



The $W_{(p,q)}^{1,2}$ -solvability for a class of fully nonlinear parabolic equations

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Abstract

The solvability in the Sobolev-Lorentz spaces $W_{(p,q)}^{1,2}(\Omega_T)$ with $p > d + 1$ and $q > 0$ is proved for a class of fully nonlinear parabolic equations with small BMO nonlinearities in (x, t) -variables over a bounded parabolic domain with $C^{1,1}$ -smooth lateral boundary. Here, we make use of the unified approach based on the Fefferman-Stein theorem in accordance with almost all pointwise estimate of the sharp functions to establish the estimates of D^2u and D_tu in the framework of Lorentz spaces.

Keywords Fully nonlinear parabolic equations · BMO nonlinearities · The Fefferman-Stein theorem · Sharp functions · Lorentz spaces

Mathematics Subject Classification Primary: 35B45, 35K61, 35R05 · Secondary: 46E30

1 Introduction

Nonlinear Calderón-Zygmund theory for elliptic and parabolic partial differential equations have been extensively studied in recent decades. In particular, an interior $W^{2,p}$ -estimate with $p > d$ for a class of fully nonlinear elliptic equations, was first obtained by Caffarelli in [6] via the Aleksandrov-Bakel'man-Pucci estimate, the covering argument and the Harnack inequality by using an improvement of Krylov-Safonov's technique. A similar interior estimate was extended by Wang in [26] to fully nonlinear parabolic equations by adapting the Aleksandrov-Bakel'man-Pucci-Krylov-Tso maximum principle and the compactness method. Later, Escauriaza in

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[11] further sharpened Caffarelli's result in [6] to the range $p > d - \varepsilon$ with a small $\varepsilon > 0$ depending on the ellipticity constants and d by considering the reverse Hölder inequality. Meanwhile, we also notice that a solvability in $W_{p,loc}^{2,1}(\Omega_T) \cap C(\overline{\Omega_T})$ for the initial-boundary problem of fully nonlinear parabolic equations was obtained by Crandall, Kocan and Świech in [7]. Winter in [27] further made use of the Alexandroff maximum principle and the weak Harnack inequality on the boundary setting to establish the boundary estimate as well as the $W^{2,p}$ -solvability for the Dirichlet boundary problem associated nonlinear elliptic equations. Recently, an interesting work from Krylov [14] provided a unified approach to show the L^p -solvability for nondivergence elliptic and parabolic equations under the regular assumption of VMO_x principle coefficients. This approach mainly relies on the pointwise estimates of sharp functions for the spatial derivatives of solutions due to the Fefferman-Stein theorem. Furthermore, such an approach was developed by Dong, Kim and Krylov to attain the $W^{2,p}$ -solvability for fully nonlinear elliptic and parabolic equations with the nonlinearities being VMO in independent variables. For examples, Krylov [15] showed the $W^{2,p}$ -solvability with $p > d$ to the Bellman's equations with the VMO nonlinearity by his unified argument. Dong-Krylov-Li [8] demonstrated an interior $W^{2,p}$ -solvability for $p > d$ and $W_p^{2,1}$ -solvability for $p > d + 1$, respectively, to fully nonlinear elliptic and parabolic equations. More generally, Krylov [16] studied the existence and uniqueness in $W^{2,p}$ for $p > d$ to the strong solution of fully nonlinear elliptic equations $H(u, Du, D^2u, x) = 0$. Dong and Krylov in [9] proved the solvability in $W_p^{2,1}$ for $p > d + 1$ to the strong solutions of fully nonlinear parabolic equations $\partial_t u(t, x) + H(u, Du, D^2u; t, x) = 0$ under some relaxed convexity assumptions instead of requiring H to be convex or concave with respect to D^2u . Very recently, Dong and Krylov [10] also derived the regularity in the mixed-norm Sobolev spaces to fully nonlinear elliptic and parabolic equations by improving the Fefferman-Stein theorem in the mixed-norm weighted Lebesgue spaces.

This paper is actually a continuation of Dong-Krylov-Li's work in [8], which extends it in two folds: the Fefferman-Stein theorem in the Lorentz spaces and Lorentz regularity of fully nonlinear parabolic equations. Here, the Fefferman-Stein inequality in the Lorentz spaces is also inspired by Dong and Krylov's paper [10] involving that in the mixed-norm form. Our main aim is to attain a global estimate in the Sobolev-Lorentz spaces by using a unified approach to fully nonlinear parabolic equations with the nonlinearities being small BMO oscillation in independent variable in a bounded $C^{1,1}$ parabolic domain. To state our problem under consideration more precisely, let us recall some related notations. Let Ω be a bounded domain of \mathbb{R}^d for $d \geq 2$ with $C^{1,1}$ -smooth boundary, and $\Omega_T := \Omega \times (0, T)$ for $T > 0$ with its usual parabolic boundary $\partial\Omega_T := (\partial\Omega \times (0, T]) \cup (\Omega \times \{0\})$. We consider the following zero Cauchy-Dirichlet problem of a fully nonlinear parabolic equation:

$$\begin{cases} \partial_t u(x, t) + H(D^2u, Du, u, x, t) = 0 & \text{in } \Omega_T, \\ u = 0 & \text{on } \partial\Omega_T, \end{cases} \quad (1.1)$$

where $D^2u = (D_{ij}u)_{d \times d}$ denotes the Hessian matrix of u , and $Du = (D_i u)_{d \times 1}$ denotes its gradient.

As usual, we use $\mathring{W}_2^{1,2}(\Omega_T)$ to denote the set of all functions $u \in L^2(\Omega_T)$ that their weak derivatives $Du, \mathring{D}^2u, D_t u$ belong to L^2 -spaces, and u vanishes on the parabolic boundary $\partial\Omega_T$ in the trace sense. Therefore, a strong solution of (1.1) which is treated throughout this paper is a Sobolev function $u \in \mathring{W}_2^{1,2}(\Omega_T)$ satisfying (1.1) almost everywhere in Ω_T , for more details to see [18, Chapter VII]. In this context, we suppose that the nonlinearity $H(D^2u, Du, u, x, t)$ can be decomposed into following two parts:

$$H(D^2u, Du, u, x, t) = F(D^2u, x, t) + G(D^2u, Du, u, x, t).$$

Let \mathcal{S} be the set of all symmetric $d \times d$ -matrices, and they satisfy the following structural assumptions:

H1. We suppose that $F(u'', x, t)$ is convex and positive homogeneous of degree one with respect to $u'' \in \mathcal{S}$, and there exist two constants $0 < \lambda \leq \Lambda < \infty$ such that

$$\lambda |\xi|^2 \leq F(u'' + \xi\xi^*, x, t) - F(u'', x, t) \leq \Lambda |\xi|^2$$

for any $u'' \in \mathcal{S}$, $\xi \in \mathbb{R}^d$ and each $(x, t) \in \Omega_T$.

H2. For any fixed $u'' \in \mathcal{S}$, $u' \in \mathbb{R}^d$ and $(x, t) \in \Omega_T$, we let $G(u'', u', u, x, t)$ be a monotone nonincreasing function in $u \in \mathbb{R}$. Moreover, there exists a positive constant K such that

$$|G(u'', u', u, x, t) - G(u'', v', v, x, t)| \leq K (|u' - v'| + |u - v|)$$

and

$$|G(u'', u', u, x, t)| \leq \psi(|u''|) |u''| + K (|u'| + |u|) + g(x, t)$$

for any $u'' \in \mathcal{S}$, $u', v' \in \mathbb{R}^d$, $u, v \in \mathbb{R}$ and each $(x, t) \in \Omega_T$. In the above, the function $\psi(\rho) : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ is bounded, monotonically decreasing and $\psi(\rho) \rightarrow 0$ as $\rho \rightarrow 0$, while $g(x, t)$ is a given function specified later.

Before stating our main result, it is necessary to introduce some notations. For any $r > 0$ and $(x, t) \in \mathbb{R}^{d+1}$, we denote

$$B_r(x) = \{y \in \mathbb{R}^d : |x - y| < r\} \quad \text{and} \quad Q_r(x, t) = B_r(x) \times (t, t + r^2).$$

For convenience, we write $B_r = B_r(0)$ and $Q_r = Q_r(0, 0)$ in the following context. An average of $f(x, t)$ over Q_r is denoted by

$$\bar{f}_{Q_r} = \int_{Q_r} f(x, t) dxdt = \frac{1}{|Q_r|} \int_{Q_r} f(x, t) dxdt,$$

where $|Q_r|$ is $(d + 1)$ -dimensional Lebesgue measure of Q_r . To our aim, we need to impose the following (δ, R_0) -vanishing assumption on the leading term in (1.1) for some δ and $R_0 > 0$.

Assumption 1.1 For any $u'' \in S$ with $|u''| = 1$, we say that $F(u'', x, t)$ is (δ, R_0) -vanishing of (x, t) . If there exists $R_0 > 0$ and a function $\tilde{F}(u'')$ (independent of $z = (x, t)$) satisfying (H1), such that for any $r \in (0, R_0]$ and $(s, y) \in \bar{\Omega}_T$, we have

$$\sup_{u'' \in S, |u''|=1} \int_{Q_r(s,y)} |F(u'', x, t) - \tilde{F}(u'')| dxdt \leq \delta,$$

where $F(u'', x, t)$ is zero extended from $Q_r(s, y) \cap \Omega_T$ to $Q_r(s, y)$ for (x, t) close to the parabolic boundary, and the parameter $\delta > 0$ will be specified later.

We refer to [8] for the same assumptions as H1, H2 and Assumption 1.1, where Dong, Krylov and Li established the L^p estimate with $p > d + 1$ for the zero Cauchy-Dirichlet problems (1.1) in a bounded $C^{1,1}$ -smooth domain. In the article, our aim is to extend it to a regularity in the scale of Lorentz space which is extremely useful when dealing with borderline existence as a refined version of Lebesgue spaces. So, it is necessary to recall the Lorentz space and Sobolev-Lorentz space since our Sobolev-Lorentz regularity for solutions is taken into consideration. The Lorentz space is a two-parameter scale of spaces which make the classical Lebesgue spaces refined in some sense, and the Sobolev-Lorentz space is similarly done from Sobolev space. More precisely, we have

Definition 1.2

- (i) Let U be an open subset in \mathbb{R}^{d+1} . For $p \in [1, \infty)$ and $q \in (0, +\infty)$, the Lorentz space $L^{(p,q)}(U)$ is the set of all measurable functions $f : U \rightarrow \mathbb{R}^{d+1}$ such that

$$\|f\|_{L^{(p,q)}(U)} := \left(p \int_0^\infty \left(\kappa^p |\{(x, t) \in U : |f(x, t)| > \kappa\}| \right)^{\frac{q}{p}} \frac{d\kappa}{\kappa} \right)^{\frac{1}{q}} < +\infty. \tag{1.2}$$

For $q = \infty$, the space $L^{(p,\infty)}(U)$ is set to be the usual Marcinkiewicz space with quasinorm

$$\|f\|_{L^{(p,\infty)}(U)} := \sup_{\kappa > 0} \left(\kappa^p |\{(x, t) \in U : |f(x, t)| > \kappa\}| \right)^{\frac{1}{p}} < +\infty. \tag{1.3}$$

- (ii) The Sobolev-Lorentz space $W_{(p,q)}^{1,2}(U)$ is defined by

$$W_{(p,q)}^{1,2}(U) = \{ f : f, \partial_j f, Df, D^2f \in L^{(p,q)}(U) \},$$

where the quasinorm of f is

$$\|f\|_{W_{(p,q)}^{1,2}(U)} := \|f\|_{L^{(p,q)}(U)} + \|\partial_j f\|_{L^{(p,q)}(U)} + \|Df\|_{L^{(p,q)}(U)} + \|D^2f\|_{L^{(p,q)}(U)}.$$

The regularity in the scale of Lorentz spaces concerning partial differential equations was originated from Talenti’s work [23] based on the symmetrization argument. In 2010, Mingione [21] developed some sharp estimates in the Lorentz spaces

for nonlinear elliptic equations in divergence form by using the duality exponent. Using the geometric argument first introduced by Byun and Wang [4], Mengesha and Phuc [20] recently obtained a global weighted Lorentz estimate to quasilinear elliptic equations with small BMO coefficients in Reifenberg flat domains. In addition, Baroni [2] showed a local Lorentz estimate for parabolic systems of p -Laplacian type by improving the classic Hardy inequality and the reverse Hölder inequality, and a similar result was also obtained in [3] for parabolic obstacle problems of p -Laplacian type by using the same technique. A global extension in the weighted Lorentz spaces was established by Tian and Zheng in [24] to the zero Cauchy-Dirichlet problems of nonlinear parabolic equations with small BMO nonlinearities in Reifenberg domains. Recently, Bui and Duong [5] got a global weighted Lorentz estimate to the weak solutions of divergence nonlinear parabolic equations with nonstandard growth over Reifenberg domains. Very recently, Tian and Zheng [25] also proved a global estimate of weak solutions in variable Lorentz spaces to parabolic obstacle problems with nonstandard growth over a quasiconvex domain based on the large- M -inequality principle originated from Acerbi and Mingione’s work [1]. Zhang and Zheng in [28] also showed local Lorentz estimates for a class of fully nonlinear parabolic and elliptic equations by using the large- M -inequality principle.

We now summarize a global Sobolev-Lorentz estimate for fully nonlinear parabolic equation (1.1) as follows.

Theorem 1.3 *Let Ω be a bounded $C^{1,1}$ -smooth domain and $T > 0$. For the Cauchy-Dirichlet problems (1.1) with assumptions H1 and H2, then there exists a small constant $\delta = \delta(d, \lambda, \Lambda, p, q, \partial\Omega) > 0$ such that if the $F(u'', x, t)$ is (δ, R_0) -vanishing shown as Assumption 1.1 for some $R_0 > 0$, then the following is true:*

- (i) *For $u \in W_{(p,q)}^{1,2}(\Omega_T)$ satisfying (1.1) with $g \in L^{(p,q)}(\Omega_T)$ for $p \in (d + 1, +\infty)$ and $q \in (0, +\infty]$, we have*

$$\|u\|_{W_{(p,q)}^{1,2}(\Omega_T)} \leq N(\|g\|_{L^{(p,q)}(\Omega_T)} + \bar{\rho}\|\psi\|_{L^\infty}) \tag{1.4}$$

for some $\bar{\rho} > 0$, where N is a positive constant depending only on $d, p, q, \lambda, \Lambda, \delta, R_0, K, T, |\Omega|$ and $\|\partial\Omega\|_{C^{1,1}}$ while $q < \infty$, and the constant is independent of q while $q = \infty$.

- (ii) *For any $g \in L^{(p,q)}(\Omega_T)$ with $p \in (d + 1, +\infty)$ and $q \in (0, +\infty]$, there is a unique solution $u \in W_{(p,q)}^{1,2}(\Omega_T)$ of (1.1). In particular, if $\psi(0) = 0$ in (H2), then there is a unique solution $u \in W_{(p,q)}^{1,2}(\Omega_T)$ of (1.1) with the estimate*

$$\|u\|_{W_{(p,q)}^{1,2}(\Omega_T)} \leq N\|g\|_{L^{(p,q)}(\Omega_T)}. \tag{1.5}$$

It is also worth noting that a solvability theorem in the space $W_{p,loc}^{1,2}(\Omega_T) \cap C(\overline{\Omega_T})$ can be found in [7] for the initial boundary value problem for a class of fully nonlinear parabolic equations. The solvability in the Sobolev spaces $W_p^{1,2}(\Omega_T)$ for $p > d + 1$ of the terminal boundary value problem is proved by Dong, Krylov and

Li in [8] for such fully nonlinear parabolic equations (1.1), in bounded cylindrical domains, in the case of VMO “coefficients”. Here, we further prove the solvability theorem 1.3 in the scale of Sobolev-Lorentz spaces under the assumptions of H1 and H2 along this line from Dong-Krylov’s argument. A key ingredient of our main proof is that we need to improve the Fefferman-Stein type theorem of sharp functions in the Lorentz spaces. We note that Krylov in [15] gave a modified Fefferman-Stein theorem in Lebesgue spaces by a filtration of partitions, which leads to that the L^p -solvability was obtained for Bellman’s equations with VMO coefficients in the whole space. Furthermore, Dong and Krylov [10] improved the Fefferman-Stein theorem in the mixed-norm weighted Lebesgue spaces, which yields a mixed-norm regularity for fully nonlinear elliptic and parabolic equations. Inspired by their works, we are going to establish a variant of the Fefferman-Stein theorem in the Lorentz spaces on the basis of the Calderón-Zygmund decomposition in the parabolic setting. Another important point is that the partitions of unity do not work for global estimates of the solutions in the Lorentz-Sobolev spaces, since the commuting transformation of integral does not hold in the Lorentz spaces. To this end, we make use of the idea of local L^p -estimate for nondivergence parabolic equations to handle the mean oscillation estimates of Hessian introduced by Krylov’s paper [15], which originated from Lin’s paper [19] for nondivergence elliptic equations.

Let us now consider the following Dirichlet problem of fully nonlinear elliptic equations:

$$\begin{cases} H(D^2u, Du, u, x) = 0 & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases} \quad (1.6)$$

As a direct consequence of main Theorem , we immediately have the following solvability regarding the elliptic setting.

Corollary 1.4 *Let Ω be a bounded $C^{1,1}$ -smooth domain. For the Dirichlet problem (1.6) with assumptions H1 and H2, then there exists a small constant $\delta = \delta(d, \lambda, \Lambda, p, q, \partial\Omega) > 0$ such that if the $F(u'', x)$ is (δ, R_0) -vanishing for some $R_0 > 0$, then the following is true:*

- (i) *For $u \in W^2_{(p,q)}(\Omega)$ satisfying (1.6) with $g \in L^{(p,q)}(\Omega)$ with $p \in (d, +\infty)$ and $q \in (0, +\infty]$, we have*

$$\|u\|_{W^2_{(p,q)}(\Omega)} \leq N(\|g\|_{L^{(p,q)}(\Omega)} + \bar{\rho}\|\psi\|_{L^\infty}) \quad (1.7)$$

for some $\bar{\rho} > 0$, where N is a positive constant depending only on $d, p, q, \lambda, \Lambda, \delta, R_0, K, |\Omega|$ and $\|\partial\Omega\|_{C^{1,1}}$ if $q < \infty$, while the constant is not dependent of q if $q = \infty$. Specifically, if $\psi(0) = 0$ in (H2), then we have

$$\|u\|_{W^2_{(p,q)}(\Omega)} \leq N\|g\|_{L^{(p,q)}(\Omega)}. \quad (1.8)$$

- (ii) For any $g \in L^{(p,q)}(\Omega)$ with $p \in (d, +\infty)$ and $q \in (0, +\infty]$, there is a unique solution $u \in W_{(p,q)}^2(\Omega)$ of (1.6).

The remainder of this paper is organized as follows. In Sect. 2, we present the Fefferman-Stein type theorem in the Lorentz spaces, and some necessary preliminary results concerning the Gagliardo-Nirenberg interpolation inequality and Alexandrov’s estimate in the Lorentz spaces. In Sect. 3 we complete the proof of Theorem 1.3. Moreover, as a corollary of main Theorem 1.3, a global Sobolev-Lorentz estimate of the solution is proved for parabolic Bellman’s equations with small BMO coefficients in a bounded $C^{1,1}$ -smooth domain.

2 Preliminary results

This section is devoted to some related notations and basic facts which will be useful in the main proof. Throughout this paper, we denote $N(d, p, q, \lambda, \Lambda, \dots)$ or $N_i(d, p, q, \lambda, \Lambda, \dots)$ for $i = 0, 1, 2, \dots$, a universal constant depending only on prescribed quantities and possibly varying from line to line. First, we give a generalized version of the Fefferman-Stein theorem of sharp functions in the Lorentz spaces. Let \mathfrak{C} be a set of all $C_{(R^2, R)}(x, t)$ for any $R > 0$ with the fixed point $(x, t) \in \mathbb{R}^{d+1}$. When f is locally integrable on \mathbb{R}^{d+1} , we define the parabolic maximal function $\mathcal{M}f$ by

$$\mathcal{M}f(x, t) = \sup_{\mathfrak{C} \in \mathfrak{C}} \int_{\mathfrak{C}} |f(s, y)| \, dyds. \tag{2.1}$$

For $\gamma \in (0, 1]$, we set

$$f_{\gamma}^{\sharp}(x, t) = \sup_{\mathfrak{C} \in \mathfrak{C}} \left(\int_{\mathfrak{C}} \int_{\mathfrak{C}} |f(s, y) - f(\tau, z)|^{\gamma} \, dydsdzd\tau \right)^{\frac{1}{\gamma}}. \tag{2.2}$$

Remark 2.1 While $\gamma = 1$, we observe that $f \in BMO$ is equivalent to $f_1^{\sharp} \in L^{\infty}$. We also note that $BMO \cong \mathcal{L}^{p,d}$ (the Campanato spaces) for any $1 \leq p < \infty$ (cf. [22, Chapter IV]), which leads to this fact that the norms $\|f_{\gamma}^{\sharp}\|_{L^{\infty}}$ and $\|f_1^{\sharp}\|_{L^{\infty}}$ are equivalent for any $\gamma \in (0, 1]$.

The Hardy-Littlewood maximal function theorem states that if $f \in L^p$ for $1 < p \leq \infty$, then $\mathcal{M}f \in L^p$ and $\|\mathcal{M}f\|_{L^p} \leq N(p, d) \|f\|_{L^p}$, for instance, see [22, Chapter I]. By interpolating argument, Mengesha and Phuc in [20, Lemma 3.11] extended the above boundedness of the Hardy-Littlewood maximal function from the Lebesgue spaces to the weighted Lorentz spaces. Similarly, if $f_{\gamma}^{\sharp} \in L^p$ for $1 < p < \infty$, then $\mathcal{M}f \in L^p$ and $\|\mathcal{M}f\|_{L^p} \leq N(d, p, \gamma) \|f_{\gamma}^{\sharp}\|_{L^p}$ as given by Krylov in [15, Theorem 5.3]. Therefore, the interpolation method is also applied to the Fefferman-Stein sharp function in the Lorentz spaces just as what Mengesha and Phuc did. Indeed, for any fixed $1 < p < \infty$, we take $p_1 = \frac{p+1}{2}$ to obtain

$$\|\mathcal{M}f\|_{L^{p_1}(\mathbb{R}^{d+1})} \leq N_1 \|f_\gamma^\#\|_{L^{p_1}(\mathbb{R}^{d+1})}$$

for any $f_\gamma^\# \in L^{p_1}(\mathbb{R}^{d+1})$; while we take $p_2 = p + 1$ to show

$$\|\mathcal{M}f\|_{L^{p_2}(\mathbb{R}^{d+1})} \leq N_2 \|f_\gamma^\#\|_{L^{p_2}(\mathbb{R}^{d+1})}$$

for any $f_\gamma^\# \in L^{p_2}(\mathbb{R}^{d+1})$. Note that for any $1 < p < \infty$, $1 < p_1 < p < p_2 < +\infty$. By taking $\vartheta = 1 - \frac{1}{p}$ such that $\frac{1}{p} = \frac{1-\vartheta}{p_1} + \frac{\vartheta}{p_2}$, we use the interpolation theorem for Lorentz spaces (cf. [12, Theorem 1.4.19]) to obtain the following conclusion

$$\|\mathcal{M}f\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \leq N_0 N_1^{1-\vartheta} N_2^\vartheta \|f_\gamma^\#\|_{L^{(p,q)}(\mathbb{R}^{d+1})}.$$

In summary, we have the following conclusion.

Lemma 2.2 *Let $1 < p < \infty$, $0 < q \leq \infty$ and $\gamma \in (0, 1]$. Then, for any $f_\gamma^\# \in L^{(p,q)}(\mathbb{R}^{d+1})$ there exists a constant N depending only on d, p, q and γ such that*

$$\|\mathcal{M}f\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \leq N \|f_\gamma^\#\|_{L^{(p,q)}(\mathbb{R}^{d+1})}.$$

Furthermore, we can also show the improved Fefferman-Stein Theorem (Lemma 2.2) in a half space via the odd or even extensions, see Lemma 5 in [8]. To this end, we write

$$\mathbb{R}_+^{d+1} = \{ (x, t) = (x, t^1, \dots, t^d) \in \mathbb{R}^{d+1} : x^1 > 0 \}.$$

Lemma 2.3 *Let $1 < p < \infty$, $0 < q \leq \infty$ and $\gamma \in (0, 1]$. Then, for any $f_\gamma^\# \in L^{(p,q)}(\mathbb{R}_+^{d+1})$ there exists a constant N depending only on d, p, q and γ such that*

$$\|\mathcal{M}f\|_{L^{(p,q)}(\mathbb{R}_+^{d+1})} \leq N \|f_\gamma^\#\|_{L^{(p,q)}(\mathbb{R}_+^{d+1})}.$$

Next let us extend the Gagliardo–Nirenberg interpolation inequality from the Lebesgue spaces to the Lorentz spaces. We first recall the conclusion by Dong and Krylov in [10, Lemma 6.5 (ii)]. For a fix $\rho > 0$ and $f(x, t) \in L^1_{loc}(\mathbb{R}^{d+1})$, we introduce the confined maximal function and the confined sharp function:

$$\begin{aligned} \mathbb{M}_\rho f(x, t) &= \sup_{\left\{ \begin{array}{l} \mathcal{C} \in \mathfrak{C} : \\ R \in [\rho, \infty) \end{array} \right\}} \int_{\mathcal{C}} |f(s, y)| \, dy ds, \\ f_{1,\rho}^\#(x, t) &= \sup_{\left\{ \begin{array}{l} \mathcal{C} \in \mathfrak{C} : \\ R \in (0, \rho] \end{array} \right\}} \int_{\mathcal{C}} \int_{\mathcal{C}} |f(s, y) - f(\tau, z)| \, dy ds dz d\tau. \end{aligned}$$

Lemma 2.4 For any $\rho > 0$, $p \in [1, \infty)$, and $f \in W_{\rho, \text{loc}}^{1,2}(\mathbb{R}^{d+1})$, we have

$$\mathbb{M}_\rho(|Df|^p) \leq N \mathbb{M}_\rho^{\frac{1}{2}}(|D^2f|^p) \mathbb{M}_\rho^{\frac{1}{2}}(|f|^p) + N \rho^{-p} \mathbb{M}_\rho(|f|^p),$$

where $N = N(d, p)$ is a positive constant.

Indeed, Lemma 2.4 yields the following interpolation inequality in the Lorentz space.

Lemma 2.5 Let $1 < p < \infty$ and $0 < q \leq \infty$. Then for any $f \in W_{(\rho,q), \text{loc}}^{1,2}(\mathbb{R}^{d+1})$, we have

$$\|Df\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \leq N \left(\varepsilon \|D^2f\|_{L^{(p,q)}(\mathbb{R}^{d+1})} + \varepsilon^{-1} \|f\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \right) \tag{2.3}$$

for any $\varepsilon > 0$, where $N = N(p, q, d)$ is a positive constant.

Proof Indeed, we note that for any $\rho > 0$,

$$(Df)_1^\#(x, t) \leq (Df)_{1,\rho}^\#(x, t) + 2\mathbb{M}_\rho(|Df|)(x, t) \quad \text{in } \mathbb{R}^{d+1}. \tag{2.4}$$

By the Poincaré inequality we get

$$(Df)_{1,\rho}^\#(x, t) \leq N\rho \mathcal{M}(|D^2f|)(x, t).$$

Following Lemma 2.4 with $p = 1$ it yields

$$\mathbb{M}_\rho(|Df|)(x, t) \leq N\rho \mathcal{M}(|D^2f|)(x, t) + N\rho^{-1} \mathcal{M}(|f|)(x, t).$$

Putting the above two estimates into (2.4) deduces

$$(Df)_1^\#(x, t) \leq N\rho \mathcal{M}(|D^2f|)(x, t) + N\rho^{-1} \mathcal{M}(|f|)(x, t)$$

for any $(x, t) \in \mathbb{R}^{d+1}$.

Now, we make use of Lemma 2.2 and the Hardy-Littlewood maximal function theorem in the Lorentz spaces to obtain the required result by taking $\varepsilon = \rho$. This completes the proof. □

In what follows, we consider the Cauchy-Dirichlet problem of linear parabolic equations in nondivergence form:

$$\begin{cases} \partial_t u + Lu := \partial_t u + a^{ij}(x, t)D_{ij}u + b^i(x, t)D_i u - c(x, t)u = \varphi(x, t) & \text{in } \Omega_T, \\ u = 0 & \text{on } \partial\Omega_T. \end{cases} \tag{2.5}$$

Here, we suppose that all coefficients $a^{ij}(x, t)$, $b^i(x, t)$ and $c(x, t)$ are bounded measurable functions defined on Ω_T and the symmetric coefficient matrix $A(x, t) = (a^{ij}(x, t))$

satisfies an uniformly ellipticity, which means that there exist positive constants $0 < \lambda \leq \Lambda < \infty$ such that

$$|b^i(x, t)| \leq \Lambda, \quad i = 1, \dots, d; \quad |c(x, t)| \leq \Lambda \quad (2.6)$$

and

$$\lambda |\xi|^2 \leq A(x, t) \xi \xi^* \leq \Lambda |\xi|^2 \quad (2.7)$$

for any $\xi \in \mathbb{R}^d$ and almost all $(x, t) \in \Omega_T$.

In the following we recall the classical Aleksandrov estimate to the Cauchy-Dirichlet problem (2.5), see [13, Lemma 3].

Lemma 2.6 *Let $p \geq d + 1$. Assume that $u \in \mathring{W}_p^{1,2}(\Omega_T)$ is a solution of (2.5) with (2.6) and (2.7). If $\varphi(x, t) \in L^p(\Omega_T)$ and $c(x, t) \geq 0$ in Ω_T , then*

$$\operatorname{ess\,sup}_{(x,t) \in \Omega_T} |u(x, t)| \leq N \left\| |c|^{\frac{d+1-p}{p}} |\det A|^{\frac{1}{p}} \varphi \right\|_{L^p(\Omega_T)},$$

where $N = N(d, p, \lambda, \Lambda, |\Omega|)$.

Combining Lemma 2.6 with the imbedding inequality in the Lorentz spaces, it yields the following $L^{(p,q)}$ -estimate to the problem (2.5).

Lemma 2.7 *For $d + 1 < p < \infty$ and $0 < q \leq \infty$, let $u \in \mathring{W}_{(p,q)}^{1,2}(\Omega_T)$ be the solution of (2.5) with (2.6) and (2.7). If $\varphi(x, t) \in L^{(p,q)}(\Omega_T)$ and $c(x, t) \geq 0$ in Ω_T , then*

$$\|u\|_{L^{(p,q)}(\Omega_T)} \leq N \|\varphi\|_{L^{(p,q)}(\Omega_T)}, \quad (2.8)$$

where $N = N(d, p, q, \lambda, \Lambda, T, |\Omega|)$.

Proof We first use Lemma 2.6 with $p = d + 1$, and the assumptions (2.6, 2.7) to obtain

$$|u(x, t)| \leq N \|\varphi\|_{L^{d+1}(\Omega_T)}$$

for almost all $(x, t) \in \Omega_T$. Furthermore, by the imbedding inequality in the Lorentz spaces (cf. [12]), for any $p > d + 1$ it yields

$$\|\varphi\|_{L^{d+1}(\Omega_T)} \leq N |\Omega_T|^{\frac{1}{d+1} - \frac{1}{p}} \|\varphi\|_{L^{(p,q)}(\Omega_T)},$$

where $N = N(p, q, d)$. Putting the above two estimates together leads to the required result. This completes the proof. \square

3 Proof of Theorem 1.3

This section is devoted to our main proof. We first consider the following parabolic equation

$$\partial_t u(x, t) + F(D^2u, x, t) = h(x, t) \quad \text{in } \mathbb{R}^{d+1}. \tag{3.1}$$

Hereafter, we suppose that Assumption 1.1 holds with \mathbb{R}^{d+1} instead of Ω_T . To obtain the Sobolev-Lorentz estimate of (3.1) in the whole space, for any fixed point $(t_0, x_0) \in \mathbb{R}^{d+1}$, we begin with establishing the following mean oscillation estimate on $Q_r(t_0, x_0)$ for $r > 0$.

Lemma 3.1 *For $\kappa \geq 2, r > 0, \beta \in (1, \infty)$ and $\alpha \in (0, 1)$, let $u \in W^{1,2}_{\beta(d+1)}(\mathbb{R}^{d+1})$ be the solution of (3.1). Then, there exists a small constant $\delta_0 = \delta_0(d, \lambda, \Lambda, \alpha, \beta, \gamma, R_0) > 0$ with some $R_0 > 0$ such that the following holds: if Assumption 1.1 holds for any $\delta \in (0, \delta_0)$, then we have*

$$\begin{aligned} & \int_{Q_r(t_0, x_0)} \int_{Q_r(t_0, x_0)} |D^2u(x, t) - D^2u(s, y)|^\gamma dxdt dyds \\ & \leq N\kappa^{d+2} \left(\int_{Q_{\kappa r}(t_0, x_0)} |h|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ & \quad + N\kappa^{d+2} \delta^{\frac{\gamma}{(d+1)\beta'}} \left(\int_{Q_{\kappa r}(t_0, x_0)} |D^2u|^{(d+1)\beta} dxdt \right)^{\frac{\gamma}{(d+1)\beta}} \\ & \quad + N\kappa^{-\alpha\gamma} \left(\int_{Q_{\kappa r}(t_0, x_0)} |D^2u|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ & \quad + N\kappa^{d+2} R_0^{-2\gamma} \left(\int_{Q_{\kappa r}(t_0, x_0)} |u|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ & \quad + N\kappa^{d+2} R_0^{-\gamma} \left(\int_{Q_{\kappa r}(t_0, x_0)} |Du|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \end{aligned} \tag{3.2}$$

for some $\gamma \in (0, 1]$, where $\beta' = \beta/(\beta - 1)$ and $N = N(d, \delta, \lambda, \Lambda, \beta)$.

Proof The proof is divided into two folds.

Case 1. If $0 < \kappa r \leq R_0$, in the Assumption 1.1 we take

$$\tilde{F}(D^2u) = \bar{F}_{Q_{\kappa r}(t_0, x_0)}(D^2u).$$

Consider

$$\partial_t u(x, t) + \tilde{F}(D^2u) = \tilde{h}(x, t) \quad \text{in } \mathbb{R}^{d+1}, \tag{3.3}$$

where

$$\tilde{h}(x, t) = h(x, t) + \tilde{F}(D^2u) - F(D^2u, x, t).$$

Let us take $\hat{u}'' = \frac{D^2u}{|D^2u|}$, then $\hat{u}'' \in S$ and $|\hat{u}''| = 1$. With (H1) we get

$$\begin{aligned} & \int_{Q_{\kappa r}(t_0, x_0)} \left| F(D^2u, x, t) - \tilde{F}(D^2u) \right|^{d+1} dxdt \\ & \leq \sup_{u'' \in S, |u''|=1} \int_{Q_{\kappa r}(t_0, x_0)} \left| F(u'', x, t) - \tilde{F}(u'') \right|^{d+1} |D^2u|^{d+1} dxdt \\ & \leq N\delta^{\frac{1}{\beta'}} \left(\int_{Q_{\kappa r}(t_0, x_0)} |D^2u|^{(d+1)\beta} dxdt \right)^{\frac{1}{\beta}}, \end{aligned}$$

where δ is a small positive constant shown as in Assumption 1.1 and $N = N(d, \lambda, \Lambda)$. It combined with Lemma 5.6 in [8] yields

$$\begin{aligned} & \int_{Q_r(t_0, x_0)} \int_{Q_r(t_0, x_0)} \left| D^2u(x, t) - D^2u(s, y) \right|^\gamma dxdt dyds \\ & \leq N\kappa^{d+2} \left(\int_{Q_{\kappa r}(t_0, x_0)} |h|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ & \quad + N\kappa^{d+2} \delta^{\frac{\gamma}{(d+1)\beta'}} \left(\int_{Q_{\kappa r}(t_0, x_0)} |D^2u|^{(d+1)\beta} dxdt \right)^{\frac{\gamma}{(d+1)\beta}} \\ & \quad + N\kappa^{-\alpha\gamma} \left(\int_{Q_{\kappa r}(t_0, x_0)} |D^2u|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \end{aligned} \tag{3.4}$$

for some $\gamma \in (0, 1]$.

Case 2. If $\kappa r > R_0$, we remark that

$$\begin{aligned} & \int_{Q_r(t_0, x_0)} \int_{Q_r(t_0, x_0)} \left| D^2u(x, t) - D^2u(s, y) \right|^\gamma dxdt dyds \\ & \leq 2 \int_{Q_r(t_0, x_0)} |D^2u(x, t)|^\gamma dxdt. \end{aligned} \tag{3.5}$$

It is a well-known fact that condition (H1) implies that there is an operator $L_0 = a^{ij}(x, t)D_{ij}$ with assumption (2.7) such that $\partial_t u(x, t) + F(D^2u, x, t) = \partial_t u(x, t) + L_0 u(x, t)$. Therefore, we use Corollary 4.2 in [15] with dilations and the standard cut-off argument to obtain

$$\begin{aligned} \int_{Q_{\kappa r/2}(t_0, x_0)} |D^2u(x, t)|^\gamma dxdt &\leq N \left(\left(\int_{Q_{\kappa r}(t_0, x_0)} |h|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \right. \\ &\quad + \frac{1}{(\kappa r)^{2\gamma}} \left(\int_{Q_{\kappa r}(t_0, x_0)} |u|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ &\quad \left. + \frac{1}{(\kappa r)^\gamma} \left(\int_{Q_{\kappa r}(t_0, x_0)} |Du|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \right) \end{aligned}$$

for some $\gamma \in (0, 1)$. With (3.5) and $\kappa \geq 2$ it leads to that

$$\begin{aligned} &\int_{Q_r(t_0, x_0)} \int_{Q_r(t_0, x_0)} |D^2u(x, t) - D^2u(s, y)|^\gamma dxdt dyds \\ &\leq N\kappa^{d+2} \left(\left(\int_{Q_{\kappa r}(t_0, x_0)} |h|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \right. \\ &\quad + R_0^{-2\gamma} \left(\int_{Q_{\kappa r}(t_0, x_0)} |u|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ &\quad \left. + R_0^{-\gamma} \left(\int_{Q_{\kappa r}(t_0, x_0)} |Du|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \right). \end{aligned} \tag{3.6}$$

Putting the above two cases together completes the proof of the lemma. □

Using Lemma 3.1 and the modified Fefferman-Stein Lemma 2.2, we get the following theorem.

Theorem 3.2 *For $p \in (d + 1, +\infty)$ and $q \in (0, +\infty]$, let $u \in W_{(p,q)}^{1,2}(\mathbb{R}^{d+1})$ satisfying (3.1) with $h \in L^{(p,q)}(\mathbb{R}^{d+1})$. Then there exists a small constant $\delta_0 = \delta_0(d, \lambda, \Lambda, R_0) > 0$ with some $R_0 > 0$ such that if Assumption 1.1 holds for any $\delta \in (0, \delta_0)$, then we have*

$$\begin{aligned} &\|D^2u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} + \|\partial_t u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \\ &\leq N \left(\|h\|_{L^{(p,q)}(\mathbb{R}^{d+1})} + \|u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \right), \end{aligned} \tag{3.7}$$

where $N = N(d, p, q, \lambda, \Lambda, \delta, R_0)$.

Proof By Lemma 3.1 with $\alpha \in (0, 1)$, we obtain

$$\begin{aligned}
 & (D^2u)_\gamma^\#(t_0, x_0) \\
 & \leq N\kappa^{\frac{d+2}{\gamma}} \mathcal{M}^{\frac{1}{d+1}}(|h|^{d+1})(t_0, x_0) \\
 & \quad + N\kappa^{\frac{d+2}{\gamma}} \delta^{\frac{1}{(d+1)\beta'}} \mathcal{M}^{\frac{1}{(d+1)\beta}}(|D^2u|^{(d+1)\beta})(t_0, x_0) \\
 & \quad + N\kappa^{-\alpha} \mathcal{M}^{\frac{1}{d+1}}(|D^2u|^{d+1})(t_0, x_0) \\
 & \quad + N\kappa^{\frac{d+2}{\gamma}} R_0^{-2} \mathcal{M}^{\frac{1}{d+1}}(|u|^{d+1})(t_0, x_0) \\
 & \quad + N\kappa^{\frac{d+2}{\gamma}} R_0^{-1} \mathcal{M}^{\frac{1}{d+1}}(|Du|^{d+1})(t_0, x_0)
 \end{aligned}$$

for some $\gamma \in (0, 1]$. Therefore, with the modified Fefferman-Stein Lemma 2.2 we get

$$\begin{aligned}
 \|D^2u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} & \leq N\left(\kappa^{\frac{d+2}{\gamma}} \delta^{\frac{1}{(d+1)\beta'}} + \kappa^{-\alpha}\right) \|D^2u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \\
 & \quad + N\kappa^{\frac{d+2}{\gamma}} (\|h\|_{L^{(p,q)}(\mathbb{R}^{d+1})}) \\
 & \quad + \|u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} + \|Du\|_{L^{(p,q)}(\mathbb{R}^{d+1})},
 \end{aligned} \tag{3.8}$$

where $N = N(d, p, q, \lambda, \Lambda, \delta, R_0)$. Now, let us first take $\kappa > 0$ being big number, and then take small $\delta > 0$ such that

$$N\left(\kappa^{\frac{d+2}{\gamma}} \delta^{\frac{1}{(d+1)\beta'}} + \kappa^{-\alpha}\right) \leq \frac{1}{2}.$$

Therefore, it follows from the Gagliardo-Nirenberg interpolation inequality (2.3) that

$$\begin{aligned}
 \|D^2u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} & \leq N(\|h\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \\
 & \quad + \|u\|_{L^{(p,q)}(\mathbb{R}^{d+1})}).
 \end{aligned} \tag{3.9}$$

On the other hand, we notice that

$$\partial_t u(x, t) = -F(D^2u, x, t) + h(x, t) \quad \text{in } \mathbb{R}^{d+1}.$$

By (H1) we have

$$\begin{aligned}
 \|\partial_t u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} & \leq N\|D^2u\|_{L^{(p,q)}(\mathbb{R}^{d+1})} + \|h\|_{L^{(p,q)}(\mathbb{R}^{d+1})} \\
 & \leq N(\|h\|_{L^{(p,q)}(\mathbb{R}^{d+1})} + \|u\|_{L^{(p,q)}(\mathbb{R}^{d+1})}).
 \end{aligned} \tag{3.10}$$

Combining (3.9) with (3.10) yields (3.7). □

In the following we consider the initial-boundary problem of (3.1) over a half space. We let

$$\partial \mathbb{R}_+^{d+1} = \{ (x, t) = (x, t^1, \dots, x^d) \in \mathbb{R}^{d+1} : x^1 = 0 \}$$

and consider

$$\begin{cases} \partial_t u(x, t) + F(D^2u, x, t) = h(x, t) & \text{in } \mathbb{R}_+^{d+1}, \\ u = 0 & \text{on } \partial \mathbb{R}_+^{d+1}. \end{cases} \tag{3.11}$$

Here, we suppose that there exists a small constant $\delta > 0$ such that Assumption 1.1 holds with \mathbb{R}_+^{d+1} instead of Ω_T . For any $r > 0$ and $(x, t) \in \mathbb{R}_+^{d+1}$, we denote

$$Q_r^+(x, t) = Q_r(x, t) \cap \overline{\mathbb{R}_+^{d+1}}, \quad Q_r^+ = Q_r^+(0, 0), \quad Q_r^+(x^1) = Q_r^+(0, x^1, 0).$$

Let us fix $(t_0, x_0) \in \overline{\mathbb{R}_+^{d+1}}$. To attain the mean oscillation estimate on $Q_r^+(t_0, x_0)$ for any $r > 0$ we mainly make use of Lemma 7.4 in [8] in terms of the odd extension technique.

Lemma 3.3 For $\kappa \geq 2$, $r > 0$, $\beta \in (1, \infty)$ and $\alpha \in (0, 1)$, let $u \in \dot{W}_{d+1}^{1,2}(\mathbb{R}_+^{d+1})$ be the solution of (3.11) and $(t_0, x_0) \in \mathbb{R}_+^{d+1}$. Then there exists a small constant $\delta_0 = \delta_0(d, \lambda, \Lambda, \beta) > 0$ such that the following holds: if Assumption 1.1 holds for any $\delta \in (0, \delta_0)$, then we have

$$\begin{aligned} & \int_{Q_r^+(t_0, x_0)} \int_{Q_r^+(t_0, x_0)} \left| D^2u(x, t) - D^2u(s, y) \right|^\gamma dxdt dyds \\ & \leq N\kappa^{d+2} \left(\int_{Q_{kr}^+(t_0, x_0)} |h|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ & \quad + N\kappa^{d+2} \delta^{\frac{\gamma}{(d+1)\beta'}} \left(\int_{Q_{kr}^+(t_0, x_0)} |D^2u|^{(d+1)\beta} dxdt \right)^{\frac{\gamma}{(d+1)\beta}} \\ & \quad + N\kappa^{-\alpha\gamma} \left(\int_{Q_{kr}^+(t_0, x_0)} |D^2u|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ & \quad + N\kappa^{d+2} R_0^{-2\gamma} \left(\int_{Q_{kr}^+(t_0, x_0)} |u|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \\ & \quad + N\kappa^{d+2} R_0^{-\gamma} \left(\int_{Q_{kr}^+(t_0, x_0)} |Du|^{d+1} dxdt \right)^{\frac{\gamma}{d+1}} \end{aligned}$$

with some $\gamma \in (0, 1]$, where $\beta' = \beta/(\beta - 1)$ and $N = N(d, \delta, \lambda, \Lambda, \beta)$.

If f is measurable function defined in \mathbb{R}_+^{d+1} , then for any $(x, t) \in \overline{\mathbb{R}_+^{d+1}}$ and $r > 0$,

$$\int_{Q_r^+(x,t)} |f(s, y)| dyds \leq 2 \int_{Q_r(x,t)} |f(s, y)I_{\mathbb{R}_+^{d+1}}| dyds \leq N(d) \mathcal{M}(f I_{\mathbb{R}_+^{d+1}})(x, t),$$

where $I_{\mathbb{R}_+^{d+1}}$ is the indicator function of \mathbb{R}_+^{d+1} . On the other hand, for the smallest $r > 0$ such that $\mathcal{C} \subset Q_r^+$, we have

$$\begin{aligned} & \left(\int_C \int_C |f(x, t) - f(s, y)|^\gamma \, dx dt dy ds \right)^{\frac{1}{\gamma}} \\ & \leq N(d, \gamma) \left(\int_{Q_r^+} \int_{Q_r^+} |f(x, t) - f(s, y)|^\gamma \, dx dt dy ds \right)^{\frac{1}{\gamma}}. \end{aligned}$$

With Lemma 3.3 we obtain the following result.

Corollary 3.4 *Under the same hypotheses as in Lemma 3.3, we have*

$$\begin{aligned} & (D^2u)_\gamma^\#(t_0, x_0) \\ & \leq N\kappa^{\frac{d+2}{\gamma}} \mathcal{M}^{\frac{1}{d+1}} \left(|h|^{d+1} I_{\mathbb{R}_+^{d+1}} \right)(t_0, x_0) \\ & \quad + N\kappa^{\frac{d+2}{\gamma}} \delta^{\frac{1}{(d+1)\beta}} \mathcal{M}^{\frac{1}{(d+1)\beta}} \left(|D^2u|^{(d+1)\beta} I_{\mathbb{R}_+^{d+1}} \right)(t_0, x_0) \\ & \quad + N\kappa^{-\alpha} \mathcal{M}^{\frac{1}{d+1}} \left(|D^2u|^{d+1} I_{\mathbb{R}_+^{d+1}} \right)(t_0, x_0) \\ & \quad + N\kappa^{\frac{d+2}{\gamma}} R_0^{-2} \mathcal{M}^{\frac{1}{d+1}} \left(|u|^{d+1} I_{\mathbb{R}_+^{d+1}} \right)(t_0, x_0) \\ & \quad + N\kappa^{\frac{d+2}{\gamma}} R_0^{-1} \mathcal{M}^{\frac{1}{d+1}} \left(|Du|^{d+1} I_{\mathbb{R}_+^{d+1}} \right)(t_0, x_0) \end{aligned}$$

for any $(t_0, x_0) \in \overline{\mathbb{R}_+^{d+1}}$, where $N = N(d, \delta, \lambda, \Lambda, \beta)$.

Hence, the following regularity theorem in $\hat{W}_{(p,q)}^{1,2}(\mathbb{R}_+^{d+1})$ can be deduced from Corollary 3.4 and Lemma 2.3 with the same way as the proof of Theorem 3.2.

Theorem 3.5 *For $p \in (d + 1, +\infty)$ and $q \in (0, +\infty]$, let $u \in \hat{W}_{(p,q)}^{1,2}(\mathbb{R}_+^{d+1})$ satisfying (3.11) with $h \in L^{(p,q)}(\mathbb{R}_+^{d+1})$. Then there exists a small constant $\delta_0 = \delta_0(d, \lambda, \Lambda, R_0) > 0$ for some $R_0 > 0$ such that if Assumption 1.1 holds for any $\delta \in (0, \delta_0)$ in \mathbb{R}_+^{d+1} , then we have*

$$\begin{aligned} & \left\| D^2u \right\|_{L^{(p,q)}(\mathbb{R}_+^{d+1})} + \left\| \partial_t u \right\|_{L^{(p,q)}(\mathbb{R}_+^{d+1})} \\ & \leq N \left(\|h\|_{L^{(p,q)}(\mathbb{R}_+^{d+1})} + \|u\|_{L^{(p,q)}(\mathbb{R}_+^{d+1})} \right), \end{aligned} \tag{3.12}$$

where $N = N(d, p, q, \lambda, \Lambda, \delta, R_0)$.

On the basis of Theorems 3.2 and 3.5, we now prove the main Theorem 1.3 to the Cauchy-Dirichlet problems (1.1) in Ω_T . Here Ω is a bounded $C^{1,1}$ domain and there exists a small constant $\delta > 0$ such that Assumption 1.1 holds in Ω_T .

Proof of Theorem 1.3 (i) Let us define

$$h(x, t) = -G(D^2u, Du, u, x, t). \tag{3.13}$$

We employ the technique of flattening the boundary and Theorems 3.2 and 3.5 which is a similar way shown as in [17]. Then, we obtain

$$\left\| D^2u \right\|_{L^{(p,q)}(\Omega_T)} + \left\| \partial_t u \right\|_{L^{(p,q)}(\Omega_T)} \leq N \left(\|h\|_{L^{(p,q)}(\Omega_T)} + \|u\|_{L^{(p,q)}(\Omega_T)} \right), \tag{3.14}$$

where $N = N(d, p, q, \lambda, \Lambda, \delta, R_0, T, C^{1,1}$ norm of $\partial\Omega$). The definition of h together with Assumption (H2) leads to that

$$\begin{aligned} \|h\|_{L^{(p,q)}(\Omega_T)} &= \left\| G(D^2u, Du, u, x, t) \right\|_{L^{(p,q)}(\Omega_T)} \\ &\leq \left\| \psi(|D^2u|) |D^2u| \right\|_{L^{(p,q)}(\Omega_T)} + K \left(\|Du\|_{L^{(p,q)}(\Omega_T)} \right. \\ &\quad \left. + \|u\|_{L^{(p,q)}(\Omega_T)} \right) + \|g\|_{L^{(p,q)}(\Omega_T)}. \end{aligned} \tag{3.15}$$

Recall that the function $\psi(\rho) : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ is bounded, monotonically decreasing and $\psi(\rho) \rightarrow 0$ as $\rho \rightarrow 0$. Therefore, for some $\bar{\rho} > 0$ we obtain

$$\begin{aligned} &\left\| \psi(|D^2u|) |D^2u| \right\|_{L^{(p,q)}(\Omega_T)} \\ &\leq \psi(\bar{\rho}) \left\| D^2u \right\|_{L^{(p,q)}(\Omega_T)} + N(p, q) \bar{\rho} \|\psi\|_{L^\infty} |\Omega_T|^{\frac{1}{p}}. \end{aligned} \tag{3.16}$$

This together with (3.15) and (3.14) yield

$$\begin{aligned} &\left\| D^2u \right\|_{L^{(p,q)}(\Omega_T)} + \left\| \partial_t u \right\|_{L^{(p,q)}(\Omega_T)} \\ &\leq N \psi(\bar{\rho}) \left\| D^2u \right\|_{L^{(p,q)}(\Omega_T)} + N \bar{\rho} \|\psi\|_{L^\infty} |\Omega_T|^{\frac{1}{p}} \\ &\quad + N \left(\|Du\|_{L^{(p,q)}(\Omega_T)} + \|u\|_{L^{(p,q)}(\Omega_T)} + \|g\|_{L^{(p,q)}(\Omega_T)} \right), \end{aligned}$$

where $N = N(d, p, q, \lambda, \Lambda, \delta, R_0, T, K, C^{1,1}$ norm of $\partial\Omega$). Taking $\bar{\rho}$ large enough that $N \psi(\bar{\rho}) \leq \frac{1}{2}$ and using again the Gagliardo-Nirenberg interpolation inequality (2.3) in the Lorentz spaces, we obtain

$$\begin{aligned} &\left\| D^2u \right\|_{L^{(p,q)}(\Omega_T)} + \left\| \partial_t u \right\|_{L^{(p,q)}(\Omega_T)} \\ &\leq N \left(\|u\|_{L^{(p,q)}(\Omega_T)} + \|g\|_{L^{(p,q)}(\Omega_T)} + \bar{\rho} \|\psi\|_{L^\infty} \right), \end{aligned} \tag{3.17}$$

where $N = N(d, p, q, \lambda, \Lambda, \delta, R_0, T, K, |\Omega|, C^{1,1}$ norm of $\partial\Omega$). Finally, to estimate the Lorentz norm of u , we rewrite the Cauchy-Dirichlet problem (1.1) as

$$\begin{cases} \partial_t u(x, t) + F(D^2u, x, t) + G(D^2u, Du, u, x, t) - G(D^2u, 0, 0, x, t) = -G(D^2u, 0, 0, x, t) & \text{in } \Omega_T, \\ u = 0 & \text{on } \partial\Omega_T. \end{cases}$$

Conditions (H1) and (H2) imply that there exists an operator $L \cdot = a^{ij} D_{ij} \cdot + b^i D_i \cdot - c \cdot$ with (2.6) and (2.7), such that

$$\begin{cases} \partial_t u(x, t) + Lu = -G(D^2u, 0, 0, x, t) & \text{in } \Omega_T, \\ u = 0 & \text{on } \partial\Omega_T. \end{cases}$$

Since G is a monotonically nonincreasing function in u , by the Alexandrov estimate in the Lorentz spaces (cf. Lemma 2.7) we obtain

$$\|u\|_{L^{(p,q)}(\Omega_T)} \leq N \left\| G(D^2u, 0, 0, x, t) \right\|_{L^{(p,q)}(\Omega_T)}, \tag{3.18}$$

where $N = N(d, p, q, \lambda, \Lambda, T, |\Omega|)$. This combined with Assumption (H2) yields

$$\|u\|_{L^{(p,q)}(\Omega_T)} \leq N \left(\left\| \psi(|D^2u|) |D^2u| \right\|_{L^{(p,q)}(\Omega_T)} + \|g\|_{L^{(p,q)}(\Omega_T)} \right).$$

Therefore, with the above estimate and (3.17), we use the same argument as in the proof of (3.16) to obtain

$$\left\| D^2u \right\|_{L^{(p,q)}(\Omega_T)} + \left\| \partial_t u \right\|_{L^{(p,q)}(\Omega_T)} \leq N \left(\|g\|_{L^{(p,q)}(\Omega_T)} + \bar{\rho} \|\psi\|_{L^\infty} \right) \tag{3.19}$$

for some $\bar{\rho} > 0$, where $N = N(d, p, q, \lambda, \Lambda, \delta, R_0, T, K, |\Omega|, \|\partial\Omega\|_{C^{1,1}})$.

(ii) With the *a priori* estimate (3.19) in hand, we make use of an approximating argument to the nonlinear function H and the domain Ω , which is just as the same proof of [8, Theorem 1.2], to lead to the existence and uniqueness claims of (1.1) in the Sobolev-Lorentz spaces $W^{1,2}_{(p,q)}(\Omega_T)$. It completes the proof. \square

Notice that parabolic Bellman’s equation is a typical prototype as fully nonlinear parabolic equations. Therefore, as a consequence of Theorem 1.3, it yields the following Sobolev-Lorentz estimate of strong solutions to parabolic Bellman’s equations under the assumption of small BMO coefficients in a bounded $C^{1,1}$ domain. More precisely, we consider the initial-boundary problem of parabolic Bellman’s equation

$$\begin{cases} \partial_t u(x, t) + \sup_{\omega \in \mathcal{A}} \left[a^{ij}(\omega, x, t) D_{ij} u(x, t) + b^i(\omega, x, t) D_i u(x, t) \right. \\ \qquad \qquad \qquad \left. - c(\omega, x, t) u(x, t) + g(\omega, x, t) \right] = 0 & \text{in } \Omega_T, \\ u(x, t) = 0 & \text{on } \partial\Omega_T, \end{cases} \tag{3.20}$$

where the set \mathcal{A} is a separable metric space, $a(\omega, x, t) = (a^{ij}(\omega, x, t))$, $b(\omega, x, t) = (b^i(\omega, x, t))$, $c(\omega, x, t) \geq 0$ and $g(\omega, x, t)$ are given functions. Here, we assume that these functions are measurable in (x, t) for each ω , continuous in ω for each (x, t) and there exist positive constants $0 < \lambda \leq \Lambda < \infty$ such that

$$|b(\omega, x, t)| \leq \Lambda, \quad c(\omega, x, t) \leq \Lambda, \quad \hat{g}(x, t) := \sup_{\omega \in \mathcal{A}} [g(\omega, x, t)] < \infty \tag{3.21}$$

and

$$\lambda|\xi|^2 \leq a^{ij}(\omega, x, t) \xi_i \xi_j \leq \Lambda|\xi|^2 \tag{3.22}$$

for any $\xi \in \mathbb{R}^d$, $\omega \in \mathcal{A}$ and almostly all $(x, t) \in \Omega_T$. Moreover, $a^{ij}(\omega, x, t)$ is (δ, R_0) -vanishing of $\overline{(x, t)}$ which means that there exists $R_0 > 0$ such that for any $r \in (0, R_0]$ and $(s, y) \in \overline{\Omega}_T$,

$$\sup_{i,j} \int_{Q_r(s,y)} \sup_{\omega \in \mathcal{A}} \left| a^{ij}(\omega, x, t) - \bar{a}_{Q_r(s,y)}^{ij}(\omega) \right| dxdt \leq \delta, \tag{3.23}$$

where the parameter $\delta > 0$ will be specified later.

One can check that **H1,H2** and Assumptions 1.1 hold in Theorem 1.3 by taking

$$F = \sup_{\omega \in \mathcal{A}} a^{ij}(\omega, x, t) D_{ij} u(x, t) \quad \text{and}$$

$$G = \sup_{\omega \in \mathcal{A}} \left[b^i(\omega, x, t) D_i u(x, t) - c(\omega, x, t) u(x, t) + g(\omega, x, t) \right],$$

which immediately leads to the following conclusion.

Corollary 3.6 For $p \in (d + 1, +\infty)$, $q \in (0, +\infty]$, $\Omega \in C^{1,1}$, let $u \in \dot{W}_{(p,q)}^{1,2}(\Omega_T)$ with $T > 0$ be the solution of the Cauchy-Dirichlet problems (3.20) with assumptions (3.21, 3.22) and $\hat{g} \in L^{(p,q)}(\Omega_T)$. Then, there exists a small constant $\delta = \delta(d, p, q, \lambda, \Lambda, R_0) > 0$ for some $R_0 > 0$ such that if the condition (3.23) holds, then we have

$$\|u\|_{W_{(p,q)}^{1,2}(\Omega_T)} \leq N \|\hat{g}\|_{L^{(p,q)}(\Omega_T)}. \tag{3.24}$$

Here, N is a positive constant depending only on $d, p, q, \lambda, \Lambda, \delta, R_0, T, |\Omega|$ and $\|\partial\Omega\|_{C^{1,1}}$.

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Compliance with ethical standards

Conflict of interest The authors declare that there is no conflict of interest regarding the publication of this paper.

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