

Optimal parameter choice of $CR-RC^m$ digital filter in nuclear pulse processing

Huai-Qiang Zhang^{1,2} · Zhuo-Dai Li¹ · Bin Tang^{1,2} · He-Xi Wu¹

Received: 2 January 2019 / Revised: 18 March 2019 / Accepted: 31 March 2019 / Published online: 5 June 2019 - China Science Publishing & Media Ltd. (Science Press), Shanghai Institute of Applied Physics, the Chinese Academy of Sciences, Chinese Nuclear Society and Springer Nature Singapore Pte Ltd. 2019

Abstract $CR-RC^m$ filters are widely used in nuclear energy spectrum measurement systems. The choice of parameters of a $CR-RC^m$ digital filter directly affects its performance in terms of energy resolution and pulse count rate in digital nuclear spectrometer systems. A numerical recursive model of a CR differential circuit and RC integration circuit is derived, which shows that the shaping result of $CR-RC^m$ is determined by the adjustment parameter (k, i) determines the shaping time of the shaper) and the integral number (m) . Furthermore, the amplitude– frequency response of $CR-RC^m$ is analyzed, which shows that it is a bandpass filter; the larger the shaping parameters $(k$ and $m)$, the narrower is the frequency band. CR–RC^m digital Gaussian shaping is performed on the actual sampled nuclear pulse signal under different shaping parameters. The energy spectrum of $137Cs$ is measured based on the $LaBr₃(Ce)$ detector under different parameters. The results show that the larger the shaping parameters (m and k), the closer the shaping result is to Gaussian shape, the wider is the shaped pulse, the higher is the energy resolution, and the lower is the pulse count rate. For the same batch of pulse signals, the energy resolution is increased from 3.8 to 3.5%, and the full energy peak area is reduced from 7815 to 6503. Thus, the optimal shaping parameters are $m = 3$ and $k = 0.95$. These research results can provide a design reference for the development of digital nuclear spectrometer measurement systems.

Keywords $CR-RC^m$ digital filter \cdot Nuclear pulse signal \cdot Gaussian shaping - Energy resolution

1 Introduction

In analog nuclear spectrometer systems, to improve the signal-to-noise ratio (SNR) of nuclear pulse processing, a $CR-RC^m$ circuit is often used to filter the nuclear pulse signal and shape it into a quasi-Gaussian shape $[1-3]$. In digital nuclear spectrometer systems, based on a highspeed analog-to-digital converter (ADC) and a high-performance field-programmable gate array (FPGA), the filtering of a nuclear pulse signal is often realized using a digital algorithm. Triangular filters, trapezoidal filters, and cusp-like filters are widely used in digital nuclear pulse processing $[4-11]$. Owing to the simple implementation and high SNR of a Gaussian filter, it is an important digital filter shaping method. The method mainly includes a time domain based on the digital Sallen–Key [[12–16](#page-4-0)] and a transfer function based on $CR-RC^m$ [[17\]](#page-4-0). The digital Gaussian filter based on $CR-RC^m$ is often used in digital nuclear spectrometer systems, owing to its superior performance and simple shaping principle and method. Based on the basic principles of CR and RC circuits, a recursive model of the digital filter shaping algorithm is derived.

This work was supported by National Natural Science Foundation of China (Nos. 11665001, 41864007), National Key R&D Project (No. 2017YFF0106503), China Scholarship Council (No. 201708360170) and One Hundred People Sail in Jiangxi Province, Open-ended Foundation from the Chinese Engineering Research Center (No. HJSJYB2014-03).

 \boxtimes Huai-Oiang Zhang zhanghq821@163.com

¹ Engineering Research Center of Nuclear Technology Application (East China University of Technology), Ministry of Education, Nanchang 330013, China

School of Nuclear Science and Engineering, East China University of Technology, Nanchang 330013, China

Furthermore, the actual sampled nuclear signal is processed using the $CR-RC^m$ filter, and the effect of the digital filter is analyzed based on the amplitude–frequency response characteristics. The performance is compared and analyzed in terms of the energy resolution and pulse count rate under different forming parameters.

2 Filter shaping principle of $CR-RC^m$

The $CR-RC^m$ filter shaping circuit used in nuclear energy spectrum measurement systems is shown in Fig. 1. In this figure, V_{in} is the input of the circuit, CR is the differential circuit (high-pass filter circuit), RC is the integral circuit (low-pass filter circuit), and V_{out} is the output of the circuit. The input nuclear pulse signal is shaped into a quasi-Gaussian shape by the circuit.

As shown in Fig. 1, $R_1C_1 = R_2C_2 = R_mC_m = \tau$. V_{in} (as a function of the time t) is the input signal. When it is a step signal, it is represented by Eq. (1) in the time and complex frequency domains, where Q is a charge and C_i is an input capacitor.

$$
V_{\text{in}}(t) = \frac{Q}{C_i} u(t) \Rightarrow V_{\text{in}}(s) = \frac{Q}{C_i} \cdot \frac{1}{s}.
$$
 (1)

The transfer function of CR differential shaping is

$$
H_{\rm CR}(s) = \frac{S\tau}{1 + S\tau}.
$$
\n(2)

The transfer function of RC integral shaping is

$$
H_{\rm RC}(s) = \frac{1}{1 + S\tau}.\tag{3}
$$

From Eqs. (2) and (3) , the output is

$$
V_{\text{out}}(s) = \frac{Q}{C_i} \frac{\tau}{(1 + S\tau)^{m+1}}
$$

\n
$$
\Rightarrow V_{\text{out}}(t) = \frac{Q}{C_i m!} \left(\frac{t}{\tau}\right)^m e^{-\frac{t}{\tau}}.
$$
\n(4)

If the input signal V_{in} is a negative exponential signal, it is necessary to add a pole-zero cancelation circuit before the differential shaping to achieve Gaussian shaping of the nuclear pulse signal [[18\]](#page-4-0).

3 Derivation of algorithm of $CR-RC^m$ digital filter

A CR–RC m filter consists of a CR high-pass filter and m RC low-pass filters. To digitize the nuclear signal more effectively, the CR and RC circuits are separately digitized [\[19](#page-4-0), [20](#page-4-0)].

1. The differential equation for CR high-pass circuits according to Kirchhoff's current law (KCL) is shown in Eq. (5).

$$
\frac{\mathrm{d}V_{\mathrm{i}}(t)}{\mathrm{d}t} = \frac{V_{\mathrm{o}}(t)}{\mathrm{CR}} + \frac{\mathrm{d}V_{\mathrm{o}}(t)}{\mathrm{d}t},\tag{5}
$$

where $V_i(t)$ and $V_0(t)$ are the input and output signals of the CR circuit, respectively. For the discrete signal sequence digitized by the high-speed ADC, $V_i(t)$ is converted into a digital sequence $X[n]$, and $V_0(t)$ is converted into a digital sequence $Y[n]$. The derivation can yield the digital recursive formula shown in Eq. (6).

$$
Y[n] = k(X[n] - X[n-1]) + k \cdot Y[n-1].
$$
 (6)

As shown in Eq. (6), $k = \frac{CR}{CR + T}$, where T is the sampling period of the ADC. Thus, the transfer function in the z-domain of the CR high-pass filter and its frequency response function are [\[20,](#page-4-0) [21\]](#page-4-0) expressed as follows:

$$
H_{\rm CR}(z) = k \frac{1 - z^{-1}}{1 - k \cdot z^{-1}}
$$

\n
$$
\Rightarrow H_{\rm CR}(e^{jw}) = k \frac{1 - e^{-jw}}{1 - ke^{-jw}}.
$$
\n(7)

In Eq. (7) , w is the angular frequency of the component signal, and Euler's formula transformation for Eq. (7) is shown in Eq. (8) .

$$
H_{\rm CR}(e^{jw}) = k \frac{1 - \cos(w) + j \cdot \sin(w)}{1 - k(\cos(w) + j \cdot \sin(w))}.
$$
 (8)

The modulo of Eq. (8) is shown in Eq. (9) .

$$
|H_{\text{CR}}(e^{iw})| = k \frac{\sqrt{2 - 2\cos(w)}}{\sqrt{1 + k^2 - 2k \cdot \cos(w)}}.
$$
 (9)

Fig. 1 Circuit of a CR-RC^m filter

2. The digital recursive equation, frequency response function, and modes of the RC low-pass circuit according to the same principle are shown in Eqs. (10) – (12) , respectively.

$$
Y[n] = (1 - k) \cdot X[n] + k \cdot Y[n - 1], \tag{10}
$$

$$
H_{\rm RC}(e^{jw}) = \frac{1-k}{1-k\cdot\cos(w)+j\cdot k\cdot\sin(w)},\tag{11}
$$

$$
|H_{\rm RC}(e^{jw})| = \frac{1-k}{\sqrt{1+k^2-2k\cdot\cos(w)}}.\tag{12}
$$

As shown in Eq. (10), $k = RC/(RC + T)$, where T is the sampling period of the ADC.

According to Eqs. (8) (8) and (11) , the frequency response function of the $CR-RC^m$ filter is shown in Eq. (13).

$$
H(e^{jw}) = \frac{(1-k)^m \cdot k \cdot (1 - k \cos(w) + j \cdot k \cdot \sin(w))}{1 - k \cdot \cos(w) + j \cdot k \cdot \sin(w)}.
$$
\n(13)

The modulo of Eq. (13) is shown in Eq. (14) .

m -

$$
|H(e^{jw})| = \frac{(1-k)^m \cdot k \cdot \sqrt{2 - 2\cos(w)}}{\left(\sqrt{1 + k^2 - 2k \cdot \cos(w)}\right)^{m+1}}.
$$
 (14)

Based on the SNR coefficient of the $CR-RC^m$ filter, $F = \frac{m!e^m}{m^m}(2m-1)^{\frac{1}{4}} \left[\frac{(2m-3)!!}{2(2m)!!} \right]^{\frac{1}{2}}$, for $m = 1$, $F = 1.36$; for $m = 2$, $F = 1.22$; for $m = 3$, $F = 1.18$; for $m = 4$, $F = 1.17$; and for $m = 5$, $F = 1.16$. Thus, with the increase in the parameter m , the change in the parameter F is small, and the filtering system becomes more complicated. Therefore, in practical applications, the parameter m is generally 3 or 4. For the value of the parameter k , the sampling period of the high-speed ADC is generally 100 ns, 50 ns, and 20 ns in the digital nuclear energy spectrum. When $RC = \tau = 1 \,\mu s$, the corresponding k values are $k = 0.91$, $k = 0.95$, and $k = 0.98$. According to

Fig. 2 (Color online) Frequency response curve of CR and RC filters (different k values)

Eqs. (9) (9) and (12) , the frequency response curve of the CR and RC filters is shown in Fig. 2.

As shown in Fig. 2, the amplitude–frequency response curve of the CR filter shows an incremental increase in frequency. As the value of k increases, the response curve rises faster, and it approaches 1 sooner. The amplitude– frequency response curve of the RC filter shows a decline. As the value of k increases, the response curve falls faster, and it approaches zero sooner. When $m = 3$, $k = 0.91$, $k = 0.95$, and $k = 0.98$, the frequency response curve of CR–RC^m is shown in Fig. 3. When $k = 0.95$, $m = 2$, $m = 3$, and $m = 4$, the frequency response curve of CR– RC^m is shown in Fig. [4.](#page-3-0)

As shown in Figs. 3 and [4,](#page-3-0) the frequency response curve of the digitized $CR-RC^m$ filter shows that it is a bandpass filter. The parameters m and k can be used to adjust the bandpass region of the filter, and the larger the values of m and k , the narrower is the bandpass region.

4 Algorithm implementation

137Cs is measured using the CH249-02 NaI detector (Hamamatsu). The high voltage is $+560$ V, and the output pulse of the signal conditioning circuit is approximately 10 μs. USB-5133 (National Instruments Corporation) is used to obtain the original nuclear pulse signal. (The sampling rate is 50 MHz.) According to the numerical recursive formula of the CR–RC^m filter, for $m = 3$, the $CR-RC^m$ shaping results under different k values are as shown in Fig. [5.](#page-3-0)

As shown in Fig. [5,](#page-3-0) with the increase of parameter k , the rising edge of the $CR-RC^m$ is gentler, and the closer to a Gaussian shape, the lower the pulse amplitude; however, for longer pulse shaping time, the generalization of the pileup occurs, and the original signal is a negative exponentially decreasing signal. After CR differential shaping, a certain degree of zero-crossing signal is present. For a

Fig. 3 Frequency response curve of the $CR-RC³$ filter (different k values)

Fig. 4 (Color online) Frequency response curve of the $CR-RC^m$ filter (different m values)

Fig. 5 (Color online) Actual nuclear pulse signal of $CR-RC^m$ shaping

pulse with a small k value or for a long interval of the original pulse, there is no pulse pileup. The undershoot signal of zero has no effect on the subsequent pulse amplitude extraction. The k value is too large, and the longer the shaping time, the higher is the pileup probability. Furthermore, the undershoot of the zero crossing directly affects the amplitude extraction of subsequent adjacent pulses [\[22](#page-4-0)].

5 Test and analysis of $CR-RC^m$ filter performance

The LaBr₃(Ce) detector (Beijing Huakailong's $2'' \times 2''$) is used to measure the $137Cs$ energy spectrum. The ADC sampling rate is 20 MHz. For the same batch of sampling pulse data, when $m = 3$, $k = 0.91$, $k = 0.95$, and $k = 0.98$, the energy spectrum is shown in Fig. 6.

As shown in Fig. 6, to demonstrate the energy spectrum results under different k values more clearly, the data of partial channels are displayed. Similarly, when $k = 0.95$, $m = 2$, $m = 3$, and $m = 4$, the statistical analysis of the energy resolution and the full energy peak area of the energy spectrum is shown in Table 1.

Fig. 6 (Color online) Energy spectrum of ^{137}Cs (partial channels)

Table 1 Performance comparison of different shaping parameters

Shaping parameter Energy resolution $(\%)$ Full energy peak area	
$CR-RC^3$ $k = 0.91$ 3.8	7815
$CR-RC^3$ $k = 0.95$ 3.6	7362
$CR-RC^3$ $k = 0.98$ 3.5	6537
$CR-RC^2 k = 0.95$ 3.7	7471
$CR-RC^4$ $k = 0.95$ 3.6	6503

As shown in Table 1, for the same batch of pulses, after the $CR-RC^m$ digital filter shaping process under different parameters, with the increase in k and m values, the energy resolution is increased. This is because the width of the shaping pulse is too large, and the full energy peak area of the energy spectrum is lowered owing to the pileup identification [[23\]](#page-4-0).

From Table 1, when $m = 3$ and $k = 0.91$, the energy resolution and the full energy peak area are, respectively, 3.8% and 7815; when $m = 3$ and $k = 0.95$, the energy resolution and the full energy peak area are, respectively, 3.6% and 7362; when $m = 3$ and $k = 0.98$, the energy resolution and the full energy peak area are, respectively, 3.5% and 6537; when $k = 0.95$ and $m = 2$, the energy resolution and pulse counting rates are, respectively, 3.7% and 7471; when $k = 0.95$ and $m = 4$, the energy resolution and the full energy peak area are, respectively, 3.6% and 6503.

6 Conclusion

 $CR-RC^m$ digital filters are widely used in digital nuclear energy spectrum measurement systems. The filter parameters directly affect the energy resolution and pulse count rate of the system. The numerical recursive algorithm of digital filtering of CR and RC is based on KCL. The measured signal is processed by the $CR-RC^m$ digital filter, and the energy resolution and maximum pulse of the system are obtained under different k and m values in the energy spectrum of ^{137}Cs . From the comparative analysis, combined with the above indicators, it is determined that if the system requires high energy resolution, the values of k and m can be larger. Moreover, when it requires a high count rate, the values of k and m can be smaller. When the system needs to balance the performance of both, the optimal parameters of the CR–RC^m digital filter are $k =$ 0.95 and $m = 3$. The relevant results can provide a design reference for further use of the $CR-RC^m$ digital filter.

Acknowledgements I would like to express my gratitude to my supervisor, Mr. Tomas Trojek, who has facilitated my research in Czech Technical University in Prague.

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