_{700}-F_A/F_B$			2.98	26 ± 0.1 [23]
$P_{700} - F_X$ core			3.35	25 ± 0.1 [24]
P_{700} -A ₁ core	150 K	3.15	3.17	25.5 ± 0.1
	280 K	3.10	3.11	25.6 ± 0.1

ESEEM signal is the sum of the contributions of two RC subensembles with separated charges in the *A* and *B* branches [see Eq. (3)]:

$$
V(\tau) = 1 - p + p * (f \cos(2\pi * 2.98 MHz * \tau) ++ (1 - f) \cos(2\pi * 3.35 MHz * \tau)),
$$
 (5)

where p is the probability of spin flip by the microwave pulse at time moment *τ* [μs].

Two cases can be recognized when the observed signal evolution in time according to (5) leads to the appearance of oscillations with the average frequency (4). Firstly, for arbitrary statistical weights of the above-described subensembles [an arbitrary value of *f* in (5)], the average frequency (4) appears if the electron can jump quickly enough between A_{1A} and A_{1B} , namely, the frequency of electron jumps should be equal to the modulation frequency difference, i.e., $3.35 - 2.98 = 0.37$ MHz. We consider this scenario unlikely.

Secondly, Eq. (5) results in the modulation frequency (4) for the case of equiprobable contribution of *A* and *B* branches to the charge separation in the RC, e.g., in the case $f = 1/2$. The modulation frequency is also equal to the average frequency (4) as a good approximation when the two branches have comparable contributions to the charge separation. We believe that this is the situation realized in the system under study. To explain the decrease in the modulation frequency with increasing temperature (see table), we must assume that the fraction *f* of the separated pairs of charges in the *A* branch should increase upon heating: $f = 0.49$ at T = 150 K, $f = 0.65$ at T = 280 K.

DISCUSSION

The distance between the spin density centers of the ion-radical pair $P_{700}^+A_1^-$ in the $P_{700}^-A_1$ core PSI complexes from the cyanobacterium *Synechocystis* sp. PCC 6803 (25.5 Å) almost coincides with the average value between the distances in the pairs $P_{700}^+A_{1A}^-$ (~26 Å) and $P_{700}^+A_{1B}^ (-25 \text{ Å})$, which corresponds to the symmetrical electron transfer from the Chl dimer P_{700} along the redox cofactor branches *A* and *B*. Analogous conclusion was made earlier on the basis of less direct experiments [12, 20]. Similar estimates of the distance in the $P_{700}^+A_{IA}^-$ pair (26 Å) were obtained earlier for intact PSI complexes and PSI complexes containing high-potential quinone derivatives in the A_1 site [16]. Similar estimates for the distances between the spin density centers of $P_{700}^+ A_1^-$ radical ion pairs (26 and 25 Å for branches *A* and *B*, respectively) were also obtained in the study of PSI mutants with the blocked transfer along one of the branches [17] for the complexes in a water–glycerin mixture at low temperature.

In [32], the spectra of the ion-radical pairs assigned by the authors to $P_{700}^+A_{1A}^-$ and $P_{700}^+A_{1B}^-$ states were recorded

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using high-frequency time-resolved EPR spectroscopy (frequency, 130 GHz) for deuterated PSI complexes isolated from the cyanobacterium *Synechococcus leopoliensis*. Similar to the earlier works conducted in PSI from the cyanobacterium *Synechocystis* sp*.* PCC 6803 [12, 16, 17, 33], the authors, who used intact $P_{700} - F_A/F_B$ complexes, observed formation of the $P_{700}^+A_{1A}^-$ pair only at low temperatures. In authors' opinion, this result could be explained by the fact that under the used experimental conditions the forward electron transfer to the 4Fe–4S clusters proceeded rapidly and predominantly along the *B* branch and the period in a series of laser pulses was too short for the backward electron transfer to occur. As a result, only the $P_{700}^+ A_{IA}^-$ pair was observed. However, the electron transfer exclusively along the *A* branch was previously found in the PSI complexes containing high-potential naphthoquinone derivatives in the A_{1A} and A_{1B} binding sites [16]. Since electron transfer to the 4Fe–4S clusters was not observed in these PSI complexes, it was concluded that the predominant electron transfer to A_{1A} was not caused by the rapid and almost irreversible transfer along the *B* branch to the 4Fe–4S clusters, but was related to a higher redox potential value of A_{1A} as compared to A_{1B} . It was shown [32] using $P_{700} - A_1$ score complexes that the resulting EPR spectrum can be approximated by the sum of two spectra characteristic of $P_{700}^{+}A_{1B}^{-}$ and $P_{700}^{+}A_{1A}^{-}$ states at a ratio of 56 : 44.

In the P₇₀₀ $-A_1$ core complexes from *Synechocystis* sp. PCC 6803, the contribution ratio of recombinations of the $P_{700}^+A_{1A}^-$ and $P_{700}^+A_{1B}^-$ pairs in solution at room temperature was 60 : 40 in favor of the *A* branch [19]. In the dried trehalose matrix, however, this ratio was inverted to 60 : 40 in favor of the *B* branch. In the $P_{700} - F_X$ core complexes lacking terminal 4Fe–4S clusters F_A/F_B and PsaC, the ratio between the branches *A* and *B* was 65 : 35 in favor of the *A* branch, while in the trehalose matrix, it was 75 : 25 in favor of the *B* branch. Charge recombination from phylloquinone acceptors A_{IA}^- and A_{IB}^- was observed in 40-45% of RCs in these samples due to the fact that the functionally active F_x cluster was present only in 55-60% of RCs in this fraction [19].

In another study [12], the temperature dependence of the recombination kinetics between A_1^- and P_{700}^+ was measured in the PSI complexes isolated from the mutant strain *RubA* that lacked the 4Fe–4S clusters and were functionally equivalent to the $P_{700} - A_1$ core preparations. It was shown that at room temperature, the ratio of the contributions of the electron transfer along the redox cofactor branches was 60 : 40 in favor of the *A* branch; when the temperature was lowered to 150 K, the contributions of the *A* and *B* branches became equal (50 : 50). In the $P_{700}-F_{X}$ core complexes, the ratio of backward electron transfer along the *A* and *B* branches was 60 : 40 at room temperature, but increased to 77 : 23 in favor of the *A* branch at 150 K [20].

Therefore, our earlier evaluation of the charge recombination kinetics of the ion-radical pairs $P_{700}^{+}A_{IA}^{-}$ and $P_{700}^+A_{IB}^-$ revealed that the backward electron transfer in the $P_{700}-A_1$ core complexes occurs along both cofactor branches with an approximately equal contributions $(50 \pm 10\%)$, at room temperature, as well as in a water– glycerol solution at 150 K and at room temperature in a dried trehalose matrix. At the same time, in the $P_{700}-F_X$ core complexes, the contribution of the *A* and *B* branches was (62.5 ± 2.5) : (37.5 ± 2.5) at room temperature in solution and increased to 75 : 25 in favor of the *A* branch at 150 K and in dried trehalose matrix. These results show that the presence of the F_x cluster promotes asymmetric electron transfer along the *A* branch, and that the degree of asymmetry increases both with decreasing temperature and upon drying in a glassy trehalose matrix, i.e., when the conformational mobility of the protein is restricted.

The question arises on why the distance between the spin density centers in the $P_{700}^+A_1^-$ ion-radical pair measured in the $P_{700}-A_1$ core complexes using ESEEM, reflects an approximately equal contribution of the *A* and *B* branches, while in the $P_{700} - F_X$ core complexes, this distance corresponds to the predominant formation of the $P_{700}^{+}A_{IB}^{-}$. There are several possible explanations for this experimental observation:

1) Electron transfer in the $P_{700} - F_X$ core complexes occurs along the *A* branch, but under conditions of limited protein mobility (low temperature, drying in a trehalose matrix) and at a laser flash frequency of 10 Hz, accumulation of the reduced acceptors $F_{\rm X}$ and $A_{\rm IA}^-$ takes place, as these acceptors do not have enough time to recombine with P_{700}^+ . At the same time, the backward electron transfer from the lower-potential A_{IB}^- to P_{700}^+ occurs faster (characteristic time, 10-20 μs); as a result, the distance characteristic of the $P_{700}^+A_{1B}^-$ pair is registered in the ESEEM experiment.

2) The balance of the redox potentials (E_m) of A_{1A} and A_{1B} changes in these complexes in the trehalose matrix, making electron transfer to quinone A_{1B} more thermodynamically preferable than to A_{1A} .

3) Drying in a trehalose matrix disrupts the cluster of intraprotein water molecules near the A_{1A} site (but not near the A_{1B} site), since this cluster is less isolated from the protein outer surface [19].

There is also the question why is the distance corresponding to the ion-radical pair $P_{700}^+ A_{1A}^-$ measured in intact $P_{700}-F_A/F_B$ complexes? We assume that the electron transfer in these complexes, as well as in the $P_{700}-F_X$ core, proceeds predominantly along the *A* branch, but the accumulation of the reduced A_{IA}^- state does not occur under experimental conditions with successive laser flashes, due to a larger capacity of the acceptors (three 4Fe–4S cluster versus one in the $P_{700} - F_X$ core complexes).

Further experiments with the PSI containing different number of 4Fe–4S clusters and exogenous quinones of various structure and redox potentials inserted into the A_{1A} and A_{1B} sites can test the validity of these hypotheses.

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Contributions. M. D. Mamedov and G. E. Milanovsky isolated PSI complexes; A. A. Sukhanov and K. M. Salikhov recorded EPR spectra; A. Yu. Semenov developed the concept, supervised the study, and wrote the article; M. D. Mamedov, G. E. Milanovsky, A. A. Sukhanov, K. M. Salikhov, and A. Yu. Semenov discussed the results and edited the manuscript.

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