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# **Design and Performance Study of Metamaterial with Quasi‑zero Stifness Characteristics Based on Human Body Structure**

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## **Abstract**

**Purpose** The design of a new meta-structure with quasi-zero stifness (QZS) characteristics based on human body structure not only expands the diversity of quasi-zero stifness structures but also demonstrates excellent performance in the feld of low-frequency vibration damping.

**Methods** Firstly, a quasi-zero stifness unit structure combining a horizontal cosine beam and vertical cosine beam is designed based on the human body structure, and multiple unit structures are combined into one meta-structure; secondly, according to the material-structural properties of the unit, a functional relationship is established for the positive stifness structure using the segmental function method, which is combined with the functional relationship of the negative stifness structure to obtain the theoretical model of the unit structure and then derives the mechanics of the whole meta-structure behavior of the whole meta-structure. Then, the unit structure model and the meta-structure model are designed and simulated, the static properties of the meta-structure are investigated, and photosensitive resin is chosen as the material in the simulation process. Finally, the transmittance comparison between the meta-structure and the linear vibration isolation structure with different damping coefficients is analyzed by theory, and the vibration isolation performance of the metastructure is simulated.

**Results** The static simulation analysis of the meta-structure verifes the theoretical solution of the meta-structure very well. According to the theoretical formulation, diferent parameters can be input to obtain the quasi-zero stifness properties of the structure. The quasi-zero stifness properties of the structure are independent of the inherent material properties. Due to the nonlinear nature of the quasi-zero stifness unit structure, it not only has a lower peak transmittance, but also a lower resonant frequency and a wider isolation range compared with the linear vibration isolator.

**Conclusions** The designed quasi-zero stifness meta-structure is suitable for low-frequency vibration isolation of precision instruments, which can be realized by diferent materials, and more complex meta-structures can be designed on this basis in the future, which has great potential for application in engineering practice.

**Keywords** Human body structure · Quasi-zero stifness · Metamaterials · Vibration isolation performance

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# **Introduction**

In many engineering practices, vibration isolation is often a critical issue [\[1](#page-14-0)]. A typical example is the optical vibration isolator, which plays an important role in the stability of the optical view axis and the accuracy of communication. The ideal optical vibration isolator achieves good vibration isolation in both the low and wide frequency ranges. To meet these diferent needs, there are many types of vibration isolators available, ranging from conventional rubber dampers to passive spring isolators to active vibration isolators [[2](#page-14-1)[–5](#page-14-2)]. Typically passive spring isolators are the most used and are known for their high reliability, ease of implementation, and



Smaller resonant frequencies are often expected in the design of vibration isolators to achieve better isolation performance. However, a smaller resonant frequency implies a lower stifness (or heavier mass), resulting in the generation of a lower load capacity (or a bulky system). Therefore, a large number of scholars have investigated some nonlinear isolation methods that possess high static stifness, while exhibiting good isolation performance in the low-frequency band [\[9,](#page-14-5) [10\]](#page-14-6).

In 1957, Molyneux proposed a three-spring QZS isolator [[11](#page-14-7)], which is considered to be the earliest QZS isolation device, and the high static-low dynamic stifness of this device opened up a new research direction for low-frequency isolation vibration. In the past 10 years, the theoretical analysis, design, and experimental research of QZS vibration isolators have been widely reported. The QZS structure is mainly composed of two parts: positive and negative stifness. Diferent types of quasi-zero stifness isolators can be constructed depending on the diferent composition structures of negative stifness (e.g., inclined springs, horizontal springs, cam springs, roller springs, magnetic springs, and Eulerian bending beams) and positive stifness [\[12](#page-14-8)[–16](#page-14-9)]. Cao et al. studied the transition from smooth to discontinuous dynamics of a three-spring system, which in the smooth region bears a clear resemblance to the Duffing oscillator, exhibiting standard dynamics controlled by a hyperbolic structure associated with the fxed state of the double trap. However, in the discontinuity limit, the dynamics are quite diferent from the standard limit [\[17](#page-14-10)]. Zhu et al. proposed a generalized inverse method for structural construction, which showed that elliptical trajectory tracking of the designed model is a sufficient condition for satisfying the zero stifness property, and that the unique geometry of the zero-stifness system provides new inspiration for the resonance-free design of metamaterial cells, and that the inverse method can even enable the more targeted application of the design to arbitrarily complex dynamical requirements [\[18](#page-14-11)]. Hao et al. proposed a dynamical model for a single-degreeof-freedom spring-mass system, which can be used as a reference for nonlinear support systems and stable-quasi-zero stifness (SQZS) vibration isolation systems for large aircraft ground vibration test (GVT) [\[19](#page-14-12)]. In addition, the dynamics analysis of the perturbed system near the optimized parameters reveals complex behaviors such as KAM structure, multiplicative period, chaotic crisis, multi-solution coexistence, intermittent chaos, and chaotic saddle, which provide a basis for better understanding of the complex dynamics of SQZS nonlinear systems. Sun et al. designed a vibration isolation platform with an n-layer scissor truss structure to achieve QZS characteristics by controlling diferent structural parameters to better achieve superior vibration isolation performance [\[20](#page-14-13)]. Zhou et al. proposed a new passive asymmetric quasi-zero stifness vibration isolator (AQZS-VI) consisting of a combination of two linear springs, cantilever plate springs and L-shaped rods (CPS-LSL). The proposed AQZS-VI has superior vibration isolation performance with lower vibration isolation onset frequency and better performance of displacement transfer capability than the linear model [[21](#page-14-14)]. Chang et al. proposed a quasi-zero stifness dynamic damper (QZS DVA) for the ultra-low-frequency absorption method. The conceptual design of the QZS DVA was carried out using a simple inclined spring model. The dynamic behavior of the QZS DVA system under harmonic excitation and random pulses was investigated and verifed by numerical analysis and experiments [[22](#page-14-15)]. Ye et al. designed an integrated fat-rotating rotating QZS isolator using a cam-roller mechanism capable of providing high static-low dynamic stifness in both directions simultaneously, and studied its nonlinear dynamic response by numerical and theoretical methods [\[23\]](#page-14-16). Xiong et al. designed a quasi-zero stiffness vibration isolation system with an additional X-shaped structure (X-QZS) by connecting the X-shaped structure to the platform. The designed X-QZS system has a wide QZS interval, small instability interval, low jump frequency, and high load carrying capacity [\[24](#page-14-17)]. Hao proposed a single- degree-of-freedom geometrically nonlinear oscillator with stable quasi-zero stifness (SQZS), modeled to include a set total mass covering an isolated object and a pair of horizontal springs that provide negative stifness parallel to a vertical linear spring to carry the load. Frequency band parameter optimization was performed using an extended averaging method to obtain the frequency response characteristics of the model. In addition, numerical simulations were performed to detect complex dynamical phenomena of periodic, chaotic motion [\[25](#page-14-18)]. Yan proposed a lever-type quasi-zero stifness vibration isolator (L-QZS-VI), where the QZS characteristic is achieved by a magnetic spring. Eddy current damping (ECD) is used to eliminate the jump phenomenon. A theoretical model of L-QZS-VI with ECD is established, and the effects of the tip mass of the lever, the lever ratio, the nonlinear stifness of the magnetic spring, and the excitation amplitude on the vibration isolation performance of L-QZS-VI are analyzed numerically and experimentally [[26\]](#page-14-19). Lu et al. investigated a coupled quasi-zero stifness (QZS) vibration isolator for axially loaded beams to improve the effectiveness of low-frequency isolation. A QZS contains two magnetic rings with negative stifness and a coil spring with positive stifness, which have high static stifness to resolve the structural instability. Analytical and numerical results show that the axially loaded beam of the parallel-coupled quasi-zero isolation system has a more signifcant suppression of power reduction at low frequencies [[27](#page-14-20)]. Hao investigated the nonlinear energy transfer in an arbitrary boundary fexible plate coupled by a high static low dynamic stifness (HSLDS) isolator. The nonlinear coupling dynamics equations are derived using the Lagrangian method, and the modifed Fourier series method and Rayleigh–Ritz method give the modal coefficients of the arbitrary boundary fexible plate with nonlinear isolators. In addition, the analytical results of the modal shape of the fexible plate are verifed by the fnite element simulation used in this paper, and the isolation performance of the nonlinear vibrator supported on the fexible plate is experimentally verifed [\[28](#page-14-21)]. Most researchers have focused on quasi-zero stifness structures for large equipment, and less on quasizero stifness vibration isolators for precision instruments.

The term "metamaterials" is commonly used to describe artifcial composite materials consisting of periodic or random arrangements of artifcial subwavelength structures. Metamaterials mainly include electromagnetic metamaterials [[29\]](#page-15-0), optical metamaterials [\[30](#page-15-1)[–33](#page-15-2)], acoustic metamaterials [[34](#page-15-3)], and mechanical metamaterials [\[35](#page-15-4)]. Mechanical metamaterials achieve idiosyncratic mechanical properties through the internal design of the unit structure. Mechanical metamaterials are not only a complement to conventional natural materials, but also they have gradually become the protagonists in the feld of new materials, endowed with richer functional properties, including reconfgurability, programmability, self-determination, and mechanical operations. In the future, the functional properties of natural materials will gradually fade away, and will more often be used as components of functional metamaterials in the form of raw materials. This paper focuses on vibration isolation research, mainly on mechanical metamaterials to manipulate the overall equivalent stifness of the structure. In recent years, the combination of quasi-zero stifness properties and mechanical metamaterials has provided a new way for the development of vibration isolation of small precision instruments.

Cai et al. proposed a new metamaterial plate structure with quasi-zero stifness (QZS) resonators, which were theoretically derived, numerically simulated, and experimentally investigated for wave propagation characteristics and showed excellent attenuation in the ultra-low-frequency band gap [\[36](#page-15-5)]. Li et al. designed a bistable hybrid symmetric laminate (BHSL) as a negative stifness element to design quasi-zero stifness vibration isolators. The theoretical model and fnite element model were validated from the static and dynamic analysis [\[37](#page-15-6)]. Lin has developed a novel two-dimensional (2D) locally resonant (LR) metamaterial with quasi-zero stifness (QZS) properties in both horizontal and vertical directions. The anisotropy of the precompression along the horizontal and vertical directions can be exploited so that the complete band gap covers diferent frequency ranges, providing a way to achieve a complete band gap at low frequencies, independent of the direction of incidence of the in-plane waves [\[38\]](#page-15-7). Ye proposed a truss-spring-based stack Miura-ori (TS-SMO) structure to provide the required stifness for high static—low dynamic requirements. A QZS vibration isolator is then designed based on specifc parameters. The static performance of the system was verifed by the force–displacement response. The dynamic analysis using the harmonic balance method (HBM) and numerical simulation was used to derive the displacement transmittance, and the isolation performance under variable viscous damping was also discussed to examine the efect of system damping [\[39](#page-15-8)]. Fan et al. designed a meta-structure combining the penetration behavior of a sinusoidal beam and the bending dominant support of two semicircular arches, which in turn achieved quasi-zero stifness properties for each cell, and analyzed the static and dynamic properties of the meta-structure. The results show that the meta-structure can achieve good quasi-zero stifness with excellent vibration isolation performance by designing the cells correctly [\[40](#page-15-9)]. The combination of quasi-zero stifness and metamaterials ofers a new approach to low-frequency vibration reduction that has an important role in future engineering.

The human body encounters several types of whole-body vibration in daily life: the more common ones are walking, running and jumping [[41\]](#page-15-10).This is when the human body acts as a damper to prevent damage to the brain from vibrations and shocks [\[42](#page-15-11)]. The resonant frequency of the human body varies for diferent body postures [[43](#page-15-12)]. Especially, when people stand with their knees bent, the peak of the transmission rate is even below 3 Hz. At the same time, the human body is able to attenuate frequencies of 20 Hz and above [[44\]](#page-15-13).

The excellent vibration isolation characteristics of the human body discussed above provide a novel vibration isolation structure design for this study, which has the potential to achieve very good ultra-low-frequency vibration isolation and excellent vibration isolation performance over a wide frequency range. In this study, a vibration-damping structure based on human body structure is proposed as shown in Fig. [1](#page-3-0) and Fig. [2](#page-3-1).The new structure consists of two parts: a double vertical cosine beam structure to simulate the human bent leg, and a horizontal cosine beam structure to imitate the human upper body arm. The layout of the proposed vibration-damping system is shown in Fig. [2.](#page-3-1)

In this paper, a new meta-structure with quasi-zero dynamic stifness is designed based on human foundation structure, and its characteristics in vibration isolation applications are investigated by combining theoretical analysis and fnite element simulation. In Section "[Human-inspired](#page-3-2) [vibration isolation structure model](#page-3-2)", a meta-structure consisting of several cells is designed and the mechanical behavior with quasi-zero stifness is derived. In Section "[Static](#page-6-0) [properties of the meta-structure](#page-6-0)", the static properties of the structure under uniaxial compression are investigated through simulations. Section "[Dynamics analysis"](#page-8-0) studies the dynamic characteristics and vibration isolation performance of the meta-structure. Section ["Conclusion"](#page-13-0) gives the





<span id="page-3-0"></span>**Fig. 1** Human body structure diagram and unit model



<span id="page-3-1"></span>**Fig. 2** Quasi-zero stifness unit cell and meta-structure for human-like structural design

conclusion. The structure not only extends the diversity of quasi-zero stifness structures, but also has great potential for application in the feld of low-frequency vibration isolation of small continuous structures.

# <span id="page-3-2"></span>**Human‑Inspired Vibration Isolation Structure Model**

To facilitate discussion and interpretation, the quasi-zero stifness structure in Fig. [2](#page-3-1) is re-arranged as Fig. [3](#page-4-0), and the structural parameters are clearly stated.



As shown in Fig. [3,](#page-4-0) the unit structure consists of horizontal and vertical cosine beams and stifer stifening walls. Since the horizontal cosine beam provides negative stifness characteristics and the vertical cosine beam provides positive stifness characteristics, the quasi-zero stifness characteristics are produced by the proper arrangement of the two structures. As shown in Fig. [3](#page-4-0), the shape of the horizontal cosine beam is given by  $y = h_1/2 \left[ 1 - \cos \left( 2 \pi x / l_1 \right) \right]$ . The length is  $l_1$ , the depth is  $b_1$ , the thickness is  $t_1$ , and the amplitude is  $t_1$ . The shape of the vertical cosine beam is then given by  $x = h_2/2[1 - \cos(2\pi y/l_2)]$ . The length is *l*<sub>2</sub>, the depth is  $b_2$ , the thickness is  $t_2$ , and the amplitude is  $h_2$ . The cosine beam is connected by a stifening wall of thickness *t*. By a



<span id="page-4-0"></span>**Fig. 3** Quasi-zero stifness structural model of human-like structural design



<span id="page-4-1"></span>**Fig. 4** Force–displacement curves of horizontal cosine beam and vertical cosine beam

reasonable arrangement of negative and positive stifnesses, the unit will have quasi-zero stifness.

## **Mechanical Model of the Unit Structure**

According to Qiu et al. [[45](#page-15-14)], when the parameters of the cosine beam satisfy  $h_1/t_1 \ge 6$ , its force–displacement curve can be approximated by three connected straight lines, as shown in Fig. [4](#page-4-1).The corresponding parameters in the force–displacement curve are shown in Eq. [1](#page-4-2), where *E* is the Young's modulus of the material and  $I_1$  is the moment of inertia of the beam.

<span id="page-4-2"></span>
$$
f_{top} \approx 740 \frac{EI_1 h_1}{l_1^3} \quad f_{bot} \approx 370 \frac{EI_1 h_1}{l_1^3} \quad d_{mid} = \frac{4}{3} h_1 \tag{1}
$$

$$
d_{top} = 0.16h_1 \ d_{bot} = 1.92h_1 \ d_{end} = 1.99h_1. \tag{2}
$$

Assuming that the transverse displacement component is constant along the thickness direction, the stifness of the cosine beam in three diferent stages is shown in Eq. [3](#page-4-3)

<span id="page-4-3"></span>
$$
k_{a1} \approx \frac{4625EI_1}{l_1^3} \quad k_{a2} \approx -\frac{13875EI_1}{22l_1^3} \quad k_{a1} \approx \frac{37000EI_1}{7l_1^3}.\tag{3}
$$

For the vertical cosine beam, which mainly provides positive stifness, the force–displacement curve under vertical displacement is shown in Fig. [4.](#page-4-1) The overall stifness of the beam shows a nonlinear relationship, but when compared with a horizontal cosine beam that provides negative stifness, the positive stifness cosine beam can be divided into three stages with approximately linear force and displacement curves in each stage, allowing the stifness to have constant values in diferent stages.

As shown in Fig. [5](#page-5-0), according to beam theory, the stifness will be proportional to the Young's modulus *E* of the material, the moment of inertia  $I_2$  of the beam and the inverse of the cube of the vertical cosine beam length *l*<sub>2</sub>.And the segmented function relationship is presented with the cut-off point of  $0.16h_2$  and  $1.92h_2$ . The method of segmented functions used in this paper can be regarded as a special form of calculus, using a straight line approximation to ft the curve. In this paper the force and displacement curves





<span id="page-5-0"></span>**Fig. 5** Dimensionless force–displacement curve of vertical cosine beam

for negative stifness are divided into three segments, so the positive stifness is also divided into three segments. Diferent analyses can be performed in subsequent applications according to their own conditions. In this paper, the variation of stifness and displacement of the vertical cosine beam can be written as:

$$
k_b = \begin{cases} \alpha \frac{EI_2}{l_2^3} & \Delta l < 0.16h_1\\ \beta \frac{EI_2}{l_2^3} & 0.16h_1 < \Delta l < 1.92h_1\\ \gamma \frac{EI_2}{l_2^3} & 1.92h_1 < \Delta l < 1.99h_1 \end{cases} \tag{4}
$$

where  $\alpha$ ,  $\beta$ ,  $\gamma$  are three dimensionless coefficients. It should be noted that the dimensionless coefficients corresponding to the diferent cross sections can be obtained by elastic analysis [[46](#page-15-15)] and the fnite element method. To refect the deformation response of the vertical cosine beam and to ensure that it is always in an elastic state when displaying quasi-zero stifness characteristics, this paper uses ABAQUS for fnite element analysis. The cosine beam is modeled with a solid unit of linear elastic material C3D20R, and the vertical end part shift is  $0.4 \times l_2$ . Multiple simulations were performed for cosine beams of diferent thicknesses, and dimensionless coefficients  $\alpha$ ,  $\beta$ ,  $\gamma$  could be obtained by fitting the curves to the stifness–thickness relationship. In this paper, we take the vertical cosine beam with the parameter  $t_2/l_2 = 0.05$  as an example. The simulation results and dimensionless force–displacement curves of the cosine beam under the action of vertical displacement are shown in Fig. [5.](#page-5-0) Since the force and displacement curves of the horizontal cosine beam are mainly divided into three segments, the vertical cosine beam force displacement curve is also divided into three segments and approximated as three straight lines, and



(a) Fitting curve of the first segment stiffness and FEA results



(b) Fitting curve of the second segment stiffness and FEA results



(c) Fitting curves of the third segment stiffness and FEA results

<span id="page-5-1"></span>**Fig. 6** Fitting curve of vertical cosine beam stifness and comparison results of fnite element analysis

the stiffness of the cosine beam can be found as  $k_b = F/\Delta L$ . The resulting fnite element result for the stifness of cosine beams with diferent thicknesses and dimensionless stifness  $k_b$ /*Eb*<sub>2</sub> under parameter  $t_2/l_2$  are shown in Fig. [6](#page-5-1).

The triple ftted curves agree well with the fnite element results, and the dimensionless coefficients  $\alpha = 150.25$ ,  $\beta$  = 50.44 and  $\gamma$  = 32.44 are obtained.



<span id="page-6-1"></span>**Fig. 7** Force–displacement curve of the unit structure

Considering that the horizontal cosine beam and the vertical cosine beam are arranged in parallel, the stifness of the unit under vertical displacement is the sum of the respective stifnesses. Therefore, the force–displacement curve of the cell can also be represented by three connected straight lines to obtain quasi-zero stifness in region II, as shown in Fig. [7.](#page-6-1) The expression for the stifness of the unit can be as follows:

$$
k_1 = k_{a1} + 2k_{b1} \ k_2 = k_{a2} + 2k_{b2}
$$
  

$$
k_3 = k_{a3} + 2k_{b3}.
$$
 (5)

By correctly selecting the parameters of the horizontal and vertical cosine beams, it is possible to design the stifness of diferent regions and obtain the appropriate static stiffness to support the isolation weight, as well as the quasi-zero dynamic stifness to ensure the vibration isolation performance.

## <span id="page-6-0"></span>**Static Properties of the Meta‑structure**

### **Unit Structure Design**

In this paper, three different cell structures  $U_1$ ,  $U_2$  and  $U_3$ are designed, as shown in Fig. [8,](#page-6-2) and only the parameters  $t_2$  of each cell are different, the rest of the parameters are the same.  $l_1 = l = 60$  mm,  $t_1 = 2$  mm,  $h_1 = 12$  mm,  $b_1 = 10$  mm, *l*<sub>2</sub> = 40 mm *h*<sub>2</sub> = 12 mm, *b*<sub>2</sub> = 10 mm, *t* = 10 mm, *h* = 50 mm. By theoretical solution, the unit structure reaches quasi-zero stiffness at  $t_2$ =2.45 *mm*, and two cases of  $t_2$ =2 *mm* and 3 *mm* are designed, respectively, with three diferent structures representing positive stifness, negative stifness and



<span id="page-6-2"></span>**Fig. 8** Unit structure model  $(U_1, U_2, U_3)$ 

quasi-zero stifness. Two samples of meta-structures with diferent parameters were also designed to study their static properties. The specifc parameters of the meta-structure are shown in Table [1.](#page-7-0) A photosensitive resin [[47\]](#page-15-16) was used as a constituent material for the preparation of the samples. Young's modulus of photosensitive resin was *E* = 1.4*Gpa*.

### **Numerical Simulation**

<span id="page-6-3"></span>Finite element analysis of the static properties of the metastructure was carried out using fnite element software. The meta-structure is modeled with a tetrahedral mesh, ten-node quadratic tetrahedral cells (C3D10), and mesh sensitivity testing is performed to ensure accuracy. Due to the large deformation behavior of the structure, geometric nonlinearities need to be taken into account. In order to simulate the actual behavior of the meta-structure under compression, all six degrees of freedom (three translations and three rotations) are constrained on the bottom surface of the structure, and the top surface of the structure is coupled to a reference point where all fve degrees of freedom are constrained, except for the displacement in the compression direction. In the static generic step, a specifed displacement is applied to the reference point to simulate the compression process of the structure. It should be clear that in the theoretical analysis, the stifened walls of the structure should be rigid. This means that the left and right sides of the structure are constrained and only translations in the displacement direction are allowed. However, in the actual compression test, the left and right edges of the structure were free, causing the stifened walls to deform. Therefore, these two diferent boundary conditions are considered simultaneously in the numerical simulation, as shown in Fig. [9.](#page-7-1) In the theoretical analysis simulation, the left and right edge of the structure are set as constrained boundaries (CB).

As an example, the cell structure  $U_2$  shows the behavior of the cell under diferent vertical displacements, as shown in Fig. [10](#page-8-1). The reaction forces of the  $U_1$ ,  $U_2$  and  $U_3$  unit structures under vertical displacement are shown in Fig. [11.](#page-8-2) The results show that the cell can exhibit positive, quasi-zero and negative stifness characteristics in region II through a reasonable design. For the unit structure  $U_1$ , the absolute



parameters

<span id="page-7-0"></span>**Table 1** Meta-structural



<span id="page-7-1"></span>**Fig. 9** Finite element model of the meta-structure

value of the positive stiffness  $2k_{b2}$  of the vertical cosine beam is smaller than the absolute value of the negative stifness  $k_{a2}$  of the horizontal cosine beam. From Eq. [5,](#page-6-3) the unit  $U_1$  is a negative stiffness of  $k<sub>2</sub> < 0$ . Similarly, for samples  $U_2$  and  $U_3$ , the absolute value of the positive stiffness  $2k_{b2}$  is almost equal to or greater than the negative stiffness  $k_{a2}$ . The  $k_2 \rightarrow 0$  quasi-zero stiffness and the positive stiffness with  $k_2$  > 0 are obtained for U<sub>2</sub> and U<sub>3</sub>, respectively. The results show that the quasi-zero stifness properties of the structure are determined due to the internal periodic cell structure rather than due by the inherent properties of the material.

By studying the response of the two meta-structures  $S_1$  and  $S_2$  under vertical displacement conditions for the structure under FB and CB conditions, the simulation results are shown in Figs. [12](#page-9-0) and [13.](#page-10-0) The results show that the deformation of the horizontal cosine beam exhibits fast passage and the deformation of the vertical cosine beam is dominated by bending in each cell under FB and CB conditions. For the internal unit of the meta-structure, the horizontal cosine beam undergoes classical symmetric buckling behavior during compression. And for the left and right edges of the meta-structure under FB conditions, the buckling behavior of the horizontal cosine beam is asymmetric. This phenomenon will change the horizontal cosine beam to be asymmetric. This phenomenon will

Metastructure Number of arrangements Structure parameters  $l_1$   $t_1$   $h_1$   $b_1$   $l_2$   $t_2$   $h_2$   $b_2$   $h$   $t$  *l S*<sup>1</sup> 3 rows and 4 columns 50 1.0 6.0 10 30 1.106 6.0 10 40 10 50 *S*<sup>2</sup> 4 rows and 5 columns 30 0.5 4.0 5.0 18 0.58 4.0 5.0 23 5.0 30

> change the horizontal cosine beam stifness properties and thus adversely afect the quasi-zero stifness properties of the meta-structure. Since the asymmetric buckling occurs only at the edges of the structure, the adverse efect on the quasi-zero stifness performance of the meta-structure will gradually decrease with the gradual increase in the number of cells.

The vertical displacements and reaction forces of the meta-structure in the numerical simulation were recorded and compared with the theoretical results, as shown in Fig. [14](#page-10-1) and Fig. [15.](#page-11-0) From the theoretical and simulation results, it can be seen that in the initial stage, the reaction force of the meta-structure increases rapidly with the increase of the compression displacement. This is because, in the initial stage, both the horizontal cosine beam and the vertical cosine beam of each unit structure provide positive stifness. When the compression displacement exceeds the initial stage, the horizontal cosine beam of the unit structure exhibits negative stifness, which cancels with the positive stifness of the vertical cosine beam, causing the positive stifness of the elemental structure to drop to quasi-zero stifness. When the vertical displacement exceeds a certain value, the stifness of the horizontal cosine beam changes back to positive stifness, and each unit structure provides positive stifness again. With the gradual increase of the compression displacement, the meta-structure will experience the transition from quasi-zero stifness to positive stifness.

Comparing the theoretical solution and simulation results of CB boundary conditions in Fig. [13](#page-10-1) and Fig. [15](#page-11-0), it can be found that both the theoretical solution and simulation results show good agreement, except for the region of the two transition points between positive and quasi-zero stifness. This is because large deformations and nonlinear responses are considered in the simulation analysis, while the theoretical model is based on the assumption of small deformations. Therefore, the simulation analysis results in a nonlinear force–displacement curve and the theoretical solution results in a linear one. As  $h_1/t_1$  of the cosine beam increases, the nonlinear efect gradually decreases, and the theoretical and simulation results become closer and closer. Even if there is a discrepancy between the theoretical and simulation results in the transition region, the theoretical solution accurately refects the quasi-zero stifness properties of the meta-structure. It is verifed that the theoretical model in this paper can efectively evaluate the mechanical



 $D=18mm$ 



<span id="page-8-1"></span>**Fig. 10** Behavior of the unitary structure under vertical displacement



<span id="page-8-2"></span>**Fig. 11** Force–displacement curves of unitary structures  $U_1$ ,  $U_2$  and  $U_3$ 

properties of quasi-zero stifness element structures when designing quasi-zero stifness element structures.

The quasi-zero stiffness properties of the simulated results in Figs. [13](#page-10-1) and [15](#page-11-0) are slightly weakened in the FB

boundary condition compared to the theoretical and simulated results in the CB boundary condition, mainly due to the asymmetric buckling of the horizontal cosine beam of the boundary cell caused by the absence of constraints on the left and right edges of the structure. However, as mentioned before, the infuence of the boundary conditions decreases as the number of unit structures gradually increases.

# <span id="page-8-0"></span>**Dynamics Analysis**

## **Power Transmission Characteristics**

Based on the static properties, the vibration isolation performance of the designed meta-structure was investigated. The results of the static force characteristics analysis show that the mechanical behavior of the meta-structure under vertical displacement contains three regions, quasi-zero stifness region II and two approximately linear stifness regions I and III. For objects placed on the meta-structure isolated from vibration, their weight is mainly supported by the approximate linear stiffness of region I. When the vibration excitation acts on the meta-structure, the object will rely on the quasi-zero stifness of region II for



<span id="page-9-0"></span>**Fig. 12** Comparison of response under FB and CB structures obtained by simulation of  $S_1$ 



 $D = 36$ mm

efective isolation. In order to study the vibration isolation characteristics of the meta-structure, the dynamic characteristics of models  $S_1$  and  $S_2$  were investigated. Assume that the upper surface of the specimen supports an object and the lower surface is excited by displacement-controlled vibration. The object with mass M is located in the metastructure II region, and the gravitational force provided by the object with mass M is replaced by the corresponding concentrated force in the simulation calculation.

The most important index to describe the performance of an isolator is transmissibility. Assume that the bottom of the meta-structure is subjected to displacement-controlled vibration excitation  $Z = Z_0 \sin \omega t$ , where  $Z_0$  is the vibration amplitude and is the  $\omega$  vibration frequency. In the case of

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<span id="page-10-0"></span>**Fig. 13** Force–displacement curves obtained from the theoretical solution and simulation of  $S_1$ 

meta-structure isolation, the vibration amplitude  $Z_1$  at the top of the meta-structure can be obtained. The transfer rate was defined as  $T = Z_1/Z_0$  to evaluate the vibration isolation performance of the meta-structure. For a conventional linear spring isolator, the transfer rate equation can be derived as:

$$
T = \sqrt{\frac{K^2 + c^2 \omega^2}{(K - M\omega^2)^2 + c^2 \omega^2}},
$$
\n(6)

where  $K$  is the linear stiffness of the spring and  $c$  is the damping of the spring material. In fact, the quasi-zero stifness is also approximately linear, having a rather small stiffness  $K_2$ . Thus, the transfer rate of the meta-structure can be obtained by Eq.  $6$ , where *K* is replaced by  $K_2$ . Since the weight of the object in the meta-structure is mainly supported by  $K_1$ , an equivalent linear spring isolator with stiffness  $K_1$  is introduced to verify the vibration isolation performance of the proposed meta-structure with QZS characteristics. Comparison with an equivalent linear vibration isolator is an efective method to verify the vibration isolation performance of the newly designed QZS structure [[48,](#page-15-17) [49](#page-15-18)].The transfer rate of the linear spring isolator can also be obtained by replacing K by  $K_1$  in Eq. [6](#page-10-2).

# **Comparison of QZS Element Structure and Linear Isolator**

Figures [16](#page-11-1) and [17](#page-11-2) plot the transmittance of the linear spring isolator and the QZS structural models  $S_1$  and  $S_2$ . All transmittance results are expressed in dB, i.e., 20log10T. The vibration excitation frequency ratio is expressed in dimensionless form as  $\Omega = \omega/\omega_n$ , where  $\omega_n = \sqrt{K_1/M}$ . The dimensionless damping coefficient of the meta-structure is defined as  $\xi = c/2\sqrt{K_1M}$ . As shown in Fig. [16](#page-11-1) and Fig. [17,](#page-11-2) for linear isolators, the efective isolation range should meet Ω *>* 1.414.And for the proposed QZS meta-structural models S<sub>1</sub> and S<sub>2</sub>, the effective isolation ranges satisfy  $\Omega > 0.116$ and  $\Omega$  > 0.114, respectively. Compared with the linear vibration isolator, the proposed QZS element structure has a wider efective vibration isolation range and can be efectively isolated at lower vibration frequencies.

<span id="page-10-2"></span>In addition to the wide efective isolation range, the QZS meta-structure also has a low transmittance.



<span id="page-10-1"></span>**Fig. 14** Comparison of response under FB and CB structures obtained by simulation of  $S_2$ 





<span id="page-11-0"></span>**Fig. 15** Force–displacement curves obtained from the theoretical solution and simulation of  $S_2$ 



<span id="page-11-1"></span>**Fig. 16** Transmittance of linear vibration isolator and  $S_1$  structure

It can be seen that the peak transmittance at the resonant frequency of the QZS superstructure corresponding to different damping coefficients is smaller than that of the linear isolator. In addition to the peak transmittance, the transmittance of the QZS element structure is also much smaller than that of the linear isolator for vibration frequencies located in the efective isolation range until the vibration frequency approaches a very large value, such as  $\Omega = 100$ .



<span id="page-11-2"></span>**Fig. 17** Transmittance of linear vibration isolator and  $S_2$  structure

For example, for the same damping factor  $\xi = 0.01$  at vibration frequency ratio  $\Omega = 3$ , the transmittance of a linear vibration isolator is  $-18.04$  dB, S<sub>1</sub> is  $-43.6$  dB and  $S_2$  is  $-$  43.46 dB.

#### **Effect of Damping Coefficient**

As shown in Fig.  $16$  and Fig. [17,](#page-11-2) the damping coefficient does no efect on the efective isolation range of the QZS element structure and the linear vibration isolator, but the transmittance at diferent vibration frequencies is afected by the damping coefficient. For the QZS element structure and linear vibration isolators, the transmittance decreases with increasing damping coefficient when the vibration frequency is not in the efective vibration isolation range. However, when the vibration frequency is in the efective vibration isolation range, the transmittance increases as the damping coefficient increases. By comparing the QZS element structure with the linear vibration isolator, it is found that the damping coefficient has a greater effect on the QZS element structure than the linear vibration isolator in the efective vibration isolation range. The results show that the QZS structure with a smaller damping coefficient has better vibration isolation performance. Although purely from the point of view of vibration isolation, the increase in damping coefficient will reduce the vibration isolation efect, in the actual working process of the machine, the external excitation, in addition to



<span id="page-12-0"></span>**Fig. 18** Finite element model of  $S_1$  vibration isolation simulation

simple harmonic vibration may also contain some irregular shocks. As the shock will cause the equipment larger amplitude of free vibration, the purpose of increasing the damping is to make the free vibration disappear quickly. Especially when the vibration isolation object through the resonance area, damping becomes more important. Therefore, in practical application, it is necessary to choose the appropriate damping coefficient for vibration isolation according to its situation.

#### **Simulation of Vibration Isolation of QZS Structure**

The fnite element model containing the gravity of the object and the meta-structure  $S_1$  is given in Fig. [18.](#page-12-0) The whole simulation process is divided into two steps. In the frst step, the meta-structure is compressed under gravity when the bottom surface of the meta-structure is fully constrained, and this step will bring the meta-structure into the quasizero stifness region. In the second step, the bottom surface is transformed from fully constrained to displacementcontrolled vibration excitation. Since the transmittance in both Fig. [16](#page-11-1) and Fig. [17](#page-11-2) is derived from the CB condition, the boundary conditions on the left and right sides of the meta-structure are defned as CB throughout the simulation. The vibration excitation amplitude  $Z_0$  is set to 5 mm. Three different excitation frequencies of  $S_1$  and  $S_2$  in the effective vibration isolation range were selected as  $\Omega = 0.5$ ,  $\Omega = 1$  and  $\Omega$  = 2. Then, the bottom and top displacements of the metastructure were recorded to evaluate the vibration isolation performance corresponding to diferent vibration frequencies and damping coefficients.

The displacements at the bottom and top of the QZS element structure for  $S_1$  and  $S_2$  at different vibration fre-quencies and damping coefficients are shown in Figs. [19](#page-12-1) and [20](#page-13-1). The simulation results in Fig. [19](#page-12-1) and Fig. [20](#page-13-1)



<span id="page-12-1"></span>**Fig. 19** Displacement at the bottom and top of the  $S_1$  quasi-zero stiffness structure

visualize the vibration isolation performance of the meta- structure. A signifcant reduction in amplitude was obtained through the meta-structure. It can also be seen that at the same vibration excitation frequency, the material with a smaller damping coefficient can obtain better vibration isolation performance, but the material with a smaller damping coefficient also takes a relatively long time to reach a smooth trend in amplitude. The main reason is that the smaller damping coefficient is weaker for the peak amplitude suppression, which makes it slower to reach the smooth state. For materials with the same damping coefficient, the vibration isolation performance of the meta-structure is better at higher vibration excitation frequencies. It can also be seen by the theoretical solution, i.e., Figs. [16](#page-11-1) and [17,](#page-11-2) that the larger the frequency ratio, the smaller the transmittance at the same damping factor.





<span id="page-13-1"></span>**Fig. 20** Displacement at the bottom and top of the  $S_2$  quasi-zero stiffness structure

That is, the further away from the peak frequency ratio of the system, the better the performance of the vibration isolation system. For each curve in Fig. [19](#page-12-1) and Fig. [20,](#page-13-1) when the behavior of the meta-structure tends to be stable, the vibration amplitude obtained is  $Z_1$ , and  $Z_1$  is much smaller than Z. It can be seen that the QZS meta-structure proposed in this paper has good vibration isolation performance. It also verifes that the theoretical results are accurate.

# <span id="page-13-0"></span>**Conclusion**

A meta-structural vibration isolation system with quasizero stifness is designed based on the bionic human structure. Combining theoretical and simulation methods, the static and dynamic properties of this meta-structure are

investigated. The vibration isolation performance of the proposed meta-structure was evaluated. The results show that the proposed QZS superstructure has good vibration isolation performance. To summarize the main research content of this paper, the following conclusions are drawn.

- (1) A unit structure with quasi-zero stifness characteristics can be realized through the rational design of horizontal cosine beams and vertical cosine beams. Based on the assumption of small deformation, the theoretical values of the mechanical properties of the meta-structure composed of numerous unit structures are derived.
- (2) Static compression simulations of the unit structure and the meta-structure model were performed, and the results showed that the designed meta-structure has ideal quasi-zero stifness characteristics, which are consistent with the theoretical results. Compared with the constrained boundary conditions, the free boundary conditions at the left and right boundaries of the meta-structure have an adverse efect on its quasi-zero stifness characteristics. However, this efect gradually becomes smaller with the gradual increase of the number of unit cells.
- (3) The dynamics analysis and the calculated transmittance are obtained that the meta-structure can obtain a larger vibration isolation range and lower transmittance compared with the linear vibration isolator. The peak value is also lower than that of the linear vibration isolator at the same damping factor. The proposed meta-structure has good vibration isolation performance.
- (4) Theoretically, in the effective vibration isolation range, the peak vibration isolation rate decreases with the increase of the damping coefficient, and after the peak, the smaller the damping coefficient, the better the vibration isolation performance of the QZS structure. Simulation results of the vibration isolation performance of this element structure show that the element structure signifcantly reduces the vibration amplitude of the vibration isolation object. At higher vibration excitation frequencies, the smaller the damping coefficient, the better the vibration isolation performance of the meta-structure. In practical application, the vibration isolation system should be selected reasonably according to its situation for use.

The structural unit designed in this paper is a basic structure, and the mechanical properties achieved depend mainly on its own structure, which can be realized by different materials. In the future, more complex two-dimensional, three-dimensional and programmable vibration

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isolation element structures can be designed on this basis, which have greater application value.

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**Data Availability** Data sets supporting the conclusions of this paper are included in the article.

#### **Declarations**

**Conflict of Interest** The authors declare that they have no known competing fnancial interests or personal relationships that could have appeared to infuence the work reported in this paper.

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