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# **Two-Grid Finite Volume Element Methods for Solving Cahn–Hilliard Equation**

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Received: 28 November 2022 / Revised: 16 March 2023 / Accepted: 21 March 2023 / Published online: 18 April 2023 © The Author(s) under exclusive licence to Iranian Mathematical Society 2023

# **Abstract**

This paper proposes a two-grid mixed finite volume element method (TGMFVE) that uses a  $\theta$  time discrete scheme to solve the Cahn–Hilliard equation. This method is separated into two steps. In the first step, the solution of the Cahn–Hilliard equation can be obtained by using a mixed  $\theta$  scheme of the finite volume element method on a coarse grid using an iterative algorithm. The second step involves using the linearized mixed  $\theta$  scheme finite volume element method to solve the equation on a fine grid. The stability analysis of the  $\theta$  scheme of the two-grid mixed finite volume element method has been performed. The priori error estimation for  $L^2$  norm and  $H^1$  norm is also analyzed. The results of theoretical analysis are confirmed by numerical experiments. The results show that the theoretical results match the actual numerical results.

**Keywords** Cahn–Hilliard equation  $\cdot \theta$  scheme  $\cdot$  A priori error estimates  $\cdot$  Stability  $\cdot$ Mixed finite volume element method · Two-grid

**Mathematics Subject Classification** 35Q35 · 74S10

# **1 Introduction**

Let  $\Omega \subset R^d$  (*d* = 2, 3) be a polygon-bounded domain. We think over the listed below phase field Cahn–Hilliard equation proposed by Cahn and Hilliard (see  $[1-3]$  $[1-3]$ ):

Communicated by Davoud Mirzaei.

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<span id="page-1-0"></span>
$$
\begin{cases}\n\frac{\partial y}{\partial t} - \Delta(-\gamma \Delta y + f(y)) = g(x, t), & t \in (0, T], x \in \Omega, \\
y(x, 0) = y_0(x), & x \in \Omega, \\
\frac{\partial y}{\partial n} = \frac{\partial (-\gamma \Delta y + f(y))}{\partial n} = 0, & t \in (0, T], x \in \partial \Omega.\n\end{cases}
$$
\n(1.1)

where *n* denotes the unit outward normal of the boundary  $\partial \Omega$ .  $\gamma$  is a prescribed positive constant, and  $g(x, t)$  is a source term,  $f(y)$  is a given nonlinear one, which satisfies  $| f(y_1) - f(y_2) |$  ≤ *C*| $y_1 - y_2$ | and  $| f(y) |$  ≤ *M*(*y*)|*y*| where *M*(*y*) is a bounded positive function.

Commonly used in fluid interfacial motion analysis, the Cahn–Hillard equation can be utilized to solve relative problems [\[4](#page-32-2)[–9\]](#page-33-0). Unfortunately, the analytical method is not always able to solve many problems related to the Cahn–Hilliard equation due to the non-linearity and fourth-order differential operator. A numerical analysis is often utilized in the study of the dynamics of the Cahn–Hillard equation. The two main types of discrete methods in the field of PDE are the finite element (FE) and the finite volume element (FVE). These two methods are very flexible and can be used in the analysis of complex computational domain (geometric region). Therefore, FE and FVE are considered the industry's first choice when it comes to engineering software. They have a wide variety of applications. In [\[10](#page-33-1)], Chen et al. derived optimal error estimates for both the first- and second-order SAV schemes with the finite element method that is a Galerkin method with standard Lagrange elements based on a mixed variational formulation in space. Ju et al. [\[11\]](#page-33-2) presented a residual-based a posteriori error estimate for the finite volume discretization of steady convectiondiffusionCreaction equations defined on surfaces in  $R<sup>3</sup>$ . In [\[12](#page-33-3)], Hu et al. proposed a finite volume solver to solve 2*D* steady Euler equations. In [\[13](#page-33-4)], Nazari and Sabzevari derived computational bases for finite element spaces  $S_2\Lambda^0(\mathcal{T}_h)$  and  $S_r\Lambda^1(\mathcal{T}_h)$  in each step of the *h*-adaptive method. Du et al. [\[14\]](#page-33-5) considered the phase separation on general surfaces by solving the nonlinear Cahn–Hilliard equation using a finite element method. In  $[15]$ , Jia et al. solved the modified Cahn–Hilliard equation via a large time-stepping mixed finite-element method. Nabet et al. [\[16](#page-33-7)] proposed a numerical scheme to solve a diphasic Cahn–Hilliard equation with dynamic boundary conditions. In [\[17](#page-33-8)], Appadu et al. constructed four finite volume methods to solve the 2*D* convective Cahn–Hilliard equation with specified initial condition and periodic boundary conditions. Thus, the finite volume element algorithm appears to be one of the optimal numerical algorithms concerning solving the Cahn–Hilliard equation as well as accurately capturing the dynamic information of phase transition. Besides ensuring the stability of a complicated system used in long-running numerical simulations, it also satisfies some unique physical properties such as mass conservation and energy decreasing progressively. Nevertheless, the traditional FVE algorithm for calculating the equations of Cahn–Hillard uses Newton's method, which also handles non-linear terms. It is very complex and causes a lot of difficulty when it comes to solving the current complicated phase field problems. For instance, it appears to be used to solve nonlinear equations with low speed and great difficulty. Therefore, it is a challenging problem to reduce the difficulty level of the traditional FVE algorithm by implementing a numerical solution of high-order differential terms and nonlinear terms. Doing so not only effectively improves the accuracy of solving problems compared with traditional FVE algorithm but also saves CPU elapsed time.

The results of numerous numerical tests have indicated that the two grid method is an effective and practical tool for solving the Cahn–Hilliard equation (see, e.g., [\[18](#page-33-9)]) when dealing with the complex nonlinear terms. For the first time, the two-grid method [\[19](#page-33-10), [20\]](#page-33-11) was proposed by professor Xu Jinchao as a discrete method about solving asymmetric indeterminate and nonlinear problems. The basic idea of the twogrid method in the calculation of the Cahn–Hilliard equation is as follows. First, a small-scale nonlinear problem is solved discretely in the coarse grid space. At the moment, the number of unknowns within the coarse grid space is few, which makes the calculations scale very small and easy to calculate. Second, the solution of the coarse grid space is projected into the fine grid space by the interpolation method. The problem of linear approximation is to be solved in a finer grid, which makes it easier to solve than the original problem. The use of this method can help reduce computation time and enhance the efficiency of the solution. The computational method is able to demonstrate its feasibility and effectiveness by accomplishing the goal of reducing the order and the computational time.

However, there is not much research on the two-grid finite volume element method. Therefore, we will use the two-grid mixed finite volume element method coupling  $\theta$ time discrete schemes to solve the Cahn–Hilliard equation problem in this paper.

The rest of the paper is organized as follows: Sect. [2](#page-2-0) develops a two-grid algorithm for solving the Cahn–Hilliard equation, as well as the corresponding time, spatial discrete schemes and a two-grid numerical solution algorithm. In Sect. [3,](#page-6-0) the theoretical analysis we provide for discrete two grid schemes includes the analysis of the stability and error. In Sect. [4,](#page-26-0) some numerical examples are summarized to corroborate the correctness of the result of theoretical derivation. In the end, conclusions are concisely summarized in Sect. [5.](#page-29-0)

Throughout this paper, we put to use standard notations for Sobolev spaces on Ω as in [\[21\]](#page-33-12). For examples,  $L^2(\Omega)$  and  $H^1(\Omega)$  are Hilbert spaces with norms  $\|\cdot\|_{L^2(\Omega)}$  $(\Vert \cdot \Vert_0)$  and  $\Vert \cdot \Vert_{H^1(\Omega)}$  ( $\Vert \cdot \Vert_1$ ). For  $\forall u \in L^2(\Omega)$ , the  $L^2$  norm for *u* could be defined as  $||u||_{L^2(\Omega)} = (\int_{\Omega} |u|^2 dx)^{\frac{1}{2}}$ . For  $\forall u \in H^1(\Omega)$ , the  $H^1$  semi-norm is defined as  $|u|_{H^1(\Omega)} = (\int_{\Omega} |\nabla u|^2 dx)^{\frac{1}{2}}$  and the  $H^1$  norm is defined as  $||u||_{H^1(\Omega)} = (||u||^2_{L^2(\Omega)} +$  $|u|_{H^1(\Omega)}^2)^{\frac{1}{2}}.$ 

## <span id="page-2-0"></span>**2 Two-Grid Algorithm for the Cahn–Hilliard Equation**

In this section, we would give the discretization of the Cahn–Hilliard equation [\(1.1\)](#page-1-0) for mixed finite volume element with  $\theta$  scheme. Let  $w = -\gamma \Delta y + f(y)$ , the mixed variational formulation of  $(1.1)$  is: find  $(y, w)$  such that

<span id="page-2-1"></span>
$$
\begin{cases} \left(\frac{\partial y}{\partial t}, v\right) + (\nabla w, \nabla v) = (g, v), & \text{for all } v \in H^1(\Omega), \\ (w, q) - \gamma(\nabla y, \nabla q) - (f(y), q) = 0, & \text{for all } q \in H^1(\Omega), \end{cases}
$$
(2.1)

where  $(\cdot, \cdot)$  is the inner product on  $\Omega$ , i.e.  $\forall u, v \in L^2(\Omega)$ ,  $(u, v) = \int_{\Omega} u(x)v(x)dx$ ,  $x \in$ Ω.

### **2.1 Temporal Discretization**

We consider the  $\theta$  scheme. Let  $t_k = k\Delta t$  ( $k = 0, 1, 2, \ldots, K$ ) be the nodes in the time interval [0, *T*], where  $t_k$  satisfy  $0 = t_0 < t_1 < t_2 < \cdots < t_K = T$  with  $\Delta t = \frac{T}{K}$ . From [\[22](#page-33-13)], we give  $\theta$  scheme approximation for the function value and the first order derivative value of function  $\phi(t) \in H^1(\Omega)$  at time  $t_{k-\theta}$  with  $\theta \in [0, \frac{1}{2}]$  as

<span id="page-3-0"></span>
$$
\phi(t_{k-\theta}) = (1 - \theta)\phi(t_k) + \theta\phi(t_{k-1}) + O(\Delta t^2) \approx (1 - \theta)\phi^k + \theta\phi^{k-1} = \phi^{k-\theta},
$$
\n(2.2)

$$
\partial_t \phi(t_{k-\theta}) = \frac{(3-2\theta)\phi(t_k) - (4-4\theta)\phi(t_{k-1}) + (1-2\theta)\phi(t_{k-2})}{2\Delta t} + O(\Delta t^2)
$$
  

$$
\approx \frac{(3-2\theta)\phi^k - (4-4\theta)\phi^{k-1} + (1-2\theta)\phi^{k-2}}{2\Delta t} =: D_t \phi^{k-\theta}.
$$
 (2.3)

Based on  $(2.2)$ ,  $(2.3)$ , we give the semi-discrete scheme for  $(2.1)$  as follows: find  $(y^{\frac{1}{2}}, w^{\frac{1}{2}}) \in H^1(\Omega) \times H^1(\Omega)$  such that

<span id="page-3-1"></span>
$$
\begin{cases}\n(D_t y^{\frac{1}{2}}, v) + (\nabla w^{\frac{1}{2}}, \nabla v) = (g(x, t_1), v), & \text{for all } v \in H^1(\Omega), \\
(w^{\frac{1}{2}}, q) - \gamma (\nabla y^{\frac{1}{2}}, \nabla q) - (f(y^{\frac{1}{2}}), q) = 0, & \text{for all } q \in H^1(\Omega),\n\end{cases}
$$
\n(2.4)

for  $k = 2, 3, ..., K$ , find  $(y^{k-\theta}, w^{k-\theta}) \in H^1(\Omega) \times H^1(\Omega)$  such that

<span id="page-3-2"></span>
$$
\begin{cases}\n(D_t y^{k-\theta}, v) + (\nabla w^{k-\theta}, \nabla v) = (g(x, t_{k-\theta}), v), & \text{for all } v \in H^1(\Omega), \\
(w^{k-\theta}, q) - \gamma (\nabla y^{k-\theta}, \nabla q) - (f(y^{k-\theta}), q) = 0, & \text{for all } q \in H^1(\Omega).\n\end{cases}
$$
\n(2.5)

#### **2.2 Spatial Discretization**

We use the mixed finite volume element methods in this paper. Let  $T_h$  be the primal quasi-uniform triangulation of  $\Omega$ , where *h* represents the largest one of the set  $h<sub>\tau</sub>$ diameters in all subdivision triangles  $\tau$ . Based on the primal triangulation, we construct the trial function space  $V_h$  which is composed of linear basis function:

$$
V_h = \{ v \in C(\Omega) : v|_{\tau} \text{ is linear, } \forall \tau \in T_h \}.
$$

Next, we begin to establish dual subdivision  $T_h^*$ . In the previous triangulation mentioned above, we make connections in each triangle  $\tau$ . Let the interior angle of any of  $T_h^*$  be no greater than 90°, take  $Z_{\tau}$  as the barycenter of  $\tau$ ,  $Z_{\tau}$  is the intersection of the mid-lines of  $\tau$  three edges, each triangle  $\tau$  can be divided into three subregions  $\tau_z$  (See Fig. [1\)](#page-4-0), where *z* represents a vertex, also known as node. Let  $\Omega_h$  be a set of the vertices of  $\tau$ . We term the new block formed of subregions  $\tau_z$  shared vertex z as control volume  $V_z$  (See Fig. [2\)](#page-4-1), where  $Z_h(\tau)$  is a set of the barycenter of  $\tau$ . Let *M<sub>h</sub>* be a set of the midpoints of all interior edges *l* of  $T_h$ . Denote by  $Z_h^0$  the set of the interior vertices. Connect the points in the sets  $Z_h(\tau)$ ,  $M_h$  in turn, we can get a polygon domain  $K_z^*$  surrounded by dotted lines in Fig. [2](#page-4-1) around *z*,  $K_z^*$  is called dual



<span id="page-4-0"></span>**Fig. 1** Triangular partition and its dual



<span id="page-4-1"></span>**Fig. 2** The control volume *Vz*

element. All dual elements form  $Ω$  new partition  $T_h^*$  donated as dual partition,  $Z_t$ is called the node of dual partition. The barycenter-type dual partition is easy to be introduced for any triangulation *Th* and will lead to relatively simple calculations. It is well known that the dual partition  $T_h^*$  is quasi-uniform since the triangulation partition *Th* is quasi-uniform. That is to say, existing positive constant *C* makes that

$$
C^{-1}h^2 \le meas(V_z) \le Ch^2, \quad \forall z \in \overline{\Omega}_h
$$

holds, where  $V_z$  is a control volume and  $meas(V_z)$  represents the area of  $V_z$ .

From [\[23](#page-33-14)], we introduce an interpolation operator  $I_h^* : H^1(\Omega) \to V_h^*$ 

<span id="page-5-1"></span>
$$
I_h^* v = \sum_{z \in Z_h^0} v(z) \psi_z,
$$
\n(2.6)

where

$$
V_h^* = \left\{ v \in L^2(\Omega) : v|_{K_z^*} \text{ is constant, } \forall K_z^* \in T_h^* \right\},\
$$

and  $\psi_z$  is characteristic function of control volume  $V_z$ . It is known that  $V_h$  is contained in  $H^1(\Omega)$ , so the interpolation operator  $I_h^*$  can also act on the function  $v_h \in V_h$ . Similarly from [\[23](#page-33-14)], it is known that

$$
\begin{cases} \|I_h^* v\|_{L^2(\Omega)} \le \|v\|_{L^2(\Omega)},\\ \|v - I_h^* v\|_{L^2(\Omega)} \le ch|v|_{H^1(\Omega)}. \end{cases}
$$
\n(2.7)

From [\[23](#page-33-14)], the definition of the bilinear form  $a_h(\cdot, \cdot)$  is as following:

<span id="page-5-2"></span>
$$
a_h(u_h, I_h^* v_h) = -\sum_{z \in Z_h^0} v_h(z) \int_{\partial V_z} \nabla u_h \cdot n \, ds, \ \ \forall u_h, v_h \in V_h,\tag{2.8}
$$

<span id="page-5-0"></span>From Lemma 2.2 in [\[24\]](#page-33-15), we have

**Lemma 2.1** *I* ∗ *<sup>h</sup> is self-adjoint in regard to the L*<sup>2</sup> *inner product,*

<span id="page-5-3"></span>
$$
\left(u_h, I_h^* v_h\right) = \left(v_h, I_h^* u_h\right), \quad \forall u_h, v_h \in V_h. \tag{2.9}
$$

*Define*

$$
\| |u_h\| |_{0} := (u_h, I_h^* u_h)^{\frac{1}{2}}.
$$
\n(2.10)

(2.11)

*Then*  $\|\cdot\|_0$  *and*  $\|\cdot\|_0$  *are equivalent. Here the equivalent constants are independent of h.*

Due to the interpolation operator  $I_h^*$ , we write the full-discrete problem of  $(2.1)$  as follows: find  $(y_h^{\frac{1}{2}}, w_h^{\frac{1}{2}}) \in V_h \times V_h$  such that

$$
\begin{cases}\n\left(D_t y_h^{\frac{1}{2}}, I_h^* v_h\right) + a_h \left(w_h^{\frac{1}{2}}, I_h^* v_h\right) = \left(g\left(x, t_{\frac{1}{2}}\right), I_h^* v_h\right), & \text{for all } v_h \in V_h, \\
\left(w_h^{\frac{1}{2}}, I_h^* q_h\right) - \gamma a_h \left(y_h^{\frac{1}{2}}, I_h^* q_h\right) - \left(f\left(y_h^{\frac{1}{2}}\right), I_h^* q_h\right) = 0, & \text{for all } q_h \in V_h,\n\end{cases}
$$

$$
\underline{\textcircled{2}}
$$
 Springer

for  $k = 2, 3, \ldots, K$ , find  $(y_h^{k-\theta}, w_h^{k-\theta}) \in V_h \times V_h$  such that

$$
\begin{cases}\n\left(D_t y_h^{k-\theta}, I_h^* v_h\right) + a_h \left(w_h^{k-\theta}, I_h^* v_h\right) \\
= \left(g\left(x, t_{k-\theta}\right), I_h^* v_h\right), & \text{for all } v_h \in V_h, \\
\left(w_h^{k-\theta}, I_h^* q_h\right) - \gamma a_h \left(y_h^{k-\theta}, I_h^* q_h\right) - \left(f\left(y_h^{k-\theta}\right), I_h^* q_h\right) \\
= 0, & \text{for all } q_h \in V_h.\n\end{cases}
$$
\n(2.12)

#### **2.3 Two-Grid Algorithm**

The above full-discrete scheme would be built for using the two-grid methods.  $T_H$  and  $T_h$  are given as two triangulations of the domain  $\Omega$  possessing different meshes size *H* and *h* and  $H > h$ ,  $T_H^*$  and  $T_h^*$  be the dual subdivision of  $T_H$  and  $T_h$  respectively. Their associated finite volume element spaces are defined as  $V_H$ ,  $V_h$ ,  $V_H^*$  and  $V_h^*$ , respectively. And the interpolation operators on  $V_H^*$  and  $V_h^*$  are denoted as  $I_H^*$  and  $I_h^*$ , respectively. The two-grid algorithm (cf. [\[25](#page-33-16)]) can be shown as follows: for a general nonlinear PDE, for example of the form  $Lu + Nu - f = 0$  where *Lu* and *Nu* are linear and nonlinear parts, respectively. *f* is the source term.

#### **Two-grid scheme 1**

1. Find  $u_H \in V_H$  such that

$$
(L(u_H),v_H)+(N(u_H),v_H)-(f,v_H)=0 \quad \forall v_H \in V_H;
$$

2. Find  $u_h \in V_h$  such that

$$
(L(u_h), v_h) + (N(u_H), v_h) + (N'(u_H)(u_h - u_H), v_h) - (f, v_h) = 0 \quad \forall v_h \in V_h;
$$

where, an "exact" coarse solver can be used for problems on the coarse grid  $T_H$  at 1, which is generally considered to be a rough area.  $u_H$ ,  $N(u_H)$  and  $N'(u_H)$  are calculated by projecting onto the fine grid  $T_h$ at 2.

Based on  $V_H$ ,  $V_H^*$  and  $V_h$ ,  $V_h^*$ , the following is a two-grid algorithm that can be used for the Cahn–Hilliard equation.

## <span id="page-6-0"></span>**3 Numerical Analysis for Two Grid Discrete Scheme**

Stability and error analysis of the two grid finite volume element with  $\theta$  scheme provided in Algorithm 1 will be shown in this section. The stability of Algorithm 1 is the first thing we shall show.

## **Algorithm** 1. Two grid finite volume element with  $\theta$  scheme.

I: Given  $y_H^0$ , which can be chosen the interpolation of  $y_0$  on the  $V_H$ . For  $k = 1$ , solve the following problem on the coarse grid *TH* ,

<span id="page-7-0"></span>
$$
\begin{cases}\n\left(D_{t} y_{H}^{\frac{1}{2}}, I_{H}^{*} v_{H}\right) + a_{H}\left(w_{H}^{\frac{1}{2}}, I_{H}^{*} v_{H}\right) \\
= \left(g\left(x, t_{\frac{1}{2}}\right), I_{H}^{*} v_{H}\right), & \text{for all } v_{H} \in V_{H}, \\
\left(w_{H}^{\frac{1}{2}}, I_{H}^{*} q_{H}\right) - \gamma a_{H}\left(y_{H}^{\frac{1}{2}}, I_{H}^{*} q_{H}\right) - \left(f\left(y_{H}^{\frac{1}{2}}\right), I_{H}^{*} q_{H}\right) \\
= 0, & \text{for all } q_{H} \in V_{H},\n\end{cases}
$$
\n(2.13)

for  $k = 2, \ldots, K$ , solve the following problem on the coarse grid  $T_H$ ,

<span id="page-7-1"></span>
$$
\begin{cases}\n\left(D_{t}y_{H}^{k-\theta}, I_{H}^{*}v_{H}\right) + a_{H}\left(w_{H}^{k-\theta}, I_{H}^{*}v_{H}\right) \\
= \left(g\left(x, t_{k-\theta}\right), I_{H}^{*}v_{H}\right), & \text{for all } v_{H} \in V_{H}, \\
\left(w_{H}^{k-\theta}, I_{H}^{*}q_{H}\right) - \gamma a_{H}\left(y_{H}^{k-\theta}, I_{H}^{*}q_{H}\right) - \left(f\left(y_{H}^{k-\theta}\right), I_{H}^{*}q_{H}\right) \\
= 0, & \text{for all } q_{H} \in V_{H}.\n\end{cases}
$$
\n(2.14)

II: Solve the linearized Cahn–Hilliard equation on the fine grid  $T_h$ , for  $k = 1$ : find  $(y_h^1, w_h^1) \in V_h \times V_h$ , such that

<span id="page-7-2"></span>
$$
\begin{cases}\n\left(D_{t} y_{h}^{\frac{1}{2}}, I_{h}^{*} v_{h}\right) + a_{h} \left(w_{h}^{\frac{1}{2}}, I_{h}^{*} v_{h}\right) \\
= \left(g\left(x, t_{\frac{1}{2}}\right), I_{h}^{*} v_{h}\right), & \text{for all } v_{h} \in V_{h}, \\
\left(w_{h}^{\frac{1}{2}}, I_{h}^{*} q_{h}\right) - \gamma a_{h} \left(y_{h}^{\frac{1}{2}}, I_{h}^{*} q_{h}\right) - \left(\mathfrak{T}\left(y_{h}^{1}, y_{H}^{1}, y_{h}^{0}\right), I_{h}^{*} q_{h}\right) \\
= 0, & \text{for all } q_{h} \in V_{h},\n\end{cases} \tag{2.15}
$$

for  $k = 2, ..., K$ : find  $(y_h^k, w_h^k) \in V_h \times V_h$ , such that

<span id="page-7-3"></span>
$$
\begin{cases}\n\left(D_{t} y_{h}^{k-\theta}, I_{h}^{*} v_{h}\right) + a_{h} \left(w_{h}^{k-\theta}, I_{h}^{*} v_{h}\right) \\
= \left(g\left(x, t_{k-\theta}\right), I_{h}^{*} v_{h}\right), \text{ for all } v_{h} \in V_{h}, \\
\left(w_{h}^{k-\theta}, I_{h}^{*} q_{h}\right) - \gamma a_{h} \left(y_{h}^{k-\theta}, I_{h}^{*} q_{h}\right) - \left(\mathfrak{T}\left(y_{h}^{k}, y_{H}^{k}, y_{h}^{k-1}\right), I_{h}^{*} q_{h}\right) \\
= 0, \text{ for all } q_{h} \in V_{h},\n\end{cases} \tag{2.16}
$$

where  $\mathfrak{T}(y_h^k, y_H^k, y_h^{k-1}) = (1 - \theta) \mathfrak{s}(y_h^k, y_H^k) + \theta f(y_h^{k-1}), \mathfrak{s}(y_h^k, y_H^k) = f(y_H^k) + f'(y_H^k)(y_h^k - y_H^k);$  $k \ge 1, 0 \le \theta \le \frac{1}{2}$ , when  $k = 1, \theta = \frac{1}{2}$ .

## **3.1 Stability**

The following are the various lemmas we will be introducing.

**Lemma 3.1** *For series*  $\{u^k_{\mathfrak{h}} \in V_{\mathfrak{h}}\}$  *(k*  $\geq$  2*) and*  $\theta \in [0, \frac{1}{2}]$ *, the following inequality holds*

<span id="page-8-2"></span><span id="page-8-0"></span>
$$
\left(D_t u_b^{k-\theta}, I_b^* u_b^{k-\theta}\right) \ge \frac{1}{4\Delta t} \left(H[u_b^k] - H[u_b^{k-1}]\right),\tag{3.1}
$$

$$
\left(D_t \nabla u_b^{k-\theta}, \nabla u_b^{k-\theta}\right) \ge \frac{1}{4\Delta t} \left(\hat{H}[u_b^k] - \hat{H}[u_b^{k-1}]\right),\tag{3.2}
$$

$$
\left(D_t u_{\mathfrak{h}}^{k-\theta}, u_{\mathfrak{h}}^{k-\theta}\right) \ge \frac{1}{4\Delta t} \left(\widetilde{H}[u_{\mathfrak{h}}^k] - \widetilde{H}[u_{\mathfrak{h}}^{k-1}]\right). \tag{3.3}
$$

*where*

$$
H[u_{\mathfrak{h}}^k] = (3 - 2\theta) ||u_{\mathfrak{h}}^k||_0^2 - (1 - 2\theta) ||u_{\mathfrak{h}}^{k-1}||_0^2 + (2 - \theta)(1 - 2\theta) ||u_{\mathfrak{h}}^k - u_{\mathfrak{h}}^{k-1}||_0^2, \quad k \ge 1,
$$
  
\n
$$
\hat{H}[u_{\mathfrak{h}}^k] = (3 - 2\theta) |u_{\mathfrak{h}}^k|_1^2 - (1 - 2\theta) |u_{\mathfrak{h}}^{k-1}|_1^2 + (2 - \theta)(1 - 2\theta) |u_{\mathfrak{h}}^k - u_{\mathfrak{h}}^{k-1}|_1^2, \quad k \ge 1,
$$
  
\n
$$
\widetilde{H}[u_{\mathfrak{h}}^k] = (3 - 2\theta) ||u_{\mathfrak{h}}^k||_0^2 - (1 - 2\theta) ||u_{\mathfrak{h}}^{k-1}||_0^2 + (2 - \theta)(1 - 2\theta) ||u_{\mathfrak{h}}^k - u_{\mathfrak{h}}^{k-1}||_0^2, \quad k \ge 1.
$$

*and*

<span id="page-8-1"></span>
$$
H[u_{\mathfrak{h}}^k] \ge \frac{1}{1-\theta} |||u_{\mathfrak{h}}^k||_0^2, \quad 0 \le \theta \le \frac{1}{2}, \tag{3.4}
$$

$$
\hat{H}[u_{\mathfrak{h}}^k] \ge \frac{1}{1-\theta} |u_{\mathfrak{h}}^k|_1^2, \quad 0 \le \theta \le \frac{1}{2}, \tag{3.5}
$$

$$
\widetilde{H}[u_{\mathfrak{h}}^k] \ge \frac{1}{1-\theta} \|u_{\mathfrak{h}}^k\|_0^2, \quad 0 \le \theta \le \frac{1}{2}.
$$
\n(3.6)

*where*  $\mathfrak{h} = h$  *or*  $H$ *.* 

*Proof* From [\(2.3\)](#page-3-0) we have

$$
D_t u_0^{k-\theta} = \frac{(3-2\theta)u_0^k - (4-4\theta)u_0^{k-1} + (1-2\theta)u_0^{k-2}}{2\Delta t}, \quad k \ge 2.
$$

The operator  $D_t u_{\mathfrak{h}}^{k-\theta}$  can be rewritten as

$$
\begin{cases} D_t u_b^{k-\theta} = (2-2\theta) \frac{u_b^{k}-u_b^{k-1}}{\Delta t} - (1-2\theta) \frac{u_b^{k}-u_b^{k-2}}{2\Delta t}, \\ D_t u_b^{k-\theta} = (\frac{3}{2}-\theta) \frac{u_b^{k}-u_b^{k-1}}{\Delta t} - (\frac{1}{2}-\theta) \frac{u_b^{k-1}-u_b^{k-2}}{\Delta t}. \end{cases}
$$

Then we have

$$
\left(D_t u_{\mathfrak{h}}^{k-\theta}, I_{\mathfrak{h}}^* u_{\mathfrak{h}}^{k-\theta}\right)
$$

$$
= (1 - \theta) \left[ (2 - 2\theta) \left( \frac{u_{\mathfrak{h}}^k - u_{\mathfrak{h}}^{k-1}}{\Delta t}, I_{\mathfrak{h}}^* u_{\mathfrak{h}}^k \right) - (1 - 2\theta) \left( \frac{u_{\mathfrak{h}}^k - u_{\mathfrak{h}}^{k-2}}{2\Delta t}, I_{\mathfrak{h}}^* u_{\mathfrak{h}}^k \right) \right] + \theta \left[ \left( \frac{3}{2} - \theta \right) \left( \frac{u_{\mathfrak{h}}^k - u_{\mathfrak{h}}^{k-1}}{\Delta t}, I_{\mathfrak{h}}^* u_{\mathfrak{h}}^{k-1} \right) - \left( \frac{1}{2} - \theta \right) \left( \frac{u_{\mathfrak{h}}^{k-1} - u_{\mathfrak{h}}^{k-2}}{\Delta t}, I_{\mathfrak{h}}^* u_{\mathfrak{h}}^{k-1} \right) \right].
$$

From Lemma [2.1,](#page-5-0) we have

$$
\left(u_{\mathfrak{h}}^{k} - u_{\mathfrak{h}}^{k-1}, I_{\mathfrak{h}}^{*} u_{\mathfrak{h}}^{k}\right) = \frac{1}{2} \left[\| |u_{\mathfrak{h}}^{k} \| \|_{0}^{2} - \| |u_{\mathfrak{h}}^{k-1} \| \|_{0}^{2} + \| |u_{\mathfrak{h}}^{k} - u_{\mathfrak{h}}^{k-1} \| \|_{0}^{2}\right],
$$
\n
$$
\left(u_{\mathfrak{h}}^{k} - u_{\mathfrak{h}}^{k-2}, I_{\mathfrak{h}}^{*} u_{\mathfrak{h}}^{k}\right) = \frac{1}{2} \left[\| |u_{\mathfrak{h}}^{k} \| \|_{0}^{2} - \| |u_{\mathfrak{h}}^{k-2} \| \|_{0}^{2} + \| |u_{\mathfrak{h}}^{k} - u_{\mathfrak{h}}^{k-2} \| \|_{0}^{2}\right],
$$
\n
$$
\left(u_{\mathfrak{h}}^{k} - u_{\mathfrak{h}}^{k-1}, I_{\mathfrak{h}}^{*} u_{\mathfrak{h}}^{k-1}\right) = \frac{1}{2} \left[\| |u_{\mathfrak{h}}^{k} \| \|_{0}^{2} - \| |u_{\mathfrak{h}}^{k-1} \| \|_{0}^{2} - \| |u_{\mathfrak{h}}^{k} - u_{\mathfrak{h}}^{k-1} \| \|_{0}^{2}\right],
$$
\n
$$
\left(u_{\mathfrak{h}}^{k-1} - u_{\mathfrak{h}}^{k-2}, I_{\mathfrak{h}}^{*} u_{\mathfrak{h}}^{k-1}\right) = \frac{1}{2} \left[\| |u_{\mathfrak{h}}^{k-1} \| \|_{0}^{2} - \| |u_{\mathfrak{h}}^{k-2} \| \|_{0}^{2} + \| |u_{\mathfrak{h}}^{k-1} - u_{\mathfrak{h}}^{k-2} \| \|_{0}^{2}\right].
$$

From the above formula, we can obtain that

$$
\left(D_{t}u_{\mathfrak{h}}^{k-\theta},I_{\mathfrak{h}}^{*}u_{\mathfrak{h}}^{k-\theta}\right) \geq \frac{3-2\theta}{4\Delta t}\left(\| |u_{\mathfrak{h}}^{k}|\|_{0}^{2}-\| |u_{\mathfrak{h}}^{k-1}|\|_{0}^{2}\right)-\frac{1-2\theta}{4\Delta t}\left(\| |u_{\mathfrak{h}}^{k-1}|\|_{0}^{2}-\| |u_{\mathfrak{h}}^{k-2}|\|_{0}^{2}\right) +\frac{2\theta^{2}-5\theta+2}{4\Delta t}\left(\| |u_{\mathfrak{h}}^{k}-u_{\mathfrak{h}}^{k-1}|\|_{0}^{2}-\| |u_{\mathfrak{h}}^{k-1}-u_{\mathfrak{h}}^{k-2}|\|_{0}^{2}\right) =\frac{1}{4\Delta t}\left(H[u_{\mathfrak{h}}^{k}]-H[u_{\mathfrak{h}}^{k-1}]\right).
$$

Further, when  $\theta = \frac{1}{2}$ , we can obviously get

$$
H[u_{\mathfrak{h}}^k] = 2||u_{\mathfrak{h}}^k||_0^2 = \frac{1}{1 - \frac{1}{2}}||u_{\mathfrak{h}}^k||_0^2.
$$

When  $\theta \in [0, \frac{1}{2})$ , from Cauchy–Schwarz inequality, we can get that

$$
H[u_{\mathfrak{h}}^{k}] = (3 - 2\theta) |||u_{\mathfrak{h}}^{k}||_{0}^{2} - (1 - 2\theta) |||u_{\mathfrak{h}}^{k-1}||_{0}^{2} + (2 - \theta)(1 - 2\theta) |||u_{\mathfrak{h}}^{k} - u_{\mathfrak{h}}^{k-1}|||_{0}^{2}
$$
  
\n
$$
\geq (2\theta^{2} - 7\theta + 5) |||u_{\mathfrak{h}}^{k}||_{0}^{2} + (2\theta^{2} - 3\theta + 1) |||u_{\mathfrak{h}}^{k-1}|||_{0}^{2}
$$
  
\n
$$
-[(2\theta^{2} - 3\theta + 1) |||u_{\mathfrak{h}}^{k-1}|||_{0}^{2} + \frac{(2\theta^{2} - 5\theta + 2)^{2}}{2\theta^{2} - 3\theta + 1} |||u_{\mathfrak{h}}^{k-1}|||_{0}^{2}]
$$
  
\n
$$
= \frac{1 - 2\theta}{2\theta^{2} - 3\theta + 1} |||u_{\mathfrak{h}}^{k}|||_{0}^{2} = \frac{1}{1 - \theta} |||u_{\mathfrak{h}}^{k}|||_{0}^{2}.
$$

Similarly we can obtain  $(3.2)$ ,  $(3.3)$ ,  $(3.5)$  and  $(3.6)$ . This completes the proof.  $\square$ 

<span id="page-9-0"></span>From Lemma 2 in [\[26\]](#page-33-17), we have

**Lemma 3.2** *Let*  $u_0$ ,  $v_0 \in V_0$ ,  $I_0^*$  *be defined in* [\(2.6\)](#page-5-1)*, for*  $a_0(\cdot, \cdot)$  *given in* [\(2.8\)](#page-5-2)*, we have* 

$$
a_{\mathfrak{h}}\left(u_{\mathfrak{h}}, I_{\mathfrak{h}}^{*}v_{\mathfrak{h}}\right) = \left(\nabla u_{\mathfrak{h}}, \nabla v_{\mathfrak{h}}\right),
$$
  

$$
a_{\mathfrak{h}}\left(u_{\mathfrak{h}}, I_{\mathfrak{h}}^{*}v_{\mathfrak{h}}\right) = a_{\mathfrak{h}}\left(v_{\mathfrak{h}}, I_{\mathfrak{h}}^{*}u_{\mathfrak{h}}\right),
$$

*where*  $\mathfrak{h} = h$  *or*  $H$ *.* 

Then we consider the following stable inequality.

**Theorem 3.3** *For the coarse solution pair*  $\{y_H^k, w_H^k\} \in V_H \times V_H$ , the stability for the *coupled system* [\(2.13\)](#page-7-0)–[\(2.14\)](#page-7-1) *holds:*

<span id="page-10-1"></span>
$$
||y_H^k||_0^2 + \Delta t \sum_{k=2}^K ||w_H^{k-\theta}||_0^2 \le C ||y_H^0||_0^2 + C \Delta t \sum_{k=0}^K ||g^k||_0^2.
$$
 (3.7)

*For the two-grid solution pair*  $\{y_h^k, w_h^k\} \in V_h \times V_h$ , the stability for the system [\(2.15\)](#page-7-2)– [\(2.16\)](#page-7-3) *holds:*

<span id="page-10-2"></span>
$$
||y_h^k||_0^2 + \Delta t \sum_{k=2}^K ||w_h^{k-\theta}||_0^2 \le C ||y_h^0||_0^2 + C ||y_H^0||_0^2 + C \Delta t \sum_{k=0}^K ||g^k||_0^2,
$$
 (3.8)  

$$
||y_h^k||_1^2 + \Delta t \sum_{k=2}^K ||w_h^{k-\theta}||_1^2 + ||y_H^k||_0^2 \le C ||y_h^0||_0^2 + C ||y_H^0||_0^2 + C \Delta t \sum_{k=0}^K ||g^k||_0^2.
$$
 (3.9)

*Proof* (I) In the coupled system, we take  $v_H = y_H^{k-\theta}$  and  $q_H = w_H^{k-\theta}$  in [\(2.14\)](#page-7-1) to get

$$
\begin{cases}\n\left(D_t y_H^{k-\theta}, I_H^* y_H^{k-\theta}\right) + a_H \left(w_H^{k-\theta}, I_H^* y_H^{k-\theta}\right) = \left(g\left(x, t_{k-\theta}\right), I_H^* y_H^{k-\theta}\right), \\
\left(w_H^{k-\theta}, I_H^* w_H^{k-\theta}\right) - \gamma a_H \left(y_H^{k-\theta}, I_H^* w_H^{k-\theta}\right) - \left(f\left(y_H^{k-\theta}\right), I_H^* w_H^{k-\theta}\right) = 0.\n\end{cases}
$$

Further, we can easily get

$$
a_H\left(y_H^{k-\theta}, I_H^* w_H^{k-\theta}\right) = \frac{1}{\gamma} \left(w_H^{k-\theta}, I_H^* w_H^{k-\theta}\right) - \frac{1}{\gamma} \left(f\left(y_H^{k-\theta}\right), I_H^* w_H^{k-\theta}\right). \tag{3.10}
$$

Based on Lemma [3.2,](#page-9-0) *f* (*y*) satisfy Lipschitz continuity and we use Cauchy– Schwarz inequality as well as Young's inequality to get

<span id="page-10-0"></span>
$$
\left(D_t y_H^{k-\theta}, I_H^* y_H^{k-\theta}\right) + \frac{1}{\gamma} \left(w_H^{k-\theta}, I_H^* w_H^{k-\theta}\right)
$$
  
=  $\frac{1}{\gamma} \left(f\left(y_H^{k-\theta}\right), I_H^* w_H^{k-\theta}\right) + \left(g\left(x, t_{k-\theta}\right), I_H^* y_H^{k-\theta}\right)$ 

$$
\leq \frac{1}{\gamma} \|f\left(y_H^{k-\theta}\right) \|_0 \|I_H^* w_H^{k-\theta} \|_0 + \|g\left(x, t_{k-\theta}\right) \|_0 \|I_H^* y_H^{k-\theta} \|_0
$$
  
\n
$$
\leq C \|y_H^k\|_0^2 + C \|y_H^{k-1}\|_0^2 + C \delta \|w_H^{k-\theta}\|_0^2 + \|g^k\|_0^2 + \|g^{k-1}\|_0^2. \tag{3.11}
$$

Combine Lemmas [2.1,](#page-5-0) [3.1](#page-8-2) with the inequality [\(3.11\)](#page-10-0), let  $\delta = \frac{1}{2C\gamma}$ , we have

<span id="page-11-0"></span>
$$
\frac{1}{4\Delta t} \left( H[y_H^k] - H[y_H^{k-1}] \right) + \frac{1}{2\gamma} ||w_H^{k-\theta}||_0^2
$$
  
\n
$$
\leq C ||y_H^k||_0^2 + C ||y_H^{k-1}||_0^2 + ||g^k||_0^2 + ||g^{k-1}||_0^2.
$$
 (3.12)

Sum [\(3.12\)](#page-11-0) with respect to k from 2 to *K* to get

<span id="page-11-2"></span>
$$
H[y_H^K] + \frac{2\Delta t}{\gamma} \sum_{k=2}^K ||w_H^{k-\theta}||_0^2 \le H[y_H^1] + 8\Delta t \sum_{k=1}^K ||g^k||_0^2 + C\Delta t \sum_{k=1}^K ||y_H^k||_0^2.
$$
\n(3.13)

In the next step, we need to estimate  $H[y_H^1]$ . We take  $v_H = y_H^{\frac{1}{2}}$  and  $q_H = w_H^{\frac{1}{2}}$  in [\(2.13\)](#page-7-0) to get

$$
\begin{cases}\n\left(D_t y_H^{\frac{1}{2}}, I_H^* y_H^{\frac{1}{2}}\right) + a_H \left(w_H^{\frac{1}{2}}, I_H^* y_H^{\frac{1}{2}}\right) = \left(g\left(x, t_{\frac{1}{2}}\right), I_H^* y_H^{\frac{1}{2}}\right), \\
\left(w_H^{\frac{1}{2}}, I_H^* w_H^{\frac{1}{2}}\right) - \gamma a_H \left(y_H^{\frac{1}{2}}, I_H^* w_H^{\frac{1}{2}}\right) - \left(f\left(y_H^{\frac{1}{2}}\right), I_H^* w_H^{\frac{1}{2}}\right) = 0.\n\end{cases}
$$

Further, we can easily get

$$
a_H\left(y_H^{\frac{1}{2}}, I_H^* w_H^{\frac{1}{2}}\right) = \frac{1}{\gamma} \left(w_H^{\frac{1}{2}}, I_H^* w_H^{\frac{1}{2}}\right) - \frac{1}{\gamma} \left(f\left(y_H^{\frac{1}{2}}\right), I_H^* w_H^{\frac{1}{2}}\right).
$$

and on the basis of Lemma [3.2,](#page-9-0)  $f(y)$  satisfy Lipschitz continuity and we use Cauchy–Schwarz inequality as well as Young's inequality to get

<span id="page-11-1"></span>
$$
\begin{aligned}\n\left(D_{t} y_{H}^{\frac{1}{2}}, I_{H}^{*} y_{H}^{\frac{1}{2}}\right) &+ \frac{1}{\gamma} \left(w_{H}^{\frac{1}{2}}, I_{H}^{*} w_{H}^{\frac{1}{2}}\right) \\
&= \frac{1}{\gamma} \left(f\left(y_{H}^{\frac{1}{2}}\right), I_{H}^{*} w_{H}^{\frac{1}{2}}\right) + \left(g\left(x, t_{\frac{1}{2}}\right), I_{H}^{*} y_{H}^{\frac{1}{2}}\right) \\
&\leq C \|y_{H}^{1}\|_{0}^{2} + C \|y_{H}^{0}\|_{0}^{2} + C\delta \|w_{H}^{\frac{1}{2}}\|_{0}^{2} + \frac{1}{2} \|g^{1}\|_{0}^{2} + \frac{1}{2} \|g^{0}\|_{0}^{2}.\n\end{aligned} \tag{3.14}
$$

Combine [\(2.2\)](#page-3-0), [\(2.3\)](#page-3-0) with the inequality [\(3.14\)](#page-11-1), let  $\delta = \frac{1}{2C\gamma}$ , we have

$$
\frac{1}{2\Delta t} \|\left|y_H^1\right\|_0^2 + \frac{1}{2\gamma} \|\left|w_H^{\frac{1}{2}}\right\|_0^2
$$

 $\hat{2}$  Springer

$$
\leq C \|y_H^1\|_0^2 + C \|y_H^0\|_0^2 + \frac{1}{2} \|g^1\|_0^2 + \frac{1}{2} \|g^0\|_0^2.
$$

In the light of  $(3.13)$ , we further get

$$
H[y_H^K] + \frac{2\Delta t}{\gamma} \sum_{k=2}^K ||\omega_H^{k-\theta}||_0^2
$$
  
\n
$$
\leq H[y_H^1] + 8\Delta t \sum_{k=1}^K ||g^k||_0^2 + C\Delta t \sum_{k=1}^K ||y_H^k||_0^2
$$
  
\n
$$
\leq C ||y_H^0||_0^2 + C\Delta t \sum_{k=0}^K ||g^k||_0^2 + C\Delta t \sum_{k=1}^K ||y_H^k||_0^2.
$$

Using Gronwall lemma for the above inequality, we complete the proof of the inequality  $(3.7)$ .

(II) For the fine grid system [\(2.16\)](#page-7-3), we take  $v_h = y_h^{k-\theta}$  and  $q_h = w_h^{k-\theta}$  to get

$$
\begin{cases}\n\left(D_t y_h^{k-\theta}, I_h^* y_h^{k-\theta}\right) + a_h \left(w_h^{k-\theta}, I_h^* y_h^{k-\theta}\right) \\
= \left(g\left(x, t_{k-\theta}\right), I_h^* y_h^{k-\theta}\right), \\
\left(w_h^{k-\theta}, I_h^* w_h^{k-\theta}\right) - \gamma a_h \left(y_h^{k-\theta}, I_h^* w_h^{k-\theta}\right) - \left((1-\theta) \mathfrak{s} \left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right), I_h^* w_h^{k-\theta}\right) \\
= 0.\n\end{cases}
$$

Further, we can easily get

$$
a_h\left(y_h^{k-\theta}, I_h^* w_h^{k-\theta}\right) = \frac{1}{\gamma} \left(w_h^{k-\theta}, I_h^* w_h^{k-\theta}\right)
$$

$$
-\frac{1}{\gamma} \left((1-\theta) \mathfrak{s}\left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right), I_h^* w_h^{k-\theta}\right). \tag{3.15}
$$

Based on Lemma [3.2,](#page-9-0) *f* (*y*) satisfy Lipschitz continuity and we use Cauchy– Schwarz inequality as well as Young's inequality to get

<span id="page-12-0"></span>
$$
\begin{split} &\left(D_{t} y_{h}^{k-\theta}, I_{h}^{*} y_{h}^{k-\theta}\right) + \frac{1}{\gamma} \left(w_{h}^{k-\theta}, I_{h}^{*} w_{h}^{k-\theta}\right) \\ &= \frac{1}{\gamma} \left( (1-\theta) \mathfrak{s} \left(y_{h}^{k}, y_{H}^{k}\right) + \theta \, f\left(y_{h}^{k-1}\right), I_{h}^{*} w_{h}^{k-\theta}\right) + \left(\mathfrak{g}\left(x, t_{k-\theta}\right), I_{h}^{*} y_{h}^{k-\theta}\right) \\ &\leq C \|y_{H}^{k}\|_{0}^{2} + C \|y_{h}^{k}\|_{0}^{2} + C \|y_{h}^{k-1}\|_{0}^{2} + C \delta \|w_{h}^{k-\theta}\|_{0}^{2} + \|g^{k}\|_{0}^{2} + \|g^{k-1}\|_{0}^{2}. \end{split} \tag{3.16}
$$

Combine Lemmas [2.1,](#page-5-0) [3.1,](#page-8-2) take and the inequality [\(3.16\)](#page-12-0), let  $\delta = \frac{1}{2C\gamma}$ , we have

<span id="page-13-0"></span>
$$
\frac{1}{4\Delta t} \left( H[y_h^k] - H[y_h^{k-1}] \right) + \frac{1}{2\gamma} ||[w_h^{k-\theta}||]_0^2
$$
\n
$$
\leq C ||y_H^k||_0^2 + C ||y_h^k||_0^2 + C ||y_h^{k-1}||_0^2 + ||g^k||_0^2 + ||g^{k-1}||_0^2. \tag{3.17}
$$

Sum [\(3.17\)](#page-13-0) with respect to *k* from 2 to *K* to get

$$
H[y_h^K] + \frac{2\Delta t}{\gamma} \sum_{k=2}^K |||w_h^{k-\theta}||_0^2
$$
  
\n
$$
\leq H[y_h^1] + 8\Delta t \sum_{k=1}^K ||g^k||_0^2 + C\Delta t \sum_{k=2}^K ||y_H^k||_0^2 + C\Delta t \sum_{k=1}^K ||y_h^k||_0^2. \quad (3.18)
$$

For  $H[y_h^1]$ , using the similar process to that as above, and applying Lemma [3.1,](#page-8-2) we have

<span id="page-13-2"></span>
$$
\frac{1}{1-\theta}||y_h^K||_0^2 + \frac{2\Delta t}{\gamma} \sum_{k=2}^K ||w_h^{k-\theta}||_0^2
$$
\n
$$
\leq C||y_h^0||_0^2 + C\Delta t \sum_{k=0}^K ||g^k||_0^2 + C\Delta t \sum_{k=1}^K ||y_H^k||_0^2 + C\Delta t \sum_{k=1}^K ||y_h^k||_0^2. \quad (3.19)
$$

Using Gronwall lemma and the conclusion of  $(3.7)$  for the above inequality, we complete the proof of the inequality  $(3.8)$ .

(III) Now we give the estimate of inequality [\(3.9\)](#page-10-2). Take  $v_h = w_h^{k-\theta}$  and  $q_h = D_t y_h^{k-\theta}$ in  $(2.16)$  to get

$$
\begin{cases}\n\left(D_{t}y_{h}^{k-\theta}, I_{h}^{*}w_{h}^{k-\theta}\right) + a_{h}\left(w_{h}^{k-\theta}, I_{h}^{*}w_{h}^{k-\theta}\right) = \left(g\left(x, t_{k-\theta}\right), I_{h}^{*}w_{h}^{k-\theta}\right), \\
\left(w_{h}^{k-\theta}, I_{h}^{*}D_{t}y_{h}^{k-\theta}\right) - \gamma a_{h}\left(y_{h}^{k-\theta}, I_{h}^{*}D_{t}y_{h}^{k-\theta}\right) - \left(\mathfrak{T}\left(y_{h}^{k}, y_{H}^{k}, y_{h}^{k-1}\right), I_{h}^{*}D_{t}y_{h}^{k-\theta}\right) = 0.\n\end{cases}
$$

Subtract the above formula, from  $(2.9)$  we have

$$
a_h\left(w_h^{k-\theta}, I_h^* w_h^{k-\theta}\right) + \gamma a_h\left(y_h^{k-\theta}, I_h^* D_t y_h^{k-\theta}\right)
$$
  
=  $\left(g\left(x, t_{k-\theta}\right), I_h^* w_h^{k-\theta}\right) - \left((1-\theta) \mathfrak{s}\left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right), I_h^* D_t y_h^{k-\theta}\right).$   
(3.20)

We use Cauchy–Schwarz inequality as well as Young's inequality for the above equation, and *f* (*y*) satisfy Lipschitz continuity

<span id="page-13-1"></span>
$$
a_h\left(w_h^{k-\theta}, I_h^* w_h^{k-\theta}\right) + \gamma a_h\left(y_h^{k-\theta}, I_h^* D_t y_h^{k-\theta}\right)
$$
  
\n
$$
\leq \|g\left(x, t_{k-\theta}\right)\|_0 \|I_h^* w_h^{k-\theta}\|_0 + \|(1-\theta)\mathfrak{s}\left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right)\|_0 \|I_h^* D_t y_h^{k-\theta}\|_0
$$

$$
\leq C \|g^k\|_0^2 + C \|g^{k-1}\|_0^2 + C \delta \|w_h^{k-\theta}\|_0^2 + C \|y_H^k\|_0^2 + C \|y_h^k\|_0^2 + C \|y_h^{k-1}\|_0^2
$$
  
+
$$
C \|D_t y_h^{k-\theta}\|_0^2.
$$
 (3.21)

Combine Lemmas [3.1,](#page-8-2) [3.2](#page-9-0) with the inequality [\(3.21\)](#page-13-1) and take  $\delta$  as suitable constant, we can get

<span id="page-14-0"></span>
$$
\frac{\gamma}{4\Delta t}(\hat{H}[y_h^k] - \hat{H}[y_h^{k-1}]) + |w_h^{k-\theta}|_1^2
$$
\n
$$
\leq C \|g^k\|_0^2 + C \|g^{k-1}\|_0^2 + C \|y_H^k\|_0^2 + C \|y_h^k\|_0^2 + C \|y_h^{k-1}\|_0^2 + C \|D_t y_h^{k-\theta}\|_0^2.
$$
\n(3.22)

Sum [\(3.22\)](#page-14-0) with respect to k from 2 to K to get

<span id="page-14-2"></span>
$$
\hat{H}[y_h^K] + \frac{4\Delta t}{\gamma} \sum_{k=2}^K |w_h^{k-\theta}|_1^2
$$
\n
$$
\leq \hat{H}[y_h^1] + C\Delta t \sum_{k=2}^K \|g^k\|_0^2 + C\Delta t \sum_{k=2}^K \|g^{k-1}\|_0^2 + C\Delta t \sum_{k=2}^K \|y_H^k\|_0^2
$$
\n
$$
+ C\Delta t \sum_{k=2}^K \|y_h^k\|_0^2 + C\Delta t \sum_{k=2}^K \|y_h^{k-1}\|_0^2 + C\Delta t \sum_{k=2}^K \|D_t y_h^{k-\theta}\|_0^2. \quad (3.23)
$$

Then we need to estimate  $||D_t y_h^{k-\theta}||_0$ . For  $k-2$ ,  $k-1$ ,  $k$ , let  $\theta = 0$  in the second formula of [\(2.16\)](#page-7-3)

$$
\begin{cases}\n\left(w_h^k, I_h^* q_h\right) - \gamma a_h \left(y_h^k, I_h^* q_h\right) = \left(\mathfrak{s}\left(y_h^k, y_H^k\right), I_h^* q_h\right), \\
\left(w_h^{k-1}, I_h^* q_h\right) - \gamma a_h \left(y_h^{k-1}, I_h^* q_h\right) = \left(\mathfrak{s}\left(y_h^{k-1}, y_H^{k-1}\right), I_h^* q_h\right), \\
\left(w_h^{k-2}, I_h^* q_h\right) - \gamma a_h \left(y_h^{k-2}, I_h^* q_h\right) = \left(\mathfrak{s}\left(y_h^{k-2}, y_H^{k-2}\right), I_h^* q_h\right).\n\end{cases}
$$

Then from [\(2.2\)](#page-3-0), we have

<span id="page-14-1"></span>
$$
\left(D_t w_h^{k-\theta}, I_h^* q_h\right) - \gamma a_h \left(D_t y_h^{k-\theta}, I_h^* q_h\right)
$$
\n
$$
= \frac{(3-2\theta)}{2\Delta t} \left(\mathfrak{s}\left(y_h^k, y_H^k\right), I_h^* q_h\right) - \frac{(4-4\theta)}{2\Delta t} \left(\mathfrak{s}\left(y_h^{k-1}, y_H^{k-1}\right), I_h^* q_h\right)
$$
\n
$$
+ \frac{(1-2\theta)}{2\Delta t} \left(\mathfrak{s}\left(y_h^{k-2}, y_H^{k-2}\right), I_h^* q_h\right). \tag{3.24}
$$

Taking  $q_h = w_h^{k-\theta}$  in [\(3.24\)](#page-14-1),  $v_h = \gamma D_t y_h^{k-\theta}$  in the first formula of [\(2.16\)](#page-7-3) and adding the resulting relations, we obtain

$$
\left(D_t w_h^{k-\theta}, I_h^* w_h^{k-\theta}\right) + \gamma \left(D_t y_h^{k-\theta}, I_h^* D_t y_h^{k-\theta}\right)
$$
  
=  $\gamma \left(g\left(x, t_{k-\theta}\right), I_h^* D_t y_h^{k-\theta}\right) + \frac{(3-2\theta)}{2\Delta t} \left(\mathfrak{s}\left(y_h^k, y_H^k\right), I_h^* w_h^{k-\theta}\right)$ 

$$
-\frac{(4-4\theta)}{2\Delta t}\left(\mathfrak{s}\left(y_h^{k-1}, y_H^{k-1}\right), I_h^* w_h^{k-\theta}\right) + \frac{(1-2\theta)}{2\Delta t}\left(\mathfrak{s}\left(y_h^{k-2}, y_H^{k-2}\right), I_h^* w_h^{k-\theta}\right). \tag{3.25}
$$

From Lemmas  $3.1$ ,  $3.2$ ,  $2.1$  and as the process from  $(3.16)$  to  $(3.19)$ , we can get

<span id="page-15-0"></span>
$$
C\Delta t \sum_{k=2}^{K} \|D_t y_h^{k-\theta}\|_0^2 \le C\Delta t \sum_{k=0}^{K} \|g^k\|_0^2 + C\|y_H^0\|_0^2 + C\|y_h^0\|_0^2. \tag{3.26}
$$

Combine  $(3.26)$  and  $(3.23)$ , we can arrive at the conclusion  $(3.9)$ .

### **3.2 Error Analysis**

<span id="page-15-2"></span>The Ritz projection operator from [\[27\]](#page-33-18) should be given to us as the first step in carrying out the error analysis.

**Lemma 3.4** *Define the Ritz projection operator*  $R_h : H^1(\Omega) \to V_h$  *as* 

<span id="page-15-1"></span>
$$
a_{\mathfrak{h}}(u - R_{\mathfrak{h}}u, I_h^* v_{\mathfrak{h}}) = 0, \quad \forall v_{\mathfrak{h}} \in V_{\mathfrak{h}}, \tag{3.27}
$$

*with the estimate inequality*

$$
||u - R_{\mathfrak{h}}u||_0 + ||(u - R_{\mathfrak{h}}u)_t||_0 \le C\mathfrak{h}^2, \quad \forall u \in H^2(\Omega),
$$
  

$$
||u - R_{\mathfrak{h}}u||_1 \le C\mathfrak{h}||u||_{3,p} \quad (p > 1),
$$

*where*  $\mathfrak{h} = h$  *or*  $H$ .

<span id="page-15-4"></span>In the following, we would show the error estimation between finite volume element solution and semi-discrete solution.

**Theorem 3.5** *Let*  $y^k$ *,*  $w^k$  *be the solution of semi-system* [\(2.4\)](#page-3-1)–[\(2.5\)](#page-3-2)*,*  $(y_H^k, w_H^k)$  *be the*  $coarse$  grid solution of system [\(2.13\)](#page-7-0)–[\(2.14\)](#page-7-1), ( $y_h^k$  ,  $w_h^k$ ) be the fine grid solution of system  $(2.15)-(2.16)$  $(2.15)-(2.16)$  $(2.15)-(2.16)$ , respectively. With  $y_h^0 = R_h y_0$ ,  $y_H^0 = R_H y_0$ , there exists a constant C *independent of h, H,*  $\Delta t$ *, such that* 

<span id="page-15-3"></span>
$$
||y^{k} - y_{H}^{k}||_{0} + \left(\Delta t \sum_{k=1}^{K} ||w^{k-\theta} - w_{H}^{k-\theta}||_{0}^{2}\right)^{\frac{1}{2}} \le CH^{2},
$$
\n(3.28)

$$
||y^{k} - y_{h}^{k}||_{0} + \left(\Delta t \sum_{k=1}^{K} ||w^{k-\theta} - w_{h}^{k-\theta}||_{0}^{2}\right)^{\frac{1}{2}} \le Ch^{2} + CH^{4}, \qquad (3.29)
$$

$$
\|y^{k} - y_{h}^{k}\|_{1} + \left(\Delta t \sum_{k=1}^{K} \|w^{k-\theta} - w_{h}^{k-\theta}\|_{1}^{2}\right)^{\frac{1}{2}} + \|y^{k} - y_{H}^{k}\|_{0} \le Ch + CH^{2}.
$$
 (3.30)

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**Proof** (I) In the coarse grid, let  $y_H^k - y^k = y_H^k - R_H y^k + R_H y^k - y^k = \sigma^{y,k,H} + \rho^{y,k,H}$ , and  $w_H^k - w^k = w_H^k - R_H w^k + R_H w^k - w^k = \sigma^{w,k,H} + \rho^{w,k,H}$ . Combine the definition of Ritz projection operator  $(3.27)$  and  $(2.14)$ , we can get

<span id="page-16-0"></span>
$$
(D_t \sigma^{y, k-\theta, H}, I_H^* v_H) + a_H(\sigma^{w, k-\theta, H}, I_H^* v_H)
$$
  
=  $(\partial_t y^{k-\theta} - D_t R_H y^{k-\theta}, I_H^* v_H),$  (3.31)

and

<span id="page-16-1"></span>
$$
(\sigma^{w,k-\theta,H}, I_H^* q_H) - \gamma a_H(\sigma^{y,k-\theta,H}, I_H^* q_H)
$$
  
=  $(f(y_H^{k-\theta}) - f(y^{k-\theta}), I_H^* q_H) + (w^{k-\theta} - R_H w^{k-\theta}, I_H^* q_H)$ . (3.32)

Let  $v_H = \sigma^{y,k-\theta,H}$  in [\(3.31\)](#page-16-0),  $q_H = \frac{\sigma^{w,k-\theta,H}}{\gamma}$  in [\(3.32\)](#page-16-1), and add two formulas, we have

1

$$
(D_t \sigma^{y,k-\theta,H}, I_H^* \sigma^{y,k-\theta,H}) + \frac{1}{\gamma} (\sigma^{w,k-\theta,H}, I_H^* \sigma^{w,k-\theta,H})
$$
  
+  $a_H (\sigma^{w,k-\theta,H}, I_H^* \sigma^{y,k-\theta,H}) - a_H (\sigma^{y,k-\theta,H}, I_H^* \sigma^{w,k-\theta,H})$   
=  $(\partial_t y^{k-\theta} - D_t R_H y^{k-\theta}, I_H^* \sigma^{y,k-\theta,H}) + \frac{1}{\gamma} (f(y_H^{k-\theta}) - f(y^{k-\theta}), I_H^* \sigma^{w,k-\theta,H})$   
+  $\frac{1}{\gamma} (w^{k-\theta} - R_H w^{k-\theta}, I_H^* \sigma^{w,k-\theta,H}).$ 

From Lemma [3.2,](#page-9-0) we use Cauchy–Schwarz inequality as well as Young's inequality to get

<span id="page-16-2"></span>
$$
(D_t \sigma^{y,k-\theta,H}, I_H^* \sigma^{y,k-\theta,H}) + \frac{1}{\gamma} (\sigma^{w,k-\theta,H}, I_H^* \sigma^{w,k-\theta,H})
$$
  
\n
$$
= (\partial_t y^{k-\theta} - D_t R_H y^{k-\theta}, I_H^* \sigma^{y,k-\theta,H}) + \frac{1}{\gamma} (f(y_H^{k-\theta}) - f(y^{k-\theta}), I_H^* \sigma^{w,k-\theta,H})
$$
  
\n
$$
+ \frac{1}{\gamma} (w^{k-\theta} - R_H w^{k-\theta}, I_H^* \sigma^{w,k-\theta,H})
$$
  
\n
$$
\leq \frac{C}{\Delta t} \int_{t_{k-2}}^{t_k} ||(y - R_H y)_t||_0^2 dt + C ||\sigma^{y,k-\theta,H}||_0^2 + C ||\rho^{w,k-\theta,H}||_0^2 + C \delta ||\sigma^{w,k-\theta,H}||_0^2
$$
  
\n
$$
+ C(||y^k - y_H^k||_0^2 + ||y^{k-1} - y_H^{k-1}||_0^2).
$$
\n(3.33)

From Lemma [3.1,](#page-8-2) we know that

<span id="page-16-3"></span>
$$
\left(D_t\sigma^{y,k-\theta,H}, I_H^*\sigma^{y,k-\theta,H}\right) \ge \frac{1}{4\Delta t} \left(H[\sigma^{y,k,H}] - H[\sigma^{y,k-1,H}]\right). \tag{3.34}
$$

Combine the inequality [\(3.33\)](#page-16-2), [\(3.34\)](#page-16-3), Lemma [2.1](#page-5-0) and let  $\delta$  as a suitable value, we can get

<span id="page-16-4"></span>
$$
\frac{1}{4\Delta t}\left(H[\sigma^{y,k,H}]-H[\sigma^{y,k-1,H}]\right)+\frac{1}{\gamma}\||\sigma^{w,k-\theta,H}\||_0^2
$$

$$
\leq \frac{C}{\Delta t} \int_{t_{k-2}}^{t_k} \| (y - R_H y)_t \|_0^2 dt + C \| \sigma^{y, k-\theta, H} \|_0^2 + C \| \rho^{w, k-\theta, H} \|_0^2
$$
  
+
$$
C \left( \| y^k - y_H^k \|_0^2 + \| y^{k-1} - y_H^{k-1} \|_0^2 \right).
$$
 (3.35)

Sum [\(3.35\)](#page-16-4) with respect to *k* from 2 to *K* to get

<span id="page-17-2"></span>
$$
H[\sigma^{y,K,H}] + \frac{4\Delta t}{\gamma} \sum_{k=2}^{K} ||\sigma^{w,k-\theta,H}||_{0}^{2} \le H[\sigma^{y,1,H}]
$$
  
+
$$
C \int_{t_{0}}^{t_{K}} ||(y - R_{H}y)_{t}||_{0}^{2} dt + C \Delta t \sum_{k=2}^{K} ||\sigma^{y,k-\theta,H}||_{0}^{2} + C \Delta t \sum_{k=2}^{K} ||\rho^{w,k-\theta,H}||_{0}^{2}
$$
  
+
$$
C \Delta t \sum_{k=2}^{K} (||y^{k} - y_{H}^{k}||_{0}^{2} + ||y^{k-1} - y_{H}^{k-1}||_{0}^{2}).
$$
(3.36)

Now we need to give the estimate of  $H[\sigma^{y,1,H}]$ . Combine the definition of Ritz projection operator [\(3.27\)](#page-15-1) and [\(2.13\)](#page-7-0), we can get

<span id="page-17-0"></span>
$$
\left(D_t \sigma^{y, \frac{1}{2}, H}, I_H^* v_H\right) + a_H \left(\sigma^{w, \frac{1}{2}, H}, I_H^* v_H\right) = \left(\partial_t y^{\frac{1}{2}} - D_t R_H y^{\frac{1}{2}}, I_H^* v_H\right),
$$
(3.37)

and

<span id="page-17-1"></span>
$$
\left(\sigma^{w,\frac{1}{2},H}, I_H^* q_H\right) - \gamma a_H \left(\sigma^{y,\frac{1}{2},H}, I_H^* q_H\right)
$$
  
=  $\left(f\left(y_H^{\frac{1}{2}}\right) - f\left(y^{\frac{1}{2}}\right), I_H^* q_H\right) + \left(w^{\frac{1}{2}} - R_H w^{\frac{1}{2}}, I_H^* q_H\right).$  (3.38)

We take  $v_H = \sigma^{y, \frac{1}{2}, H}$  in [\(3.37\)](#page-17-0) and  $q_H = \frac{\sigma^{w, \frac{1}{2}, H}}{\gamma}$  in [\(3.38\)](#page-17-1), and use a similar derivation to the one of inequality  $(3.36)$  to get

<span id="page-17-3"></span>
$$
\|\|\sigma^{y,1,H}\|\|_{0}^{2} + \frac{2\Delta t}{\gamma} \|\|\sigma^{w,\frac{1}{2},H}\|\|_{0}^{2}
$$
  
\n
$$
\leq C \int_{t_{0}}^{t_{1}} \|\left(y - R_{H}y\right)_{t}\|_{0}^{2} dt + C \Delta t \|\sigma^{y,\frac{1}{2},H}\|_{0}^{2} + C \Delta t \|\rho^{w,\frac{1}{2},H}\|_{0}^{2} + \|\sigma^{y,0,H}\|\|_{0}^{2}
$$
  
\n
$$
+ C \Delta t \left(\|y^{1} - y_{H}^{1}\|_{0}^{2} + \|y^{0} - y_{H}^{0}\|_{0}^{2}\right). \tag{3.39}
$$

From [\(3.39\)](#page-17-3) and Lemmas [2.1,](#page-5-0) [3.1,](#page-8-2) we easily know that

<span id="page-17-4"></span>
$$
H[\sigma^{y,1,H}] = 2||\sigma^{y,1,H}||^2_0.
$$
 (3.40)

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Combine [\(3.36\)](#page-17-2), [\(3.39\)](#page-17-3) with [\(3.40\)](#page-17-4) to get

$$
H[\sigma^{y,K,H}] + \frac{2\Delta t}{\gamma} \sum_{k=1}^{K} ||\sigma^{w,k-\theta,H}||_{0}^{2} \le C ||\sigma^{y,0,H}||_{0}^{2}
$$
  
+
$$
C \int_{t_{0}}^{t_{K}} ||(y - R_{H}y)_{t}||_{0}^{2} dt + C \Delta t \sum_{k=1}^{K} ||\sigma^{y,k-\theta,H}||_{0}^{2} + C \Delta t \sum_{k=1}^{K} ||\rho^{w,k-\theta,H}||_{0}^{2}
$$
  
+
$$
C \Delta t \sum_{k=1}^{K} (||y^{k} - y_{H}^{k}||_{0}^{2} + ||y^{k-1} - y_{H}^{k-1}||_{0}^{2}),
$$
(3.41)

which is combined with Gronwall lemma, triangle inequality and Lemmas [2.1,](#page-5-0) [3.4](#page-15-2) to arrive at  $(3.28)$ .

(II) In the fine grid, let  $y_h^k - y^k = y_h^k - R_h y^k + R_h y^k - y^k = \sigma^{y,k,h} + \rho^{y,k,h}$ , and  $w_h^k - w^k = w_h^k - R_h w^k + R_h w^k - w^k = \sigma^{w,k,h} + \rho^{w,k,h}$ . Firstly we would give the estimation of  $(3.29)$ . Combine the definition of Ritz projection operator  $(3.27)$ and  $(2.16)$ , we can get

<span id="page-18-0"></span>
$$
\left(D_t \sigma^{y,k-\theta,h}, I_h^* v_h\right) + a_h \left(\sigma^{w,k-\theta,h}, I_h^* v_h\right)
$$

$$
= \left(\partial_t y^{k-\theta} - D_t R_h y^{k-\theta}, I_h^* v_h\right), \tag{3.42}
$$

and

<span id="page-18-1"></span>
$$
(\sigma^{w,k-\theta,h}, I_h^* q_h) - \gamma a_h(\sigma^{y,k-\theta,h}, I_h^* q_h)
$$
  
=  $((1-\theta)\mathfrak{s}(y_h^k, y_H^k) + \theta f(y_h^{k-1}) - f(y^{k-\theta}), I_h^* q_h) + (w^{k-\theta} - R_h w^{k-\theta}, I_h^* q_h).$  (3.43)

Let  $v_h = \sigma^{y,k-\theta,h}$  in [\(3.42\)](#page-18-0) and  $q_h = \frac{\sigma^{w,k-\theta,h}}{\gamma}$  in [\(3.43\)](#page-18-1), and add two formula, we have

$$
\left(D_t \sigma^{y,k-\theta,h}, I_h^* \sigma^{y,k-\theta,h}\right) + \frac{1}{\gamma} \left(\sigma^{w,k-\theta,h}, I_h^* \sigma^{w,k-\theta,h}\right)
$$
  
+  $a_h \left(\sigma^{w,k-\theta,h}, I_h^* \sigma^{y,k-\theta,h}\right) - a_h \left(\sigma^{y,k-\theta,h}, I_h^* \sigma^{w,k-\theta,h}\right)$   
=  $\left(\partial_t y^{k-\theta} - D_t R_h y^{k-\theta}, I_h^* \sigma^{y,k-\theta,h}\right) + \frac{1}{\gamma} \left(w^{k-\theta} - R_h w^{k-\theta}, I_h^* \sigma^{w,k-\theta,h}\right)$   
+  $\frac{1}{\gamma} \left((1-\theta) \mathfrak{s}\left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right) - f\left(y^{k-\theta}\right), I_h^* \sigma^{w,k-\theta,h}\right)$ . (3.44)

From Lemma [3.2,](#page-9-0) we use Cauchy–Schwarz inequality as well as Young's inequality to get

<span id="page-18-2"></span>
$$
\left(D_t\sigma^{y,k-\theta,h},I_h^*\sigma^{y,k-\theta,h}\right)+\frac{1}{\gamma}\left(\sigma^{w,k-\theta,h},I_h^*\sigma^{w,k-\theta,h}\right)
$$

$$
= \left(\partial_t y^{k-\theta} - D_t R_h y^{k-\theta}, I_h^* \sigma^{y, k-\theta, h}\right) + \frac{1}{\gamma} \left(w^{k-\theta} - R_h w^{k-\theta}, I_h^* \sigma^{w, k-\theta, h}\right) + \frac{1}{\gamma} \left((1-\theta) \mathfrak{s}\left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right) - f\left(y^{k-\theta}\right), I_h^* \sigma^{w, k-\theta, h}\right) \leq \frac{C}{\Delta t} \int_{t_{k-2}}^{t_k} \| (y - R_h y)_t \|_0^2 dt + C \| \sigma^{y, k-\theta, h} \|_0^2 + C \| \rho^{w, k-\theta, h} \|_0^2 + C \delta \| \sigma^{w, k-\theta, h} \|_0^2 + C \left( \|y^k - y_h^k\|_0^2 + \| \left(y^k - y_H^k\right)^2 \|_0^2 + \|y^{k-1} - y_h^{k-1}\|_0^2 \right). \tag{3.45}
$$

From lemma [3.1,](#page-8-2) we know that

<span id="page-19-0"></span>
$$
\left(D_t \sigma^{y,k-\theta,h}, I_h^* \sigma^{y,k-\theta,h}\right) \ge \frac{1}{4\Delta t} \left(H[\sigma^{y,k,h}] - H[\sigma^{y,k-1,h}]\right). \quad (3.46)
$$

Combine the inequality [\(3.45\)](#page-18-2), [\(3.46\)](#page-19-0) with Lemma [2.1](#page-5-0) and let  $\delta$  as a suitable value, we can get

<span id="page-19-1"></span>
$$
\frac{1}{4\Delta t} \left( H[\sigma^{y,k,h}] - H[\sigma^{y,k-1,h}] \right) + \frac{1}{\gamma} ||\sigma^{w,k-\theta,h}||_0^2
$$
\n
$$
\leq \frac{C}{\Delta t} \int_{t_{k-2}}^{t_k} ||(y - R_h y)_t||_0^2 dt + C ||\sigma^{y,k-\theta,h}||_0^2 + C ||\rho^{w,k-\theta,h}||_0^2
$$
\n
$$
+ C \left( ||y^k - y_h^k||_0^2 + ||(y^k - y_H^k)|_0^2 + ||y^{k-1} - y_h^{k-1}||_0^2 \right). \quad (3.47)
$$

Sum [\(3.47\)](#page-19-1) with respect to *k* from 2 to *K* to get

<span id="page-19-4"></span>
$$
H[\sigma^{y,K,h}] + \frac{4\Delta t}{\gamma} \sum_{k=2}^{K} ||\sigma^{w,k-\theta,h}||_{0}^{2} \le H[\sigma^{y,1,h}]
$$
  
+
$$
C \int_{t_{0}}^{t_{K}} ||(y - R_{h}y)_{t}||_{0}^{2} dt + C \Delta t \sum_{k=2}^{K} ||\sigma^{y,k-\theta,h}||_{0}^{2} + C \Delta t \sum_{k=2}^{K} ||\rho^{w,k-\theta,h}||_{0}^{2}
$$
  
+
$$
C \Delta t \sum_{k=2}^{K} (||y^{k} - y_{h}^{k}||_{0}^{2} + ||(y^{k} - y_{H}^{k})^{2}||_{0}^{2} + ||y^{k-1} - y_{h}^{k-1}||_{0}^{2}). \quad (3.48)
$$

Now we need to give the estimation of  $H[\sigma^{y,1,h}]$ . Combine the definition of Ritz projection operator [\(3.27\)](#page-15-1) with [\(2.15\)](#page-7-2), we can get

<span id="page-19-2"></span>
$$
\left(D_t \sigma^{y, \frac{1}{2}, h}, I_h^* v_h\right) + a_h \left(\sigma^{w, \frac{1}{2}, h}, I_h^* v_h\right) = \left(\partial_t y^{\frac{1}{2}} - D_t R_h y^{\frac{1}{2}}, I_h^* v_h\right),
$$
\n(3.49)

and

<span id="page-19-3"></span>
$$
\left(\sigma^{w,\frac{1}{2},h}, I_h^*q_h\right) - \gamma a_h\left(\sigma^{y,\frac{1}{2},h}, I_h^*q_h\right)
$$

$$
= \left(\frac{1}{2}\mathfrak{s}(y_h^1, y_H^1) + \frac{1}{2}f(y_h^0) - f\left(y^{\frac{1}{2}}\right), I_h^*q_h\right) + \left(w^{\frac{1}{2}} - R_hw^{\frac{1}{2}}, I_h^*q_h\right). \tag{3.50}
$$

We take  $v_h = \sigma^{y, \frac{1}{2}, h}$  in [\(3.49\)](#page-19-2) and  $q_h = \frac{\sigma^{w, \frac{1}{2}, h}}{\gamma}$  in [\(3.50\)](#page-19-3), and use a similar derivation to the one of inequality [\(3.48\)](#page-19-4) to get

<span id="page-20-0"></span>
$$
\|\sigma^{y,1,h}\|_{0}^{2} + \frac{2\Delta t}{\gamma} \|\sigma^{w,\frac{1}{2},h}\|_{0}^{2}
$$
  
\n
$$
\leq C \int_{t_{0}}^{t_{1}} \|(y - R_{h}y)_{t}\|_{0}^{2} dt + C \Delta t \|\sigma^{y,\frac{1}{2},h}\|_{0}^{2} + C \Delta t \|\rho^{w,\frac{1}{2},h}\|_{0}^{2} + \|\sigma^{y,0,h}\|_{0}^{2}
$$
  
\n
$$
+ C \Delta t \left(\|y^{1} - y_{h}^{1}\|_{0}^{2} + \|\left(y^{1} - y_{H}^{1}\right)^{2}\|_{0}^{2} + \|y^{0} - y_{h}^{0}\|_{0}^{2}\right).
$$
\n(3.51)

From [\(3.51\)](#page-20-0) and Lemmas [2.1,](#page-5-0) [3.1,](#page-8-2) we easily know that

<span id="page-20-1"></span>
$$
H[\sigma^{y,1,h}] = 2||\sigma^{y,1,h}||_0^2.
$$
 (3.52)

Combine  $(3.48)$ ,  $(3.51)$  with  $(3.52)$  to get

$$
H[\sigma^{y,K,h}] + \frac{2\Delta t}{\gamma} \sum_{k=1}^{K} ||\sigma^{w,k-\theta,h}||_{0}^{2} \le C ||\sigma^{y,0,h}||_{0}^{2}
$$
  
+
$$
C \int_{t_{0}}^{t_{K}} ||(y - R_{h}y)_{t}||_{0}^{2} dt + C \Delta t \sum_{k=1}^{K} ||\sigma^{y,k-\theta,h}||_{0}^{2} + C \Delta t \sum_{k=1}^{K} ||\rho^{w,k-\theta,h}||_{0}^{2}
$$
  
+
$$
C \Delta t \sum_{k=1}^{K} \left( ||y^{k} - y_{h}^{k}||_{0}^{2} + ||(y^{k} - y_{H}^{k})^{2}||_{0}^{2} + ||y^{k-1} - y_{h}^{k-1}||_{0}^{2} \right),
$$

which is combined with Gronwall lemma, triangle inequality with Lemmas [2.1,](#page-5-0) [3.4](#page-15-2) to arrive at  $(3.29)$ .

(III) Then we would give the estimation of  $(3.30)$ . Combine the definition of Ritz projection operator  $(3.27)$  with  $(2.16)$ , we can get

<span id="page-20-2"></span>
$$
\left(D_t \sigma^{y,k-\theta,h}, I_h^* v_h\right) + a_h \left(\sigma^{w,k-\theta,h}, I_h^* v_h\right)
$$

$$
= \left(\partial_t y^{k-\theta} - D_t R_h y^{k-\theta}, I_h^* v_h\right),
$$
(3.53)

and

<span id="page-20-3"></span>
$$
\begin{aligned} &\left(\sigma^{w,k-\theta,h}, I_h^* q_h\right) - \gamma a_h \left(\sigma^{y,k-\theta,h}, I_h^* q_h\right) \\ &= \left((1-\theta) \mathfrak{s}\left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right) - f\left(y^{k-\theta}\right), I_h^* q_h\right) \end{aligned}
$$

$$
+\left(w^{k-\theta}-R_hw^{k-\theta},I_h^*q_h\right). \tag{3.54}
$$

Let  $v_h = \sigma^{w,k-\theta,h}$  in [\(3.53\)](#page-20-2) and  $q_h = D_t \sigma^{y,k-\theta,h}$  in [\(3.54\)](#page-20-3) and subtract the resulting equations following [\(2.9\)](#page-5-3), we have

<span id="page-21-0"></span>
$$
a_h\left(\sigma^{w,k-\theta,h}, I_h^*\sigma^{w,k-\theta,h}\right) + \gamma a_h\left(\sigma^{y,k-\theta,h}, I_h^*D_t\sigma^{y,k-\theta,h}\right)
$$
  
\n
$$
= \left(\partial_t y^{k-\theta} - D_t R_h y^{k-\theta}, I_h^*\sigma^{w,k-\theta,h}\right) - \left(w^{k-\theta} - R_h w^{k-\theta}, I_h^*D_t\sigma^{y,k-\theta,h}\right)
$$
  
\n
$$
- \left((1-\theta) \mathfrak{s}\left(y_h^k, y_H^k\right) + \theta f\left(y_h^{k-1}\right) - f\left(y^{k-\theta}\right), I_h^*D_t\sigma^{y,k-\theta,h}\right)
$$
  
\n
$$
= I + II + III. \tag{3.55}
$$

For I and II, from the Cauchy–Schwarz inequality as well as Young's inequality, we have

$$
I + II
$$
  
\n
$$
\leq \|\partial_t y^{k-\theta} - D_t R_h y^{k-\theta}\|_{0} \|I_h^* \sigma^{w, k-\theta, h}\|_{0} + \|w^{k-\theta} - R_h w^{k-\theta}\|_{0} \|I_h^* D_t \sigma^{y, k-\theta, h}\|_{0}
$$
  
\n
$$
\leq C \|D_t \left(\rho^{y, k-\theta, h}\right)\|_{0}^{2} + C \delta \|\sigma^{w, k-\theta, h}\|_{0}^{2} + \|\rho^{w, k-\theta, h}\|_{0}^{2} + C \|D_t \sigma^{y, k-\theta, h}\|_{0}^{2}
$$
  
\n
$$
\leq C \| \frac{(3-2\theta) [\rho^{y, k, h} - \rho^{y, k-1, h}] - (1-2\theta) [\rho^{y, k-1, h} - \rho^{y, k-2, h}]}{2\Delta t} \|_{0}^{2}
$$
  
\n
$$
+ \|\rho^{w, k-\theta, h}\|_{0}^{2} + C \|D_t \sigma^{y, k-\theta, h}\|_{0}^{2} + C \delta \|\sigma^{w, k-\theta, h}\|_{0}^{2}
$$
  
\n
$$
\leq \frac{C}{\Delta t} \int_{t_{k-2}}^{t_k} \| (y - R_h y)_t \|_{0}^{2} dt + \|\rho^{w, k-\theta, h}\|_{0}^{2} + C \|D_t \sigma^{y, k-\theta, h}\|_{0}^{2} + C \delta \|\sigma^{w, k-\theta, h}\|_{0}^{2}.
$$
  
\n(3.56)

Use Taylor formula and Cauchy–Schwarz inequality, we can get

<span id="page-21-1"></span>III  
\n
$$
\leq ||(1-\theta)\mathfrak{s}(y_h^k, y_H^k) + \theta f(y_h^{k-1}) - f(y^{k-\theta})||_0||I_h^*D_t\sigma^{y,k-\theta,h}||_0
$$
\n
$$
\leq C||y^k - y_h^k||_0^2 + C||(y^k - y_H^k)^2||_0^2 + C||y^{k-1} - y_h^{k-1}||_0^2 + ||D_t\sigma^{y,k-\theta,h}||_0^2.
$$
\n(3.57)

From Lemmas [3.1,](#page-8-2) [3.2](#page-9-0) and  $(3.55)$ – $(3.57)$ , we can obtain that

<span id="page-21-2"></span>
$$
\frac{1}{4\Delta t} \left( \hat{H}[\sigma^{y,k,h}] - \hat{H}[\sigma^{y,k-1,h}] \right) + |\sigma^{w,k-\theta,h}|_1^2
$$
\n
$$
\leq \frac{C}{\Delta t} \int_{t_{k-2}}^{t_k} ||(y - R_h y)_t||_0^2 dt + ||\rho^{w,k-\theta,h}||_0^2 + C ||D_t \sigma^{y,k-\theta,h}||_0^2
$$
\n
$$
+ C ||y^k - y_h^k||_0^2 + C ||(y^k - y_H^k)^2||_0^2 + C ||y^{k-1} - y_h^{k-1}||_0^2. \tag{3.58}
$$

Sum  $(3.58)$  with respect to *k* from 2 to *K* to get

<span id="page-22-2"></span>
$$
\hat{H}[\sigma^{y,K,h}] + 4\Delta t \sum_{k=2}^{K} |\sigma^{w,k-\theta,h}|_{1}^{2} \leq \hat{H}[\sigma^{y,1,h}] + C \int_{t_{0}}^{t_{K}} ||(y - R_{h}y)_{t}||_{0}^{2} dt
$$

$$
+ C\Delta t \sum_{k=2}^{K} ||D_{t}\sigma^{y,k-\theta,h}||_{0}^{2} + C\Delta t \sum_{k=2}^{K} ||\rho^{w,k-\theta,h}||_{0}^{2}
$$

$$
+ C\Delta t \sum_{k=2}^{K} (||y^{k} - y_{h}^{k}||_{0}^{2} + ||(y^{k} - y_{H}^{k})^{2}||_{0}^{2} + ||y^{k-1} - y_{h}^{k-1}||_{0}^{2}).
$$
(3.59)

Now we need to give the estimate of  $\hat{H}[\sigma^{y,1,h}]$ . Combine the definition of Ritz projection operator  $(3.27)$  and  $(2.15)$ , we can get

<span id="page-22-0"></span>
$$
\left(D_t \sigma^{y, \frac{1}{2}, h}, I_h^* v_h\right) + a_h \left(\sigma^{w, \frac{1}{2}, h}, I_h^* v_h\right) = \left(\partial_t y^{\frac{1}{2}} - D_t R_h y^{\frac{1}{2}}, I_h^* v_h\right),
$$
(3.60)

and

<span id="page-22-1"></span>
$$
\begin{aligned} \left(\sigma^{w,\frac{1}{2},h}, I_h^* q_h\right) - \gamma a_h \left(\sigma^{y,\frac{1}{2},h}, I_h^* q_h\right) \\ &= \left(\frac{1}{2}\mathfrak{s}(y_h^1, y_H^1) + \frac{1}{2}f(y_h^0) - f\left(y^{\frac{1}{2}}\right), I_h^* q_h\right) + \left(w^{\frac{1}{2}} - R_h w^{\frac{1}{2}}, I_h^* q_h\right). \end{aligned} \tag{3.61}
$$

We take  $v_h = \sigma^{w, \frac{1}{2}, h}$  in [\(3.60\)](#page-22-0) and  $q_h = D_t \sigma^{y, \frac{1}{2}, h}$  in [\(3.61\)](#page-22-1), and use the similar derivation to the one of inequality  $(3.59)$  to get

$$
\frac{\gamma}{2\Delta t} |\sigma^{y,1,h}|_1^2 + |\sigma^{w,\frac{1}{2},h}|_1^2
$$
\n
$$
= \frac{\gamma}{2\Delta t} |\sigma^{y,0,h}|_1^2 + \frac{C}{\Delta t} \int_{t_0}^{t_1} ||(y - R_h y)_t||_0^2 dt + ||\rho^{w,\frac{1}{2},h}||_0^2 + C||D_t \sigma^{y,\frac{1}{2},h}||_0^2
$$
\n
$$
+ C||y^1 - y_h^1||_0^2 + C||(y^1 - y_H^1)|_0^2 + C||y^0 - y_h^0||_0^2. \tag{3.62}
$$

Multiply  $2\Delta t$  on both sides of the above formula

<span id="page-22-3"></span>
$$
\gamma |\sigma^{y,1,h}|_{1}^{2} + 2\Delta t |\sigma^{w,\frac{1}{2},h}|_{1}^{2}
$$
  
=  $\gamma |\sigma^{y,0,h}|_{1}^{2} + C \int_{t_{0}}^{t_{1}} ||(y - R_{h}y)_{t}||_{0}^{2} dt + 2\Delta t ||\rho^{w,\frac{1}{2},h}||_{0}^{2} + C\Delta t ||D_{t}\sigma^{y,\frac{1}{2},h}||_{0}^{2}$   
+ $C\Delta t ||y^{1} - y_{h}^{1}||_{0}^{2} + C\Delta t ||(y^{1} - y_{H}^{1})^{2}||_{0}^{2} + C\Delta t ||y^{0} - y_{h}^{0}||_{0}^{2}.$  (3.63)

From [\(3.63\)](#page-22-3) and Lemmas [2.1,](#page-5-0) [3.1,](#page-8-2) we easily know that

<span id="page-23-0"></span>
$$
\hat{H}[\sigma^{y,1,h}] = 2|\sigma^{y,1,h}|_1^2.
$$
\n(3.64)

Combine [\(3.59\)](#page-22-2), [\(3.63\)](#page-22-3) with [\(3.64\)](#page-23-0) to get

<span id="page-23-2"></span>
$$
\hat{H}[\sigma^{y,K,h}] + 2\Delta t \sum_{k=1}^{K} |\sigma^{w,k-\theta,h}|_{1}^{2}
$$
\n
$$
\leq C \int_{t_{0}}^{t_{K}} \|(y - R_{h}y)_{t}\|_{0}^{2} dt
$$
\n
$$
+ C \Delta t \sum_{k=2}^{K} \|D_{t}\sigma^{y,k-\theta,h}\|_{0}^{2} + C \Delta t \|D_{t}\sigma^{y,\frac{1}{2},h}\|_{0}^{2} + C \Delta t \sum_{k=1}^{K} \|\rho^{w,k-\theta,h}\|_{0}^{2}
$$
\n
$$
+ C \Delta t \sum_{k=1}^{K} (\|y^{k} - y_{h}^{k}\|_{0}^{2} + \|(y^{k} - y_{H}^{k})^{2}\|_{0}^{2} + \|y^{k-1} - y_{h}^{k-1}\|_{0}^{2}) + C|\sigma^{y,0,h}|_{1}^{2}.
$$
\n(3.65)

Then we need to estimate  $||D_t \sigma^{y,k-\theta,h}||_0$ . Let  $\theta = 0$  in [\(3.54\)](#page-20-3), we have

$$
\begin{split}\n&\left(\sigma^{w,k,h}, I_h^* q_h\right) - \gamma a_h \left(\sigma^{y,k,h}, I_h^* q_h\right) \\
&= \left(\mathfrak{s}\left(y_h^k, y_H^k\right) - f\left(y^k\right), I_h^* q_h\right) + \left(w^k - R_h w^k, I_h^* q_h\right), \\
&\left(\sigma^{w,k-1,h}, I_h^* q_h\right) - \gamma a_h \left(\sigma^{y,k-1,h}, I_h^* q_h\right) \\
&= \left(\mathfrak{s}\left(y_h^{k-1}, y_H^{k-1}\right) - f\left(y^{k-1}\right), I_h^* q_h\right) + \left(w^{k-1} - R_h w^{k-1}, I_h^* q_h\right), \\
&\left(\sigma^{w,k-2,h}, I_h^* q_h\right) - \gamma a_h \left(\sigma^{y,k-2,h}, I_h^* q_h\right) \\
&= \left(\mathfrak{s}\left(y_h^{k-2}, y_H^{k-2}\right) - f\left(y^{k-2}\right), I_h^* q_h\right) + \left(w^{k-2} - R_h w^{k-2}, I_h^* q_h\right).\n\end{split}
$$

Then from  $(2.3)$ , we have

<span id="page-23-1"></span>
$$
\left(D_{t}\sigma^{w,k-\theta,h}, I_{h}^{*}q_{h}\right) - \gamma a_{h}\left(D_{t}\sigma^{y,k-\theta,h}, I_{h}^{*}q_{h}\right)
$$
\n
$$
= \left(\frac{(3-2\theta)\sigma^{w,k,h} - (4-4\theta)\sigma^{w,k-1,h} + (1-2\theta)\sigma^{w,k-2,h}}{2\Delta t}, I_{h}^{*}q_{h}\right)
$$
\n
$$
- \gamma a_{h}\left(\frac{(3-2\theta)\sigma^{y,k,h} - (4-4\theta)\sigma^{y,k-1,h} + (1-2\theta)\sigma^{y,k-2,h}}{2\Delta t}, I_{h}^{*}q_{h}\right)
$$
\n
$$
= \left(\frac{(3-2\theta)\epsilon\left(y_{h}^{k}, y_{H}^{k}\right) - (4-4\theta)\epsilon\left(y_{h}^{k-1}, y_{H}^{k-1}\right) + (1-2\theta)\epsilon\left(y_{h}^{k-2}, y_{H}^{k-2}\right)}{2\Delta t}, I_{h}^{*}q_{h}\right)
$$
\n
$$
- \left(\frac{(3-2\theta)\ f\left(y^{k}\right) - (4-4\theta)\ f\left(y^{k-1}\right) + (1-2\theta)\ f\left(y^{k-2}\right)}{2\Delta t}, I_{h}^{*}q_{h}\right)
$$

$$
+\left(D_t\left(w^{k-\theta}-R_hw^{k-\theta}\right),I_h^*q_h\right). \hspace{1cm} (3.66)
$$

Taking  $q_h = \sigma^{w,k-\theta,h}$  in [\(3.66\)](#page-23-1),  $v_h = \gamma D_t \sigma^{y,k-\theta,h}$  in [\(3.53\)](#page-20-2) and adding the resulting relations we obtain

$$
\left(D_{t}\sigma^{w,k-\theta,h}, I_{h}^{*}\sigma^{w,k-\theta,h}\right) + \gamma \left(D_{t}\sigma^{y,k-\theta,h}, I_{h}^{*}D_{t}\sigma^{y,k-\theta,h}\right) \n= \gamma \left(\partial_{t} y^{k-\theta} - D_{t}R_{h} y^{k-\theta}, I_{h}^{*}D_{t}\sigma^{y,k-\theta,h}\right) + \left(D_{t}\left(w^{k-\theta} - R_{h} w^{k-\theta}\right), I_{h}^{*}\sigma^{w,k-\theta,h}\right) \n+ \left(\frac{(3-2\theta) \mathfrak{s}\left(y_{h}^{k}, y_{H}^{k}\right) - (4-4\theta) \mathfrak{s}\left(y_{h}^{k-1}, y_{H}^{k-1}\right)}{2\Delta t}, I_{h}^{*}\sigma^{w,k-\theta,h}\right) \n+ \left(\frac{(1-2\theta) \mathfrak{s}\left(y_{h}^{k-2}, y_{H}^{k-2}\right)}{2\Delta t}, I_{h}^{*}\sigma^{w,k-\theta,h}\right) \n- \left(\frac{(3-2\theta) f\left(y^{k}\right) - (4-4\theta) f\left(y^{k-1}\right) + (1-2\theta) f\left(y^{k-2}\right)}{2\Delta t}, I_{h}^{*}\sigma^{w,k-\theta,h}\right).
$$
\n(3.67)

From Lemmas [3.1,](#page-8-2) [2.1,](#page-5-0) [3.4,](#page-15-2) [3.2](#page-9-0) and as the process from [\(3.45\)](#page-18-2) to [\(3.52\)](#page-20-1), we can get

<span id="page-24-0"></span>
$$
C\Delta t \sum_{k=2}^{K} \|D_t \sigma^{y,k-\theta,h}\|_0^2 + C\Delta t \|D_t \sigma^{y,\frac{1}{2},h}\|_0^2 \le C(h^2 + H^4). \tag{3.68}
$$

Then from Lemma [3.4,](#page-15-2) triangle inequality, [\(3.68\)](#page-24-0) and [\(3.65\)](#page-23-2), we can arrive at  $(3.30)$ .

<span id="page-24-1"></span>Next, we would show the main error estimation of this paper.

**Theorem 3.6** *Let y, w be the solution of system* [\(2.1\)](#page-2-1),  $(y_H^k, w_H^k)$  *be the coarse grid solution of the system*  $(2.13)$ – $(2.14)$ ,  $(y_h^k, w_h^k)$  *be the fine grid solution of the system*  $(2.15)-(2.16)$  $(2.15)-(2.16)$  $(2.15)-(2.16)$ , respectively. With  $y_h^0 = R_h y_0$ ,  $y_H^0 = R_H y_0$ , there exists a constant C *independent of h, H,*  $\Delta t$ *, such that* 

$$
\|y(t_k) - y_H^k\|_0 + \left(\Delta t \sum_{k=1}^K \|w(t_{k-\theta}) - w_H^{k-\theta}\|_0^2\right)^{\frac{1}{2}} \le C\left(\Delta t^2 + H^2\right), \quad (3.69)
$$

$$
||y(t_k) - y_h^k||_0 + \left(\Delta t \sum_{k=1}^K ||w(t_{k-\theta}) - w_h^{k-\theta}||_0^2\right)^{\frac{1}{2}} \leq C\left(\Delta t^2 + h^2 + H^4\right), \quad (3.70)
$$

$$
||y(t_k) - y_h^k||_1 + \left(\Delta t \sum_{k=1}^K ||w(t_{k-\theta}) - w_h^{k-\theta}||_1^2\right)^{\frac{1}{2}}
$$
  
+ 
$$
||y(t_k) - y_H^k||_0 \le C\left(\Delta t^2 + h + H^2\right).
$$
 (3.71)

*Proof* Let  $E_y^{k-\theta} = y(t_{k-\theta}) - y^{k-\theta}$  and  $E_w^{k-\theta} = w(t_{k-\theta}) - w^{k-\theta}$  and we easily know from systems [\(2.5\)](#page-3-2) and the mixed variational formulation of [\(1.1\)](#page-1-0) that for  $k \ge 2$ , we can obviously get

$$
[(D_t y(t_{k-\theta}), v) + (\nabla w(t_{k-\theta}), \nabla v)] - [(D_t y^{k-\theta}, v) + (\nabla w^{k-\theta}, \nabla v)]
$$
  
= 
$$
[(D_t y(t_{k-\theta}), v) - (D_t y^{k-\theta}, v)] + [(\nabla w(t_{k-\theta}), \nabla v) - (\nabla w^{k-\theta}, \nabla v)]
$$
  
= 
$$
(D_t (y(t_{k-\theta}) - y^{k-\theta}), v) + (\nabla (w(t_{k-\theta}) - w^{k-\theta}), \nabla v)
$$
  
= 
$$
(D_t E_y^{k-\theta}, v) + (\nabla E_w^{k-\theta}, \nabla v)
$$
  
= 
$$
(R_1^{k-\theta}, v),
$$

and

$$
\begin{aligned} & \left( w \left( t_{k-\theta} \right) - w^{k-\theta}, q \right) - \gamma \left( \nabla \left( y \left( t_{k-\theta} \right) - y^{k-\theta} \right), \nabla q \right) \\ & - \left( f \left( y \left( t_{k-\theta} \right) \right) - f \left( y^{k-\theta} \right), q \right) \\ & = \left( E_w^{k-\theta}, q \right) - \gamma \left( \nabla E_y^{k-\theta}, \nabla q \right) - \left( f \left( y \left( t_{k-\theta} \right) \right) - f \left( y^{k-\theta} \right), q \right) \\ & = \left( R_2^{k-\theta}, q \right), \end{aligned}
$$

where

$$
R_1^{k-\theta} = D_t \left( y(t_{k-\theta}) - y^{k-\theta} \right) - \Delta \left( w(t_{k-\theta}) - w^{k-\theta} \right) = O \left( \Delta t^2 \right),
$$
  
\n
$$
R_2^{k-\theta} = \left( w(t_{k-\theta}) - w^{k-\theta} \right) + \gamma \Delta \left( y(t_{k-\theta}) - y^{k-\theta} \right) - \left( f \left( y(t_{k-\theta}) \right) - f \left( y^{k-\theta} \right) \right)
$$
  
\n
$$
= O \left( \Delta t^2 \right).
$$

Similarly, we can follow that

$$
[(D_t y(t_1), v) + (\nabla w(t_1), \nabla v)] - [(D_t y^{\frac{1}{2}}, v) + (\nabla w^{\frac{1}{2}}, \nabla v)]
$$
  
\n
$$
= [(D_t y(t_1), v) - (D_t y^{\frac{1}{2}}, v)] + [(\nabla w(t_1), \nabla v) - (\nabla w^{\frac{1}{2}}, \nabla v)]
$$
  
\n
$$
= (D_t (y(t_1) - y^{\frac{1}{2}}), v) + (\nabla (w(t_1) - w^{\frac{1}{2}}), \nabla v)
$$
  
\n
$$
= (D_t E_y^{\frac{1}{2}}, v) + (\nabla E_w^{\frac{1}{2}}, \nabla v)
$$
  
\n
$$
= (R_3^{\frac{1}{2}}, v),
$$

and

$$
(w(t_{\frac{1}{2}}) - w^{\frac{1}{2}}, q) - \gamma (\nabla(y(t_{\frac{1}{2}}) - y^{\frac{1}{2}}), \nabla q) - (f(y(t_{\frac{1}{2}})) - f(y^{\frac{1}{2}}), q)
$$
  
= 
$$
(E_w^{\frac{1}{2}}, q) - \gamma (\nabla E_y^{\frac{1}{2}}, \nabla q) - (f(y(t_{\frac{1}{2}})) - f(y^{\frac{1}{2}}), q) = (R_4^{\frac{1}{2}}, q),
$$

where

$$
R_3^{\frac{1}{2}} = D_t(y(t_{\frac{1}{2}}) - y^{\frac{1}{2}}) - \Delta(w(t_{\frac{1}{2}}) - w^{\frac{1}{2}}) = O(\Delta t^2),
$$
  
\n
$$
R_4^{\frac{1}{2}} = (w(t_{\frac{1}{2}}) - w^{\frac{1}{2}}) + \gamma \Delta(y(t_{\frac{1}{2}}) - y^{\frac{1}{2}}) - (f(y(t_{\frac{1}{2}})) - f(y^{\frac{1}{2}})) = O(\Delta t^2).
$$

We use the similar analysis as the ones in the Theorem [3.5](#page-15-4) to simply get

<span id="page-26-1"></span>
$$
\| (y(t_k) - y^k) \|_0 + \left( \Delta t \sum_{k=1}^K \| w(t_{k-\theta}) - w^{k-\theta} \|_0^2 \right)^{\frac{1}{2}} \le C \Delta t^2, \qquad (3.72)
$$

and

<span id="page-26-2"></span>
$$
||y(t_k) - y^k||_1 + \left(\Delta t \sum_{k=1}^K ||w(t_{k-\theta}) - w^{k-\theta}||_1^2\right)^{\frac{1}{2}} \le C \Delta t^2.
$$
 (3.73)

Finally, we combine [\(3.72\)](#page-26-1) and [\(3.73\)](#page-26-2) with the result in Theorem [3.5](#page-15-4) and use triangle inequality to get the conclusion of Theorem [3.6.](#page-24-1)

*Remark 3.7* From the Theorem [3.6](#page-24-1) one also can see that the coarse grid can be much coarser than the fine grid and achieve asymptotically optimal approximation as long as the mesh sizes satisfy  $H \le O(h^{\frac{1}{2}})$ , this means that the convergence rate of the space is not lowered. From the second and third inequalities of Theorem [3.6,](#page-24-1) it can be obtained that the spatial convergence order of error  $y - y_h ||_0$  and  $y - y_h ||_1$  is 2 and 1, respectively. The time convergence order of errors  $||y - y_h||_0$  and  $||y - y_h||_1$  is 2.

### <span id="page-26-0"></span>**4 Numerical Examples**

The two numerical examples given in this section will help us test the efficiency of our computation carried out with respect to the two-grid mixed finite volume element method united with the  $\theta$ -scheme. For the implementation of numerical computations in two-dimensional cases, we take triangle segmentation for spatial domain  $\overline{\Omega}$ . In the following description, the mesh length of the coarse grid is taken as *H* and *h* is taken as the mesh length of the fine grid. Then we choose *H* to satisfy  $h < H \leq h^{\frac{1}{2}}$ . Therefore, only the error order of *h* is shown in our example.

*Example 4.1* In [\(1.1\)](#page-1-0), on the space domain  $\overline{\Omega} = [0, 1] \times [0, 1]$  and the time interval [0, 1], we take a term  $f(y) = y^3 - y$  that is nonlinear and the exact solution  $y(t, x_1, x_2) = \cos(\pi t) \cos(\pi x_1) \cos(\pi x_2)$ , then we give the source term  $g(t, x_1, x_2) =$  $\frac{\partial y}{\partial t}$  –  $\Delta(-\gamma \Delta y + f(y)).$ 

In Tables [1,](#page-27-0) [2](#page-27-1) and [3,](#page-27-2) by taking  $θ = 0.5$ ,  $Δt = 1/250$ , changed  $γ = 1, 0.1, 10$ . For testing CPU time, we arrive at the CPU time comparison between two-grid mixed

<b>TGMFVE</b>				<b>MFE</b>				
H	h	$  y - y_h  _1$	$rac{CPU \ time}{Seconds}$	h	$  y - y_h  _1$	CPU time Seconds	Ratio	
0.0884	0.0221	0.1090	6.0740	0.0221	0.1091	8.1911	1.3486	
0.0442	0.0055	0.0273	69.7599	0.0055	0.0273	167.4619	2.4005	
0.0221	0.0014	0.0068	$1.6526e+03$	0.0014	0.0068	$6.5479e+03$	3.9622	

<span id="page-27-0"></span>**Table 1**  $\theta = 0.5$ ,  $\Delta t = 1/250$ ,  $\nu = 1$ 

<span id="page-27-1"></span>**Table 2**  $\theta = 0.5$ ,  $\Delta t = 1/250$ ,  $\gamma = 0.1$ 

<b>TGMFVE</b>				<b>MFE</b>				
H	h	$  y - y_h  _1$	CPU time Seconds	h	$  y - y_h  _1$	CPU time Seconds	Ratio	
0.0884	0.0221	0.1092	6.0679	0.0221	0.1090	9.7017	1.5989	
0.0442	0.0055	0.0273	71.9671	0.0055	0.0273	185.4141	2.5764	
0.0221	0.0014	0.0068	$1.6526e+03$	0.0014	0.0068	$4.6761e+03$	2.8295	

<span id="page-27-2"></span>**Table 3**  $\theta = 0.5$ ,  $\Delta t = 1/250$ ,  $\gamma = 10$ 



finite volume element (TGMFVE) and mixed finite element (MFE) under same mesh and same order  $H^1$  error result. The radio in Tables [1,](#page-27-0) [2](#page-27-1) and [3](#page-27-2) is  $\frac{CPU \text{ time of MFE}}{CPU \text{ time of TGMFVE}}$ . By the contrast between two-grid mixed finite volume element method and mixed finite element method, we see that two-grid mixed finite volume element method can not only economize the CPU time to a great extent, but also get the better convergence rate.

In Tables [4,](#page-28-0) [5](#page-28-1) and [6,](#page-29-1) with the parameter  $\theta = 0.5$ , altered  $\gamma = 1, 0.1, 10$ . We arrive at the CPU time comparison between two-grid mixed finite volume element (TGMFVE) and mixed finite element (MFE) under same mesh, same time step and same order *H*<sup>1</sup> error result. The ratio is also the specific value of CPU time. Compared with a mixed finite element method in these examples, we can see distinctly from the calculated data in Tables [1,](#page-27-0) [2,](#page-27-1) [3,](#page-27-2) [4,](#page-28-0) [5](#page-28-1) and [6](#page-29-1) shows clearly the advantages of utilizing the two-grid mixed finite volume element method. It not only effectively reduce calculation error but also achieve the preferable convergence rate.

In Table [7,](#page-29-2) by taking  $\theta = 0.5$ , changed  $\gamma = 1, 0.1, 10$ . For testing the order of spatial convergence, we keep  $\Delta t = 1/250$  unchanged. We arrive at two-grid mixed finite volume element method error estimates  $(L^2 \text{ norm})$  with second-order convergence rate and error estimates ( $H^1$  norm) with first-order convergence rate.



<span id="page-28-0"></span>**Table 4**  $\theta = 0.5$ ,  $\gamma = 1$ 

<span id="page-28-1"></span>**Table 5**  $\theta = 0.5$ ,  $\gamma = 0.1$ 

$\Delta t$	H	$\boldsymbol{h}$	$  y - y_h  _1$	Order	CPU time Seconds	Ratio
<b>TGMFVE</b>						
1/80	0.0884	0.0221	0.1091	1.9987	2.0568	1.6846
1/160	0.0442	0.0055	0.0273	2.0053	47.1549	2.6680
1/320	0.0221	0.0014	0.0068		$2.1912e+03$	2.6604
<b>MFE</b>						
1/80		0.0221	0.1090	0.9954	3.4648	
1/160		0.0055	0.0273	2.0053	125.8101	
1/320		0.0014	0.0068		$5.8294e+03$	

In Table [8,](#page-30-0) with the parameter  $\theta = 0.5$ , altered  $\gamma = 1, 0.1, 10$ . We inspect the rate at which temporal convergence of  $\theta$ -scheme. Here we use the same order reduction method, that is to say, in estimates  $L^2$  norm, while  $\Delta t$  becomes 1/2 of the previous value, *H* and *h* change  $1/2$  of the previous value, we get the  $L^2$  norm convergence order of two-grid mixed finite volume element method is close to 2. Similarly, in estimates  $H^1$  norm, while  $\Delta t$  becomes 1/2 of the previous value, *H* and *h* change 1/4 of the previous value, we also attain convergence rate of the second-order  $(H^1$  norm). Our method can reach the calculation accuracy of the second-order convergence rate of time.

Further, we show the figures of numerical solutions. In Figs. [3](#page-31-0) and [4,](#page-31-1) based on the parameters  $\theta = 0.5$ ,  $\gamma = 1$  and  $\Delta t = 1/250$ , under the parameter  $h = 0.0221$  the figures of numerical solution  $y_h$  is given at  $t = 0.25$  and  $t = 1$ .

**Example 4.2** We illustrate the typical phase separation phenomena of the Cahn– Hilliard equation through a numerical example from [\[28\]](#page-33-19). The space domains are all the unit square  $\overline{\Omega} = [0, 1] \times [0, 1]$ , with uniform triangulation thereon. The scale of the coarse triangulation is  $H = 0.0625$ , the scale of the fine triangulation is  $h = 0.0078$ , the stepsize is  $\Delta t = 0.001$ ,  $\gamma = 4 \times 10^{-4}$ , taking the source term  $g(t, x_1, x_2) = 0$ , the nonlinear term  $f(y) = y^3 - y$ , we choose the numerical example only including

$\Delta t$	Η	$\boldsymbol{h}$	$  y - y_h  _1$	Order	CPU time Seconds	Ratio
<b>TGMFVE</b>						
1/80	0.0884	0.0221	0.1090	0.9954	2.0819	1.3056
1/160	0.0442	0.0055	0.0273	2.0053	47.1750	2.5748
1/320	0.0221	0.0014	0.0068		$2.1615e + 03$	85.4731
<b>MFE</b>						
1/80		0.0221	0.1091	1.9987	2.7181	
1/160		0.0055	0.0273	2.0053	121.4641	
1/320		0.0014	0.0068		$1.8475e+05$	

<span id="page-29-1"></span>**Table 6**  $\theta = 0.5$ ,  $\nu = 10$ 

<span id="page-29-2"></span>**Table 7**  $\theta = 0.5$ ,  $\Delta t = 1/250$ 

Mesh		$\nu = 1$		$\nu = 0.1$		$\nu = 10$	
H	$\boldsymbol{h}$	$  y - y_h  _0$	Order	$  y - y_h  _0$	Order	$  y-y_h  _0$	Order
0.0884	0.0221	0.0017	2.0495	0.0011	2.0894	0.0018	2.0397
0.0442	0.0055	$9.8280e - 05$	2.3477	$6.0161e - 0.5$	1.6497	$1.0549e - 04$	2.4796
0.0221	0.0014	$3.9572e - 06$		$6.2951e - 06$		$3.5460e - 06$	
H	$\boldsymbol{h}$	$  y - y_h  _1$	Order	$  y - y_h  _1$	Order	$  y - y_h  _1$	Order
0.0884	0.0221	0.1090	0.9954	0.1092	0.9967	0.1090	0.9954
0.0442	0.0055	0.0273	1.0159	0.0273	1.0159	0.0273	1.0159
0.0221	0.0014	0.0068		0.0068		0.0068	

the initial value

$$
y(0, x_1, x_2) = 10^{-3} \sin^3\left(\frac{2\pi}{0.0624}(x_1 - 0.5)\right) \sin^3\left(\frac{2\pi}{0.0624}(x_2 - 0.5)\right)
$$

Following the computation of our method, we can observe the typical phase transition phenomena: phase separation-coarsening see Figure. [5.](#page-32-3)

## <span id="page-29-0"></span>**5 Conclusion**

In this thesis, we develop the two-grid mixed finite-volume element method with  $\theta$  schemes that can solve the Cahn–Hilliard equation. The theoretical conclusions encompassing stability analysis and a priori error estimation in *L*<sup>2</sup> norm and *H*<sup>1</sup> norm for the  $\theta$  scheme with two grid mixed finite volume element method have been given, the numerical experiments results exhibited during the verification process are then used to demonstrate the theoretical correctness of the proposed study.

<span id="page-30-0"></span>



<span id="page-31-0"></span>**Fig. 3** *y<sub>h</sub>* with  $h = 0.0221$ ,  $\Delta t = \frac{1}{250}$ ,  $\theta = 0.5$  and  $\gamma = 1$ 



<span id="page-31-1"></span>**Fig. 4** *y<sub>h</sub>* with  $h = 0.0221$ ,  $\Delta t = \frac{1}{250}$ ,  $\theta = 0.5$  and  $\gamma = 1$ 



<span id="page-32-3"></span>**Fig. 5** The phase evolution of Example II

**Acknowledgements** This work is supported by Shandong Provincial Natural Science Foundation (No. ZR2022MA005).

**Data Availability** The datasets generated during and/or analysed during the current study are available from the corresponding author on reasonable request.

## **Declarations**

**Conflict of interest** We declare that we have no conflict of interest in this work. We declare that we do not have any commercial or associative interest that represents a conflict of interest in connection with the work submitted.

# **References**

- <span id="page-32-0"></span>1. Cahn, J., Hilliard, J.: Free energy of a nonuniform system. I. Interfacial free energy. J. Chem. Phys. **28**, 258–267 (1958)
- 2. Cahn, J., Hilliard, J.: Free energy of a nonuniform system. II. Thermodynamic basis. J. Chem. Phys. **30**, 1121–2214 (1959)
- <span id="page-32-1"></span>3. Cahn, J., Hilliard, J.: Free energy of a nonuniform system. III. Nucleation in a two-component incompressible fluid. J. Chem. Phys. **31**, 688–699 (1959)
- <span id="page-32-2"></span>4. Colturato, M.: Sliding mode control for a diffuse interface tumor growth model coupling a Cahn– Hilliard equation with a reaction–diffusion equation. Math. Methods Appl. Sci. **43**, 6598–6626 (2020)
- 5. Menshov, I.S., Zhang, C.: Interface capturing method based on the Cahn–Hilliard equation for twophase flows. Comput. Math. Math. Phys. **60**, 472–483 (2020)
- 6. Barclay, P., Lukes, J.: Cahn–Hilliard mobility of fluid–fluid interfaces from molecular dynamics. Phys. Fluids **31**, 092107 (2019)
- 7. Wu, X., Zwieten, G., Zee, K.: Stabilized second-order convex splitting schemes for Cahn–Hilliard models with application to diffuse-interface tumor-growth models. Int. J. Numer. Methods Biomed. Eng. **30**, 180–203 (2014)
- 8. Ye, Q., Ouyang, Z., Chen, C., Yang, X.: Efficient decoupled second-order numerical scheme for the flow-coupled Cahn–Hilliard phase-field model of two-phase flows. J. Comput. Appl.Math. **405**, 113875 (2022)
- <span id="page-33-0"></span>9. Wang, X., Li, K., Jia, H.: A linear unconditionally stable scheme for the incompressible Cahn–Hilliard– Navier–Stokes phase-field model. Bull. Iran. Math. Soc. **48**, 1991–2017 (2022)
- <span id="page-33-1"></span>10. Chen, H., Mao, J., Shen, J.: Optimal error estimates for the scalar auxiliary variable finite-element schemes for gradient flows. Numer. Math. **145**, 1–30 (2020)
- <span id="page-33-2"></span>11. Ju, L., Tian, L., Wang, D.: A posteriori error estimates for finite volume approximations of elliptic equations on general surfaces. Comput. Methods Appl. Mech. Eng. **198**(5/8), 716–726 (2009)
- <span id="page-33-3"></span>12. Hu, G., Li, R., Tang, T.: A robust WENO type finite volume solver for steady Euler equations on unstructured grids. Commun. Comput. Phys, **9**(3), 627–648 (2011)
- <span id="page-33-4"></span>13. Nazari, S., Sabzevari, M.: Computational bases for  $S_r\Lambda^1(\mathbb{R}^2)$  and their application in mixed finite element method. Bull. Iran. Math. Soc. **44**, 1141–1153 (2018)
- <span id="page-33-5"></span>14. Du, Q., Ju, L., Tian, L.: Finite element approximation of the Cahn–Hilliard equation on surfaces. Comput. Methods Appl. Mech. Eng. **200**(29), 2458–2470 (2011)
- <span id="page-33-6"></span>15. Jia, H., Hu, H., Meng, L.: A large time-stepping mixed finite method of the modified Cahn–Hilliard equation. Bull. Iran. Math. Soc. **46**, 1551–1569 (2020)
- <span id="page-33-7"></span>16. Nabet, F.: Finite-volume method for the Cahn–Hilliard equation with dynamic boundary conditions. ESAIM: Proc. Surv. **45**(1), 502–511 (2014)
- <span id="page-33-8"></span>17. Appadu, A., Djoko, J., Gidey, H., Lubuma, J.: Analysis of multilevel finite volume approximation of 2D convective Cahn–Hilliard equation. Jpn. J. Ind. Appl. Math. **34**, 253–304 (2017)
- <span id="page-33-9"></span>18. Zhou, J., Chen, L., Huang, Y., Wang, W.: An efficient two-grid scheme for the Cahn–Hilliard equation. Commun. Comput. Phys. **17**, 127–145 (2015)
- <span id="page-33-10"></span>19. Liu, S., Chen, Y., Huang, Y., Zhou, J.: An efficient two grid method for miscible displacement problem approximated by mixed finite element methods. Comput. Math. Appl. **77**, 752–764 (2019)
- <span id="page-33-11"></span>20. Tian, Z., Chen, Y., Wang, J.: A two-grid discretization method for nonlinear SchrÖdinger equation by mixed finite element methods. Comput. Math. Appl. **130**, 10–20 (2023)
- <span id="page-33-12"></span>21. Adams, R.: Sobolev Spaces. Academic, New York (1975)
- <span id="page-33-13"></span>22. Yin, B., Liu, Y., Li, H., He, S., Wang, J.: TGMFE algorithm combined with some time second-order schemes for nonlinear fourth-order reaction–diffusion system. Results Appl. Math. **4**, 100080 (2019)
- <span id="page-33-14"></span>23. Chou, S., Kwak, D., Li, Q.:  $L^p$  error estimates and superconvergence for covolume or finite volume element methods. Numer. Methods Partial Differ. Equ. **19**, 463–486 (2003)
- <span id="page-33-15"></span>24. Chou, S., Li, Q.: Error estimates in  $L^2$ ,  $H^1$  and  $L^\infty$  in covolume methods for elliptic and parabolic problems: a unified approach. Math. Comput. **69**, 103–120 (1999)
- <span id="page-33-16"></span>25. Xu, J.: Two-grid discretization techniques for linear and nonlinear PDEs. SIAM J. Numer. Anal. **33**(5), 1 (2018)
- <span id="page-33-17"></span>26. Xu, J., Zou, Q.: Analysis of linear and quadratic simplicial finite volume methods for elliptic equations. Numer. Math. **111**, 469–492 (2009)
- <span id="page-33-18"></span>27. Thomée, V.: Galerkin Finite Element Methods for Parabolic Problems. Springer, Berlin (1984)
- <span id="page-33-19"></span>28. Zhang, S., Wang, M.: A nonconforming finite element method for the Cahn–Hilliard equation. J. Comput. Phys. **229**, 7361–7372 (2010)

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