#### **TECHNICAL PAPER**



# **A novel design of nanochannel structure in a micro–nanofuidic preconcentrator for electrokinetic ion enrichment**

**Wenbo Han1 · Xueye Chen[1](http://orcid.org/0000-0002-2377-352X)**

Received: 14 July 2019 / Accepted: 6 December 2019 / Published online: 16 December 2019 © The Brazilian Society of Mechanical Sciences and Engineering 2019

#### **Abstract**

Since the concentration of target molecules contained in most reagents in analytical chemistry experiments is lower than the minimum concentration required for subsequent detection, it is necessary for sample enrichment. Electrokinetic trapping that utilizes the principle of ion concentration polarization to achieve biomacromolecular enrichment is the most efficient. In this paper, based on the Poisson–Nernst–Plank equation, a novel design of nanochannel structure in a micro–nanofuidic preconcentrator for electrokinetic ion enrichment is carried out by numerical simulation. The results show that the enrichment process can be divided into three stages: enrichment generation, enrichment promotion and enrichment breakdown when the applied voltage is increased. Importantly, by comparing the six diferent structure nanochannels (straight line type, square wave type, rectangular type, circle type, zigzag type and multi-wave type), we found that the enrichment ratio produced by the rectangular nanochannel is the highest of the six nanochannels, which is 20.7 times. Rectangular nanochannel requires the lowest applied voltage to achieve the highest enrichment ratio. This work provides a novel design method and theoretical basis for the design of detection equipment in ultra-low concentration molecular detection tasks.

**Keywords** Nanochannel structure · Micro–nanofuidics · Electrokinetic ion enrichment · Ion concentration polarization · Poisson–Nernst–Plank

## **1 Introduction**

With the rapid development of micro-nano-manufacturing technology, micro–nanofuidic technology has been widely developed. The functions of preparation, enrichment, separation, purifcation, mixing, detection and reaction of the sample have been realized on the micro–nanofuidic chips [\[1](#page-8-0)–[3\]](#page-8-1). Therefore, miniaturization and personalized analysis experiments are completed on the micro–nanofuidic chip [\[4](#page-8-2), [5](#page-8-3)]. The main advantages of the micro–nanofuidic chips are small sample consumption, rapid response, low manufacturing cost and so on. So micro–nanofuidic chips have been widely used in life sciences, chemical engineering and environmental engineering [\[6](#page-8-4)[–10\]](#page-8-5).

Technical Editor: Erick de Moraes Franklin, Ph.D.

 $\boxtimes$  Xueye Chen xueye\_chen@126.com

Enrichment, separation, purification and detection are involved in the analysis of chemical experiments on micro–nanofuidic chips. The enrichment of the samples is mainly for molecular detection tasks with very low concentration. There are many ways to enrich, such as feld amplifed sample stacking (FASS) [\[11](#page-8-6)], isoelectric focusing (IEF) [\[12](#page-8-7)], dielectrophoretic trapping (DEPT) [\[13](#page-8-8), [14\]](#page-8-9) and electrokinetic trapping (EKT) [[15](#page-8-10)[–18\]](#page-8-11). Among them, EKT of enriching charged molecules by using the principle of ion concentration polarization (ICP) is the most efficient [\[19](#page-8-12)]. The principle of ICP mainly applies the characteristics of nanochannels for selective transport of ions [\[20](#page-8-13)]. As the pioneers of ICP theory, Pu et al. [[21](#page-8-14)] frst demonstrated the enrichment and consumption of ions at both ends of a negatively charged glass nanochannel and theoretically explained the ICP phenomenon. Wang et al. [\[22\]](#page-8-15) used a micro–nanofuidic chip to achieve an enrichment rate of fuorescein isothiocyanate that increased to 600 times in 120 s at 120 V. Jain et al. [[23](#page-8-16)] shown that the sub-2 nm pores separated in graphene are opposite to the larger pores and exhibit diferent transport behaviors consistent with ion transport under the free energy barrier caused by ion dehydration

 $1$  Faculty of Mechanical Engineering and Automation, Liaoning University of Technology, Jinzhou 121001, China

and electrostatic interaction. Yang et al. [[24\]](#page-8-17) found that the voltage-controlled surface charge pattern and ion current rectifcation are caused by the enhanced cation concentration in the bipolar bilayer and the negatively charged silicon nitride nanopore. Prakash et al. [[25\]](#page-8-18) indicated that while maintaining the overall system electrical neutrality, the addition of trace amounts of divalent  $Mg^{2+}$  ions to the nitric oxide (NaCl) electrolyte solution increased the maximum electroosmotic rate of the solution by nearly two times. Through investigating nanopore electrode arrays (NEA) with diferent geometries, Ma et al. [\[26](#page-8-19)] determined the individual contribution of ion accumulation and ion migration to the entire additional current amplifcation (AFad), achieving a large AFad value in NEA with smaller pores.

Most of the above studies focused on ion transport in straight line nanochannels or nanopores. However, few studies conducted on nanochannels with diferent structures. In the previous research work, we used numerical simulation to study the ion enrichment of high-viscosity fuids and the Cantor fractal nanochannel. In previous work, we found that the wall structure of nanochannels has a signifcant impact on electro-nanoparticle enrichment [[27,](#page-8-20) [28](#page-8-21)]. In view of the important guiding position of electrokinetic ion enrichment theory in ultra-low concentration molecular detection tasks, it is necessary to carry out research on ion transport in different structures of nanochannels.

In this paper, based on the Poisson–Nernst–Plank (PNP) equation, the ion enrichment of six diferent structure nanochannels in the micro–nanofuidic preconcentrator was studied by numerical simulation. Analysis of the principle of electrokinetic ion enrichment is discussed in detail. The results show that the enrichment process can be divided into three stages. The rectangular type has the highest enrichment rate of 20.7 times. This work provides a novel design for equipment in ultra-low concentration tasks.

#### **2 Methodology**

We have established a mathematical model and conducted numerical simulation analysis. The numerical simulation of this work is based on the fnite element software *COMSOL Multiphysics®*. In order to study the nanochannel structure in a micro–nanofuidic preconcentrator for nano-electrokinetic ion enrichment, a set of coupled partial diferential equations are solved, including the Poisson equation, the Nernst–Planck equation and the Navier–Stokes equation [\[29\]](#page-8-22).

*Electric feld* The Poisson equation (Eq. [1\)](#page-1-0) is applied to solve the distribution of the electrical potential. The distribution of the electrical potential depends mainly on the applied voltage and surface charge.

<span id="page-1-0"></span>
$$
\nabla \cdot (\varepsilon_r \varepsilon_0 \nabla \phi) = -\sum_{k=1}^N e z_k n_k \tag{1}
$$

In the above equation,  $\varepsilon_r$  is relative permittivity,  $\varepsilon_0$  denotes electrical permittivity of the vacuum,  $\phi$  represents potential, *e* is elementary charge,  $z_k$  is the valence of the *k*th ion, *k* is the *k*th species and  $n_k$  is concentration of the *k*th ion.

*Ionic concentration field* As shown in Eq. ([2](#page-1-1)), the Nernst–Planck equation was used to simulate ion transport and ion concentration distribution which depended on the electrophoresis, electroosmosis and difusion in the micro-nano-composite channel.

<span id="page-1-1"></span>
$$
\frac{\partial n_k}{\partial t} + \nabla \cdot \left( -\omega_k z_k n_k \nabla \phi \right) - \nabla \cdot \left( D_k \nabla n_k \right) + \nabla \cdot \left( n_k u \right) = 0 \tag{2}
$$

where *u* is fluid velocity,  $w_k$  and  $D_k$  are electrophoretic mobility and ion diffusion coefficient.

*Flow feld* The incompressible Navier–Stokes equation (Eq. [3](#page-1-2)) with electrical force  $(-\rho_e \nabla \phi)$  and the continuity equation (Eq. [4\)](#page-1-3) are used to describe the fow feld. This means that the density of the fuid is always constant.

<span id="page-1-2"></span>
$$
\rho_0(\partial_t u + u \cdot \nabla u) = -\nabla p + \mu \nabla^2 u - \rho_e \nabla \phi \tag{3}
$$

<span id="page-1-3"></span>
$$
\nabla \cdot u = 0 \tag{4}
$$

where  $\rho_0$  is fluid density, *p* is pressure,  $\mu$  is dynamic viscosity of fluid and  $\rho_e$  is volumetric charge density.

A two-dimensional numerical model as shown in Fig. [1a](#page-2-0) is designed. The preconcentrator consists of two microchannels and one nanochannel. The length and width of the microchannels are a and b, respectively. The width of the nanochannel is 20 nm. The vertical distance between the two microchannels is 1 um. Then six diferent nanochannels as shown in Fig. [1b](#page-2-0) were designed to connect the two microchannels. S1 is a straight line type. S2 is a square wave type. S3 is a rectangular type. S4 is a circle type. S5 is a zigzag type. S6 is a multi-wave type. Their dimensions are shown in Fig. [1](#page-2-0)b. In this paper, potassium chloride (KCl) can be completely disassociated into K+ ions and Cl− ions in aqueous solution. Detailed ion physical parameters are shown in Table [1](#page-2-1). The boundary conditions of the two-dimensional geometric model are shown in Table [2](#page-3-0).

#### **3 Mesh independence**

We use the model of the S1 structure to divide four different numbers of grids to test the grid independence. The number of grids is 8784, 10,724, 11,936 and 15,184. The physical parameters and boundary conditions of the fuid are set as the parameters in Tables [1](#page-2-1) and [2.](#page-3-0) The applied electric <span id="page-2-0"></span>**Fig. 1** The two-dimensional design of the micro–nanofuidic preconcentrator. **a** Twodimensional geometric model of the preconcentrator with micro-nano-composite channels. **b** Six diferent structures of nanochannels: (S1) straight line type, (S2) square wave type, (S3) rectangular type, (S4) circle type, (S5) zigzag type, (S6) multi-wave type



feld voltage is fxed at 100 V. We obtained the velocity profle at the centerline of the nanochannel under diferent grids as shown in Fig. [2.](#page-3-1) From Fig. [2,](#page-3-1) it is found that the

<span id="page-2-1"></span>**Table 1** The value of the ion

velocity value is within the error tolerance range when the number of grids is 11,936. Therefore, in order to reduce the amount of calculation and maintain the accuracy of the

<span id="page-3-0"></span>**Table 2** Boundary conditions of the model in the simulation

	Inlet and outlet	The wall of the microchannel The wall of the nanochannel	
	Electric field $\phi_{\text{inlet}} = 1-200 \text{ V}, \phi_{\text{outlet}} = 0 \text{ V}$ $\sigma_{\text{micro}} = -8e^{-4}C/m^2$		$\sigma_{\text{nano}} = -8e^{-4}C/m^2$
Ionic field	$cinlet = coutlet = 0.1 mol/m3$	$i_1 = 0$	$j_1 = 0$
Flow field	$p=0$ Pa	$u=0$	$u=0$



<span id="page-3-1"></span>**Fig. 2** Velocity at the centerline of the nanochannel under diferent grids

calculation results, we use a meshing scheme with a grid number 11,936. The mesh in the nanochannel is 5 nm, and the mesh in the microchannel is 50 nm.

# **4 Results and discussion**

Ion enrichment through straight line structure nanochannel (S1) was investigated. It can be clearly seen from Fig. [3a](#page-4-0) that the process of ion enrichment has a peak voltage. The enrichment ratio is gradually increased as the applied voltage is increased. When the applied voltage reaches a peak, the enrichment ratio reaches a maximum value. It can be seen from Fig. [3a](#page-4-0) that the enrichment ratio is 10.86 and the peak voltage is 56 V. When the applied voltage continues to increase, the enrichment ratio begins to decrease. When the applied voltage is between 56 and 140 V, the enrichment ratio drops rapidly. When the applied voltage is between 140 and 200 V, the enrichment rate decreases at a slower rate. Therefore, the process of ion enrichment can be divided into three stages: enrichment generation, enrichment promotion and enrichment breakdown. To further explain the reasons, the ion concentration image at the applied voltages of 16 V, 36 V, 56 V and 76 V, the concentration curve and the velocity curve on the centerline of the micro-nanochannel were plotted. As can be seen from Fig. [3](#page-4-0)b–d, when the applied voltage rises from 16 to 36 V, the concentration near the entrance of the nanochannel decreases. However, the concentration near the exit of the nanochannel increases. This is due to the selective transport properties of the nanochannels. When the applied voltage has not reached the peak voltage, even if the applied voltage is increased, the electrophoretic force generated by the electrophoretic effect is not sufficient to balance the repulsive force generated by the overlap of the electric double layer. The dominant repulsive force causes ions to pass through the nanochannels to the anode region, thus creating an enrichment at the exit of the nanochannel. As the applied voltage increases, the electroosmotic fow in the nanochannels strengthens. Therefore, the fow rate in the nanochannel is increased. When the applied voltage reaches a peak voltage of 56 V, the ion concentration of the enrichment region reaches a peak. The velocity in the nanochannel also reaches the peak velocity. As the applied voltage continues to increase, the ion concentration in the enrichment zone decreases. However, the ion concentration in the depletion region increases. The velocity within the nanochannel decreases. This is because the electrophoretic force is generated by the electrophoretic efect when the applied voltage is increased exceeds the repulsive force generated by the overlap of the electric double layer. Therefore, the ion selective permeability of the nanochannel is broken. The ions reach the anode region through the nanochannel due to the electrophoretic effect.

In order to study the efect of nanochannel structure on ion enrichment, we added the same boundary conditions as S1 on the remaining fve nanochannels. The curve of the ion enrichment ratio of the six diferent structural nanochannels and the concentration image corresponding to the peak voltage were plotted. It can be seen from Fig. [4](#page-5-0)a that the rates of increase in the ion enrichment ratio of the four channels of S1, S2, S5 and S6 are uniform when the applied voltage is increased. The rates of increase in S2 and S3 ion enrichment ratio are signifcantly higher than that of the other four structures. This is due to the relatively complicated shuntconvergence structure of S3 and S4. The fluid needs to flow longer in S3 and S4 structures. It is found from Fig. [4](#page-5-0)a that when the nanochannel has such a complicated structure, the enrichment performance of the nanochannel is enhanced. In order to clearly see the ion concentration distribution at the peak voltage, ion concentration diagrams of the six structures at the peak voltage were obtained, respectively. Figure [4b](#page-5-0) shows that the S3 structure has a peak voltage of 97 V consistent with the S5 structure. However, the enrichment



<span id="page-4-0"></span>**Fig. 3 a** The curve of the enrichment ratio of the straight line type with the applied voltage. **b** Anion concentration image when applied voltages are 16 V, 36 V, 56 V and 76 V, respectively. **c** The concentration curve of anions on the centerline of the micro-nano-composite

channel when the applied voltages are 16 V, 36 V, 56 V and 76 V, respectively. **d** The velocity curve at the centerline of the micro-nanocomposite channel when the applied voltages are 16 V, 36 V, 56 V and 76 V

<span id="page-5-0"></span>**Fig. 4 a** The curve of the ion enrichment ratio of the six diferent structure nanochannels. **b** The concentration image corresponding to the peak voltage for six diferent structure nanochannels



performance of S3 structure is much higher than that of S5. The enrichment ratio of S3 is higher than that of S5 by fve times. It can be known from the peak voltages of the six structures that the peak voltage of the nanochannel having the shunt-convergence structure is lower than the peak voltage of the three structures (S2, S5 and S6). In terms of enrichment ratio, the enrichment ratio of S2 and S3 is 20.7 times. However, the applied voltage of S3 is signifcantly lower than S2. This means that the advantages of the S3 structure are particularly evident in some experiments that require lower voltage. Then, four kinds of S1, S2, S5 and S6 without complex shunt-convergence structure were studied. The study found that their enrichment ratio and peak voltage are ranked as S2>S6>S5>S1. It is revealed that anions do not easily pass through such nanochannel due to the presence of a right-angled nanochannel in the S2 structure. Therefore, the S2 structure requires a larger applied voltage to provide a stronger electrophoretic force to overcome this obstacle due to the right-angled nanochannel.

In order to further study the efect of the nanochannels of this shunt-convergence structure on electrokinetic ion enrichment, we have conducted in-depth research on the two confuent structures of S3 and S4. Previous studies have shown that ion fux is mainly composed of electrophoresis fux, electroosmotic fux and difusion fux [\[30](#page-8-23)]. Since the diffusion coefficient of ions is very small, the diffusion flux



is ignored in this paper. The transport forms of the anion and cation in the micro-nanochannel caused by the applied voltage mainly include electrophoresis migration of ions and ion transport caused by fuid electroosmotic fow. The direction of electrophoresis of the anion is from the cathode to the anode. The electroosmotic fow direction of the fuid is from the anode to the cathode. Because of the opposite ion transport directions caused by electrophoresis and electroosmotic flow, in order to facilitate the comparison of the effects of electrophoresis and electroosmotic fow on ion enrichment,

a relative fux is defned. The relative fux is equal to the value of the electrophoresis fux minus the electroosmotic fux as follows,

$$
\Delta EP - EOF = \int_{\Omega} (\omega_k z_k n_k E - u n_k) d\Omega \tag{5}
$$

In this way, the enrichment magnifcation can be quantitatively calculated, and the integration region is as shown in Fig. [5](#page-7-0)a.

As can be seen from Fig. [5b](#page-7-0), c, the electrophoresis fux of anions during the ion enrichment process plays a decisive role. The anion's electroosmotic fux is very weak. It can be seen from Fig. [5d](#page-7-0), e that the anion relative fux of the rectangular nanochannel is higher than that of the circular structure in the case of the same applied voltage. Therefore, it also explains why the enrichment ratio of the rectangular nanochannel is higher than the enrichment ratio of the circular nanochannel when the applied voltage is the same.



<span id="page-7-0"></span>**Fig. 5 a** Schematic diagram of the integration zone, **b** electroosmotic fux of anions in circle and rectangular nanochannels, **c** electrophoresis fux of anions in circle and rectangular nanochannels, **d** relative

fux (ΔEP-EOF) of anions in circle and rectangular nanochannels, **e** ion enrichment ratio in circle and rectangular nanochannels

## **5 Conclusion**

In this paper, the ion enrichment of six diferent structure nanochannels in the micro–nanofluidic preconcentrator was studied by numerical simulation. The results show that with the increase in applied voltage, ion enrichment can be divided into three stages: enrichment generation, enrichment promotion and enrichment breakdown. In addition, the study found that when the applied voltage is the same, the enrichment ratio of the rectangular nanochannel is the highest among the six structural species. Rectangular channels have a maximum enrichment rate of up to 20.7 times. The research results in this paper lay a theoretical foundation for the design of detection equipment in ultra-low concentration molecular detection tasks.

**Acknowledgements** This work was supported by Liaoning Natural Science Foundation (2019-MS-169), the Key Project of Department of Education of Liaoning Province (JZL201715401) and Liaoning BaiQianWan Talents Program. We sincerely thank Prof. Chong Liu for his kind guidance.

#### **Compliance with ethical standards**

**Conflict of interest** The authors declare no confict of interest.

#### **References**

- <span id="page-8-0"></span>1. Han W, Chen X, Hu Z, Yang K (2018) Three-dimensional numerical simulation of a droplet generation in a double T-junction microchannel. J Micro/Nanolithogr MEMS MOEMS 17(2):025502
- 2. Han W, Chen X (2019) Numerical simulation of the droplet formation in a T-junction microchannel by a level-set method. Aust J Chem 71(12):957–964
- <span id="page-8-1"></span>3. Han W, Chen X (2019) New insights into the pressure during the merged droplet formation in the squeezing time. Chem Eng Res Des 145:213–225
- <span id="page-8-2"></span>4. Zhang S, Chen X (2019) Secondary bonding of PMMA micromixer with high-pressure. Microchem J 144:339–344
- <span id="page-8-3"></span>5. Zhang S, Chen X, Wu Z, Zheng Y (2019) Numerical study on stagger Koch fractal baffles micromixer. Int J Heat Mass Transf 133:1065–1073
- <span id="page-8-4"></span>6. Wang X, Zhu C, Fu T, Ma Y (2014) Critical lengths for the transition of bubble breakup in microfuidic T-junctions. Chem Eng Sci 111:244–254
- 7. Timung S, Chaudhuri J, Borthakur MP, Mandal TK, Biswas G, Bandyopadhyay D (2017) Electric feld mediated spraying of miniaturized droplets inside microchannel. Electrophoresis 38(11):1450–1457
- 8. Santra S, Das S, Das SS, Chakraborty S (2018) Surfactant-induced retardation in lateral migration of droplets in a microfuidic confinement. Microfluid Nanofluid 22(8):88
- 9. Anna SL, Bontoux N, Stone HA (2003) Formation of dispersions using "fow focusing" in microchannels. Appl Phys Lett 82(3):364–366
- <span id="page-8-5"></span>10. Kim SH, Kim JW, Kim DH, Han SH, Weitz DA (2013) Enhancedthroughput production of polymersomes using a parallelized capillary microfuidic device. Microfuid Nanofuid 14(3–4):509–514
- <span id="page-8-6"></span>11. Burgi DS, Chien RL (1991) Optimization in sample stacking for high-performance capillary electrophoresis. Anal Chem 63(18):2042–2047
- <span id="page-8-7"></span>12. Cui H, Horiuchi K, Dutta P, Ivory CF (2005) Multistage isoelectric focusing in a polymeric microfuidic chip. Anal Chem 77(24):7878–7886
- <span id="page-8-8"></span>13. Gebauer P, Boček P (2000) Recent progress in capillary isotachophoresis. Electrophoresis Int J 21(18):3898–3904
- <span id="page-8-9"></span>14. Asbury CL, Van Den Engh G (1998) Trapping of DNA in nonuniform oscillating electric felds. Biophys J 74(2):1024–1030
- <span id="page-8-10"></span>15. Kim SJ, Ko SH, Kang KH, Han J (2010) Direct seawater desalination by ion concentration polarization. Nat Nanotechnol 5(4):297
- 16. Yeh LH, Zhang M, Hu N, Joo SW, Qian S, Hsu JP (2012) Electrokinetic ion and fuid transport in nanopores functionalized by polyelectrolyte brushes. Nanoscale 4(16):5169–5177
- 17. Choi E, Kwon K, Lee SJ, Kim D, Park J (2015) Non-equilibrium electrokinetic micromixer with 3D nanochannel networks. Lab Chip 15(8):1794–1798
- <span id="page-8-11"></span>18. Gong L, Ouyang W, Li Z, Han J (2018) Direct numerical simulation of continuous lithium extraction from high  $Mg^{2+}/Li^{+}$  ratio brines using microfuidic channels with ion concentration polarization. J Membr Sci 556:34–41
- <span id="page-8-12"></span>19. MacDonald BD, Gong MM, Zhang P, Sinton D (2014) Out-ofplane ion concentration polarization for scalable water desalination. Lab Chip 14(4):681–685
- <span id="page-8-13"></span>20. Li M, Anand RK (2016) Recent advancements in ion concentration polarization. Analyst 141(12):3496–3510
- <span id="page-8-14"></span>21. Pu Q, Yun J, Temkin H, Liu S (2004) Ion-enrichment and ion-depletion effect of nanochannel structures. Nano Lett 4(6):1099–1103
- <span id="page-8-15"></span>22. Wang JY, Xu Z, Li YK, Liu C, Liu JS, Chen L, Du L, Wang LD (2013) Nanopore density efect of polyacrylamide gel plug on electrokinetic ion enrichment in a micro-nanofuidic chip. Appl Phys Lett 103(4):043103
- <span id="page-8-16"></span>23. Jain T, Rasera BC, Guerrero RJS, Boutilier MS, O'hern SC, Idrobo JC, Karnik R (2015) Heterogeneous sub-continuum ionic transport in statistically isolated graphene nanopores. Nat Nanotechnol 10(12):1053
- <span id="page-8-17"></span>24. Yang C, Hinkle P, Menestrina J, Vlassiouk IV, Siwy ZS (2016) Polarization of gold in nanopores leads to ion current rectifcation. J Phys Chem Lett 7(20):4152–4158
- <span id="page-8-18"></span>25. Prakash S, Zambrano HA, Rangharajan KK, Rosenthal-Kim E, Vasquez N, Conlisk AT (2016) Electrokinetic transport of monovalent and divalent cations in silica nanochannels. Microfuid Nanofluid 20(1):8
- <span id="page-8-19"></span>26. Ma C, Xu W, Wichert WR, Bohn PW (2016) Ion accumulation and migration efects on redox cycling in nanopore electrode arrays at low ionic strength. ACS Nano 10(3):3658–3664
- <span id="page-8-20"></span>27. Han W, Chen X (2019) Nano-electrokinetic ion enrichment in a micro-nanofuidic preconcentrator with nanochannel's Cantor fractal wall structure. Appl Nanosci. [https://doi.org/10.1007/s1320](https://doi.org/10.1007/s13204-019-01049-7) [4-019-01049-7](https://doi.org/10.1007/s13204-019-01049-7)
- <span id="page-8-21"></span>28. Han W, Chen X (2019) Nano-electrokinetic ion enrichment of highly viscous fuids in micro-nanochannel. Chem Eng Process Process Intensif 143:107626
- <span id="page-8-22"></span>29. Wang Y, Pant K, Chen Z, Wang G, Difey WF, Ashley P, Sundaram S (2009) Numerical analysis of electrokinetic transport in micro-nanofuidic interconnect preconcentrator in hydrodynamic flow. Microfluid Nanofluid 7(5):683
- <span id="page-8-23"></span>30. Wang J, Liu C, Xu Z (2016) Electrokinetic ion transport in confned micro-nanochannel. Electrophoresis 37(5–6):769–774

**Publisher's Note** Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.