#### RESEARCH



# Constructing core-shell carbon fiber/polypyrrole/CoFe<sub>2</sub>O<sub>4</sub> nanocomposite with optimized conductive loss and polarization loss toward efficient electromagnetic absorption

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#### Abstract

The inadequate impedance matching and weak attenuation capability for incident electromagnetic waves exhibited by carbon fibers (CF) are critical factors limiting their application served as absorbing materials. Constructing a nanocomposite system that simultaneously exhibits both dielectric loss and magnetic loss characteristics is a feasible strategy to overcome these limitations. In the present study, a core-shell CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is fabricated through electrodeposition and subsequent hydrothermal methods to enhance the attenuation capacity and impedance matching of bare CF. Under the synergistic effects of diverse components and a peculiar network structure, the nanocomposite demonstrates optimized conductive loss and polarization loss, which results in a remarkable electromagnetic wave absorption performance with a minimum reflection loss ( $RL_{min}$ ) of -55.33 dB and an effective absorption bandwidth (EAB) of 6.48 GHz (12 ~ 18 GHz) at optimal thicknesses of 2.11 and 2.42 mm, respectively, suggesting its promising application as a candidate absorber. More importantly, the exploration concerning the absorption mechanism provides significant insights into the attenuation modes of the dielectric-magnetic loss hetero-junction, which is beneficial for developing similar absorbing materials.

**Keywords** Conductive loss  $\cdot$  Polarization loss  $\cdot$  Electromagnetic absorption  $\cdot$  The absorption mechanism  $\cdot$  Dielectricmagnetic loss hetetero-junction  $\cdot$  The synergistic effects

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## **1** Introduction

Over the last decade, a variety of electronic devices that use electromagnetic wave as a medium have been extensively deployed in various domains, including telecommunications, aerospace, and automobiles [1]. However, the widespread presence of these electronic products has heightened public concerns about electromagnetic radiation pollution and potential health risks [2, 3]. To mitigate these issues, electromagnetic wave absorbing materials, which possess the capability to effectively dissipate electromagnetic wave and convert it into other energy forms, have increasingly captured the attention of researchers [4, 5].

Currently, due to their high electrical conductivity, excellent chemical stability, tunable dielectric properties, costeffectiveness, and wide availability of raw materials, carbonbased materials (such as carbon nanotubes, porous carbon, graphene, and carbon fibers, etc.) have emerged as the preferred choice for producing electromagnetic wave absorbers [6]. Undoubtedly, carbon fiber (CF) [7], serving as a outstanding reinforcing material, has exhibited exceptional performance in the realm of high-strength structural composites [8, 9]. As the concept of integrated structure and function in carbon fiber materials gradually takes root, research into their novel functional characteristics becomes increasingly pressing. This exploration is beneficial for expanding their application range and meeting the practical requirements of electromagnetic protection in various future industries [10].

It is widely acknowledged that a single carbon component struggles to achieve optimal impedance matching conditions [11]. Moreover, the simplistic mechanism by which CF interacts with incident electromagnetic wave often results in unsatisfactory attenuation effects. Constructing CF-based composite is a feasible and frequently utilized strategy, often employed by numerous researchers to enhance the electromagnetic wave absorption performance of individual CF [12]. Wang et al. [13] successfully developed flexible CF@ ZnO composites, which achieve an effective absorption bandwidth (EAB) of 10.6 GHz, covering both the whole X and Ku bands. In another study, Wang and colleagues [14] fabricated a CF@MXene@MoS<sub>2</sub> nanocomposite with a three-dimensional layered structure, and the product achieves a minimum reflection loss ( $RL_{min}$ ) of -61.51 dB at a matching thickness of 3.5 mm, while demonstrating an EAB value of 7.6 GHz at 2.1 mm. Zhang et al. [15] investigated the wave absorption performance of the CF@MoS<sub>2</sub> composite, which exhibited an EAB value of 10.85 GHz at a matching thickness of 3.8 mm. These research achievements have not only contributed significantly towards constructing CF-based composite absorbing materials, but also highlighted the tremendous application potential as effective electromagnetic wave absorbers. However, the pursuit of superior absorption performance at thinner matching thickness inspires ongoing research in this field.

Conductive polymers [16], represented by polyurethane, polypyrrole, and polythiophene, fall under the category of dielectric loss type absorbing materials, which also exhibits remarkable compatibility with CF [17]. The composite made of CF and conductive polymer retains the native conductive loss of CF. It also enhances the interface loss capability when encountering incident electromagnetic wave, achieving superior absorption performance compared to CF alone. On the other hand, the magnetic components based on iron, cobalt, nickel can generate strong magnetic loss against incident electromagnetic wave. By compounding them with dielectric loss components, it ensures that the composite material obtains more appropriate impedance matching condition and improved electromagnetic wave attenuation performance. However, when relying solely on dielectric loss and interface polarization, these composites struggle to effectively attenuate incident electromagnetic wave [18]. Moreover, the impedance matching has not yet reached optimal condition, resulting in that the EAB value, matching thickness, and RL<sub>min</sub> cannot meet the requirements for "thinness, lightness, width, and strength" [19-21]. To address this issue, magnetic components are typically introduced into the absorbing material system [22]. While realizing the dual action of dielectric-magnetic losses [23], it further improves the impedance matching of the composites, thus achieving superior absorptive performance [24, 25]. Ding et al. [26] fabricated a core-shell nanomaterial (TiO<sub>2</sub>@ Fe<sub>3</sub>O<sub>4</sub>@PPy) via a sequential solvent heat treatment and polymerization process. By manipulating the thickness of the polypyrrole layer, the electromagnetic wave absorption properties can be regulated. This achieves a RL<sub>min</sub> of -61.8 dB at a thickness of 3.2 mm and an EAB value of 6 GHz at an optimal thickness of 2.2 mm. Zhang and collaborators [27] successfully synthesized CF@PANI@Fe<sub>3</sub>O<sub>4</sub> hybrid nanocomposite and also proposed a synergistic mechanism of multiple reflection/scattering, dipole polarization, interfacial polarization, eddy current, and magnetic resonance, thus demonstrating a multiple attenuation method of incident electromagnetic wave, and the RL<sub>min</sub> of 46.86 dB was achieved at a matching thickness of 2.7 mm. Liu's group [28] successfully synthesized structured  $Fe_3O_4@SiO_2@$ PPy microspheres using a microemulsion polymerization technique. The product demonstrated a RL<sub>min</sub> of -40.9 dB at a matching thickness of 5 mm and an EAB value of approximately 6.88 GHz (11.12~18 GHz), which encapsulates the entire K-band (12~18 GHz).

In this paper, a nanocomposite electromagnetic wave absorber, CF@PPy@CoFe2O4, is developed by successively depositing a PPy layer and magnetic  $CoFe_2O_4$  particles onto the CF surface using the electrodeposition and hydrothermal methods, respectively. The PPy layer enhances the impedance matching condition and dielectric loss capability, while the presence of CoFe<sub>2</sub>O<sub>4</sub> nanoparticle provides magnetic loss capability to the bare CF, further amplifying its attenuation capability. Through the use of a large number of heterogeneous interfaces, the coaxial core-shell structures generate robust interfacial polarization, meanwhile, the structural defects and adsorbed functional groups produce dipole polarization, owing to the synergistic effect of these multiple loss mechanisms, the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite exhibits optimal impedance matching and excellent attenuation capabilities.

#### 2 Experimental methods

#### 2.1 Materials and chemicals

The CF substrate used in this study was provided by Shanghai Huayu Instrument Co., Ltd. Various chemical reagents, including hydrochloric acid (HCl), sodium chloride (NaCl), iron chloride hexahydrate (FeCl<sub>3</sub>·6H<sub>2</sub>O), cobalt chloride hexahydrate (CoCl<sub>2</sub>·6H<sub>2</sub>O), and ethylene glycol (EG) were all purchased from Sinopharm Chemical Reagent Co., Ltd. Pyrrole (C<sub>4</sub>H<sub>5</sub>N) and anhydrous sodium acetate (CH<sub>3</sub>COONa) were acquired from Aladdin Reagent. The pyrrole used was of analytical purity grade, whereas the other reagents were of chemical grade and did not require further purification.

#### 2.2 Preparation process of CF@PPy nanocomposite

In this study, a CF@PPy composite was prepared by electrochemical deposition using CF as the substrate. Initially, the CF was segmented into 2 cm x 2 cm sections, then immersed in 1 mol/L hydrochloric acid for 3 h, followed by repeated rinsing with deionized water. In the electrochemical workstation setup, the CF served as the working electrode, a platinum plate electrode was employed as the counter electrode, and the Hg/Hg<sub>2</sub>Cl<sub>2</sub> was used as reference electrode. The electrolyte was prepared with a pyrrole monomer concentration of 0.1 mol/L and a sodium chloride solution concentration of 1 mol/L. Cyclic voltammetry mode (CV) was conducted within a voltage window of  $0 \sim 0.8$  V at a scan rate of 10 mV/s. At a predetermined cycles number, the sample was rinsed with deionized water and dried at 60 °C to obtain the CF@ PPy sample. By varying the number of cycles to 30, 60, 90, and 120, the corresponding resultant samples were labeled as CP-30, CP-60, CP-90, and CP-120, respectively.

The formation mechanism of PPy molecules can be summarized as follows: Firstly, under the influence of an electric field, pyrrole monomer molecules lose electrons at the surface of the electrode, resulting in the formation of cationic free radicals. Then, these pyrrole free radicals interact with another monomer to create a pyrrole dimer. Subsequently, an extensive PPy molecular chain is produced through a series of similar reactions.

## 2.3 Preparation of CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite

A solution was prepared by dissolving 5 mmol FeCl<sub>3</sub>·6H<sub>2</sub>O and 2.5 mmol CoCl<sub>2</sub>·6H<sub>2</sub>O in 70 mL ethylene glycol, followed by vigorous stirring for 30 min. Subsequently, 7.3 g anhydrous sodium acetate was added to the solution and stirred for an additional hour. The resulting mixture was then transferred into a reaction vessel, and the pre-determined CF@PPy sample was also introduced. The reaction was maintained at a temperature of 180 °C for a period of 8 h. The sample was subsequently rinsed with deionized water and dried at 60 °C to yield the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite.

## 2.4 Characterization

The electrochemical deposition process was conducted using a CHI660E electrochemical workstation. The morphology of the

products was examined using scanning electron microscopy (SEM, Hitachi S-8400) and transmission electron microscopy (TEM, Hitachi H-8100), while their phase composition was investigated by a powder X-ray diffraction (XRD, Bruker D8) conducted at room temperature. Fourier infrared spectroscopy (FT-IR, Nicolet-6700) was used to analyze the functional groups of the products, and Raman spectroscopy (Renishaw 2000) was employed to evaluate the graphitization degree and defect content within the samples. The magnetization hysteresis loops of CF@PPy@CoFe2O4 hybrid nanomaterial was recorded using a vibrating sample magnetometer (VSM, Quantum Design, PPMS-9). X-ray photoelectron spectrometer (XPS, Thermo Scientific, ESCALAB 250XI) was used to detect the composition and chemical bonding state of the samples. The N<sub>2</sub> adsorption/desorption isotherm was determined by automatic specific surface and porosity analyzer (Micromeritics ASAP 2460), and the specific surface area of the sample was calculated by Brunauer-Emmett-Teller (BET) the model. Thermogravimetric analysis (TGA) of CF@PPy@ CoFe<sub>2</sub>O<sub>4</sub> nanocomposite was carried out in an argon atmosphere with a heating rate of 10 °C/min. The electromagnetic wave absorption performance of the samples was gauged by measuring electromagnetic parameters using a vector network analyzer (Agilent, N5230A) after compressing a ring with an internal diameter of 3.04 mm, an external diameter of 7.00 mm, and a thickness of 3.00 mm, which was mixed with epoxy resin under a packing load of 2 wt%.

## **3** Results and discussion

Figure 1a depicts a schematic representation of the synthesis process of CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite. The PPy shell was initially wrapped on the surface of CF via an electrodeposited process, resulting in the creation of CF@PPy nanocomposite. Subsequently, magnetic CoFe<sub>2</sub>O<sub>4</sub> nanoparticles were integrated onto the nanocomposite using a straightforward hydrothermal method. To provide a more detailed understanding of the microscopic morphology, SEM was employed for characterization. Figure 1b displays a representative SEM image of a parallel CF sample with a diameter of approximately 8 µm. Notably, numerous subtle ditches aligned along the axis direction of CF provide a wealth of functional groups and active sites, which is beneficial for the subsequent growth of PPy shell. Figure 1c-f display typical SEM images of the CF@PPy nanocomposite, with the PPy shell densely coated on the CF surface, without disturbing the original orientation of the CF. Significantly, an increase in the nanocomposites diameter is observed in line with the number of electrodeposition cycles. Figure S1 provides crosssectional SEM images of various CF@PPy nanocomposites. It can be clearly observed that the variation in the diameter of different nanocomposites can be clearly attributed to the



**Fig. 1** a Schematic illustration of synthesis process of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite; SEM images of the **b** CF, **c** CP-30, **d** CP-60, **e** CP-90, **f** CP-120; **g** Comparison of the diameter values of the CF and the as-prepared CF@PPy samples

increasing thickness of the PPy shell, which correlates directly with the number of electrodeposition cycles. Furthermore, the PPy shell demonstrates an irregular profile, which potentially causes incident electromagnetic wave to undergo multiple reflections, thereby enhancing the attenuation characteristics of the nanocomposites. Crucially, an effective interface is established between the CF and the PPy shell, fostering increased interfacial polarization and promoting the loss of incident electromagnetic waves at the interface.

The phase structure information of CF and CF@PPy samples was analyzed using XRD, with results displayed in Fig. 2a. The XRD patterns of the pure CF sample reveal



two broad diffraction peaks at approximately 25.8° and 44.6°, characteristic of disordered carbon. The diffraction peaks in the CF@PPy sample are notably less intense than in the CF, suggesting the formation of amorphous PPy. To better understand the microstructure of PPy on CF, Raman and FT-IR spectra were captured, as depicted in Fig. 2b-e. Figure 2b shows the Raman spectra of the two samples with characteristic PPy peaks delineated in the highlighted vellow box (Fig. 2c). The peaks at 931 cm<sup>-1</sup> and 970 cm<sup>-1</sup> correspond to the polarization and bipolarization properties of the quinone-type structure, while the peak near 1044 cm<sup>-1</sup> can be assigned to the in-plane deformation of the C-H bonds [29]. Moreover, the peaks observed near 1355 cm<sup>-1</sup> and 1568 cm<sup>-1</sup> can be associated with the ring stretching of C=C bonds and the  $\pi$ -bonded conjugated structure of the polymer backbone [30, 31]. Both samples, due to the presence of carbon materials, exhibit two characteristic peaks around 1340 cm<sup>-1</sup> and 1591 cm<sup>-1</sup>, representing the D-band and G-band, respectively [32]. The D-band arises from the vibrations of defects in the disordered graphitic carbon and sp<sup>3</sup> hybridization, while the G-band is induced by the planar vibrations of sp<sup>2</sup> hybridized graphitic carbon within the carbocyclic ring system. Generally, the relative intensity ratio  $(I_D/I_G)$  of the D and G bands is representative of the degree of graphitization or defects within the carbon material. In the given figure, the  $I_D/I_G$  values of the two materials are 1.17 and 0.52, respectively. The decrease in the  $I_D/I_G$  value of CF@PPy nanocomposite is mainly due to the incorporation of PPy, which results in a greater degreee of  $\pi$ -bond conjugation within the polymer backbone [33]. This leads to an enhanced level of graphitization and a subsequent decrease in the  $I_D/I_G$  value [34].

Figure 2d shows the FT-IR spectra of the sample, with the sections enclosed in the boxes of varying colors representing the distinct peaks of the CF. The peaks located near  $3646 \text{ cm}^{-1}$  and  $1664 \text{ cm}^{-1}$  correspond to the O-H stretching vibration of the hydrogen bond, while the features located near 2912 cm<sup>-1</sup> and 1021 cm<sup>-1</sup> are attributed to the in-plane vibration of the -C-H bond [35]. The characteristic peak appearing near 1385 cm<sup>-1</sup> is due to the C-O bending vibration. Some characteristic peaks of PPy can be observed in the locally zoomed-in image within the yellow box [36]. These encompass the stretching vibration of the C-N bond located near 1382 cm<sup>-1</sup>, the intra-plane deformation vibration of the NH<sup>+</sup> detected near 1094 cm<sup>-1</sup>, and the intra-plane deformation vibration of the N-H bond as well as C-H bond (1048 cm<sup>-1</sup>). The characteristic peaks positioned at 960 cm<sup>-1</sup> and 880 cm<sup>-1</sup> can be respectively ascribed to the ring out-of-plane deformation vibrations of the C-C bond and the ring out-of-plane deformation vibrations of the C-H bond.

XPS was employed to investigate the surface elemental valence states and chemical bonding configurations of CF@PPy. The comprehensive XPS spectrum of the CF@ PPy sample (Fig. 2f) distinctly demonstrates the presence of C, N, and O elements. Compared with bare CF characterization results (Fig. S2) [37], the detection of N can be reasonably attributed to the introduction of PPy [38]. The energy spectrum of O1s (Fig. 2g) reveals a characteristic peak at a binding energy of 531.79 eV, correlating with the surface adsorption of water (OH-) [39]. The C1s energy spectrum of CF@PPy (Fig. 2h) can be deconvoluted into four splitting peaks, corresponding to C-C (284.8 eV), C-N (285.6 eV), C-O (286.3 eV), and C=O (288.2 eV) bonds. The emergence of C-N is most likely due to the existence of PPy [40]. Furthermore, the N1s energy spectrum (Fig. 2i) can be divided into three splitting peaks, indicating that the N atom primarily exists in the form of benzenemethanamine (-NH-), with a lesser amount present as quinoneimine (-N=) and a positively charged nitrogen atom  $(-NH^+=)$  [41]. These findings further corroborate the successful introduction and encapsulation of PPy on the CF surface.

For evaluating the response characteristics of an absorber to electromagnetic waves, two factors are significant, relative complex permittivity ( $\varepsilon_r = \varepsilon' - i\varepsilon''$ ) and complex permeability  $(\mu_r = \mu' - i\mu'')$  [42]. These factors reflect the capability to store and dissipate electric field energy and magnetic field energy, and further determine the reflection coefficient and transmission coefficient of the absorber. Figure 3a-c depict the trends of  $\varepsilon'$ ,  $\varepsilon''$ , and tan $\delta_{\varepsilon}$  curves of both CF and CF@PPy nanocomposites within the frequency range of 2~18 GHz [43]. It is clearly evident from the figure that the CF@PPy nanocomposites exhibit superior values of  $\varepsilon'$ ,  $\varepsilon''$ , and  $\tan \delta_{\varepsilon}$ as compared to bare CF, signifying an enhanced dielectric loss capacity. Moreover, Fig. 3a demonstrates that with an increase in frequency, the curves tend to stabilize in the frequency range of 2~12 GHz, indicating a consistent capability of the absorber to store electrical energy. However, the  $\varepsilon$ ' curves of all samples exhibit fluctuations within the range of 12~18 GHz, likely due to inherent defects. In the  $\varepsilon''$  curve (Fig. 3b), multiple resonance peaks are observed in the CF@ PPy samples, which can be attributed to the existence of numerous heterogeneous interfaces. These interfaces induce interfacial polarization, leading to the appearance of resonance peaks. The dielectric loss capability of an absorber is usually assessed by its dielectric loss tangent, calculated as  $\tan \delta_{\varepsilon} = \varepsilon''/\varepsilon'$ . As depicted in Fig. 3c, the CP-90 sample also exhibits elevated  $\varepsilon''$  values and superior dielectric loss capability, suggesting that the complex permittivity of the sample can be effectively manipulated by adjusting the number of electrodeposition cycles, thereby optimizing its dielectric loss capability [44].

$$Z_{\rm in} = Z_0 \sqrt{\frac{\mu_{\rm r}}{\epsilon_{\rm r}}} \tanh[j(\frac{2\pi\,{\rm fd}}{\rm c})\sqrt{\mu_{\rm r}\epsilon_{\rm r}}] \tag{1}$$

$$RL(dB)=20lg \left| \frac{Z_{in} - Z_0}{Z_{in} + Z_0} \right|$$
(2)

In above formula, the normalized input impedance is denoted as  $Z_{in}$ , while  $Z_0$  represents the impedance in free space [45]. The complex permittivity and complex permeability are  $\varepsilon_r$  and  $\mu_r$ , respectively. The frequency of the incident electromagnetic wave is denoted as f, while the matching thickness of the absorber is symbolized as d, and c denotes the propagation speed of the electromagnetic wave in free space.

Figure 3d illustrates the RL projection curve of the CF sample, which reveals a poor electromagnetic wave absorption performance is revealed with a RL<sub>min</sub> value of -7.93 dB at a matching thickness of 5.50 mm. The absorption performance of the CF@PPy nanocomposites is markedly enhanced with the application of a PPy shell to the CF core. As the thickness of the PPy shell increases, the absorption performance first increases but subsequently decreases, as indicated in Fig. S3. Figure S3a displays RL values versus frequency of the CP-30 nanocomposite, displaying an  $RL_{min}$ value of -10.07 dB at a matching thickness of 2.90 mm. For the CP-60 nanocomposite (Fig. S3b), the RL<sub>min</sub> value is -13.21 dB at a matching thickness of 5.50 mm, and the optimal EAB value is 1.12 GHz. Figure 3e illustrates the performance of the CP-90 nanocomposite with a RL<sub>min</sub> value of -27.98 dB at a matching thickness of 3.05 mm and an EAB value of 2.72 GHz at a matching thickness of 5.25 mm. Figure S3c shows RL values versus frequency of CP-120 nanocomposite, from which it can be seen that the RL<sub>min</sub> value is -13.21 dB when the matching thickness is 2.57 mm, and the optimal EAB value is 1.12 GHz when the matching thickness is 2.60 mm Fig. 3f provides a more visual comparison of the absorption performance of the various CF@ PPy samples with a bar graph of the RL and EAB values. Unquestionably, the CP-90 nanocomposite outperforms the others in terms of absorption performance.

It is widely acknowledged that dipole polarization significantly contributes to the dielectric loss absorbing material, which can be explained by the Debye relaxation phenomenon. The relaxation process of electromagnetic wave



**Fig.3 a** Real part of complex permittivity, **b** Imaginary part of complex permittivity, **c**  $\tan \delta_e$  of CF, CP-30, CP -60, CP -90, CP -120; **d** RL values versus frequency of CF; **e** RL values versus frequency

absorbing materials can be represented by the following equation:

$$\left(\varepsilon' - \frac{\left(\varepsilon_s + \varepsilon_{\infty}\right)}{2}\right)^2 + \left(\varepsilon''\right)^2 = \left(\frac{\varepsilon_s - \varepsilon_{\infty}}{2}\right)^2 \tag{3}$$

herein,  $\varepsilon_s$  and  $\varepsilon_{\infty}$  represent the static and high-frequency optical permittivity, respectively. According to above equation, the relationship between  $\varepsilon'$  and  $\varepsilon''$  should form a semicircle, and each semicircle corresponds to a Debye relaxation

of CP -90; **f** Comparison plots of RL and EAB; **g** Cole-Cole curve of CP -90; **h**  $|Z_{in}/Z_0|$  of CP -90; **i** Attenuation constant of CF, CP -30, CP -60, CP -90, CP -120

process. As illustrated in Figs. 3g and S4, several semicircles can be identified in both the CF and CF@PPy samples, indicating the presence of multiple dielectric relaxations. The CP-90 nanocomposites display more semicircles, which may stem from the interfacial polarization instigated by the hetero-interfaces between CF and PPy, as well as from the inherent dipole polarization of the absorber, suggesting its good electromagnetic wave absorption capacity.

It is universally acknowledged that the impedance matching and attenuation coefficient are crucial when evaluating the efficiency of an absorber designed for electromagnetic wave absorption. An optimal absorber should permit the maximum penetration of electromagnetic waves, necessitating effective impedance matching, which is typically evaluated using the delta function, as delineated in the following formula:

$$|\Delta| = \left|\sinh^2\left(kfd\right) - M\right| \tag{4}$$

in which, the values of K and M can be determined by the relationship between complex permittivity and complex permeability, as expressed in the following equation.

$$K = \frac{4\pi \sqrt{\mu_r \varepsilon_r} \sin\left(\frac{\delta_e + \delta_m}{2}\right)}{c \cos \delta_e \cos \delta_m} \tag{5}$$

$$M = \frac{4\mu' \cos\delta_e \varepsilon' \cos\delta_m}{(\mu' \cos\delta_e - \varepsilon' \cos\delta_m)^2 + [\tan(\frac{\delta_m - \delta_e}{2})]^2 (\mu' \cos\delta_e + \varepsilon' \cos\delta_m)^2}$$
(6)

In general, a material possesses optimal impedance matching when the  $|\Delta|$  value tends to zero. When the  $|\Delta|$ value is less than 0.4, it can be considered that the material exhibits effective impedance matching. Figures 3h and S10 illustrate the impedance matching and  $|\Delta|$  values of various CF@PPy nanocomposites under their optimal conditions, respectively. It can be observed that the CP-90 nanocomposite exhibits a larger mapping area ( $|\Delta| < 0.4$ ) compared to other materials. These findings indicate that CP-90 nanocomposites demonstrate optimal impedance matching, enabling a higher penetration of incident electromagnetic waves into the absorber. The attenuation coefficient is used to assess the dissipation ability of incident electromagnetic waves within the absorbing material. When an absorber exhibits good impedance matching, it is necessary to maximize the dissipation of the incident electromagnetic waves to achieve effective electromagnetic wave absorption. This can be quantified using the following equation.

$$\alpha = \frac{\sqrt{2}\pi f}{c} \sqrt{\left(\mu''\varepsilon'' - \mu'\varepsilon'\right) + \sqrt{\left(\mu''\varepsilon'' - \mu'\varepsilon'\right)^2 + \left(\mu'\varepsilon'' + \mu''\varepsilon'\right)^2}}$$
(7)

As depicted in Fig. 3i, the CF@PPy nanocomposite exhibits a significantly higher attenuation coefficient  $\alpha$  compared to the CF sample. Notably, CP-90 demonstrates the highest  $\alpha$  value (especially in Ku band), indicating its superior capability to dissipate electromagnetic waves energy. These findings highlight the importance of synergistically considering both impedance matching and attenuation capability to achieve outstanding microwave absorption performance. Therefore, it can be observed that the incorporation of the conductive polymer PPy effectively enhances the impedance matching and electromagnetic wave loss capability of CF sample.

To induce magnetic loss within the CF@PPy nanocomposite, CoFe<sub>2</sub>O<sub>4</sub> magnetic nanoparticles are fabricated using CF@PPy as substrate, resulting in the construction of multi-layer heterogeneous CF@PPy@CoFe2O4 nanocomposite. As illustrated in Fig. 4a, b, the CoFe<sub>2</sub>O<sub>4</sub> nanoparticles have been successfully dispersed across the surface of the CF@PPy nanocomposites without any discernible agglomeration. The diameter of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is approximately 18 µm, larger than CF@ PPy nanocomposites. The typical XRD spectrum analysis of  $CF@PPy@CoFe_2O_4$  is presented in Fig. 4c. The broad diffraction peak observed near 26° is attributed to the presence of CF and PPy. Moreover, other distinct diffraction peaks located at 30.08°, 35.50°, 43.05°, 56.98°, and 65.68° correspond to the (220), (311), (400), (511), and (440) crystallographic facets of the CoFe<sub>2</sub>O<sub>4</sub> phase (JCPDS No 22-1086) [46], thereby verifies the successful preparation of  $CoFe_2O_4$ nanoparticles. As PPy and CoFe<sub>2</sub>O<sub>4</sub> were coated on the CF surface step by step, and the CF had good stability under the given reaction conditions, the mass of CF would keep remain unchanged before and after the reaction [47, 48]. To obtain the actual content of CF and PPy in the CF@PPy@ CoFe<sub>2</sub>O<sub>4</sub> nanocomposite, we compared the mass changes of the products before and after the reaction and designed several contrast experiments. The actual mass fractions of various components were obtained, as shown in Table S1. The employed CF sample was 0.055 g. After being coated with the PPy layer, its mass increased to 0.092 g. The final mass of the obtained CF@PPy@CoFe2O4 nanocomposite was 0.104 g. Therefore, it can be calculated that the mass fraction of CF in the product is approximately 53.9 wt%, and the mass fraction of PPy in the product is approximately 34.6 wt%. To further investigate the thermal stability and decomposition behavior of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite, thermogravimetric analysis (TG) coupled with differential scanning calorimetry (DSC) measurements were conducted [9]. The TG-DSC curves are presented in Fig. S6 Within the test temperature range, the mass change of the product can be roughly divided into three distinct stages. In the first stage, from room temperature to 120 °C, a mass loss of 3.59% was observed, which can be attributed to the decomposition of physically adsorbed water and the removal of functional groups on the surface of the sample. The second stage, from 120 to 840 °C, exhibited a significant mass loss of 32.16%, corresponding to the decomposition of PPy. This value is in good agreement with the calculated mass fraction of PPv in the nanocomposites (Table S1). During this stage, the decomposition of PPy involves the evolution of various small-molecule substances, such as water, carbon monoxide, carbon dioxide, and ammonium chloride, in the



**Fig. 4** a, b SEM images, c XRD pattern, d, e Typical TEM images and f-h HRTEM images of  $CoFe_2O_4$  coating capture from CF@PPy@  $CoFe_2O_4$ , i Raman spectra, j hysteresis loop, XPS spectra k full scan, l Co 2p, m Fe 2p, n O 1s of CF@PPy@CoFe\_2O\_4

form of gas. The third stage, above 840 °C, showed a relatively slow and steady mass loss, which can be ascribed to the gradual decomposition of CoFe<sub>2</sub>O<sub>4</sub> and the remaining carbon residue. It is important to note that the mass loss observed in the third stage does not solely correspond to the decomposition of  $CoFe_2O_4$ . Therefore, the value of 1.43% reported in the previous version of the manuscript does not accurately represent the CoFe<sub>2</sub>O<sub>4</sub> content in the nanocomposites. Figure S7a illustrates typical type IV nitrogen adsorption-desorption isotherms of the CF@PPy@CoFe2O4 nanocomposites in the relative pressure  $(P/P_0)$  range of  $0.5 \sim 1.0$ . The BET specific surface area of the sample is evaluated to be around 56.59 m<sup>2</sup>/g, which is larger than that of the employed commercial CF carrier (about 1.31 m<sup>2</sup>/g, provided by the manufacturer). Figure S7b shows the pore size and their distribution obtained from the adsorption branch by the Barrett-Joyner-Halenda (BJH) method, demonstrating that the pore diameter is 3.818 nm.

Systematic characterization results revealed that the CF carrier used in the experiment and the electrodeposited PPy coating are amorphous, while the CoFe<sub>2</sub>O<sub>4</sub> synthesized via the hydrothermal method is crystalline. To further investigate the microstructure of the product, TEM characterization was performed. Figure 4d shows a representative image of the product, with a size of several microns, corresponding to the micro protrusions observed on the fiber surface in Fig. 4b. Figure 4e shows a magnified view of the region marked by the light blue block in Fig. 4d, revealing distinct contrast differences. Figure 4f-h are high-resolution TEM images taken from the corresponding areas in Fig. 4e, respectively. Numerous nanoscale crystal particles are randomly distributed on the product surface, exhibiting clearly visible lattice stripes. Their lattice stripes are clearly visible, and the interplanar crystal spacing corresponds to the (220), (311), (222), (400), (422), (511) planers of cobalt ferrite, respectively.

Additionally, to determine the physical phase composition of the composites, Raman and FT-IR analyses were conducted. The Raman energy spectrum of the CF@PPy@ CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is displayed in Fig. 4i. The D and G bands are discernible around 1350 cm<sup>-1</sup> and 1554 cm<sup>-1</sup>, respectively. The  $I_D/I_G$  ratio value increases to 0.87, indicating the presence of more defect in the CF and polypyrrole due to the introduction of CoFe<sub>2</sub>O<sub>4</sub> nanoparticles. Significantly, the characteristic peaks of PPy remain visible in the magnified image (highlighted in yellow). Additionally, the peaks at 614 cm<sup>-1</sup> and 661 cm<sup>-1</sup> are associated with the  $T_{1g}^{-1}$  and  $A_{1g}^{-1}$  vibrational modes of the CoFe<sub>2</sub>O<sub>4</sub> phase respectively. Figure S8 depicts the FT-IR spectra of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite [49]. No new distinctive peaks emerge when compared with CF@PPy, signifying that the PPy continues to be encapsulated on the CF and that the integration of CoFe<sub>2</sub>O<sub>4</sub> nanoparticles has not compromised its structural integrity. The magnetization capacity of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is characterized by room-temperature hysteresis loops [50], as depicted in Fig. 4j. The saturation magnetization intensity is 26.86 emu/g, while the remanent magnetization intensity (Mr) stands at 4.76 emu/g, as seen in the magnified image. The coercivity (Hc) is measured at 163.3 Oe, denoting the ferromagnetic nature of the CoFe<sub>2</sub>O<sub>4</sub> nanoparticles. This magnetic attribute permits the composite to exhibit magnetic loss contribution, thus enhancing its impedance matching potential and loss capacity.

Figures 4k and S9 presents the XPS full scan of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite, with the emergence of characteristic peaks associated with the elements Fe and Co, signifying the successful incorporation of  $CoFe_2O_4$ nanoparticles into CF@PPy [51]. Figure 41 illustrates the Gaussian-fitted Co 2p spectrum, comprising two primary peaks aligned with the Co  $2p_{1/2}$  and Co  $2p_{3/2}$  energy levels, supplemented by two vibrational satellite peaks. The binding energies of the primary Co 2p<sub>3/2</sub> peaks are observed at 780.5 eV and 781.6 eV, while the peak at 796.4 eV/781.9 eV corresponds to Co  $2p_{1/2}$ , indicating the existence of two distinct valence states of Co. The peaks at 785.7 eV and 803.2 eV are vibrational peaks originating from the coupling of the d-electron orbitals of Co atoms with spin orbita. In the Fe 2p spectrum (Fig. 4m), four main peaks are observed, with additional vibrational satellite peaks at 719.5 eV and 733.7 eV. The peaks situated at 713.7 eV and 727.5 eV mainly derive from Fe<sup>3+</sup> in the tetrahedral sites, while the peaks at 711.6 eV and 723.6 eV are mainly contributed by  $Fe^{2+}$  in the octahedral sites, indicating the presence of two distinct valence states of the Fe elements. Figure 4n depicts the O 1s spectrum, which can be deconvoluted into two distinctive peaks. The peak located at 530.6 eV corresponds to the M-O bond, typical of  $CoFe_2O_4$  (Co single bond O single bond Fe), while the peak at approximately 531.8 eV can be attributed to the O-H bond, originating from surface oxygenated functional groups and adsorbed water molecules [52].

In order to further elucidate the role of  $CoFe_2O_4$  magnetic nanoparticles, the electromagnetic parameters of the  $CF@PPy@CoFe_2O_4$  nanocomposite were examined, as illustrated in Fig. 5a–c. In Fig. 5a, a decrease  $\varepsilon'$  values of the  $CF@PPy@CoFe_2O_4$  nanocomposite is observed with increasing frequency. This trend aligns with the frequency dispersion behavior, mainly attributed to the delayed response rate of dipole polarization to the rapidly alternating electromagnetic field. This emergence of frequency dispersion behavior enables a broader absorption band of incident electromagnetic waves. The enhanced  $\varepsilon'$  value of  $CF@PPy@CoFe_2O_4$  nanocomposite may be ascribed to the incorporation of magnetic  $CoFe_2O_4$  nanoparticles, which creates a greater number of heterogeneous interfaces between the three materials, thereby prompting



**Fig. 5** a Real part of complex permittivity, **b** imaginary part of complex permittivity, **c**  $\tan \delta_e$  of CF, CF@PPy and CF@PPy@CoFe<sub>2</sub>O<sub>4</sub>; **d** RL values versus frequency; **e** conduction loss and **f** polarization loss of CF, CF@PPy and CF@PPy@CoFe<sub>2</sub>O<sub>4</sub>; **g** Cole-Cole curve of CF@PPy@CoFe<sub>2</sub>O<sub>4</sub>; **h** impedance matching of CF@PPy@CoFe<sub>2</sub>O<sub>4</sub>;

i attenuation constant of CF,CF@PPy and CF@PPy@CoFe<sub>2</sub>O<sub>4</sub>; j schematic illustration of conductive loss and polarization loss of CF, CF@PPy and CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> samples; k  $RL_{min}$  versus thickness of microwave absorbing materials reported recently

stronger interfacial polarization and distinct polarization relaxation phenomena. Figure 5b presents the variation curves of  $\varepsilon$ " values with frequency for the three materials. Notably, the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite displays significantly higher values than the other two samples, suggesting its superior microwave absorption properties. The existence of resonance peaks in the high-frequency range also indicates the occurrence of more polarization phenomena. As depicted in Fig. 5c, the CF@PPy@ CoFe<sub>2</sub>O<sub>4</sub> nanocomposite exhibits an enhanced capacity for dielectric loss, particularly within the Ku band. Figure S10a illustrates the variation of  $\mu'$  and  $\mu''$  values with the varying frequency of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite. The  $\mu'$  values range from 0.8 to 1.2, while  $\mu''$ values span from 0 to 0.8. Notably, the complex permeability exhibits resonance peaks in 12~16 GHz, indicating its effective attenuation of incident microwave energy magnetic field. Generally, magnetic loss tangent  $(tan\delta_{\mu})$ , a measure of a material's magnetic loss ability, is characterized by higher values indicating greater dissipation of magnetic field energy. The  $tan \delta_{\mu}$  curve of the CF@ PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite (Fig. S10b) mirrors the trend observed in the imaginary part of the permeability (as shown in Fig. S10a). Notably, in the range of 2~18 GHz, the dielectric loss tangent values overshadow those of the magnetic loss tangent, suggesting that the dielectric loss mechanism plays a dominant role in the attenuation of electromagnetic waves. Figure S10c presents the C<sub>0</sub> curve for the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite. The C<sub>0</sub> value remains low, while in the 2~7 GHz range, it gradually decreases, indicating that magnetic loss is primarily caused by natural ferromagnetic resonance. Additionally, slight fluctuations in  $C_0$  values in the 12 ~ 16 GHz range suggest magnetic losses due to exchange resonance. Figure 5d presents the RL projection drawing of the CF@  $PPy@CoFe_2O_4$  nanocomposite, shows their exceptional microwave absorption properties. Notably, the RL<sub>min</sub> reaches – 55.33 dB at a matching thickness of 2.11 mm, indicating efficient absorption of electromagnetic waves. Furthermore, the EAB spans a broad range of 6.48 GHz at an optimal thickness of 2.42 mm, demonstrating the material's wide absorption band. Figure S11 shows RLfrequency curves, and the relationship between the simulated thickness and the peak frequency for the samples. The quarter-wavelength formula is shown below:

$$t_{\rm m} = \frac{n\lambda}{4} = \frac{nc}{4f_{\rm m}\sqrt{|\mu_{\rm r}||\epsilon_{\rm r}|}} (n = 1, 3, 5 \dots )$$
(8)

where tm is the matching thickness of the absorber,  $\lambda$  is the wavelength of the incident electromagnetic wave, c represents the speed of light in vacuum, and fm is the corresponding matching frequency. It can be clearly seen that the vertical line of the  $RL_{min}$  peak of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is on the 1/4 $\lambda$  curve at the corresponding matching thickness.

To further analyze the dominant role of the loss mechanism in dielectric loss,  $\varepsilon''$  can be divided into two parts, as represented by the following equations:

$$\epsilon'' = \epsilon_p'' + \epsilon_c'' = (\epsilon_s - \epsilon_\infty) \frac{2\pi f \tau}{1 + (2\pi f)^2 \tau^2} + \frac{\sigma}{2\pi f \epsilon_0}$$
(9)

where  $\sigma$  denotes the conductivity of the absorbing material,  $\varepsilon_{p}^{\ \prime\prime}$  and  $\varepsilon_{c}^{\ \prime\prime}$  represent the polarization loss and conductive loss of the absorber, respectively. The values of  $\varepsilon_n^{\prime\prime}$  and  $\varepsilon_{\rm c}{}''$  are obtained using the nonlinear least-squares method of fitting. Typically,  $\varepsilon_c''$  can be equated to the parallel circuit of resistance and capacitance, while  $\varepsilon_{p}^{"}$  represents the series circuit of resistance and capacitance. Since both CF and PPy are dielectric absorbing materials, the main loss mechanism is dielectric loss. Furthermore, the introduction of CoFe<sub>2</sub>O<sub>4</sub> magnetic nanoparticles plays a role in optimizing the impedance matching as well as providing a certain amount of magnetic loss. The contribution of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite to the loss mechanism, including polarization loss and conductive loss, was analyzed and fitted, and the results are depicted in Fig. 5e, f. As revealed in Fig. 5e, the conductive loss curves of all three samples exhibit a gradual decrease with the increasing frequency. The polarization loss of CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is about 1.5 in the high frequency region. Furthermore, the polarization loss curves of the three samples show similar variations to the imaginary part of the complex permittivity, displaying several resonance peaks at high frequency (Fig. 5b). Figure S12 reveals that the polarization loss ratio of the sample is greater than 0.5, indicating that the polarization loss is the dominant factor in the dielectric loss. However, it is crucial to realize that neither of these loss mechanisms can be ignored over the entire range of measured frequencies. In contrast, dielectric loss is the result of the combined action of conductive loss and polarization loss. The dielectric relaxation phenomenon of CF@PPy@CoFe2O4 nanocomposite is investigated using Cole-Cole semicircular curves, as shown in Fig. 5g. The appearance of semicircles indicates the occurrence of Debye relaxation phenomenon, which positively affects the dielectric relaxation of the absorber. Furthermore, the linear tails observed in the Cole-Cole curves suggest that the dominant loss mechanism in the material is conductive loss in the low-frequency range. Furthermore, the  $|\Delta|$  value of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite (Fig. 5h) indicates a significantly larger area ( $|\Delta| < 0.4$ ) compared to the other two reference samples, CF and CF@PPy, suggesting that the introduction of CoFe<sub>2</sub>O<sub>4</sub> magnetic nanoparticles greatly enhances the impedance matching, enabling more efficient penetration of electromagnetic waves into the interior of the absorbing material. In Fig. 5i, it can be observed that the attenuation coefficient of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is higher compared to the other two samples. It suggests that the incident electromagnetic wave can be absorbed to a greater extent after entering the interior of the material.

To gain a deeper understanding of the loss mechanisms of the three materials, a comprehensive comparison of their respective conduction and polarization losses was undertaken, as illustrated in Fig. 5j. It becomes apparent that the integration of PPy shells and CoFe<sub>2</sub>O<sub>4</sub> nanoparticles into the CF matrix creates multiple heterogeneous interfaces, leading to interfacial polarization. Furthermore, these diverse interfaces give rise to a multiple of dipoles that act as polarization centers, facilitating effective absorption of incident electromagnetic waves. Additionally, the introduction of magnetic CoFe<sub>2</sub>O<sub>4</sub> nanoparticles not only enhances the polarization loss but also enriches the magnetic loss mechanisms of the nanocomposite. Consequently, the superior electromagnetic wave absorption capacity of CF@ PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite is governed by the combined effect of conductive loss in the low frequency range and the polarization loss in the high frequency range. In Fig. 5k, we present a comprehensive comparison of the electromagnetic wave absorption properties of the CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite with those of recently reported similar nanocomposites containing conductive polymers, and magnetic particle components [26, 29, 33, 53-59]. Notably, the CF@ PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite stands out with its broader EAB and superior RL<sub>min</sub> compared to the other materials.



Fig. 6 Schematic diagrams of electromagnetic wave attenuation approaches of  $CF@PPy@CoFe_2O_4$  nanocomposite

These remarkable findings highlight the exceptional electromagnetic wave absorption performance of the CF@PPy@ CoFe<sub>2</sub>O<sub>4</sub> nanocomposite, making it a highly promising candidate for applications demanding thin, width band, lightweight, and mechanically robust absorbers.

Figure 6 exhibits the electromagnetic loss mechanism diagram of the CF@PPy@CoFe2O4 nanocomposite, and its outstanding microwave absorption properties can be summarized to the following reasons. Firstly, the highly conductive CF constructs a 3D conductive network, providing efficient pathways for electrons hopping and migration, thereby amplifying the conduction loss of the sample. Additionally, the 3D conductive network triggers a multiple scattering effect, generating several reflections and scattering for the incident electromagnetic wave reaching the absorber surface. Secondly, during the fabrication process of CF@PPy@CoFe2O4 nanocomposite, an abundance of defects and surface functional groups are introduced, resulting in the congregation of dipoles within the product. Influenced by the alternating electromagnetic field, the motion of these dipoles cannot keep up with the periodic changes of the electromagnetic wave, leading to the polarization dissipation of the electromagnetic wave. Additionally, the CoFe<sub>2</sub>O<sub>4</sub> nanoparticles within the nanocomposites present eddy current effects and natural resonance, contributing to the magnetic loss capability of the sample and enriching its microwave absorption performance. Finally, the core-shell structure of CF@PPy@  $CoFe_2O_4$  nanocomposite leads to multiple heterogeneous structures, inducing interfacial polarization phenomena, further enhancing the electromagnetic attenuation capability of the material. In conclusion, these results provide an in-depth exploration of the electromagnetic wave attenuation mechanism of the CF@PPy@CoFe2O4 nanocomposite and highlight the reasons behind their superior microwave absorption properties.

### 4 Conclusions

In this study, a core-shell heterogeneous structure by integrating a dielectric CF core, a conductive polymer PPy interlayer, and magnetic  $CoFe_2O_4$  nanoparticles have been innovatively developed via electrodeposition and hydrothermal methodologies. The synthesized CF@PPy@CoFe<sub>2</sub>O<sub>4</sub> nanocomposite displays excellent electromagnetic wave attenuation ability and optimal impedance matching properties, which can be attributed to the synergistic interplay of the various components. Notably, at a matching thickness of 2.11 mm, the nanocomposite attains a RL<sub>min</sub> of -55.33 dB and an EAB value of 6.48 GHz at a matching thickness of 2.42 mm, covering the entire Ku-band. This superior performance is credited to multiple loss mechanisms, including conduction loss, interfacial/dipole polarization, and natural resonance, which interact with each other and contribute to the electromagnetic wave attenuation of the absorber. Importantly, the introduction of the conductive polymer PPy enhances the impedance matching of the sample, while the employment of  $CoFe_2O_4$  nanoparticles enriches the magnetic loss mechanisms. In conclusion, the  $CF@PPy@CoFe_2O_4$  nanocomposite provides significant insights into the microwave absorption mechanism of core-shell structured nanocomposites and serve as a robust foundation for future investigations in this arena.

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Data availability The data will be available upon proper request.

### Declarations

Competing interests The authors declare no competing interests.

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