

# Dynamic High-Temperature Tensile Characterization of an Iridium Alloy with Kolsky Tension Bar Techniques

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Abstract Conventional Kolsky tension bar techniques were modified to characterize an iridium alloy in tension at elevated strain rates and temperatures. The specimen was heated to elevated temperatures with an induction coil heater before dynamic loading; whereas, a cooling system was applied to keep the bars at room temperature during heating. A preload system was developed to generate a small pretension load in the bar system during heating in order to compensate for the effect of thermal expansion generated in the high-temperature tensile specimen. A laser system was applied to directly measure the displacements at both ends of the tensile specimen in order to calculate the strain in the specimen. A pair of high-sensitivity semiconductor strain gages was used to measure the weak transmitted force due to the low flow stress in the thin specimen at elevated temperatures. The dynamic hightemperature tensile stress–strain curves of a DOP-26 iridium alloy were experimentally obtained at two different strain rates ( $\sim$ 1000 and 3000 s<sup>-1</sup>) and temperatures ( $\sim$ 750 and 1030 °C). The effects of strain rate and temperature on the tensile stress–strain response of the iridium alloy were determined. The iridium alloy exhibited high ductility in stress–strain response that strongly depended on strain-rate and temperature.

Keywords Kolsky bar · High temperature · Dynamic tension - Iridium alloy - Stress–strain response

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#### Introduction

Iridium alloys possess unique property combinations of high melting temperature, high-temperature strength and ductility, and excellent oxidation and corrosion resistance [\[1](#page-8-0)], making them ideal for high-temperature applications. In some applications, the high-temperature impact response of the material must be fully understood in order to meet the safety requirements for the design of components. High-strain-rate and high-temperature stress– strain data are thus needed to develop strain-rate and temperature dependent material models for safety analysis. However, current mechanical characterization of iridium alloys has been limited to relatively low strain rates (below 50 s<sup>-1</sup>) [2-5], which are not sufficiently high for their potential applications. Song et al. [[6,](#page-8-0) [7](#page-8-0)] recently employed Kolsky compression bar testing, also known as split Hopkinson pressure bar testing, to characterize the compressive stress–strain properties of a DOP-26 iridium alloy at high strain rates  $(300-10,000 s^{-1})$  and high temperatures (750 and 1030 °C). The iridium alloy showed significant strain-rate and temperature effects on the compressive stress–strain response. However, the compression tests did not provide ductility and failure information at high strain rates and temperatures. Therefore, dynamic tensile characterization of iridium alloys at high temperatures was needed.

The Kolsky bar was originally developed for compression testing [\[8](#page-8-0)] and subsequently modified for use in tension and torsion testing [\[9](#page-8-0), [10](#page-8-0)]. A variety of Kolsky tensile bar techniques have been developed since the 1960s [\[11](#page-8-0)]. The most commonly used method is direct-tension Kolsky bars [[12,](#page-8-0) [13\]](#page-8-0). Recently, Song et al. [\[14](#page-8-0)] developed a directtension Kolsky bar that launches a solid striker to an end cap on the open end of the gun barrel to generate a dynamic

tensile load in the bar system. In contrast to the compression tests, more attention must be paid to the Kolsky tensile tests due to the complication of specimen attachment to the bar ends. A dog-bone-shaped cylinder with threads on both ends is a typical specimen design for Kolsky tension bar tests. However, special fixtures are needed when characterizing thin sheet specimens like the iridium specimen investigated in this study. In addition, the complex specimen fixtures may modify the stress wave propagation. For instance, a stress wave reflection may be generated at each interface within the joints/fixtures between the specimen and the bars. In this case, the reflected wave, which is usually used to calculate the strain rate and strain in the specimen, is no longer reliable. Direct displacement measurements, i.e. non-contact optical measurements, on both ends of the specimen are preferred.

It has been very challenging to conduct hightemperature Kolsky tension bar experiments. For hightemperature Kolsky bar tests, which are typically in compression, it is a common practice to heat the specimen individually while the whole bar system is kept at room temperature before dynamic loading. The high-temperature specimen is dynamically loaded immediately upon contact with the room-temperature bars. Cold contact time (CCT) has been defined as the time during which the high-temperature specimen stays in contact with the room-temperature bars until being dynamically loaded [\[11](#page-8-0)]. The CCT is usually required to be as short as several milliseconds, and even within 1 ms when the specimen is thin [\[7](#page-8-0)]. Appropriate modifications have been made to Kolsky compression bars for generating short CCTs. However, it has been difficult to directly apply these modifications to Kolsky tension bars. In a Kolsky tension bar test, the specimen has to be firmly attached to the bar ends before dynamic loading, which makes it nearly impossible to heat the specimen individually. Su et al. [\[15](#page-8-0)] applied a layer of thermal-protective coating with a very low heat transfer coefficient to the bar surface to mitigate the heat transfer to the bars. Using the same method, the temperature in the bars was reduced to below 300 $\degree$ C when the specimen was heated to 527 °C  $[16]$  $[16]$ . This means the heat was still transferred from the high-temperature specimen to the bars through the threads, which generated a thermal gradient in the bars. Such a thermal gradient may become more significant when the testing temperature increases. A significant thermal gradient will result in erroneous stress and strain measurements in the specimen, particularly when the test temperature is over  $600\text{ °C}$  where the Young's modulus of steel, a typical bar material, decreases significantly. These challenges limit current Kolsky tension bar tests to temperatures below 600 °C  $[15–19]$  $[15–19]$ , where the effects of temperature gradients on the steel bars can be neglected [\[18](#page-8-0)]. However, special experimental design considerations

are required for Kolsky tension bar experiments at higher temperatures, i.e., 750 and 1030  $^{\circ}$ C in this study.

In this study, the conventional Kolsky tension bar was modified to characterize dynamic tensile stress–strain response of a thin-sheet DOP-26 iridium alloy at elevated temperatures. The iridium alloy was characterized at two different strain rates ( $\sim$ 1000 and 3000 s<sup>-1</sup>) and temperatures ( $\sim$ 750 and 1030 °C) in order to determine the effects of strain rate and temperature on the tensile stress– strain response.

### Modified High-Temperature Kolsky Tension Bar System

The Kolsky tension bar described in [[14\]](#page-8-0) was modified in this study for dynamic high-temperature characterization of the DOP-26 iridium alloy. Figure [1](#page-2-0) shows a schematic of the modified high-temperature Kolsky tension bar system. Since it is not possible to directly thread a thin and flat iridium specimen into the bar ends, a pair of specimen fixtures was designed, as shown in Fig. [2a](#page-2-0). This design used a flat dog-bone shaped tensile specimen which is similar to the one used in Ref. [[20\]](#page-8-0). The fixture was machined with a slot with the same dimensions as the nongage section of the specimen such that the whole non-gage section of the specimen was placed into the slot of the fixture. This design allows the specimen shoulder take the load and transfer it to the specimen gage section. In addition, the area of the shoulder was much larger than that of the specimen gage section (Fig. [2](#page-2-0)a). The deformation in the non-gage section (shoulder) was therefore minimized during dynamic loading. The specimen was then covered with a semicircular cap. The depth of the fixture was made the same as the specimen thickness such that the semicircular cap did not provide additional perpendicular force on the specimen but retained the specimen during dynamic loading. Both the fixtures and the semicircular caps were made of Inconel 718 steel, which has a relatively high strength at elevated temperatures.

An induction coil heater, which was first applied by Rosenberg et al. [\[21](#page-8-0)] for high-temperature Kolsky bar experiments, was installed on the testing section. Due to the small size of the specimen under investigation, the induction coil was set to heat the relatively large fixtures and then transfer the heat to the specimen (Fig. [2b](#page-2-0)). This experimental design enabled direct measurement of displacements at the specimen ends using a laser extensometer. When the fixtures were heated with the induction coil, the heat was transferred to both the specimen and the bars simultaneously. In order to prevent heating of the bars, a pair of hollow water-cooled pillow blocks was installed on the bar ends, as shown in Fig. [2c](#page-2-0). This design is similar

<span id="page-2-0"></span>





Fig. 2 Details of the high-temperature Kolsky tension bar design. a Specimen and fixture design; b induction coil design; c cooling system design; d spring pre-tension load design

<span id="page-3-0"></span>in principle to that developed by Scapin et al. [[22\]](#page-8-0). The difference is that Scapin et al. [[22\]](#page-8-0) applied a Cortex-tubebased air cooling system to cool down the bars for testing up to 400  $^{\circ}$ C. The water cooling system used in this study was proven to be capable of cooling the bars below room temperature when the testing temperature was as high as 1030 °C.

As shown in Fig. [2](#page-2-0)b, the high-temperature fixtures attached to the bar ends still generated a thermal gradient between the bars and the specimen even though the bars were kept at room temperature, which might modify the wave propagation. In addition, the complicated design of the fixtures themselves exhibited many interfaces that might also modify the wave propagation. Both resulted in an unreliable reflected pulse for specimen strain calculation. As mentioned earlier, we employed a laser system to directly measure the displacement histories at the specimen ends of the fixtures. Non-contact laser systems have been recently used for direct displacement measurements in Kolsky bar experiments [[23,](#page-8-0) [24\]](#page-8-0). The splitting-beam laser extensometer presented in Ref. [[24\]](#page-8-0) was employed in this study to independently measure the displacements at both ends of the specimen.

The working principle of the laser system is illustrated in Fig. 3. A uniform laser line generator was used as a light source and then split into two independent beams. The movements of the incident and transmission bars, which consequently stretch the specimen, generate the intensity changes of the laser beams that are independently detected with two separate laser detectors. The laser detector used in this study had a bandwidth of 100 kHz or even higher, depending on the resolution, which is sufficiently high for Kolsky bar experiments. The resolutions of the laser detectors were able to be set independently to further increase the accuracy of the specimen strain measurement. In this study, the laser detectors have been calibrated to exhibit a perfect linearity but with different factors: 0.732 mm/v at

the incident bar side and 0.258 mm/v at the transmission bar side. The specimen strain can be calculated as

$$
\varepsilon = \frac{L_1 - L_2}{L_s},\tag{1}
$$

where  $L_1$  and  $L_2$  are displacements of the specimen ends attached to the incident and transmission bars, respectively;  $L<sub>s</sub>$  is the gage length of the specimen. As shown in Fig. [1](#page-2-0)a, the non-gage sections of the specimen were enclosed in the fixtures. Equation (1) maximizes the representation of the actual deformation of the gage section without the need for correction with respect to the deformation in the non-gage sections.

When the thin iridium specimen was heated, i.e. to 750 and  $1030$  °C in this study, the specimen became longer due to thermal expansion. However, the force generated by such a thermal expansion was not sufficiently high to overcome the friction between the bars and the bar supports and to push the incident and transmission bars back. As a consequence, the thin iridium specimen was buckled, which produced an erroneous stress–strain response and was difficult to correct numerically. A spring-loaded pretension system, as shown in Fig. [2](#page-2-0)d, was developed to prevent the specimen from buckling during heating. The spring was placed between a rigid mass and a flange (Fig. [2d](#page-2-0)). The flange was screwed toward the rigid mass to compress the spring, which in turn generated a tension load in the tension bar system including the tensile specimen. Another rigid mass was placed against the gun barrel (Fig. [1\)](#page-2-0) to prevent the bar system from moving backwards when the whole bar system was pre-loaded in tension [\[25](#page-8-0)]. In this study, the spring was set to generate a pre-tension load of approximately 18 N which is sufficient to straighten the iridium specimen during heating but insufficient to produce further stretch on the iridium specimen.

Another high temperature testing issue is thermal softening of the specimen. In general, the flow stress in



<span id="page-4-0"></span>metallic materials decreases significantly at elevated temperatures, which results in a very weak transmitted signal in Kolsky bar experiments. In order to measure the weak transmitted signal with relatively high resolution, a pair of semiconductor strain gages was used to replace the regular resistor strain gages on the transmission bar. The semiconductor strain gages had a gage factor of 139, which is approximately 70 times more sensitive than the regular resistor strain gages. The specimen stress is calculated as

$$
\sigma = \frac{E_0 A_0 \varepsilon_t}{A_s},\tag{2}
$$

where  $E_0$  is Young's modulus of the bar material;  $A_0$  is the cross-sectional area of the transmission bar;  $\varepsilon_t$  is the transmitted strain; and  $A_s$  is the cross-sectional area of the specimen.

Combining the measurements of the semiconductor strain gages (Eq. 2) and the laser system (Eq. [1\)](#page-3-0) for specimen stress and strain histories, respectively, yields the stress–strain curve of the specimen under investigation.

## Dynamic High-Temperature Tensile Characterization of Iridium Alloy

The material investigated in this study was DOP-26 iridium alloy which contains iridium as the base material alloyed with nominally 0.3 wt% tungsten, 60 wppm thorium, and 50 wppm aluminum  $[3, 4]$  $[3, 4]$  $[3, 4]$  $[3, 4]$ . The iridium tensile specimens were removed from prime DOP-26 alloy blanks using electrical discharge machining (EDM) with zinc-coated brass wire. The specimens were ground to remove the residual EDM layer, and then deburred and polished. All specimens were acid cleaned and then annealed at 1375 °C  $\pm$  25 °C for 1 h  $\pm$  10 min in vacuum (1  $\times$  10<sup>-4</sup> torr), producing an average grain size of approximately  $25 \mu m$ . Typical annealed DOP-26 iridium alloy microstructures may be found in Ref. [[5\]](#page-8-0). The tensile specimens had a thickness of 0.66 mm, a width of 2.54 mm, and a gage length of 6.35 mm. The specimen thickness was specified to match the thickness of interest for the applications. The detailed dimensions of the iridium specimens used in this study are shown in Fig. 4. In this study, the DOP-26 iridium alloy was dynamically characterized at 750 and 1030 °C, which are approximately 38 and 48  $\%$  of the melting temperature  $(2446 \degree C)$  of pure iridium, respectively.

In order to check the temperature uniformity in the specimen during heating, three thermocouples were attached to the gage section of the iridium specimen: one was placed in the center and the other two were placed close to each end of the specimen, respectively, as shown in Fig. [5](#page-5-0)a. Figure [5b](#page-5-0) shows the temperature histories from the



Fig. 4 Iridium tensile specimen

three thermocouples when the induction heater was programmed to heat the specimen to  $750^{\circ}$ C. The thermocouple signals were noisy with high frequencies caused by the electromagnetic field generated by the induction coil. In order to provide clearer temperature readings, a 10 Hz digital filter was applied to the thermocouple signals, the results of which are also shown in Fig. [5b](#page-5-0). It is noted that Fig. [5](#page-5-0)b records the whole heating process of the specimen. However, it takes much less time, usually 3–4 s (called "time window of testing" here), to complete the dynamic testing procedure, as marked in Fig. [5c](#page-5-0). The temperatures were very consistent during the heating process as indicated by the thermocouples in these three different locations. The results demonstrate the reasonable uniformity of the temperature across the whole specimen gage section (756  $\pm$  17 °C). In actual iridium alloy testing, only one thermocouple was attached to the fixture surface under the specimen, as shown in Fig. [2a](#page-2-0), to avoid spot welding the thermocouple to the surface and causing specimen microstructure changes.

Figure [6](#page-5-0) shows typical strain gage signals on the incident and transmission bars for the incident, reflected, and transmitted waves at the test temperature of 1030 °C. As shown in Fig. [6](#page-5-0), the transmission bar strain was low (only 15 microstrains) due to the low strength of the iridium specimen at such a high temperature. High resolution measurement of the low transmitted signal was achieved by using high sensitivity semiconductor strain gages. Again, the reflected pulse was not reliable and could not be used for specimen strain measurement. Instead, we used the laser system (Fig. [3](#page-3-0)) to directly track the movements of the specimen ends that were attached to the incident and transmission bars. The laser outputs are shown in Fig. [7.](#page-5-0) Figure [7](#page-5-0) clearly shows significant change in the laser

<span id="page-5-0"></span>

Fig. 5 Temperature uniformity during induction heating. a Induction coil and thermocouple locations; b temperatures recorded at three different locations; c temperatures over the duration of dynamic test

output for the front end (on the incident bar side) but no significant change for the back end (on the transmission bar side). This is because the transmitted force was too small to generate significant displacement on the transmission bar side.



Fig. 6 Typical incident, reflected, and transmitted signals



Fig. 7 Laser outputs for bar-end displacement measurements

The engineering stress and strain histories in the specimen, which were calculated with Eqs. ([2\)](#page-4-0) and ([1\)](#page-3-0), respectively, are shown in Fig. [8.](#page-6-0) The strain rate was then calculated with the slope of the strain history as a nearly constant  $860 s^{-1}$ . It is noted that when the reflected pulse becomes unreliable, it is difficult to compare the force histories at both ends of the specimen for force/stress equilibrium checks. In this study, we conducted preliminary tests with high-rate high-temperature digital image correlation (DIC) approaches to qualitatively check the deformation uniformity across the whole specimen gage section. The DIC results showed that, under the same loading condition, the specimen with a gage length of 6.35 mm was deformed uniformly until necking occurred. However, due to the interaction between the high-intensity light for the DIC measurement and the laser system used

<span id="page-6-0"></span>

Fig. 8 Engineering stress and strain histories

for the direct strain measurement in the specimen, we were not able to apply DIC in each test. Based on the stress and strain histories shown in Fig. 8, the tensile stress–strain curves at 860  $s^{-1}$  and 1030 °C were able to be obtained.

Following the same procedure, the iridium alloy was characterized in tension at two different strain rates  $({\sim}1000$  and 3000 s<sup>-1</sup>) and temperatures ( ${\sim}750$  and 1030 °C). At each condition, three experiments were repeated and the results were consistent (within 10 %) at the same testing condition. Figure 9 shows the mean tensile stress–strain curves of the iridium alloy at different strain rates and temperatures. It is noted that, due to the superior ductility of the iridium alloy at elevated temperatures, the specimens did not fail during the first dynamic tensile load except for the testing condition of 735  $\degree$ C/3450 s<sup>-1</sup>. At the condition of 735 °C/3450 s<sup>-1</sup>, the specimens possessed engineering failure strains varying between 0.5 and 0.7. It



Fig. 9 Engineering stress–strain curves at different strain rates and temperatures

is noted that the engineering failure strains were not representative of the actual failure strains since significant strain localization and necking occurred in the specimens before failure. Therefore, the stress–strain curves are plotted up to a strain of 0.5 in Fig. 9. It is noted that the stress–strain curve at 1030 °C/850 s<sup>-1</sup> ends at the strain of 0.4 due to unloading, not failure of the specimen. The dynamic high-temperature stress–strain curves of the iridium alloy show different profiles than quasi-static curves [\[3](#page-8-0)]. All stress–strain curves show an initial elasticity followed by significant work hardening behavior when the strain is below 10 %. This phenomenon may be related to a change in deformation mechanism at high strain rates and elevated temperatures, which is still under investigation. In addition, the stress–strain curve at 735 °C/3450 s<sup>-1</sup> had more oscillations than the others. The reason for this is still under investigation. When the strain is below 10 %, the stress–strain curves show neither strain-rate nor



Fig. 10 Strain rate effect on the tensile stress–strain response at different temperatures. **a** 750 °C; **b** 1030 or 1090 °C



Fig. 11 Temperature effect on the tensile stress–strain response at different strain rates.  $\mathbf{a} \sim 1000 \text{ s}^{-1}$ ;  $\mathbf{b} \sim 3000 \text{ s}^{-1}$ 

temperature effects. When the strain increases, the stress– strain curves show plastic flow with significant strain-rate and temperature effects. The insignificant work hardening behavior of the DOP-26 iridium alloy at high strain rates shown in Fig. [9](#page-6-0) is consistent with the quasi-static data shown in Ref. [[3\]](#page-8-0) at similar high temperatures. At strains greater than 10 % and similar strain rates the flow stresses decrease when the temperature increases from 750 to 1030 °C, showing significant thermal-softening behavior. At the same temperature, the flow stresses increase when the strain rate increases from  $\sim$  1000 to 3500 s<sup>-1</sup>, showing positive strain-rate sensitivity. Detailed strain-rate effects on the tensile stress–strain response including the quasi-static data at different strain rates and temperatures presented in Ref. [[3\]](#page-8-0) are shown in Fig. [10](#page-6-0) at two differ-ent temperatures (750 °C (Fig. [10](#page-6-0)a) and 1030/1090 °C (Fig. [10](#page-6-0)b), respectively. The data obtained from this study

were boxed and the rest are from the Ref. [\[3](#page-8-0)] in both figures. Figure [10](#page-6-0) clearly shows a significant strain-rate effect on the tensile flow stress of the iridium alloy at both temperatures. The strain-rate sensitivities are slightly different at the two temperatures and are also dependent on the level of strain. Figure 11 shows the temperature effect on the flow stresses at the strains of 0.15 and 0.30 and two different strain rates ( $\sim$  1000 and 3000 s<sup>-1</sup>), respectively. The flow stresses decreased when the temperature increased from 750 to 1030  $\degree$ C, showing significant thermal softening behavior.

#### Conclusion

The conventional direct-tension Kolsky bar was modified for high-temperature tensile characterization of the DOP-26 iridium alloy. An induction coil was applied to heat the iridium specimen to elevated temperatures up to  $1030$  °C while the specimen ends of the incident and transmission bars were cooled to reduce the thermal gradient in the bars. A pair of semiconductor strain gages on the transmission bar was used to directly measure the force/stress in the specimen during dynamic loading. A laser system was developed to independently measure the displacements at the specimen ends on the incident- and transmission-bar sides so that the specimen strain could be calculated. A spring-loaded pretension system was installed on the free end of the transmission bar to prevent the high-temperature specimen from buckling during heating. Dynamic tensile stress–strain curves of the iridium alloy were obtained at two elevated temperatures (750 and 1030  $^{\circ}$ C) and strain rates ( $\sim$  1000 and 3000 s<sup>-1</sup>). The iridium alloy shows high ductility at elevated temperatures and strain rates. The effects of strain rate and temperature on the tensile stress– strain response of the iridium alloy were also determined. The iridium alloy exhibited little sensitivity to strain rate or temperature when the strain was below 10 % but strong sensitivities to both strain rate and temperature when the strain was greater than 10 %. The DOP-26 iridium alloy exhibited significant strain-rate hardening behavior at the same temperatures and thermal softening behavior at the same strain rates.

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