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Computational investigation of magnetohydrodynamics convective heat transfer in I‑shaped wavy enclosure considering various shapes of inner bodies flled with nanofuid–porous layers

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Abstract

The present work examines numerically the inclined magnetic feld on thermogravitional heat transfer in a novel I-shaped enclosure flled partially with nanofuid in the left layers and flled by partially by porous medium saturated by the same nanofuid using fnite element method. Three diferent shapes of inner bodies had been embedded in the enclosure. The enclosure is partially wavy from its vertical walls with four diferent cases of multi-inner bodies of various shapes such as case 1, 2, 3 and 4 represent circular, square, rhombus and triangular in order to examine their impact on heat transfer and fuid fow. Also, the influence of nanofluid loading, Rayleigh number $(10^4 \leq Ra \leq 10^6)$, Darcy number $(10^{-5} \leq Da \leq 0.1)$, Hartmann number ($0 \le Ha \le 60$), MHD angle ($0^\circ \le \gamma \le 90^\circ$) along with the number ($1 \le No \le 3$) and position ($0.3 \le Y \le 1.3$) of inner hot bodies had been examined in terms of streamlines, isotherms and Nusselt number. The results indicate that the number of inner body and its position along with its shape infuence on the heat transfer rate. It is obtained that Nusselt number for *Case*1 *> Case*3 *> Case*2 *> Case*4. Also, movement the inner hot body from bottom to the top leads to an obvious reduction in the Nusselt number. The increasing of magnetic field angle from $\gamma = 0^\circ$ into $\gamma = 30^\circ$ leads to decreases the heat transfer rate while more increasing of magnetic feld angle augments the rate of heat transfer. Finally, increasing the number of inner hot bodies leads to reduce the total Nusselt number. Thus, for better heat transfer augmentation it is recommended to locate the inner hot body at $Y = 0.3$ and $N\sigma = 1$.

Keywords Magnetic feld · Natural convection · Nanofuid · Porous medium · Wavy enclosure

Abbreviations

- C_p Specific heat at constant pressure (KJ/kg K)
- *g* Gravitational acceleration (m/s^2)
- *k* Thermal conductivity (W/m K)
- R Radius differences of inner and outer cylinder cavity (m)
- R_0 Base circle (m)
- P Dimensionless pressure

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- *p* Pressure (Pa)
- Pr Prandtl number (ν_f/α_f)
- Ra Rayleigh number $(g\beta_f L^3 \Delta T / v_f \alpha_f)$
T Temperature (K)
- Temperature (K)
- T_c Temperature of the cold surface (K)
T_h Temperature of the hot surface (K)
- T_h Temperature of the hot surface (K)
No Number of inner hot bodies
- Number of inner hot bodies
- Nu Local Nusselt number on the hot inner cylinder
- AR Aspect ratio
- U Dimensionless velocity component in x-direction
- *u* Velocity component in x-direction (m/s)
- *V* Dimensionless velocity component in y-direction
- *v* Velocity component in y-direction (m/s)
- X Dimensionless coordinate in horizontal direction
- x Cartesian coordinates in horizontal direction (m)

Y Dimensionless coordinate in vertical direction
- Dimensionless coordinate in vertical direction
- y Cartesian coordinate in vertical direction (m)
- Gr Grashof number

Greek symbols

- α Thermal diffusivity (m²/s)
- *θ* Dimensionless temperature (T-Tc/ΔT)
- Ψ Dimensional stream function (m²/s)
- ψ Dimensionless stream function
- ϕ Nanofluid volume fraction
- *μ* Dynamic viscosity (kg s/m)
- ν Kinematic viscosity (μ/ρ)(Pa s)
- β Volumetric coefficient of thermal expansion (1/K)
- ρ Density (kg/m³)

Subscripts

- c Cold
- bf Base Fluid
- *γ* Inclination angle of magnetic field
- h Hot
- na Nanofuid

Introduction

The natural convection within confned enclosures of different shapes flled partly by fuid and partly with porous medium saturated by the same fuids had been take a lot of numerical and experimental investigation for its importance applications such as heat exchangers, fuel cell, solar collectors, cooling of electronic equipment in addition to the nuclear thermal systems, drying processes, fuid fow of geophysics, pollution of ground water, etc. There are serious problem in augmentation of the heat transfer using of the traditional fuids which is due to low thermal conductivities so the researchers introduced the concepts of nanofuid which is a promoting techniques in the augmentation of heat transfer. Besides, corrugating the wall of the enclosure is recommending technique in improving the thermal performance of the system. So, here a full—description to the most important studies that covers these techniques will be reviewed comprehensively.

The nanofuid flled regular and simple shapes of enclosure had been studied abundantly. For example, Ghasemi et al. ([2011](#page-19-0)) studied the infuence of magneto-hydrodynamics applied uniformly in a horizontal direction on natural convection in a square enclosure. They examined the infuence of Hartmann number, Rayleigh number in addition to the nanofuid loading on heat transfer rate and behavior of fuid fows. The results indicate that there would be an improving in the rate of heat transfer as the Hartmann number decreases along with increasing of Rayleigh number. While the heat transfer may be increases or decreases when the nanofuid is adding will be highly infuenced by Rayleigh and Hartmann number. Izadi et al. [\(2014](#page-19-1)) utilized fnite volume method to examine the infuence of diferent heat source and sink position on the mixed convection within a square

enclosure considering Al_2O_3 -water nanofluid filled the space. Control volume fnite element formulation had been utilized by Bouhalleb and Abbassi [\(2015\)](#page-19-2) to analyze the transient thermogravitional fow in an inclined rectangular enclosure. It heated sinusoidally from its left sidewall while the right sidewall kept at cold temperature. The infuence of aspect ratio, angle of enclosure inclination in addition to the nanofuid loading had been discussed in full-details. They found that increasing or decreasing the aspect ratio leads to diferent behaviors on Nusselt number which the latter increases as nanofuid loading increases. Izadi ([2020\)](#page-19-3) and Izadi et al. [\(2020\)](#page-19-4) studied the natural convection within isosceles triangular enclosure filled by Al_2O_3 –Cu/water hybrid nanofuid with porous medium. The two-phase mixture approach and Darcy–Brinkman model had been used to simulate the nanofuid and porous medium, respectively. The natural convection within triangular enclosure considering the Brownian motion along with magnetic feld had been studied by Ghasemi and Aminossadati [\(2010\)](#page-19-5), Mahmoudi et al. ([2012\)](#page-20-0), Ghasemi and Aminossadati ([2010](#page-19-5)) and Mahmoudi et al. ([2012\)](#page-20-0). Also, Saleh et al. ([2011\)](#page-20-1) examined the problem of natural convection within trapezoidal enclosure filled by Al_2O_3 -water and Cu–water nanofluids using fnite diference method. They developed new correlations of Nusselt number. The results showed that Cu–water nanofuid is recommended in heat transfer more than Al_2O_3 . Mehryan et al. ([2020a](#page-20-2), [b\)](#page-20-3) studied the natural convection within a trapezoidal enclosure with fexible partition using fnite element formulation. It had been indicated that the best heat transfer rate was at zero inclination angle and reduces with increasing the inclination angle of the trapezoidal enclosure. Hussein & Mustafa (2017) studied the thermally driven flow in a parallelogrammic enclosure open from its top wall and flled with Cu–water nanofuid using fnite volume method. The enclosure is heated partially from its bottom wall while the tilted wall is maintained at cold temperature. The selected dimensionless parameters of this study are Rayleigh number, position of the heat source, inclination angle of the two cold walls of the enclosure in addition to the nanofluid concentration. The results show that increasing nanofuid loading along with the dimensionless value of Rayleigh number improves noticeably the Nusselt number. Also, for better heat transfer it is better to locate the heat source closed to the left sidewall with tilting angle of 60°.

The nanofuid flled complex shapes of enclosure investigated by many researchers. For example, Cho et al. ([2013\)](#page-19-7) utilized fnite volume formulation to examine the infuence of different types of nanofluids particles such as $Cu, A₂O₃$ and $TiO₂$ filled complex wavy enclosure on natural convection as well as entropy generation. Additionally, the authors investigated the infuence of wavy surface amplitude and length in addition to Rayleigh number on fluid flow and heat transfer. The results indicate that Cu nanoparticle reveals the best performance in terms of heat transfer and lowest generation of entropy. Sadeghi et al. ([2020](#page-20-4)) studied the natural convection in a novel-complex enclosure shape filled with Al_2O_3 nanofuid with internal heat generation along with inclined magnetic feld using fnite element scheme. The enclosure includes trapezoidal heater and wavy cold wall. The infuence of Hartmann and Rayleigh numbers, concentration of nanoparticle and its shapes, trapezoidal heater location and parameter of internal heat generation on fluid flow and heat transfer were examined. The results show that the impact of heat generation at low value of Rayleigh number was so signifcant. Also, applying of horizontal magnetic feld can suppress the natural convection heat transfer.

The nanofuid saturated with porous medium studied a lot using various models like LTE, LTNE, Darcy model, Darcy–Brinkmann model etc. For example, Ahmed and Rashed [\(2019\)](#page-19-8) illustrated the MHD buoyancy driven fow in a wavy enclosure flled by nanofuid and porous medium along with internal heat generation using fnite diference approach. Izadi et al. ([2019\)](#page-19-9) used fnite element method to study the non-uniform magnetic feld applied from two semi-circular hot cylinders. The gap flled by hybrid nanofuid and porous medium considering the local thermal non-equilibrium model under various selected parameters like the magnetic source power ratio, coefficient of porosity, Hartmann and Rayleigh number. Kadhim et al. [\(2020](#page-19-10)) studied the natural convection in a wavy enclosure using fnite element method. The enclosure was flled by two diferent layers, the right layer filled by $Cu - Al₂O₃$ hybrid nanofluid while the left layer flled by porous medium saturated by the same hybrid nanofuid. Darcy–Brinkmann model had been used to model the porous layer. The researchers examined the infuence of Rayleigh number, porous layer width, Darcy number, number of undulations, inclination angle and the nanofuid loading. The results showed that using of hubrid nanofuid reveals better augmentation in Nusselt number in a comparison with Al_2O_3 -water nanofluid. Besides, Rayleigh and Darcy number were strongly infuenced on the Nusselt number. Mehryan et al. ([2020a,](#page-20-2) [b](#page-20-3)) used LTNE to simulate the natural convection within Ag–MgO/water nanofuid in porous square enclosure. The porous medium had been treated via Darcy model.

The existence of inner body considering various shapes within diferent shapes of enclosure had been studied considering the infuence of inner body size, position, number and shapes etc. For example, the infuence of inner body locations that moved in vertical, horizontal and diagonal direction had been studied by Abdulkadhim ([2019](#page-19-11)), Hussain and Hussein ([2010\)](#page-19-12), Kim et al. [\(2008](#page-20-5)), Lee et al. [\(2010\)](#page-20-6) and Majdi et al. ([2019\)](#page-20-7). Also, Yoon et al. ([2012\)](#page-20-8) studied utilizing immersed boundary scheme which is based upon formulation of fnite volume approach the natural convection between two inner circular cylinder immersed in a square enclosure considering the infuence of various size and Rayleigh number. The results showed that reduction of the size of the inner circular cylinders leads to increasing of the dependency on Rayleigh number. Kefayati and Tang [\(2018](#page-19-13)) used lattice Boltzmann scheme to study the infuence of one cylinder of circular and elliptical shape on natural convection within square enclosure under various inclination angle. They examined the infuence of various inner body position and size under diverse Rayleigh number values. Roy ([2018](#page-20-9)) studied using fnite diference method the natural convection between three diferent shapes of inner shapes (circular, elliptical and rectangular) immersed in a nanofuid square enclosure. The results indicate that using of inner circular cylinder reveals better rate in thermal heat transfer more than rectangular and elliptical shapes. Wang et al. [\(2021\)](#page-20-10) studied the transient convection considering various inner circular body positions moved in the vertical direction immersed in square enclosure. The results indicate that the inner body position is highly afects on the fuid fow behavior and heat transfer feld. Bhowmick et al. ([2020](#page-19-14)) studied the infuence of magnetic feld on entropy generation and natural convection of square enclosure with pair of circular cylinders located immersed in a porous medium. They illustrate the impact of distance between the inner cylinders along with Darcy, Hartmann and Rayleigh number. The results displayed that the thermal rate of heat transfer augmented as there was an increasing in the dimensionless distance between the inner circular cylinders. Yan et al. [\(2020](#page-20-11)) used fnite element method to study the natural convection within square enclosure flled by nanofuid and porous medium considering inner hot elliptical body under inclined magnetic feld. Alsabery et al. ([2020](#page-19-15)) studied the location and the size of inner solid body of squared shape on the rate of heat transfer utilizing two-phase nanofuid model. Shehzad et al. ([2021\)](#page-20-12) studied the infuence of diferent inner hot and cold pipes with a fin attached to them on the natural convection within square heat exchangers.

The natural convection within more complex shapes like C-shape, H-shape, U-shape and especially I-shaped had been received little attention despite their importance in engineering applications (Armaghani et al. [2020;](#page-19-16) Keramat et al. [2020;](#page-20-13) Ma et al. [2019a](#page-20-14), [b,](#page-20-15) [2020;](#page-20-16) Malekpour et al. [2018](#page-20-17); Mohebbi et al. [2017](#page-20-18)). For example, Armaghani et al. ([2020\)](#page-19-16) studied the infuence of inclined magnetic feld along with heat generation in I-shaped enclosure flled by nanofuid and porous medium. Also, Ma et al. [\(2020](#page-20-16)) studied the inclined magnetic feld within I-shaped enclosure with multi-square inner bodies considering nanofuid flled the gap.

Thus, it can be seen and according to the best authors' knowledge that there are serious limitations in the complex shapes of enclosures especially the I-shapes enclosure despite its applications in the shell and tube heat exchangers industry. The I-shaped enclosure is a combination of more than one enclosure which is the real case in the industrial engineering projects. Besides, there is no study up to date that discuss the inclined magnetic field in I-shaped enclosure with wavy walls flled by multi-layers of nanofuid and porous medium considering four diferent cases of inner bodies (circular, square, rhombus and triangular) using fnite element method. In this way, this paper will be the frst attempt to solve this problem under wide range of Rayleigh number, Darcy number, Hartmann number in addition to the infuence of diferent thermal management of the inner bodies.

Mathematical model

The computational CFD model used in the present study will be presented here considering many assumptions. Firstly, the physical model is presented in Fig. [1](#page-3-0) which consists of a three-dimensional view for the enclosure with the internal circular bodies along with the two-dimensional view of the enclosure that contain the two layers. The right layer consists of nanofuid while the left layer consists of porous medium saturated with the same nanofuid. Four diferent cases of inner body shapes had been included in the present work. Case 1, 2, 3 and 4 represent circular, square, rhombus and triangular shapes, respectively. With respect to the

physical properties

thermal boundary conditions, the top and bottom wall of the enclosure are fully insulated. The right and left wall included the wavy, horizontal and the vertical parts which form the I-shape are kept at isotherm hot and cold temperature, respectively. Also, diferent boundary conditions of the inner body had been selected in order to examine the infuence of the inner body conditions and its position on the heat transfer and fuid fows.

Governing equations

The governing equations of the nanofluid and porous medium under the assumptions inserted below;

- The fuid fow of the two layers within the enclosure are considered to be laminar,
- The left layer which consists of the porous medium had been modeled via Darcy–Brinkman model.
- Local thermal equilibrium between the each layer had been considered
- The infuence of thermal radiation and the internal heat generation, Soret and Dufour are ignoring.
- Only the density is considered to be changed in the momentum equations and treated via Boussinesq approximation while the thermophysical nanofuid properties are assumed to be constant.

In this way and based upon the above mentioned conditions, the fnal form of the governing equations may be written as indicated below (Al-Zamily [2014;](#page-19-17) Hussain and Rahomey [2019](#page-19-18); Sheikholeslami et al. [2013](#page-20-19));

$$
\frac{\partial U}{\partial X} + \frac{\partial V}{\partial Y} = 0\tag{1}
$$

X − *momentum*

$$
U\frac{\partial U}{\partial X} + V\frac{\partial V}{\partial Y} = -C_1\frac{\partial P}{\partial X} + C_2\frac{\mu_{na}}{\rho_{na}\alpha_{bf}} \left(\frac{\partial^2 U}{\partial X^2} + \frac{\partial^2 U}{\partial Y^2}\right) - C_3\frac{\mu_{na}}{\rho_{na}\alpha_{bf}}\frac{U}{Da} + C_1PrHa^2\left(V\sin(\gamma)\cos(\gamma) - U\sin^2(\gamma)\right)
$$
\n(2)

Y − *momentum*

$$
U\frac{\partial V}{\partial X} + V\frac{\partial V}{\partial Y} = -C_1\frac{\partial P}{\partial Y} + C_2\frac{\mu_{na}}{\rho_{na}\alpha_{bf}} \left(\frac{\partial^2 V}{\partial X^2} + \frac{\partial^2 V}{\partial Y^2}\right) - C_3\frac{\mu_{nf}}{\rho_{nf}\alpha_f}\frac{V}{Da}
$$

+ $C_1\frac{(\rho\beta)_{na}}{\rho_{bf}\beta_{na}}Ra\theta + C_1Ha^2Pr\left(Usin(\gamma)cos(\gamma) - Vcos^2(\gamma)\right)$ (3)

Energy

$$
U\frac{\partial\theta}{\partial X} + V\frac{\partial\theta}{\partial Y} = C_4 \left(\frac{\partial^2\theta}{\partial X^2} + \frac{\partial^2\theta}{\partial Y^2}\right)
$$
(4)

Considering the following dimensionless parameter

$$
X = \frac{x}{L}, Y = \frac{y}{L}, U = \frac{uL}{\alpha_f}, V = \frac{vL}{\alpha_f}, P = \frac{pL^2}{n f \alpha_f^2},
$$

$$
= \frac{T - T_c}{T_h - T_c}, Ra = \frac{g_f L^3 \Delta T}{f \alpha_f}, Da = \frac{k_p}{L^2}, Pr = \frac{f}{\alpha_f}
$$
(5)

The coefficients in the Eqs. $2-4$ $2-4$ are as indicated below; Right Layer (Nanofluid Layer): $C_1 = 1, C_2 = 1, C_3 = 0$

and $C_4 = \alpha_{\text{nf}} / \alpha_f$ Left Layer (Nano-Porous Layer): $C_1 = \varepsilon^2$, $C_2 = \varepsilon$, $C_3 = \varepsilon^2$ *and* $C_4 = \alpha_{\text{eff}} / \alpha_f$

The fluid flow distributions within the wavy enclosure are represented in term of stream function as indicated below;

$$
U = \frac{\partial \Psi}{\partial Y}, V = -\frac{\partial \Psi}{\partial X}
$$
 (6)

$$
U_{na} = \frac{\partial \Psi_{na}}{\partial Y}, V_{na} = \frac{\partial \Psi_{na}}{\partial X}
$$
 (7)

Which leads to the dimensionless equation;

$$
\frac{\partial^2 \Psi}{\partial X^2} + \frac{\partial^2 \Psi}{\partial Y^2} = \frac{\partial U}{\partial Y} - \frac{\partial V}{\partial X}
$$
(8)

$$
\frac{\partial^2 \Psi_{na}}{\partial X^2} + \frac{\partial^2 \Psi_{na}}{\partial Y^2} = \frac{\partial U_{na}}{\partial Y} - \frac{\partial V_{na}}{\partial X}
$$
(9)

The thermophysical nanofuid properties are inserted below (Table [1\)](#page-4-2);

$$
\alpha_{na} = \frac{k_{na}}{\left(\rho C_p\right)_{na}}\tag{10}
$$

$$
\rho_{na} = (1 - \phi)\rho_{bf} + \phi\rho_{sp} \tag{11}
$$

$$
\left(\rho c_p\right)_{na} = (1 - \phi)\left(\rho c_p\right)_{bf} + \phi\left(\rho c_p\right)_{sp} \tag{12}
$$

$$
(\rho \beta)_{na} = (1 - \phi)(\rho \beta)_{bf} + \phi(\rho \beta)_{sp}
$$
\n(13)

Table 1 Thermo-physical properties of water and nanoparticles (Al_2O_3) at T = 25 °C (Motlagh and Soltanipour [2017\)](#page-20-20)

Properties	ρ (kg/m ³)	Cp (J/kg K)	k (W/m K)	β (1/K)
Water	997	4180	0.614	2.1×10^{-4}
Copper	3970	765	40	0.85×10^{-5}

As regards the dynamic viscosity, Brinkman correlations have been taken into account in the numerical simulation (Chamkha and Ismael [2013](#page-19-19); Garoosi et al. [2015](#page-19-20); Oztop and Abu-Nada [2008\)](#page-20-21)

$$
\mu_{na} = \frac{\mu_{bf}}{(1 - \phi)^{2.5}}
$$
 (14)

Maxwell approach considered for thermal conductivity (Abdelmalek et al. [2020](#page-19-21); Bessaïh et al. [2017](#page-19-22); Sheremet et al. [2016\)](#page-20-22)

$$
\frac{k_{nf}}{k_{bf}} = \frac{k_{sp} + 2k_{bf} - 2\phi(k_{bf} - k_{sp})}{(k_{sp} + 2k_{bf}) + \phi(k_{bf} - k_{sp})}
$$
\n(15)

The electrical conductivity of the mixture of nanofuid are expressed as indicated below;

$$
\frac{\sigma_{na}}{\sigma_{bf}} = 1 + \frac{3(\frac{\sigma_{sp}}{\sigma_{bf}} - 1)\phi}{(\frac{\sigma_{sp}}{\sigma_{bf}} + 2) - (\frac{\sigma_{sp}}{\sigma_{bf}} - 1)\phi}
$$
(16)

The thermal diffusivity α_{eff} in addition to the thermal conductivity effectiveness k_{eff} are given by:

$$
\alpha_{\text{eff}} = \frac{k_{\text{eff}}}{\left(\rho c_p\right)_{na}}\tag{17}
$$

$$
k_{\text{eff}} = (1 - \varepsilon)k_{\text{sp}} + \varepsilon k_{\text{na}} \tag{18}
$$

Boundary conditions

The boundary conditions in this study are presented in shown in Fig. [1](#page-3-0) and indicated below;

The left walls including the horizontal, vertical and wavy parts are kept at hot temperature:

$$
T_h = 1, U = V = 0
$$

The right wall including the horizontal, vertical and wavy parts are kept at cold temperature:

$$
T_c = 0, U = V = 0
$$

The top and bottom walls are insulated $\frac{\partial T}{\partial Y} = 0, U = V = 0$ The inner bodies are located at hot and cold temperature as indicated in the results.

Also, with respect to the permeable surface's boundary conditions that separate the nanofuid and porous medium may be defned as indicated below;

$$
\theta_{po} = \theta_{na}, \frac{\partial \theta_{na}}{\partial X} = \frac{K_{eff}}{K_{na}} \frac{\partial \theta_{po}}{\partial X}, \Psi_{po} = \Psi_{na}, \frac{\partial \Psi_{po}}{\partial X} = \frac{\partial \Psi_{na}}{\partial X}, \Omega_{po} = \Omega_{na},
$$

$$
\frac{\partial \Omega_{na}}{\partial X} = \frac{\partial \Omega_{po}}{\partial X}, \mu_{po} \left(\frac{\partial U_{po}}{\partial Y} + \frac{\partial V_{po}}{\partial X} \right) = \mu_{na} \left(\frac{\partial U_{na}}{\partial Y} + \frac{\partial V_{na}}{\partial X} \right),
$$

$$
P_{po} = P_{na}, \frac{\partial P_{na}}{\partial X} = \frac{\partial P_{po}}{\partial Y}
$$

Nusset number

Nusselt number had been determined on the inner hot body from the equation inserted below;

$$
Nu_L = \frac{k_{nf}}{k_f} \frac{\partial \theta}{\partial n}, \overline{Nu} = \frac{1}{l} \int_0^l Nu_L(\varphi) d\varphi \tag{19}
$$

With respect to the determination of total Nusselt number for the case of multi-inner hot bodies, it had been predicted by summing all of the Nusselt number at each hot inner body.

Validation and mesh independent study

In order to check the accuracy of the present numerical CFD results, validations with signifcant researchers are presented. First validation is with Malekpour et al. ([2018\)](#page-20-17) results in terms of Nusselt number for I-shaped enclosure flled with nanofuid along with applied magnetic feld uniformly in the horizontal direction. The results are presented in Table [2](#page-5-0) and graphed in Fig. [2](#page-6-0) under wide range of Rayleigh and Hartmann numbers.

It can be seen that the present good agreement with their results. Besides that two diferent cases of validation in terms of streamlines and isotherms had been presented in Figs. [3](#page-6-1) and [4.](#page-7-0) The frst validation presented, in Fig. [3](#page-6-1), with Ma et al. ([2020\)](#page-20-16) for I-shaped nanofuid enclosure with multi-inner body of square shape considering

Table 2 Validation of the present work with Malekpour et al. ([2018\)](#page-20-17)

	Rayleigh number	Present work	Malekpour and Karimi
$Ha=0$	10.000	6.60545	6.988
	$1.00E + 0.5$	13.48916	13.072
	$1.00E + 06$	31.64713	30.599
$Ha = 40$	10.000	6.12548	6.628
	$1.00E + 0.5$	8.61114	8.747
	$1.00E + 06$	25.26759	21.724
$Ha = 80$	10.000	6.08909	6.599
	$1.00E + 0.5$	6.95519	7.329
	$1.00E + 06$	17.73851	17.706

The present work

Fig. 3 Validation in terms of streamlines and isotherms with Ma et al. Ma et al. ([2020\)](#page-20-16)

applied magnetic field under $\phi = 0$, $Ha = 0$, $Ra = 10^6$. Also, in Fig. [4](#page-7-0), A validation with results found by Hussain and Rahomey ([2019\)](#page-19-18) for inner circular body embedded in a square enclosure flled partially by nanofuid and partially by porous medium with the same nanofuid at $Ra = 10^6$, $Da = 10^{-3}$, $\phi = 0$ was presented. Again it can be seen that there is excellence agreement with each study.

With respect to the numerical grid generation and independent study that presented in Fig. [5](#page-8-0). Due to the CPU memory, various number of numerical element generation had been test for high Rayleigh number to check the error that may occur at $Ra = 10^6$, $Da = 10^{-3}$, $\phi = 0.02$, $Ha = 20$. It can be seen that when the number of element are 28,146 leads to the same Nusselt number at number of element are 32,234 so there is no necessity in increasing the number of

Hussain & Rahomey (2019)

The present work

Fig. 4 Validation in terms of streamlines and isotherms with Hussain and Rahomey Hussain & Rahomey ([2019\)](#page-19-18)

elements and the time of numerical solution. That's why we use, in our all calculations, 28,146 elements.

Results and discussion

The present section discuss the numerical results of natural convection heat transfer within I-shaped wavy-walled enclosure flled by nanofuid saturated with porous medium (left layer) and the same nanofuid (right layer) under wide range of Rayleigh number ($10^4 \leq Ra \leq 10^6$), Darcy number (10−⁵ ≤ *Da* ≤ 10−¹), Hartmann number (0 ≤ *Ha* ≤ 60) considering three diferent positions of the hot inner body along with four diferent cases of the shapes of inner body (circle, square, rhombus and triangle) for three diferent number of the hot source inner body ($1 \leq No \leq 3$) and three different positions (0.3 \leq *Y* \leq 1.3). The influence of these mentioned dimensionless parameters on heat transfer and fuid fow are presented in terms of streamlines, isotherms in addition to Nusselt number.

The infuence of dimensionless numbers on fuid fow and heat transfer

Since many dimensionless parameters under investigated, we start by illustrating the infuence of Rayleigh, Darcy and Hartmann numbers along with four diferent shapes of inner body on fuid fow and heat transfer. First of all, the infuence of Rayleigh number and the inner body shape is presented in terms of streamlines and isotherms in Figs. [6](#page-9-0)a, b, respectively.

With respect to Fig. [6](#page-9-0)a, at low Rayleigh number $(Ra = 10⁴)$, where the power intensity of natural convection is low and the conductive mode of heat transfer is the dominant mode, it can be seen that there are just two inner vortices formed closed to each shape of the inner body on the nanofuid region (right layer) while the streamlines on the nano-porous region (left layer) does not contain any circulation which refect the lower strength of the fluid flow. This is due to the porous medium which plays as additional resistance to the fuid fow in addition to the low Rayleigh number. This can be observed from the low **Fig. 5** Mesh generation for the enclosure (top image) and mesh independent study in terms of Nusselt number (bottom image)

stream function value which is influenced by the inner body shape. For example, the strength of fluid flow for Case4 *> Case*3 *>* Case2 *> Case*1 which makes the triangular shape (Case 4) recorded the better case in terms of the best fuid fow strength while the inner circular body gives the lowest strength of the fuid fow.

With increasing the Rayleigh number into $(Ra = 10^6)$, which refect the changing of the heat transfer mode from the conduction mode into the thermogravitational (natural convection) mode leading to increasing the fuid fow strength for each case of the inner body shape. For example, the strength for Case3 *> Case*1 *>* Case2 *> Case*4 which makes the triangular shape reveals the lowest case in terms of fuid fow strength while the rhombus shape give the best fluid flow strength. It is worthy to mention that increasing the Rayleigh number helps more nanofuid to penetrate from the right layer into the left layer due to increasing the permeability of the porous medium which leads to the formation of the inner vortices in many regions within the enclosure at high Rayleigh number.

With respect to Fig. [6b](#page-9-0), which reveals the isotherms under various dimensionless values of Rayleigh number considering four diferent cases of inner body, it can be seen that the isotherms lines at low Rayleigh number $(Ra = 10^4)$ are vertical which refect the conduction heat transfer mode is dominant while at high Rayleigh number $(Ra = 10^6)$, the shapes of isotherms lines obviously changed into the non-uniform curved lines which is an indication of the thermogravitional mode is dominated. It can be seen that the behavior of isotherms are similar regardless the inner body shape.

The infuence of Darcy number and the inner body shape is presented in terms of streamlines and isotherms

in Fig. [6b](#page-9-0) and [7](#page-10-0)a, respectively. Firstly, with respect to the Fig. [7](#page-10-0)a which illustrates the infuence of various Darcy number and diferent four cases of inner body shape on the fuid fow strength. By constriction our analysis on the left layer of the enclosure which contains (porous medium saturated with nanofuid) and following the infuence of Darcy number along with the inner body shape, it can be noted that at low Darcy number $(Da = 10^{-5})$, there is no inner vortices formed while as the Darcy number increases into $Da = 10^{-1}$, more inner vortices formed which means that the fluid flow strength is increasing as Darcy number increases. The physical reason behind this, is that increasing the dimensionless value of Darcy number leads to an obvious increment in the permeability of the porous medium which helps nanofuid to penetrate into the left layer leading to increases the convection heat transfer mode because the conductive mode is dominant at low Darcy number and this is obvious also from the shape of the streamlines in the left layer as there is no inner cells which refect low fuid fow strength.

Also, the shape of inner body had great impact on the fluid flow strength. For example, at low Darcy number $(Da = 10^{-5})$, the maximum stream function of the left layer for Case1 *> Case*2 *>* Case3 *> Case*4, while at high Darcy number $(Da = 10^{-5})$, the maximum stream function of the left layer for Case2 *> Case*4 *>* Case3 *> Case*1. Thus, at low Darcy number Case 1 which represents the inner circular shape had the highest fuid fow strength while the lowest fuid fow strength was when the inner shape is triangular. Actually, at low Darcy number the infuence

of inner body shape had great infuence on the fuid fow characteristics. On the other hand, at high Darcy number Case 1 which is the inner circular shapes recorded the lowest case in the fuid fow strength while the square shape is better in the enhancing the fuid fow strength.

With respect to the Fig. [7](#page-10-0)b, which examines the isotherms under various Darcy number and diferent cases of inner body shape, it can be seen obviously that increasing Darcy number leads to change the isotherm lines from the vertical lines as the conduction heat transfer is dominant at low

Darcy number into curved-horizontal lines at high Darcy number.

With respect to the magnetic field influence on fluid flow and heat transfer which is presented in Fig. [8a](#page-11-0), b. Firstly, the infuence of Hartmann number considering diferent shapes of inner body had been presented in Fig. [8a](#page-11-0). It can be seen that the Hartmann number infuence strongly on the fuid flow contours because increasing Hartmann number leads to reduce the fuid fow strength. For example, for case 1, it can be seen that maximum stream function reduces from $Ψ_{max}$ = 29.6179 at *Ha* = 0, into Ψ_{*max*} = 13.2388 at *Ha* = 60 which makes this behavior is quite reverse to the infuence of Rayleigh number on the strength of the fuid fow. This behavior is similar for all of the other cases of inner body shapes but at diferent rate. For example, in the absence of magnetic field $(Ha = 0)$, it can be seen that the maximum

stream function for diferent cases of inner body shape is Case3 *> Case*1 *>* Case2 *> Case*4 while increasing the Hartmann number into $(Ha = 60)$, leads for further decrement in the maximum stream function which makes it for Case2 *> Case*4 *>* Case1 *> Case*3. The physical reason behind this is that increasing the Hartmann number is an indication of increasing the electromagnetic force which its role is to reduce the natural convection buoyancy force and reduces the fuid fow strength. With respect to Fig. [8b](#page-11-0) which shows that increasing Hartmann number slightly reduces on the isotherms.

The infuence of inner hot body number

The infuence of number of inner hot pipe for each of the four selected cases of the inner body shape (circle, square, rhombus and triangular) on streamlines and isotherm are presented in Fig. [9a](#page-13-0), b, respectively.

With respect to Fig. [9a](#page-13-0), it may be noted that regardless the inner body shape, increasing the number of inner hot body from $N = 1$ into $N = 2$, leads to increasing the fluid flow strength by increasing the maximum stream function. Also, further increasing the number of hot body for $No = 3$, leads to slight reduction the fluid flow strength. For example, regarding Case 1 the maximum stream function increases from $\Psi_{max} = 18.3070$ into $\Psi_{max} = 27.9309$ and then reduces into $\Psi_{max} = 25.0989$ when the number of inner hot body increases from $N_0 = 1, 2$, and 3, respectively. It is worthy to mention, that the inner body shape along with the number of inner hot body had great impact of the strength of fuid fow. For example, when the number of inner hot body is $N_0 = 1$, the maximum stream function for $Case4 > Case2 > Case1 > Case3$ which make the fluid flow strength of Case 4 which represents the triangular shape had the highest fuid fow strength. The infuence of inner body shape for further increasing of inner hot body number is approximately negligible with slight better augmentation in fluid flow strength for Case 1 at $N_0 = 1$.

With respect to the isotherms as shown in Fig. [9](#page-13-0)-b it can be seen that increasing the number of hot inner body leads to increase the hot temperature lines in the left layer (nanofuid saturated with the porous medium) as the hot lines resulted from the inner hot bodies will combine with the hot line resulted from the left wall of the enclosure.

The infuence of position of the inner hot body

The infuence of three diferent inner body position which they are based on the position of the inner hot body are bottom when the distance between the position of the inner hot body and the base of the insulated enclosure wall is $Y = 0.3$, in the same manner, center $Y = 0.8$ and top $Y = 1.3$ on streamlines and isotherms are presented in Fig. [10a](#page-14-0),

b, respectively. Firstly, with respect to the streamlines which is presented in Fig. [10a](#page-14-0), so for Case 1 which indicate the circular shape of inner body by following its position when moved the inner hot body up and down, that the fuid fow strength because the maximum stream function increases from $\Psi_{max} = 17.0935$ when it located at $Y = 0.3$ into $\Psi_{max} = 24.1271$ when it located at $Y = 0.8$. However, more upward movement leads to a reduction in the fuid flow strength into $\Psi_{max} = 16.1245$ at $Y = 1.3$. This behavior is repeated to Case 2 and Case 3 except Case 4 because the fuid fow strength decreases as the inner hot triangular body moved upward. It worthy to mention that the infuence of inner hot body shape is negligible on fuid fow strength when it is located at the center and top of the enclosure while the infuence of inner body location when it located at the bottom of the enclosure is strong especially for the Case 4 which is the triangular shape.

With respect to Fig. [10b](#page-14-0) which shows the isotherms clearly infuenced by the position of the inner body shape.

Nusselt number

The Nusselt number is the most important parameters in the augmentation of heat transfer rate. So the infuence of the mentioned parameters on Nusselt number will be discussed in full details to draw the major conclusions.

The infuence of Rayleigh number considering diferent inner body shape on average Nusselt number along the inner hot body had been displayed in Fig. [11.](#page-15-0) It can be seen that increasing Rayleigh number leads to an increasing in the Nusselt number because of increasing the fuid flow strength and natural convection rate which increasing the rate impact between the fuid molecules which increasing the thermal energy transferred from the hot inner body into the enclosure area. This increasing of Rayleigh number changed the mode of heat transfer into the thermogravitional (natural convection) mode. It is also noted that the inner body shape plays an important role in the augmentation of Nusselt number which leads to improving the thermal rate of heat transfer. For example, Nusselt number for *Case*1 *> Case*3 *> Case*2 *> Case*4.

With respect to the infuence of Darcy number as illustrated in Fig. [12](#page-15-1) which shows obviously that increasing Darcy number improves Nusselt number because increasing Darcy number leads to augmentation of natural convection and buoyancy force which changed the mode of heat transfer from the conductive mode into the convection. Again, Nusselt number for *Case*1 *> Case*3 *> Case*2 *> Case*4.

Besides, increasing Hartmann number leads to express more resistance of the MHD force to the movement of the fuid fow which leads to a reduction of Nusselt number as shown in Fig. [13](#page-15-2) which makes its infuence is quite reversible to the infuence of Rayleigh and Darcy number.

Fig. 9 a Streamlines contours considering diferent number of inner hot shapes at $Ra = 10^6$, $\varphi = 0.02$, $Da = 0.001$, $Ha = 20$ **b** Isotherms contours considering diferent number of inner hot shapes at $Ra = 10^6$, $\varphi = 0.02$, $Da = 0.001$, $Ha = 20$

a

 $\Psi_{max} = 18.3070$ $\Psi_{max} = 27.9309$ $\Psi_{max} = 25.0989$

 $\Psi_{max} = 26.5025 \qquad \qquad \Psi_{max} = 27.5546 \qquad \qquad \Psi_{max} = 25.7497$

b

Fig. 10 a Streamlines contours considering diferent position of inner hot shapes at $Ra = 10^6$, $\varphi = 0.02$, $Da = 0.001$, $Ha = 20$ **b** Isotherms contours considering diferent position of inner hot shapes at $Ra = 10^6$, $\varphi = 0.02$, $Da = 0.001$, $Ha = 20$ **a**

b

 $\Psi_{max} = 26.5024$ $\Psi_{max} = 23.8868$ $\Psi_{max} = 16.0134$

Fig. 11 Nu with respect to Ra for diferent shapes of inner body at $\varphi = 0.02, Da = 0.001, Ha = 0$

Fig. 12 Nu with respect to Da for diferent shapes of inner body at $\varphi = 0.02, Ra = 10^5, Ha = 0$

It may be noted that infuence of Hartmann number on Nusselt number is highly afected by the inner body shape of each case. For example, the reduction in Nusselt number for *Case*1 *> Case*3 *> Case*2 *> Case*4 which makes Case 4 is the best solution to reduce the heat transfer reduction in a comparison with the other cases.

Also, the infuence of MHD angle on the Nusselt number is presented in Fig. [14](#page-15-3). Diferent behaviors of Nusselt

Fig. 13 Nu with respect to Ha for diferent shapes of inner body at $\varphi = 0.02, Ra = 10^5, Da = 0.001$

Fig. 14 Nu with respect to MHD Angle for diferent shapes of inner body at $\varphi = 0.02$, $Ra = 10^5$, $Da = 0.001$, $Ha = 60$

number is obtained based upon the angle of magnetic feld inclination angle. When $0° \le \gamma \le 30°$, it is observed that increasing the inclination angle of magnetic feld leads to decreasing in the Nusselt number for all of the inner body shapes. When $\gamma \geq 30^{\circ}$, it can be seen that the influence of angle of MHD is positive on the Nusselt number as the latter

Fig. 15 Nu with respect to nanofuid loading for diferent shapes of inner body at $\varphi = 0.02$, $Ra = 10^5$, $Da = 0.001$

increases as the magnetic feld angle increases for all of the inner body shape. It can be seen that Nusselt number for *Case*1 *> Case*3 *> Case*2 *> Case*4 for the entire magnetic feld angle of inclination.

The influence of nanofluid loading is showed in Fig. [15](#page-16-0) for various inner body shapes. It can be seen that Nusselt number increases as the nanofluid loading increases. It can be seen again that the inner body shape plays an important role in the enhancement of Nusselt number which Case 1 indicate the best case in the augmentation of heat transfer while the Case 4 record the lowest improving in Nusselt number.

The impact of number of inner hot bodies on the Nusselt number is presented in Fig. [16](#page-17-0). It can be seen that for all numbers of inner bodies and various shapes, there would be an increasing of Nusselt number with increasing of Rayleigh number. However, it can be seen that Nusselt number is in its highest value when the number of inner heater bodies is $N_0 = 1$ for all shapes (circle, square, rhombus and triangle). It is observed that when the number of inner bodies are $N_0 = 3$, leads to reduce Nusselt number for all of the four cases of inner body shapes. It is noted that at low Rayleigh there is no change on Nusselt when the number of inner hot bodies for all cases are $N_0 = 1, 2$. While the deviation (difference) in Nusselt number for $N_0 = 1, 2$ increases as Rayleigh number increases.it is obtained that Case 1 reveals better augmentation of Nusselt number for all cases.

Finally, with respect to the influence of the inner hot body location on Nusselt number is discussed in Figs. [17](#page-18-0) and [18.](#page-19-23) It can be seen that when the inner hot body start moving upwards, there would be a reduction in Nusselt number for all of the shapes of inner bodies. Besides that, the Nusselt number for *Case*1 *> Case*3 *> Case*2 *> Case*4 in any position of inner bodies.

Conclusion

- 1. Increasing Ra, Da, nanofuid loading and reducing the Ha leads to augmentation of heat transfer
- 2. Best location in terms of best strength of fuid fow was for triangular shapes when its located at the bottom of the enclosure
- 3. Increasing the number of inner heated body into $N = 2$, leads to highest fuid fow strength. However, further increasing of number of the heated body leads to reduce the fuid fow strength
- 4. For various values of Ra, Da, Ha, γ_{MHD} , ϕ , Case 1 which represent the circular shapes reveals the highest augmentation in Nusselt number followed by Case 3 (rhombus), Case 2 (square) and Case 4 (triangular) reveals the lowest improving in heat transfer rate.
- 5. The magnetic feld inclination angle reveals diferent behavior on heat transfer rate where it was concluded that increasing its value from $0°$ into $30°$ leads to reduces the Nusselt number. further increasing of magnetic feld angle more than 30° leads to increasing of Nusselt number

Fig. 16 Nu with respect to Rayleigh number for various number of inner hot pipes considering different shapes at $\varphi = 0.02$, $Ra = 10^6$, $Da = 0.001$, $Ha = 20$

Fig. 17 Nu with respect to Rayleigh number under various location of inner hot body considering different shapes at $\varphi = 0.02$, $Ra = 10^6$, $Da = 0.001$, $Ha = 20$

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Fig. 18 Nu with respect to various position of inner hot body considering different shapes at $\varphi = 0.02$, $Ra = 10^6$, $Da = 0.001$, $Ha = 20$

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