



# Approximations to the Euler–Mascheroni Constant

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## Abstract

In this paper, we establish an asymptotic expansion for the Euler–Mascheroni constant. Based on this expansion, we establish a two-sided inequality and a continued fraction approximation for the Euler–Mascheroni constant.

**Keywords** Euler–Mascheroni constant · Psi (or Digamma) function · Polygamma functions · Asymptotic expansion · Inequality

**Mathematics Subject Classification** 41A60; 11Y60, 40A05

## 1 Introduction

Throughout this paper,  $\mathbb{N}$  represents the set of positive integers and  $\mathbb{N}_0 := \mathbb{N} \cup \{0\}$ . The Euler–Mascheroni constant  $\gamma = 0.577215664 \dots$  is defined as the limit of the sequence

$$D_n = \sum_{k=1}^n \frac{1}{k} - \ln n \quad (n \in \mathbb{N}). \quad (1.1)$$

The Euler–Mascheroni constant is a number that appears in analysis and number theory. It is not known yet whether the number is irrational or transcendental [20, 25]. The following two-sided inequality was presented in [24, 28]:

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$$\frac{1}{2(n+1)} < D_n - \gamma < \frac{1}{2n} \quad \text{for } n \in \mathbb{N}, \tag{1.2}$$

which shows that the convergence of the sequence  $D_n$  to  $\gamma$  is very slow (like  $n^{-1}$ ). By changing the logarithmic term in (1.1), faster approximation formulas to the Euler–Mascheroni constant were presented in [7, 10, 16, 17, 23]. For example, Chen and Mortici [7] established the following approximation formula:

$$\sum_{k=1}^n \frac{1}{k} - \ln \left( n + \frac{1}{2} + \frac{1}{24n} - \frac{1}{48n^2} + \frac{23}{5760n^3} \right) = \gamma + O(n^{-5}) \quad \text{as } n \rightarrow \infty \tag{1.3}$$

and posed an open question as follows: For a given  $p \in \mathbb{N}_0$ , find the constants  $a_i$  ( $i = 0, 1, 2, \dots, p$ ), such that

$$\sum_{k=1}^n \frac{1}{k} - \ln \left( n + \sum_{i=0}^p \frac{a_i}{n^i} \right)$$

is the fastest sequence which would converge to  $\gamma$ . This open problem has been considered by Yang [27], Gavrea and Ivan [18], and Lin [21]. Recently, Chen [5] determined the coefficients  $a_j$  and  $b_j$ , such that

$$\sum_{k=1}^n \frac{1}{k} - \ln \frac{n^p + \sum_{j=1}^p a_j n^{p-j}}{n^q + \sum_{j=1}^q b_j n^{q-j}} = \gamma + O\left(\frac{1}{n^{p+q+1}}\right) \quad (n \rightarrow \infty),$$

where  $p \in \mathbb{N}$ ,  $q \in \mathbb{N}_0$  and  $p = q + 1$ . This solves an open problem of Mortici [22]. There are a lot of formulas expressing  $\gamma$  as series, integrals, or products [4, 8, 9, 13, 14, 20]. For more information on the Euler–Mascheroni constant  $\gamma$ , please refer to survey papers [15, 26] and expository book [19].

In this paper, we consider the sequence  $(A_n)_{n \in \mathbb{N}}$  defined by

$$A_n = \sum_{i=1}^n \sum_{j=0}^n \frac{1}{(i+j)^2} - 1 + \ln 2 - \ln n. \tag{1.4}$$

Using the computer program MAPLE 13, we find, as  $n \rightarrow \infty$

$$\begin{aligned} A_n - \gamma \sim & \frac{1}{2n} - \frac{11}{48n^2} + \frac{7}{48n^3} - \frac{31}{640n^4} - \frac{31}{960n^5} + \frac{635}{16128n^6} + \frac{127}{5376n^7} \\ & - \frac{3577}{61440n^8} - \frac{511}{15360n^9} + \dots \end{aligned} \tag{1.5}$$

We here give a formula for determining these coefficients in the right-hand side of (1.5). Then, we establish a two-sided inequality and a continued-fraction approximation for the Euler–Mascheroni constant.

The numerical values given in this paper have been calculated via the computer program MAPLE 13.

### 2 Lemmas and Preliminaries

The gamma function may be defined by  $\Gamma(x) = \int_0^\infty t^{x-1} e^{-t} dt$  ( $x > 0$ ). The logarithmic derivative of the gamma function  $\psi(x) = \Gamma'(x) / \Gamma(x)$  is known as the psi (or digamma) function. The derivatives of the psi function  $\psi^{(n)}(x)$  ( $n \in \mathbb{N}$ ) are called the polygamma functions. It is well known that the psi function has the duplication formula [1, p. 259, Eq. (6.3.8)]:

$$\psi(2x) = \frac{1}{2} \left[ \psi(x) + \psi\left(x + \frac{1}{2}\right) \right] + \ln 2. \tag{2.1}$$

The following series expansion and asymptotic formula hold (see [1, p. 260]):

$$\psi'(x) = \sum_{k=0}^\infty \frac{1}{(x+k)^2} \tag{2.2}$$

and

$$\psi'(x) \sim \frac{1}{x} + \frac{1}{2x^2} + \sum_{k=1}^\infty \frac{B_{2k}}{x^{2k+1}} = \frac{1}{x} + \frac{1}{2x^2} + \sum_{j=2}^\infty \frac{B_j}{x^{j+1}} \quad (x \rightarrow \infty), \tag{2.3}$$

where  $B_n$  ( $n \in \mathbb{N}_0$ ) are the Bernoulli numbers defined by

$$\frac{t}{e^t - 1} = \sum_{n=0}^\infty B_n \frac{t^n}{n!}.$$

It follows from the known result (see [6, Eq. (3.26)]) that:

$$\psi'\left(x + \frac{1}{2}\right) \sim \frac{1}{x} - \sum_{k=2}^\infty \frac{(1 - 2^{1-k})B_k}{x^{k+1}} = \frac{1}{x} - \sum_{k=3}^\infty \frac{(1 - 2^{2-k})B_{k-1}}{x^k} \quad (x \rightarrow \infty). \tag{2.4}$$

**Lemma 2.1** (see [3, Theorem 9]) *Let  $k \geq 1$  and  $n \geq 0$  be integers. Then, for all real numbers  $x > 0$*

$$S_k(2n; x) < (-1)^{k+1} \psi^{(k)}(x) < S_k(2n + 1; x), \tag{2.5}$$

where

$$S_k(p; x) = \frac{(k - 1)!}{x^k} + \frac{k!}{2x^{k+1}} + \sum_{i=1}^p \left[ B_{2i} \prod_{j=1}^{k-1} (2i + j) \right] \frac{1}{x^{2i+k}},$$

$B_n$  are the Bernoulli numbers.

It follows from (2.5) that for  $x > 0$ :

$$\frac{1}{x} + \frac{1}{2x^2} + \frac{1}{6x^3} - \frac{1}{30x^5} < \psi'(x) < \frac{1}{x} + \frac{1}{2x^2} + \frac{1}{6x^3} - \frac{1}{30x^5} + \frac{1}{42x^7}. \tag{2.6}$$

Using the recurrence formula

$$\psi'(x + 1) = \psi'(x) - \frac{1}{x^2}, \tag{2.7}$$

we deduce from (2.6) that for  $x > 0$

$$\frac{1}{x} - \frac{1}{2x^2} + \frac{1}{6x^3} - \frac{1}{30x^5} < \psi'(x + 1) < \frac{1}{x} - \frac{1}{2x^2} + \frac{1}{6x^3} - \frac{1}{30x^5} + \frac{1}{42x^7}. \tag{2.8}$$

**Lemma 2.2** (see [2]) For  $x > \frac{1}{2}$  and  $N \in \mathbb{N}_0$

$$\begin{aligned} \frac{(n - 1)!}{(x - \frac{1}{2})^n} + \sum_{k=1}^{2N+1} \frac{B_{2k}(1/2)}{(2k)!} \frac{(n + 2k - 1)!}{(x - \frac{1}{2})^{n+2k}} &< (-1)^{n+1} \psi^{(n)}(x) \\ &< \frac{(n - 1)!}{(x - \frac{1}{2})^n} + \sum_{k=1}^{2N} \frac{B_{2k}(1/2)}{(2k)!} \frac{(n + 2k - 1)!}{(x - \frac{1}{2})^{n+2k}} \quad (n = 1, 2, \dots), \end{aligned} \tag{2.9}$$

where  $B_n(t)$  are the Bernoulli polynomials defined by the following generating function:

$$\frac{x e^{tx}}{e^x - 1} = \sum_{n=0}^{\infty} B_n(t) \frac{x^n}{n!}.$$

Noting that

$$B_k(1/2) = - \left( 1 - \frac{1}{2^{k-1}} \right) B_k \quad (k \in \mathbb{N}_0),$$

we obtain from (2.9) that for  $x > 0$

$$\frac{1}{x} - \frac{1}{12x^3} + \frac{7}{240x^5} - \frac{31}{1344x^7} < \psi' \left( x + \frac{1}{2} \right) < \frac{1}{x} - \frac{1}{12x^3} + \frac{7}{240x^5}. \tag{2.10}$$

**Lemma 2.3** (see [11, 12]) *Let  $a_1 \neq 0$  and  $A(x) \sim \sum_{j=1}^{\infty} \frac{a_j}{x^j}$  ( $x \rightarrow \infty$ ) be a given asymptotic expansion. Then, the function  $B(x) := a_1/A(x)$  has asymptotic expansion of the following form:*

$$B(x) \sim x + \sum_{j=0}^{\infty} \frac{b_j}{x^j} \quad (x \rightarrow \infty),$$

where

$$b_0 = -\frac{a_2}{a_1}, \quad b_j = -\frac{1}{a_1} \left( a_{j+2} + \sum_{k=1}^j a_{k+1} b_{j-k} \right) \quad (j \geq 1).$$

**Remark 2.4** Lemma 2.3 provides a method to construct a continued-fraction approximation based on a given asymptotic expansion. The details of this method are given below.

Let  $a_1 \neq 0$  and

$$A(x) \sim \sum_{j=1}^{\infty} \frac{a_j}{x^j} \quad (x \rightarrow \infty) \tag{2.11}$$

be a given asymptotic expansion. Then, the asymptotic expansion (2.11) can be transformed into the continued-fraction approximation of the form

$$A(x) \approx \frac{a_1}{x + b_0 + \frac{b_1}{x + c_0 + \frac{c_1}{x + d_0 + \ddots}}} \quad (x \rightarrow \infty), \tag{2.12}$$

wherein the constants are given by the following recurrence relations:

$$\left\{ \begin{array}{l} b_0 = -\frac{a_2}{a_1} \quad \text{and} \quad b_j = -\frac{1}{a_1} \left( a_{j+2} + \sum_{k=1}^j a_{k+1} b_{j-k} \right) \\ c_0 = -\frac{b_2}{b_1} \quad \text{and} \quad c_j = -\frac{1}{b_1} \left( b_{j+2} + \sum_{k=1}^j b_{k+1} c_{j-k} \right) \\ d_0 = -\frac{c_2}{c_1} \quad \text{and} \quad d_j = -\frac{1}{c_1} \left( c_{j+2} + \sum_{k=1}^j c_{k+1} d_{j-k} \right) \\ \dots \quad \dots \end{array} \right. \tag{2.13}$$

Clearly, since  $a_j \implies b_j \implies c_j \implies d_j \implies \dots$ , the asymptotic expansion (2.11) is transformed into the continued-fraction approximation (2.12), and the constants in the right-hand side of (2.12) are determined by the system (2.13).

### 3 Asymptotic Expansion

We provide a recurrence relation for successively determining the coefficients of  $\frac{1}{n^k}$  ( $k \in \mathbb{N}$ ) in expansion (1.5) given by Theorem 3.1.

**Theorem 3.1** *The sequence  $(A_n)_{n \in \mathbb{N}}$ , defined by (1.4), has the asymptotic expansion*

$$A_n - \gamma \sim \sum_{k=1}^{\infty} \frac{a_k}{n^k}, \quad n \rightarrow \infty, \quad (3.1)$$

with the coefficients  $a_k$  given by the recurrence relation

$$a_1 = \frac{1}{2}, \quad a_k = \frac{1}{k} \left\{ - \sum_{j=1}^{k-1} a_j (-1)^{k-j} \binom{k}{k-j+1} - \frac{(4-2^{1-k})B_k}{2} + (-1)^k \left( \frac{k}{2^k} - \frac{3k}{4} + \frac{1}{k+1} \right) \right\} \quad (k \geq 2), \quad (3.2)$$

where  $B_n$  are the Bernoulli numbers.

**Proof** Denote

$$I_n = A_n - \gamma \quad \text{and} \quad J_n = \sum_{k=1}^{\infty} \frac{a_k}{n^k}.$$

In view of (1.5), we can let  $I_n \sim J_n$  and

$$\Delta I_n := I_{n+1} - I_n \sim J_{n+1} - J_n =: \Delta J_n$$

as  $n \rightarrow \infty$ , where  $a_k$  are real numbers to be determined.

We obtain by (2.2) that

$$\sum_{j=0}^n \frac{1}{(i+j)^2} = \psi'(i) - \psi'(i+n+1). \quad (3.3)$$

Using (3.3) and (2.7), we have

$$\begin{aligned} \Delta I_n &= \sum_{i=1}^{n+1} \sum_{j=0}^{n+1} \frac{1}{(i+j)^2} - \sum_{i=1}^n \sum_{j=0}^n \frac{1}{(i+j)^2} - \ln \left( 1 + \frac{1}{n} \right) \\ &= \sum_{i=0}^n \frac{1}{(i+n+2)^2} + \sum_{j=0}^n \frac{1}{(j+n+1)^2} - \ln \left( 1 + \frac{1}{n} \right) \end{aligned}$$

$$\begin{aligned}
 &= \psi'(n + 2) - \psi'(2n + 3) + \psi'(n + 1) - \psi'(2n + 2) - \ln\left(1 + \frac{1}{n}\right) \\
 &= 2\psi'(n) - \frac{2}{n^2} - 2\psi'(2(n + 1)) - \frac{3}{4(n + 1)^2} - \ln\left(1 + \frac{1}{n}\right). \tag{3.4}
 \end{aligned}$$

We obtain from (2.1) that

$$2\psi'(2x) = \frac{1}{2}\psi'(x) + \frac{1}{2}\psi'\left(x + \frac{1}{2}\right). \tag{3.5}$$

We obtain by (3.5) and (2.7) that

$$2\psi'(2(n + 1)) = \frac{1}{2}\psi'(n) - \frac{1}{2n^2} + \frac{1}{2}\psi'\left(n + \frac{1}{2}\right) - \frac{1}{2(n + \frac{1}{2})^2}. \tag{3.6}$$

Substituting (3.6) into (3.4), we obtain

$$\begin{aligned}
 \Delta I_n &= \frac{3}{2}\psi'(n) - \frac{3}{2n^2} - \frac{1}{2}\psi'\left(n + \frac{1}{2}\right) + \frac{1}{2(n + \frac{1}{2})^2} - \frac{3}{4(n + 1)^2} - \ln\left(1 + \frac{1}{n}\right). \\
 &\tag{3.7}
 \end{aligned}$$

Using (2.3) and (2.4), we find as  $n \rightarrow \infty$

$$\begin{aligned}
 \Delta I_n &\sim \frac{3}{2}\left(\frac{1}{n} + \frac{1}{2n^2} + \sum_{j=2}^{\infty} \frac{B_j}{n^{j+1}}\right) - \frac{3}{2n^2} - \frac{1}{2}\left(\frac{1}{n} - \sum_{k=3}^{\infty} \frac{(1 - 2^{2-k})B_{k-1}}{n^k}\right) \\
 &+ \sum_{k=2}^{\infty} \frac{(-1)^k(k-1)}{2^{k-1}n^k} - \sum_{k=2}^{\infty} \frac{(-1)^k 3(k-1)}{4n^k} - \sum_{k=1}^{\infty} \frac{(-1)^{k-1}}{kn^k},
 \end{aligned}$$

which can be written as

$$\begin{aligned}
 \Delta I_n &\sim \sum_{k=2}^{\infty} \left\{ \frac{(4 - 2^{2-k})B_{k-1}}{2} + (-1)^k \left( \frac{k-1}{2^{k-1}} - \frac{3(k-1)}{4} + \frac{1}{k} \right) \right\} \frac{1}{n^k} \quad (n \rightarrow \infty). \\
 &\tag{3.8}
 \end{aligned}$$

Direct computation yields

$$\begin{aligned}
 \sum_{k=1}^{\infty} \frac{a_k}{(n+1)^k} &= \sum_{k=1}^{\infty} \frac{a_k}{n^k} \left(1 + \frac{1}{n}\right)^{-k} = \sum_{k=1}^{\infty} \frac{a_k}{n^k} \sum_{j=0}^{\infty} \binom{-k}{j} \frac{1}{n^j} \\
 &= \sum_{k=1}^{\infty} \frac{a_k}{n^k} \sum_{j=0}^{\infty} (-1)^j \binom{k+j-1}{j} \frac{1}{n^j} = \sum_{k=1}^{\infty} \sum_{j=1}^k a_j (-1)^{k-j} \binom{k-1}{k-j} \frac{1}{n^k}.
 \end{aligned}$$

We then obtain

$$\Delta J_n = \sum_{k=2}^{\infty} \left\{ \sum_{j=1}^k a_j (-1)^{k-j} \binom{k-1}{k-j} - a_k \right\} \frac{1}{n^k}. \tag{3.9}$$

Equating coefficients of the term  $n^{-k}$  on the right-hand sides of (3.8) and (3.9) yields

$$\frac{(4 - 2^{2-k})B_{k-1}}{2} + (-1)^k \left( \frac{k-1}{2^{k-1}} - \frac{3(k-1)}{4} + \frac{1}{k} \right) = \sum_{j=1}^k a_j (-1)^{k-j} \binom{k-1}{k-j} - a_k$$

for  $k \geq 2$ . For  $k = 2$  we obtain  $a_1 = \frac{1}{2}$ , and for  $k \geq 3$ , we have

$$\begin{aligned} & \frac{(4 - 2^{2-k})B_{k-1}}{2} + (-1)^k \left( \frac{k-1}{2^{k-1}} - \frac{3(k-1)}{4} + \frac{1}{k} \right) \\ &= \sum_{j=1}^{k-2} a_j (-1)^{k-j} \binom{k-1}{k-j} - (k-1)a_{k-1}, \end{aligned}$$

which gives the desired result (3.2). The proof of Theorem 3.1 is complete. □

### 4 A Two-Sided Inequality

Motivated by (1.5), we establish a two-sided inequality for the Euler–Mascheroni constant given by Theorem 4.1.

**Theorem 4.1** *Let the sequence  $(A_n)_{n \in \mathbb{N}}$  be defined by (1.4). Then, for  $n \geq 1$*

$$\frac{1}{2n} - \frac{11}{48n^2} < A_n - \gamma < \frac{1}{2n} - \frac{11}{48n^2} + \frac{7}{48n^3}. \tag{4.1}$$

**Proof** We consider the sequences  $(x_n)_{n \in \mathbb{N}}$  and  $(y_n)_{n \in \mathbb{N}}$  defined by

$$x_n := A_n - \gamma - \frac{1}{2n} + \frac{11}{48n^2} \quad \text{and} \quad y_n := A_n - \gamma - \frac{1}{2n} + \frac{11}{48n^2} - \frac{7}{48n^3}.$$

Clearly

$$\lim_{n \rightarrow \infty} x_n = 0 \quad \text{and} \quad \lim_{n \rightarrow \infty} y_n = 0.$$

Noting that (3.7), we obtain by (2.8) and (2.10) that

$$x_{n+1} - x_n = \Delta I_n - \frac{1}{2(n+1)} + \frac{11}{48(n+1)^2} + \frac{1}{2n} - \frac{11}{48n^2}$$



$$\begin{aligned}
 &= \frac{3}{2}\psi'(n+1) - \frac{1}{2}\psi'\left(n + \frac{1}{2}\right) + \frac{1}{2(n+\frac{1}{2})^2} - \frac{3}{4(n+1)^2} - \ln\left(1 + \frac{1}{n}\right) \\
 &\quad - \frac{1}{2(n+1)} + \frac{11}{48(n+1)^2} + \frac{1}{2n} - \frac{11}{48n^2} \\
 &< \frac{3}{2}\left(\frac{1}{n} - \frac{1}{2n^2} + \frac{1}{6n^3} - \frac{1}{30n^5} + \frac{1}{42n^7}\right) \\
 &\quad - \frac{1}{2}\left(\frac{1}{n} - \frac{1}{12n^3} + \frac{7}{240n^5} - \frac{31}{1344n^7}\right) \\
 &\quad + \frac{1}{2(n+\frac{1}{2})^2} - \frac{3}{4(n+1)^2} - \left(\frac{1}{n} - \frac{1}{2n^2} + \frac{1}{3n^3} - \frac{1}{4n^4}\right) \\
 &\quad - \frac{1}{2(n+1)} + \frac{11}{48(n+1)^2} + \frac{1}{2n} - \frac{11}{48n^2} \\
 &= -\frac{P_7(n-2)}{13440n^7(2n+1)^2(n+1)^2}
 \end{aligned}$$

and

$$\begin{aligned}
 y_{n+1} - y_n &= \Delta I_n - \frac{1}{2(n+1)} + \frac{11}{48(n+1)^2} - \frac{7}{48(n+1)^3} + \frac{1}{2n} - \frac{11}{48n^2} + \frac{7}{48n^3} \\
 &= \frac{3}{2}\psi'(n+1) - \frac{1}{2}\psi'\left(n + \frac{1}{2}\right) + \frac{1}{2(n+\frac{1}{2})^2} - \frac{3}{4(n+1)^2} - \ln\left(1 + \frac{1}{n}\right) \\
 &\quad - \frac{1}{2(n+1)} + \frac{11}{48(n+1)^2} - \frac{7}{48(n+1)^3} + \frac{1}{2n} - \frac{11}{48n^2} + \frac{7}{48n^3} \\
 &> \frac{3}{2}\left(\frac{1}{n} - \frac{1}{2n^2} + \frac{1}{6n^3} - \frac{1}{30n^5}\right) - \frac{1}{2}\left(\frac{1}{n} - \frac{1}{12n^3} + \frac{7}{240n^5}\right) \\
 &\quad + \frac{1}{2(n+\frac{1}{2})^2} - \frac{3}{4(n+1)^2} - \left(\frac{1}{n} - \frac{1}{2n^2} + \frac{1}{3n^3} - \frac{1}{4n^4} + \frac{1}{5n^5}\right) \\
 &\quad - \frac{1}{2(n+1)} + \frac{11}{48(n+1)^2} - \frac{7}{48(n+1)^3} + \frac{1}{2n} - \frac{11}{48n^2} + \frac{7}{48n^3} \\
 &= \frac{P_5(n-2)}{480n^5(2n+1)^2(n+1)^3},
 \end{aligned}$$

where

$$\begin{aligned}
 P_7(n) &= 2152565 + 8661410n + 14190237n^2 + 12516340n^3 \\
 &\quad + 6481224n^4 + 1980440n^5 + 331632n^6 + 23520n^7
 \end{aligned}$$

and

$$P_5(n) = 6715 + 31135n + 36739n^2 + 18489n^3 + 4268n^4 + 372n^5.$$

We then obtain  $x_{n+1} < x_n$  and  $y_{n+1} > y_n$  for  $n \geq 2$ . Direct computations give

$$x_1 = -\frac{1}{48} + \ln 2 - \gamma = 0.095098\dots, \quad x_2 = \frac{341}{576} - \gamma = 0.014798\dots,$$

$$y_1 = -\frac{1}{6} + \ln 2 - \gamma = -0.050735\dots, \quad y_2 = \frac{661}{1152} - \gamma = -0.003430\dots$$

We see that the sequence  $(x_n)$  is strictly decreasing and  $(y_n)$  is strictly increasing for  $n \geq 1$ , and we have

$$x_n = A_n - \gamma - \frac{1}{2n} + \frac{11}{48n^2} > \lim_{m \rightarrow \infty} x_m = 0 \quad (n \in \mathbb{N})$$

and

$$y_n = A_n - \gamma - \frac{1}{2n} + \frac{11}{48n^2} - \frac{7}{48n^3} < \lim_{m \rightarrow \infty} y_m = 0 \quad (n \in \mathbb{N}).$$

The proof of Theorem 4.1 is complete. □

### 5 Continued-Fraction Approximation

We convert the asymptotic expansion (3.1) into a continued-fraction approximation given by Theorem 5.1.

**Theorem 5.1** *It is asserted that*

$$A_n - \gamma \approx \frac{\frac{1}{2}}{n + \frac{11}{24} + \frac{-\frac{47}{576}}{n - \frac{5129}{5640} + \frac{\frac{112521}{55225}}{n + \frac{2058785551}{1480776360} + \ddots}}} \quad (n \rightarrow \infty). \quad (5.1)$$

**Proof** By Remark 2.4, we can convert (3.1) into a continued-fraction approximation of the form

$$A_n - \gamma \approx \frac{a_1}{n + b_0 + \frac{b_1}{n + c_0 + \frac{c_1}{n + d_0 + \ddots}}} \quad (n \rightarrow \infty),$$

where the constants in the right-hand side can be determined by using the system (2.13). We see from (1.5) that

$$a_1 = \frac{1}{2}, \quad a_2 = -\frac{11}{48}, \quad a_3 = \frac{7}{48}, \quad a_4 = -\frac{31}{640}, \quad a_5 = -\frac{31}{960}, \quad a_6 = \frac{635}{16128}, \dots$$

We obtain from the first recurrence relation in (2.13) that

$$b_0 = -\frac{a_2}{a_1} = \frac{11}{24},$$

$$\begin{aligned}
 b_1 &= -\frac{a_3 + a_2b_0}{a_1} = -\frac{47}{576}, \\
 b_2 &= -\frac{a_4 + a_2b_1 + a_3b_0}{a_1} = -\frac{5129}{69120}, \\
 b_3 &= -\frac{a_5 + a_2b_2 + a_3b_1 + a_4b_0}{a_1} = \frac{163853}{1658880}, \\
 b_4 &= -\frac{a_6 + a_2b_3 + a_3b_2 + a_4b_1 + a_5b_0}{a_1} = \frac{2749273}{278691840}, \dots
 \end{aligned}$$

We obtain from the second recurrence relation in (2.13) that

$$\begin{aligned}
 c_0 &= -\frac{b_2}{b_1} = -\frac{5129}{5640}, \\
 c_1 &= -\frac{b_3 + b_2c_0}{b_1} = \frac{112521}{55225}, \\
 c_2 &= -\frac{b_4 + b_2c_1 + b_3c_0}{b_1} = -\frac{2058785551}{726761000}, \dots
 \end{aligned}$$

Continuing the above process, we get

$$d_0 = -\frac{c_2}{c_1} = \frac{2058785551}{1480776360}, \dots$$

The proof of Theorem 5.1 is thus completed. □

**Remark 5.2** Based on (5.1), we find the following two-sided inequality:

$$\frac{\frac{1}{2}}{n + \frac{11}{24} + \frac{-\frac{47}{576}}{n - \frac{5129}{5640} + \frac{\frac{112521}{55225}}{n + \frac{2058785551}{1480776360}}}} < A_n - \gamma < \frac{\frac{1}{2}}{n + \frac{11}{24} + \frac{-\frac{47}{576}}{n - \frac{5129}{5640}}} \quad (n \in \mathbb{N}).$$

(5.2)

Following the same method as was used in the proof of Theorem 4.1, we can prove (5.2). We here omit the proof. Elementary calculations show that

$$\begin{aligned}
 &\frac{\frac{1}{2}}{n + \frac{11}{24} + \frac{-\frac{47}{576}}{n - \frac{5129}{5640} + \frac{\frac{112521}{55225}}{n + \frac{2058785551}{1480776360}}}} - \left( \frac{1}{2n} - \frac{11}{48n^2} \right) \\
 &= \frac{1764329280n^2 + 1071186592n + 667905887}{48n^2(252047040n^3 + 236742280n^2 + 229857800n + 60718717)} > 0 (n \geq 1)
 \end{aligned}$$

and

$$\frac{\frac{1}{2}}{n + \frac{11}{24} + \frac{-\frac{47}{576}}{n - \frac{5129}{5640}}} - \left( \frac{1}{2n} - \frac{11}{48n^2} + \frac{7}{48n^3} \right)$$

$$= -\frac{4371n - 6559}{48n^3(1880n^2 - 848n - 937)} < 0 \quad (n \geq 2).$$

Hence, for  $n \geq 2$ , the two-sided inequality (5.2) is more accurate than the two-sided inequality (4.1).

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## Declarations

**Conflict of interest** The authors declare that they have no conflicts of interest.

## References

1. Abramowitz, M., Stegun, I.A. (eds.): Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables, Applied Mathematics Series 55. Ninth Printing, National Bureau of Standards, Washington, D.C. (1972)
2. Allasia, G., Giordano, C., Pečarić, J.: Inequalities for the gamma function relating to asymptotic expansions. *Math. Inequal. Appl.* **5**(3), 543–555 (2002)
3. Alzer, H.: On some inequalities for the gamma and psi functions. *Math. Comput.* **66**, 373–389 (1997)
4. Chen, C.-P.: Inequalities for the Lugo and Euler-Mascheroni constants. *Appl. Math. Lett.* **25**(4), 787–792 (2012)
5. Chen, C.-P.: Approximation formulas and inequalities for the Euler-Mascheroni constant, *Revista de la Real Academia de Ciencias Exactas, Físicas y Naturales. Serie A. Matemáticas*, **115**(2), Article 56 (2021). <https://doi.org/10.1007/s13398-021-00999-4>
6. Chen, C.-P., Choi, J.: Inequalities and asymptotic expansions for the constants of Landau and Lebesgue. *Appl. Math. Comput.* **248**, 610–624 (2014)
7. Chen, C.-P., Mortici, C.: New sequence converging towards the Euler-Mascheroni constant. *Comput. Math. Appl.* **64**, 391–398 (2012)
8. Chen, C.-P., Srivastava, H.M.: New representations for the Lugo and Euler-Mascheroni constants. *Appl. Math. Lett.* **24**(7), 1239–1244 (2011)
9. Chen, C.-P., Srivastava, H.M.: New representations for the Lugo and Euler-Mascheroni constants, II. *Appl. Math. Lett.* **25**(3), 333–338 (2012)
10. Chen, C.-P., Srivastava, H.M., Li, L., Manyama, S.: Inequalities and monotonicity properties for the psi (or digamma) function and estimates for the Euler-Mascheroni constant. *Integral Transforms Spec. Funct.* **22**, 681–693 (2011)
11. Chen, C.-P., Srivastava, H.M., Wang, Q.: A method to construct continued fraction approximations and its applications, *Revista de la Real Academia de Ciencias Exactas, Físicas y Naturales. Serie A. Matemáticas*, **115**(3), Article 97 (2021)
12. Chen, C.-P., Wang, Q.: Asymptotic expansions and continued fraction approximations for the harmonic numbers. *Appl. Anal. Discrete Math.* **13**(2), 569–582 (2019)
13. Choi, J.: Some mathematical constants. *Appl. Math. Comput.* **187**, 122–140 (2007)
14. Choi, J., Srivastava, H.M.: Integral representations for the Euler-Mascheroni constant  $\gamma$ . *Integral Transforms Spec. Funct.* **21**(9), 675–690 (2010)
15. Dence, T.P., Dence, J.B.: A survey of Euler's constant. *Math. Mag.* **82**, 255–265 (2009)

16. DeTemple, D.W.: The non-integer property of sums of reciprocals of consecutive integers. *Math. Gaz.* **75**, 193–194 (1991)
17. DeTemple, D.W.: A quicker convergence to Euler's constant. *Am. Math. Mon.* **100**, 468–470 (1993)
18. Gavrea, I., Ivan, M.: Optimal rate of convergence for sequences of a prescribed form. *J. Math. Anal. Appl.* **402**, 35–43 (2013)
19. Havil, J.: *Gamma: exploring Euler's constant*. Princeton University Press, Princeton (2003)
20. Lagarias, J.C.: Euler's constant: Euler's work and modern developments. *Bull. Am. Math. Soc.* **50**(4), 527–628 (2013)
21. Lin, L.: Asymptotic formulas associated with psi function with applications. *J. Math. Anal. Appl.* **405**, 52–56 (2013)
22. Mortici, C.: On new sequences converging towards the Euler-Mascheroni constant. *Comput. Math. Appl.* **59**, 2610–2614 (2010)
23. Negoi, T.: A faster convergence to the constant of Euler. *Gazeta Matematică, seria A* **15**, 111–113 (1997). ((in Romanian))
24. Rippon, P.J.: Convergence with pictures. *Am. Math. Mon.* **93**, 476–478 (1986)
25. Sondow, J.: Criteria for irrationality of Euler's constant. *Proc. Am. Math. Soc.* **131**(11), 3335–3345 (2003)
26. Srivastava, H.M.: A survey of some recent developments on higher transcendental functions of analytic number theory and applied mathematics. *Symmetry* **13** **2294**, 1–22 (2021)
27. Yang, S.: On an open problem of Chen and Mortici concerning the Euler-Mascheroni constant. *J. Math. Anal. Appl.* **396**, 689–693 (2012)
28. Young, R.M.: Euler's constant. *Math. Gaz.* **75**, 187–190 (1991)

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