

A new effective coherent numerical technique based on shifted Vieta–Fibonacci polynomials for solving stochastic fractional integro-differential equation

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Abstract

In this article, an operational matrix method based on shifted Vieta–Fibonacci polynomials is utilised to find the numerical solution of fractional order stochastic integro-differential equations. In this method, the operational matrices are developed by using the shifted Vieta– Fibonacci polynomials for the fractional order Caputo differential operator in order to solve the present concerned problem. Using Newton cotes nodes as collocation points, operational matrices are employed to convert the above-mentioned equation into a system of linear algebraic equations. The coherent procedure for the appropriate numerical technique is described in this article. Additionally, the convergence analysis and error bound of the suggested method are well established. In order to illustrate the effectiveness, consistency, plausibility, and reliability of the proposed technique, three numerical examples are given. Moreover, the results obtained by the proposed method have been compared with those obtained by the Chelyshkov operational matrix method.

Keywords Fractional stochastic integro-differential equation · Itô integral · Brownian motion · Vieta–Fibonacci polynomial · Convergence analysis

Mathematics Subject Classification 60H20 · 34A08 · 97N50 · 65D30 · 41A15

1 Introduction

It is generally known that fractional derivatives may characterise the memory and heredity properties of certain materials and processes in ways that integer order derivatives can not. Recently, many applications across a wide range of fields, including viscoelastic materials (Meral et al[.](#page-24-0) [2010\)](#page-24-0), signal processing (Machado and Lope[s](#page-24-1) [2015\)](#page-24-1), meteorology, earthquake

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(El-Misiery and Ahme[d](#page-24-2) [2006\)](#page-24-2), optimal control (Sahu and Saha Ra[y](#page-24-3) [2018](#page-24-3)), fluid-dynamic (Momani and Odiba[t](#page-24-4) [2006\)](#page-24-4), qua[n](#page-24-5)tum mechanics (Atman and Şirin [2020](#page-24-5)), finance (Scalas et al[.](#page-24-6) [2000](#page-24-6)) and in other fields of science and engineering (Sun et al[.](#page-24-7) [2018](#page-24-7)) have been remodeled using fractional calculus. Integro-differential equations have a strong physical foundation and are widely used in fields of study including polymer rheology (Lodge et al[.](#page-24-8) [1978](#page-24-8)) and population model (Yzbaı et al[.](#page-24-9) [2013\)](#page-24-9). Deterministic fractional equations such as fractional order pantograph Volterra delay-integro-differential equations (Behera and Saha Ra[y](#page-24-10) [2022\)](#page-24-10), Riemann–Liouville fractional integro-differential equations (Ahmad and Niet[o](#page-23-0) [2011](#page-23-0)), fractional integro-differential equations (Arikoglu and Ozko[l](#page-24-11) [2009\)](#page-24-11) are used to represent real physical problems, which often rely on a noise source that is disregarded owing to the absence of sophisticated computational tools. As computational power has increased recently, real world phenomena can now be more effectively modeled using stochastic fractional equations such as stochastic fractional differential equations, stochastic fractional integral equations, stochastic fractional integro-differential equations (SFIDE).

This article investigates the numerical solution of the following SFIDE:

$$
{}^{C}D_{\eta}^{\alpha}z(\eta) = g(\eta) + \lambda_{1} \int_{0}^{\eta} \kappa_{1}(\eta, \zeta)z(\zeta)d\zeta + \lambda_{2} \int_{0}^{\eta} \kappa_{2}(\eta, \zeta)z(\zeta)dB(\zeta),
$$

\n
$$
\eta \in [0, 1], \ z(0) = z_{0},
$$
\n(1.1)

where λ_1 and λ_2 are constant numbers and ${}^CD_{\eta}^{\alpha}$ is the Caputo fractional differential operator of order $0 < \alpha < 1$. In Eq. [\(1.1\)](#page-1-0), $g(\eta)$ and $\kappa_i(\eta, \zeta)$ for $i = 1, 2$ are known smooth functions, and $z(\eta)$ is an unknown function. Brownian motion process is defined as $B = \{B(t); t \geq 0\}$ and $z(\eta)$ is a stochastic process defined on the probability space $(\Omega, \mathcal{F}, \mathbb{P})$. This is referred to as the solution of SFIDE.

There are several numerical methods to solve SFIDE, such as block pulse approximation (Mirzaee et al[.](#page-24-12) [2019\)](#page-24-12), cubic B spline approximation (Mirzaee and Alipou[r](#page-24-13) [2020\)](#page-24-13), meshless discrete collocation method based on radial basis functions (Mirzaee and Samadya[r](#page-24-14) [2019\)](#page-24-14), shifted Legendre spectral collocation method (Taheri et al[.](#page-24-15) [2017\)](#page-24-15), Bernstein polynomials approximation (Mirzaee and Samadya[r](#page-24-16) [2017](#page-24-16)), Galerkin method (Kamran[i](#page-24-17) [2016\)](#page-24-17), explicit finite difference method (Saha Ray and Patr[a](#page-24-18) [2013\)](#page-24-18) and different other methods that have been implemented to solve SFIDE.

The main motivation of this study is to solve the SFIDE (Eq. (1.1)) using shifted Vieta– Fibonacci polynomials. These kinds of equations may be found in many different fields, including physics, biology, physiology, optics, and climatology. Explicitly solving SFIDE can be difficult and time-consuming. So, here, the operational matrix method is implemented to solve these equations. Using shifted Vieta–Fibonacci polynomials, a new stochastic operational matrix has been derived for the first time in this paper. The proposed method is effective, applicable, and consistent.

In this study, the numerical results of Eq. (1.1) obtained by the shifted Vieta–Fibonacci operational matrix (SVFOM) method are further compared with the orthonormal Chelyshkov operational matrix (OCOM) method and actual solutions. Equation [\(1.1\)](#page-1-0) can be transformed into a system of algebraic equations by using operational matrices along with suitable collocation points. The resultant equations can be easily solved to get the desired approximate solution.

This article is organised as follows:

A few fundamental concepts about shifted Vieta–Fibonacci polynomials (SVFPs), stochastic calculus, and fractional calculus have been introduced in Sect. [2.](#page-2-0) In Sect. [3,](#page-4-0) operational matrices (OMs) for product, integral, fractional, and stochastic integrals have been

constructed using shifted Vieta–Fibonacci polynomial. The SFIDE problem is handled using the suggested operational matrix approach in Sect. [4,](#page-7-0) which also provides a review of the collocation technique. In Sect. [5,](#page-8-0) theorems relating to error estimation and convergence analysis are covered. Section [6](#page-13-0) represents the reliability and efficiency of the suggested numerical method using a few illustrative examples, and a brief overview is provided in Sect. [7.](#page-19-0)

2 Preliminaries

This section covers the properties of SVFPs as well as some fundamental stochastic calculus concepts.

2.1 Stochastic calculus

Definition 1 (*Itô Integra[l](#page-24-19)* (Øksendal [2003](#page-24-19))) Let $V = V(U, V)$ be the class functions $g(\gamma, \delta)$: $[0, \infty) \times \Omega \to \mathbb{R}$ and $g \in V(U, V)$. Thus, the definition of the Itô integral of *g* is given by

$$
\int_{U}^{V} g(\gamma, \delta) dB_{\gamma}(\delta) = \lim_{m \to \infty} \int_{U}^{V} \psi_{m}(\gamma, \delta) dB_{\gamma}(\delta) \quad (\lim_{n \to \infty} L^{2}(\mathbb{P})), \tag{2.1}
$$

where ψ_m is a sequence of elementary functions such that

$$
E\left[\int_{U}^{V}(g(\gamma,\delta)-\psi_m(\gamma,\delta))^2d\gamma\right]\to 0 \text{ as } m\to\infty.
$$
 (2.2)

Theorem 2.1.1 (The Itô isometry (Øksenda[l](#page-24-19) [2003\)](#page-24-19)). Let $g \in V(U, V)$, be elementary and *bounded functions. Then*

$$
E\left[\left(\int_{U}^{V} g(\gamma,\delta) dB_{\gamma}(\delta)\right)^{2}\right] = E\left[\int_{U}^{V} g^{2}(\gamma,\delta) d\gamma\right].
$$
 (2.3)

2.2 Fractional calculus

Definition 2 Consider $p - 1 < \alpha < p, \alpha > 0, \eta > 0, \alpha, \eta \in \mathbb{R}$, then the Caputo fractional differential operator ${}^CD_{\eta}^{\alpha}z(\eta)$ of order α is defined as (Saha Ra[y](#page-24-20) [2015](#page-24-20))

$$
{}^{C}D_{\eta}^{\alpha}z(\eta) = \frac{1}{\Gamma(p-\alpha)} \int_{0}^{\eta} (\eta - \zeta)^{p-\alpha-1} z^{(p)}(\zeta) d\zeta.
$$
 (2.4)

Also, the Riemann–Liouville (RL) fractional integral operator J_{η}^{α} of order α is defined as

$$
J_{\eta}^{\alpha}z(\eta) = \frac{1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \zeta)^{\alpha - 1} z(\zeta) d\zeta, \quad J_{\eta}^0 z(\eta) = z(\eta). \tag{2.5}
$$

The operators ${}^C D^{\alpha}_{\eta}$ and J^{α}_{η} has the following characteristics:

- 1. $J_{\eta}^{\alpha}(\delta_1 z(\eta) + \delta_2 z(\eta)) = \delta_1 J_{\eta}^{\alpha}(z(\eta)) + \delta_2 J_{\eta}^{\alpha}(z(\eta))$, $\alpha \ge 0$.
- 2. $J_{\eta}^{\beta_1} J_{\eta}^{\beta_2} z(\eta) = J_{\eta}^{\beta_1 + \beta_2} z(\eta), \quad \beta_1, \beta_2 \ge 0.$
- 3. ${}^CD_\eta^\alpha J_\eta^\alpha z(\eta) = z(\eta), \quad \alpha \ge 0.$
- 4. $J_{\eta}^{\alpha}({}^C D_{\eta}^{\alpha})z(\eta) = z(\eta) \sum_{i=0}^{p-1} z^i(0) \frac{\eta^i}{i!}, \quad p-1 < \alpha < p, \quad p \in \mathbb{N}.$

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2.3 Shifted Vieta–Fibonacci polynomials and its characteristics

Many problems in mathematical physics have been solved using SVFPs, such as the Lane– Emden equation, reaction-advection–diffusion, Emden–Fowler equation, etc.

Vieta–Fibonacci polynomials

According to the following relation, the Vieta–Fibonacci polynomials $V_{\tau_m}(\eta)$ of degree *m* in η are defined on the interval $[-2, 2]$.

$$
\mathcal{VF}_m(\eta) = \frac{\sin(m\theta)}{\sin\theta},
$$

where $\eta = 2 \cos \theta$ and $\theta \in [0, \pi]$.

These polynomials can also be generated by the following recurrence relation:

$$
\mathcal{VF}_m(\eta) = \eta \mathcal{VF}_{m-1}(\eta) - \mathcal{VF}_{m-2}(\eta), \quad m = 2, 3, \dots,
$$

with the initial values $\mathcal{VF}_0(\eta) = 0$, $\mathcal{VF}_1(\eta) = 1$. **Shifted Vieta–Fibonacci polynomials**

Definition 5 The shifted Vieta–Fibonacci polynomials $V\mathcal{F}_{m}^{*}(\eta)$, of degree *m* in η on [0, 1] are defined as follows (Sadri et al[.](#page-24-21) [2022](#page-24-21))

$$
\mathcal{VF}_m^*(\eta) = \mathcal{VF}_m(4\eta - 2).
$$

Also, these polynomials, can be generated via the following recurrence relation:

$$
\mathcal{VF}_{m}^{*}(\eta) = (4\eta - 2)\mathcal{VF}_{m-1}^{*}(\eta) - \mathcal{VF}_{m-2}^{*}(\eta), \quad m = 2, 3, ..., \tag{2.6}
$$

using the initial values $V\mathcal{F}_0^*(\eta) = 0$, $V\mathcal{F}_1^*(\eta) = 1$.

The SVFPs are also defined by using the following series:

$$
\mathcal{VF}_{m}^{*}(\eta) = \sum_{l=0}^{m-1} \frac{(-1)^{m-l-1} 2^{2l} \Gamma(m+l+1)}{\Gamma(m-l) \Gamma(2l+2)} \eta^{l}, \quad m = 2, 3, \dots, \tag{2.7}
$$

These polynomials are orthogonal with respect to the weight function $w(\eta) = \sqrt{\eta - \eta^2}$, i.e.

$$
\int_0^1 \mathcal{VF}_m^*(\eta) \mathcal{VF}_n^*(\eta) w(\eta) d\eta = \begin{cases} \frac{\pi}{8} & m = n \neq 0 \\ 0 & m \neq n \end{cases} \tag{2.8}
$$

2.4 Function approximation by SVFPs

Let $H = L^2_{\omega}(I), I = [0, 1],$ and $S = \text{span}\{ \mathcal{VF}_1^*(\eta), \mathcal{VF}_2^*(\eta), \dots, \mathcal{VF}_{m+1}^*(\eta) \}$. Then for any $y(\eta) \in H$, $y_m(\eta) \in S$ is a best approximation; that is

$$
y(\eta) \simeq y_m(\eta) = \sum_{i=1}^{m+1} a_i \mathcal{V} \mathcal{F}_i^*(\eta) = A^T V_F^*(\eta), \tag{2.9}
$$

where $A = [a_1, a_2, \dots, a_{m+1}]^T$ and $V_F^*(\eta) = [\mathcal{VF}_1^*(\eta), \mathcal{VF}_2^*(\eta), \dots, \mathcal{VF}_{m+1}^*(\eta)]^T$. Furthermore,

$$
a_i = \left(\frac{8}{\pi}\right) \int_0^1 y(\eta) \mathcal{V} \mathcal{F}_i^*(\eta) w(\eta) d\eta, \quad i = 1, 2, \dots, m+1. \tag{2.10}
$$

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In order to approximate the two-dimensional kernel function $\kappa(\eta, \zeta)$, following approximation is used.

$$
\kappa(\eta,\zeta) \simeq \kappa_{m,m}(\eta,\zeta) = \sum_{i=1}^{m+1} \sum_{j=1}^{m+1} \kappa_{ij} \mathcal{V} \mathcal{F}_i^*(\eta) \mathcal{V} \mathcal{F}_j^*(\zeta) = V_F^* T(\eta) \mathbb{K} V_F^*(\zeta), \qquad (2.11)
$$

where K is a $(m + 1) \times (m + 1)$ order kernal matrix.

Here, the orthogonality property of the SVFPs, together with the weight function $w(\eta)$ in Eq. [\(2.8\)](#page-3-0), is used to generate the kernel matrix.

It follows

$$
\mathbb{K} = Q^{-1} \left(\int_0^1 w(\eta) V_F^*(\eta) \left(\int_0^1 \kappa(\eta, \zeta) V_F^{*T}(\zeta) w(\zeta) d\zeta \right) d\eta \right) Q^{-1}, \tag{2.12}
$$

where

$$
Q = \left\langle V_F^*(.) , V_F^{*T}(.) \right\rangle_w.
$$

The matrix form for these SVFPs is as follows:

$$
V_F^*(\eta) = \tilde{A}L_m(\eta),\tag{2.13}
$$

where

$$
V_F^*(\eta) = [\mathcal{V}\mathcal{F}_1^*(\eta), \mathcal{V}\mathcal{F}_2^*(\eta), \dots, \mathcal{V}\mathcal{F}_{m+1}^*(\eta)]^T, \quad L_m(\eta) = [1, \eta, \dots, \eta^m]^T. \tag{2.14}
$$

2.4.1 *A***˜ matrix**

Using Eq. [\(2.7\)](#page-3-1)

$$
\tilde{A} = \begin{pmatrix}\n1 & 0 & \cdots & 0 \\
\frac{(-1)^{1} \Gamma(2+1)}{\Gamma(2) \Gamma(2)} & \frac{(-1)^{2-2} 2^{2} \Gamma(2+2)}{\Gamma(2-1) \Gamma(2+2)} & \cdots & 0 \\
\vdots & \vdots & \ddots & \vdots \\
\frac{(-1)^{m} \Gamma(m+2)}{\Gamma(m+1) \Gamma(2)} & \frac{(-1)^{m-1} 2^{2} \Gamma(m+3)}{\Gamma(m) \Gamma(2+2)} & \cdots & \frac{(-1)^{0} 2^{2m}}{\Gamma(1)}\n\end{pmatrix}_{(m+1)\times(m+1)},
$$

where \tilde{A} is lower triangular non singular matrix, hence \tilde{A}^{-1} exists.

Therefore,

$$
L_m(\eta) = \tilde{A}^{-1} V_F^*(\eta). \tag{2.15}
$$

3 Operational matrix for SVFPs

To solve the SFIDE by operational matrix method, it is necessary to evaluate the following OMs:

3.1 Product operational matrix

The OM for the product is determined in this section.

$$
V_F^*(\eta)V_F^{*T}(\eta)P \simeq \hat{P}V_F^*(\eta),\tag{3.1}
$$

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where \hat{P} is an OM of $(m + 1) \times (m + 1)$ order that is found by applying the orthogonality property of SVFPs with the weight function $w(\eta)$.

$$
\hat{P} = \left\langle V_F^*(\eta) V_F^{*T}(\eta) P, V_F^{*T}(\eta) \right\rangle_{w(\eta)} Q^{-1}.
$$
\n(3.2)

3.2 Integral operational matrix

In terms of OM, the integration of vector $V_F^*(\eta)$ can be described as follows:

$$
\int_0^\eta V_F^*(\zeta)d\zeta \simeq \tilde{P}V_F^*(\eta),\tag{3.3}
$$

where \tilde{P} is an integral OM with a $(m + 1) \times (m + 1)$ dimension that can be found by utilising the orthogonality property of SVFPs with the weight function $w(\eta)$.

Using Eq. [\(3.3\)](#page-5-0), we obtain

$$
\tilde{P} = \left\langle \left(\int_0^\eta V_F^*(\zeta) d\zeta \right), V_F^{*T}(\eta) \right\rangle_{w(\eta)} Q^{-1}.
$$
\n(3.4)

3.3 Stochastic operational matrix

Here, the stochastic OM can be used to approximate the Itô integral of the vector $V_F^*(\eta)$ as follows:

$$
\int_0^\eta V_F^*(\zeta)dB(\zeta) \simeq H_s V_F^*(\eta),\tag{3.5}
$$

where H_s is a stochastic OM with a $(m + 1) \times (m + 1)$ dimension that can be found by utilising the orthogonality property of SVFPs with the weight function $w(\eta)$.

From Eq. (3.5) , we have

$$
H_s = \left\langle \left(\int_0^{\eta} V_F^*(\zeta) dB(\zeta) \right), V_F^{*T}(\eta) \right\rangle_{w(\eta)} Q^{-1}.
$$
 (3.6)

3.3.1 Calculation for *Hs* **matrix**

From Eq. [\(2.13\)](#page-4-1)

$$
V_F^*(\eta) = \tilde{A}L_m(\eta).
$$

Now,

$$
\int_0^{\eta} V_F^*(\zeta) dB(\zeta) = \int_0^{\eta} \tilde{A} L_m(\eta) dB(t) = \tilde{A} \int_0^{\eta} L_m(\zeta) dB(\zeta), \tag{3.7}
$$

$$
\int_0^\eta L_m(\zeta)dB(\zeta) = \left[\int_0^\eta dB(\zeta), \int_0^\eta \zeta dB(\zeta), \dots, \int_0^\eta \zeta^m dB(\zeta)\right]^T. \tag{3.8}
$$

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Thus,

$$
\begin{bmatrix}\n\int_0^{\eta} dB(\zeta) \\
\int_0^{\eta} \zeta dB(\zeta) \\
\int_0^{\eta} \zeta^2 dB(\zeta)\n\end{bmatrix} = \begin{bmatrix}\nB(\eta) \\
\eta B(\eta) - \int_0^{\eta} B(\zeta) d\zeta \\
\eta^2 B(\eta) - 2 \int_0^{\eta} \zeta B(\zeta) d\zeta\n\end{bmatrix} = Y_m(\eta) = [y_j]_{(m+1)\times 1}, (3.9)
$$
\n
$$
\vdots \qquad \vdots \qquad \vdots \qquad \vdots
$$
\n
$$
\eta^m B(\eta) - m \int_0^{\eta} \zeta^{(m-1)} B(\zeta) d\zeta
$$

where,

$$
y_j = \eta^j B(\eta) - j \int_0^{\eta} \zeta^{j-1} B(\zeta) d\zeta
$$
 and $j = 0, 1, 2, ..., m$.

The Simpson's $\frac{1}{3}$ rule is used to evaluate the integrals in Eq. [\(3.9\)](#page-6-0), resulting in

$$
y_j = \left(1 - \frac{j}{6}\right) \eta^j B(\eta) - \frac{j}{3 \times 2^{j-2}} \eta^j B\left(\frac{\eta}{2}\right), \quad j = 0, 1, 2, \dots, m. \tag{3.10}
$$

Now, $B\left(\frac{\eta}{2}\right)$ 2), $B(\eta)$ in Eq. [\(3.10\)](#page-6-1) are approximated by $B(0.25)$ and $B(0.5)$, respectively. Thus,

$$
\int_0^\eta L_m(\zeta) dB(\zeta) \simeq \Lambda L_m(\eta),\tag{3.11}
$$

where

$$
\Lambda = \begin{bmatrix} B(0.5) & 0 & \cdots & 0 \\ 0 & -\frac{2}{3}B(0.25) + \frac{5}{6}B(0.5) & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & -\frac{m}{3 \times 2^{m-2}}B(0.25) + \left(1 - \frac{m}{6}\right)B(0.5) \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & -\frac{m}{3 \times 2^{m-2}}B(0.25) + \left(1 - \frac{m}{6}\right)B(0.5) \end{bmatrix}_{(m+1) \times (m+1)},
$$

and $L_m(\eta) = [1, \eta, \dots, \eta^m]_{(m+1)\times 1}^T$.
Using Eqs. [\(2.15\)](#page-4-2) and [\(3.11\)](#page-6-2), we get

$$
\int_0^\eta V_F^*(\zeta)dB(\zeta) = \tilde{A}\Lambda L_m(\eta) = \tilde{A}\Lambda \tilde{A}^{-1}V_F^*(\eta) = H_s V_F^*(\eta). \tag{3.12}
$$

Hence,

$$
H_s = \tilde{A}\Lambda \tilde{A}^{-1}.
$$
\n(3.13)

3.4 Fractional integral operational matrix

The OM for fractional integrals is discussed in this section.

$$
J_{\eta}^{\alpha} V_F^*(\eta) \simeq F^{\alpha} V_F^*(\eta), \tag{3.14}
$$

where F^{α} is a fractional OM with a $(m + 1) \times (m + 1)$ dimension that can be found by utilising the orthogonality property of SVFPs with the weight function $w(\eta)$.

From Eq. [\(3.14\)](#page-6-3), we obtain

$$
F^{\alpha} = \left\langle J_q^{\alpha} V_F^*(\eta), V_F^{*T}(\eta) \right\rangle_{w(\eta)} Q^{-1}, \tag{3.15}
$$

where J_{η}^{α} is defined in Eq. [\(2.5\)](#page-2-1).

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4 Numerical method

In the operational matrix technique, SVFPs are used to approximate each term in Eq. (1.1) . Let,

$$
{}^{C}D_{\eta}^{\alpha} \simeq V_{F}^{*T}(\eta)A_{1} = A_{1}^{T}V_{F}^{*}(\eta), \qquad (4.1)
$$

$$
z(0) = z_0 \simeq A_2^T V_F^*(\eta), \tag{4.2}
$$

$$
g(\eta) \simeq A_3^T V_F^*(\eta), \tag{4.3}
$$

where A_1 , A_2 , and A_3 are vectors of order $(m+1) \times 1$, which can be defined in the following manner as in Eqs. (2.9) and (2.10) .

By using the RL operator properties

$$
J_{\eta}^{\alpha}({}^C D_{\eta}^{\alpha}) \simeq A_1^T J_{\eta}^{\alpha} V_F^*(\eta), \qquad (4.4)
$$

then, applying Eq. (3.14) into Eq. (4.4) ,

$$
z(\eta) - z_0 \simeq A_1^T F^\alpha V_F^*(\eta), \qquad (4.5)
$$

by using Eq. (4.2) in Eq. (4.5)

$$
z(\eta) \simeq z_m(\eta) = (A_2^T + A_1^T F^{\alpha}) V_F^*(\eta) = \Delta^T V_F^*(\eta) = V_F^{*T}(\eta) \Delta, \qquad (4.6)
$$

where $\Delta = A_2 + (F^{\alpha})^T A_1$ and F^{α} is defined in Eq. [\(3.15\)](#page-6-4).

Now, by substituting Eqs. (2.11) , (4.1) , (4.3) , and (4.6) into Eq. (1.1) , the following is obtained:

$$
A_1^T V_F^*(\eta) = A_3^T V_F^*(\eta) + \lambda_1 \int_0^{\eta} (V_F^{*T}(\eta) \mathbb{K}_1 V_F^*(\zeta) V_F^{*T}(\zeta) \Delta) d\zeta + \lambda_2 \int_0^{\eta} (V_F^{*T}(\eta) \mathbb{K}_2 V_F^*(\zeta) V_F^{*T}(\zeta) \Delta) dB(\zeta) = A_3^T V_F^*(\eta) + \lambda_1 V_F^{*T}(\eta) \mathbb{K}_1 \int_0^{\eta} (V_F^*(\zeta) V_F^{*T}(\zeta) \Delta) d\zeta + \lambda_2 V_F^{*T}(\eta) \mathbb{K}_2 \int_0^{\eta} (V_F^*(\zeta) V_F^{*T}(\zeta) \Delta) dB(\zeta).
$$
 (4.7)

By using Eq. [\(3.1\)](#page-4-4) in Eq. [\(4.7\)](#page-7-5),

$$
A_1^T V_F^*(\eta) = A_3^T V_F^*(\eta) + \lambda_1 V_F^{*T}(\eta) \mathbb{K}_1 \hat{\Delta} \int_0^{\eta} V_F^*(\zeta) d\zeta + \lambda_2 V_F^{*T}(\eta) \mathbb{K}_2 \hat{\Delta} \int_0^{\eta} V_F^*(\zeta) d\zeta,
$$
\n(4.8)

where $\hat{\Delta} = \langle V_F^*(\eta) V_F^{*T}(\eta) \Delta, V_F^{*T}(\eta) \rangle_{w(\eta)} Q^{-1}$. By substituting Eqs. (3.3) and (3.5) in Eq. (4.8) .

$$
A_1^T V_F^*(\eta) = A_3^T V_F^*(\eta) + \lambda_1 V_F^{*T}(\eta) \mathbb{K}_1 \hat{\Delta} \tilde{P} V_F^*(\eta) + \lambda_2 V_F^{*T}(\eta) \mathbb{K}_2 \hat{\Delta} H_s V_F^*(\eta). \tag{4.9}
$$

An algebraic system of equations is created by collocating Eq. [\(4.9\)](#page-7-7) at the Newton cotes nodes provided by $\eta_r = \frac{2r-1}{2(m+1)}$, $r = 1, 2, ..., m+1$. After solving this system of algebraic equations, the coefficient vector A_1 is generated. Now, calculate $\Delta^T = A_2^T + A_1^T F^{\alpha}$. After that the final approximate solution by the SVFPs method is obtained by the equation $z(\eta) \simeq z_m(\eta) = \Delta^T V_F^*(\eta).$

5 Error bound and convergence analysis

5.1 Error bound

Theorem 5[.](#page-23-1)1.1 (Agarwal et al. [2021\)](#page-23-1) *Suppose that* $z(\eta) \in C^{m+1}[0, 1]$ *and* $z_m(\eta)$ *be the approximate solution of z*(η) *defined in Eq.* [\(4.6\)](#page-7-4)*, then*

$$
||z(\eta) - z_m(\eta)|| \le \frac{\hat{E}\lambda^{m+1}}{2(m+1)!} \sqrt{\frac{\pi}{2}},
$$
\n(5.1)

where

$$
\hat{E} = \max_{\eta \in [0,1]} z^{m+1}(\eta) \quad \text{and} \quad \lambda = \max\{\eta_0, 1 - \eta_0\}.
$$

Theorem 5.1.2 *Let* $k(\eta, \zeta)$ *be the sufficiently smooth function in* Ω *such that* $k(\eta, \zeta) \in$ $L^2(\Omega) \cap C^{\infty}(\Omega)$, where $\Omega = ([0, L] \times [0, T])$. Suppose that $k_{m,n}(\eta, \zeta)$ is the best approxi*mation to* $k(\eta, \zeta)$ *out of the linear span* $\Pi_{m,n}(\Omega)$ *. Now assume*

$$
\sup_{(\eta,\zeta)\in\Omega} \left| \frac{\partial^{m+1}k(\eta,\zeta)}{\partial \eta^{m+1}} \right| \le b_1,
$$

\n
$$
\sup_{(\eta,\zeta)\in\Omega} \left| \frac{\partial^{n+1}k(\eta,\zeta)}{\partial \zeta^{n+1}} \right| \le b_2,
$$

\n
$$
\sup_{(\eta,\zeta)\in\Omega} \left| \frac{\partial^{m+n+2}k(\eta,\zeta)}{\partial \eta^{m+1}\partial \zeta^{n+1}} \right| \le b_3,
$$

then there exists $R > 0$ *such that*

$$
||k(\eta,\zeta) - k_{m,n}(\eta,\zeta)||_2 \leq \mathcal{R}\left[\frac{1}{2^m(m+1)!} + \frac{1}{2^n(n+1)!} + \frac{1}{2^{m+n}(m+1)!(n+1)!}\right]\sqrt{\mathcal{C}},\tag{5.2}
$$

where $\mathcal{R} = \max\{b_1, b_2, b_3\}$ *and* $\mathcal{C} = \int_0^L \int_0^T w(\eta)w(\zeta) d\eta d\zeta$.

According the concept of interpolation, which is similar as Saha Ray and Sing[h](#page-24-22) [\(2021\)](#page-24-22)*, we obtain the following desired results.*

Theorem 5.1.2 *Let* $z_m(\eta) = \Delta^T V_F^*(\eta)$ *be the approximate solution and* $z(\eta)$ *be the exact solution of Eq.* [\(1.1\)](#page-1-0)*. Furthermore, suppose that if*

1. $|z(\eta)| \leq M, \forall \eta \in [0, 1],$ 2. $|\kappa_i(\eta, \zeta)| \leq K_i, i = 1, 2, \forall (\eta, \zeta) \in [0, 1] \times [0, 1],$ 3. $\frac{4}{\Gamma(\alpha)^2} [\lambda_1^2 (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) + \lambda_2^2 (2\mathcal{K}_2^2 + 4\mathcal{S}_2^2(m))] < 1.$

Then,

$$
||e(\eta)|| \leq \sqrt{\frac{\frac{4}{(\Gamma(\alpha))^2}\mathcal{P}^2(m) + 8\mathcal{M}^2\frac{(\lambda_1^2\mathcal{S}_1^2(m) + \lambda_2^2\mathcal{S}_2^2(m))}{\Gamma(\alpha))^2}}{1 - \frac{4}{\Gamma(\alpha))^2}[\lambda_1^2(2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) + \lambda_2^2(2\mathcal{K}_2^2 + 4\mathcal{S}_2^2(m))]},
$$

and by the Theorems [5.1.1](#page-8-1) *and* [5.1.2](#page-8-2)

$$
||g(\eta) - g_m(\eta)|| \le \mathcal{P}(m),\tag{5.3}
$$

$$
||\kappa_i(\eta, \zeta) - P_{(m,m)}[\kappa_i](\eta, \zeta)|| \le S_i(m), i = 1, 2.
$$
 (5.4)

where, $g_m(\eta)$ *and* $P_{(m,m)}[k_i](\eta, \zeta)$ *are the approximate polynomials using SVFPs and*

$$
\mathcal{P}(m) = \frac{\hat{E}\lambda^{m+1}}{2(m+1)!} \sqrt{\frac{\pi}{2}} \quad \text{and} \quad \mathcal{S}_i(m) = \mathcal{R} \left[\frac{1}{2^{m-1}(m+1)!} + \frac{1}{2^{2m}(m+1)!^2} \right] \sqrt{\mathcal{C}}.
$$

Proof Let $z_m(\eta)$ be the approximate solution of Eq. [\(1.1\)](#page-1-0).

$$
{}^{c}D_{\eta}^{\alpha}z_{m}(\eta) = g_{m}(\eta) + \lambda_{1} \int_{0}^{\eta} P_{(m,m)}[\kappa_{1}](\eta,\zeta)z_{m}(\zeta)d\zeta
$$

+
$$
\lambda_{2} \int_{0}^{\eta} P_{(m,m)}[\kappa_{2}](\eta,\zeta)z_{m}(\zeta)dB(\zeta), \quad \eta \in [0,1],
$$
(5.5)

Let $|z(\eta) - z_m(\eta)|$ be an error function, then by Eq. [\(1.1\)](#page-1-0) and [\(5.5\)](#page-9-0),

$$
{}^{C}D_{\eta}^{\alpha}(z(\eta)-z_{m}(\eta)) = g(\eta) - g_{m}(\eta) + \lambda_{1} \int_{0}^{\eta} (\kappa_{1}(\eta,\zeta)z(\zeta) - P_{(m,m)}[\kappa_{1}](\eta,\zeta)z_{m}(\zeta))d\zeta
$$

$$
+ \lambda_{2} \int_{0}^{\eta} (\kappa_{2}(\eta,\zeta)z(\zeta) - P_{(m,m)}[\kappa_{2}](\eta,\zeta)z_{m}(\zeta))dB(\zeta). \tag{5.6}
$$

Now, applying the RL operator (J_{η}^{α}) on both sides of the Eq. [\(5.6\)](#page-9-1),

$$
J_{\eta}^{\alpha}({^{C}D_{\eta}^{\alpha}})(z(\eta) - z_{m}(\eta)) = J_{\eta}^{\alpha}(g(\eta) - g_{m}(\eta)) + J_{\eta}^{\alpha} \left(\lambda_{1} \int_{0}^{\eta} (\kappa_{1}(\eta, \zeta)z(\zeta) - P_{(m,m)}[\kappa_{1}](\eta, \zeta)z_{m}(\zeta))d\zeta \right) + J_{\eta}^{\alpha} \left(\lambda_{2} \int_{0}^{\eta} (\kappa_{2}(\eta, \zeta)z(\zeta) - P_{(m,m)}[\kappa_{2}](\eta, \zeta)z_{m}(\zeta))dB\zeta \right),
$$
(5.7)

where J_{η}^{α} is defined in Eq. [\(2.5\)](#page-2-1).

By using the properties of J_{η}^{α} which are given in Sect. [2.2,](#page-2-2) Eq. [\(5.7\)](#page-9-2) can be written as

$$
z(\eta) - z_m(\eta) = \frac{1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \zeta)^{\alpha - 1} (g(\zeta) - g_m(\zeta)) d\zeta
$$

+
$$
\frac{\lambda_1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \zeta)^{\alpha - 1} \left(\int_0^{\gamma} (\kappa_1(\gamma, \zeta) z(\zeta) - P_{(m,m)}[\kappa_1](\eta, \zeta) z_m(\zeta)) d\zeta \right) d\gamma
$$

+
$$
\frac{\lambda_2}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \zeta)^{\alpha - 1} \left(\int_0^{\gamma} (\kappa_2(\gamma, \zeta) z(\zeta) - P_{(m,m)}[\kappa_2](\eta, \zeta) z_m(\zeta)) dB(\zeta) \right) d\gamma.
$$
(5.8)

Using inequality $(c_1 + c_2 + c_3)^2 \le 4(c_1^2 + c_2^2 + c_3^2)$, we obtain $||e(\eta)||^2 = ||z(\eta) - z_m(\eta)||^2$

$$
\leq 4 \left\| \frac{1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \zeta)^{\alpha - 1} (g(\zeta) - g_m(\zeta)) d\zeta \right\|^2
$$

+4
$$
\left\| \frac{\lambda_1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \gamma)^{\alpha - 1} \left(\int_0^{\gamma} (k_1(\gamma, \zeta)z(\zeta) - P_{(m,m)}[k_1](\eta, \zeta)z_m(\zeta)) d\zeta \right) d\gamma \right\|^2
$$

+4
$$
\left\| \frac{\lambda_2}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \gamma)^{\alpha - 1} \left(\int_0^{\gamma} (k_2(\gamma, \zeta)z(\zeta) - P_{(m,m)}[k_2](\eta, \zeta)z_m(\zeta)) d\zeta(\zeta) \right) d\gamma \right\|^2.
$$
(5.9)

Let Eq. (5.9) be written as

$$
||e(\eta)||^2 \le T_2 + T_3 + T_4. \tag{5.10}
$$

Now,

$$
T_2 = 4 \left\| \frac{1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \zeta)^{\alpha - 1} (g(\zeta) - g_m(\zeta)) d\zeta \right\|^2
$$

=
$$
4E \left[\left| \frac{1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \zeta)^{\alpha - 1} (g(\zeta) - g_m(\zeta)) d\zeta \right|^2 \right].
$$
 (5.11)

Since, $0 < \zeta < 1$ and $0 < \alpha < 1$, $0 < \zeta < \eta < 1$. It implies $0 < \eta - \zeta < 1 - \zeta < 1$. Now, using the Cauchy–Schwarz inequality in Eq. [\(5.11\)](#page-10-0), the following is obtained:

$$
\begin{split} T_2 &\leq \frac{4\eta}{(\Gamma(\alpha))^2} E\left[\int_0^\eta |g(\zeta) - g_m(\zeta)|^2 d\zeta\right] \\ &= \frac{4\eta}{(\Gamma(\alpha))^2} ||g(\zeta) - g_m(\zeta)||^2 \\ &\leq \frac{4\eta}{(\Gamma(\alpha))^2} \mathcal{P}^2(m). \end{split} \tag{5.12}
$$

Again,

$$
T_3 = 4 \left\| \frac{\lambda_1}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \gamma)^{\alpha - 1} \left(\int_0^{\gamma} (\kappa_1(\gamma, \zeta)z(\zeta) - P_{(m,m)}[\kappa_1](\eta, \zeta)z_m(\zeta)) d\zeta \right) d\gamma \right\|^2
$$

=
$$
4 \left\| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right\| E\left[\left| \int_0^{\eta} (\eta - \gamma)^{\alpha - 1} \left(\int_0^{\gamma} (\kappa_1(\gamma, \zeta)z(\zeta) - P_{(m,m)}[\kappa_1](\eta, \zeta)z_m(\zeta)) d\zeta \right) d\gamma \right\|^2 \right].
$$

(5.13)

Since, $|\eta - \gamma| \le 1$, then by using the Cauchy–Schwarz inequality

$$
T_3 \le 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E \left[\eta \int_0^{\eta} \left(\int_0^{\gamma} (k_1(\gamma, \zeta)z(\zeta) - P_{(m,m)}[\kappa_1](\eta, \zeta)z_m(\zeta))d\zeta \right)^2 d\gamma \right] \n\le 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E \left[\eta \int_0^{\eta} \left(\gamma \int_0^{\gamma} (k_1(\gamma, \zeta)z(\zeta) - P_{(m,m)}[\kappa_1](\eta, \zeta)z_m(\zeta)) \right)^2 d\zeta \right) d\gamma \right] \n= 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E \left[\eta \int_0^{\eta} (\gamma \int_0^{\gamma} |\kappa_1(\gamma, \zeta)(z(\zeta) - z_m(\zeta)) - \zeta(\zeta)z_m(\zeta)) \right)^2 d\zeta d\gamma \right] \n+ (\kappa_1(\gamma, \zeta) - P_{(m,m)}[\kappa_1](\eta, \zeta)) \times (z_m(\zeta)) - z(\zeta) + z(\zeta)) \Big|^2 d\zeta d\gamma \Big]
$$
\n
$$
\le 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E \left[\eta \int_0^{\eta} (\gamma \int_0^{\gamma} (2|\kappa_1(\gamma, \zeta)(z(\zeta) - z_m(\zeta))|^2 - \zeta(\zeta)) \right)^2 d\zeta d\gamma \Big]
$$
\n
$$
+ 2 |(\kappa_1(\gamma, \zeta) - P_{(m,m)}[\kappa_1](\eta, \zeta)) (z_m(\zeta)) - z(\zeta) + z(\zeta))|^2 d\zeta d\gamma \Big]
$$
\n
$$
\le 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E \left[\eta^2 \int_0^{\eta} (\int_0^{\gamma} (2K_1^2 |\epsilon(\zeta)|^2 + 4S_1^2(m) |\epsilon(\zeta)|^2 + 4M^2 S_1^2(m)) d\zeta \right) d\gamma \right]
$$
\n
$$
= 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E \left[\eta^2 (2K_1^2 + 4S_1
$$

By changing the order of integration, the following is obtained:

$$
\mathcal{T}_3 \le 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E\left[\eta^2 (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) \int_0^\eta |e(\zeta)|^2 \left(\int_\zeta^\eta d\gamma \right) d\zeta + 4\eta^2 \mathcal{M}^2 \mathcal{S}_1^2(m) \int_0^\eta \int_0^\gamma d\zeta d\gamma \right]
$$

$$
\text{Springer } \text{DNN}
$$

$$
\leq 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| E \left[\eta^2 (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) \int_0^{\eta} |e(\zeta)|^2 d\zeta + 2\eta^4 \mathcal{M}^2 \mathcal{S}_1^2(m) \right]
$$

\n
$$
= 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| [\eta^2 (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) ||e(\zeta)||^2 + 2\eta^4 \mathcal{M}^2 \mathcal{S}_1^2(m)]
$$

\n
$$
= 4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| \eta^2 (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) ||e(\zeta)||^2 + 8 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| \eta^4 \mathcal{M}^2 \mathcal{S}_1^2(m).
$$
 (5.15)

Now,

$$
\mathcal{T}_4 = 4 \left\| \frac{\lambda_2}{\Gamma(\alpha)} \int_0^{\eta} (\eta - \gamma)^{\alpha - 1} \left(\int_0^{\gamma} (\kappa_2(\gamma, \zeta) z(\zeta) - P_{(m,m)}[\kappa_2](\eta, \zeta) z_m(\zeta)) dB \zeta \right) d\gamma \right\|^2
$$

=
$$
4 \left\| \frac{\lambda_2^2}{(\Gamma(\alpha))^2} \right\| E \left[\left| \int_0^{\eta} (\eta - \gamma)^{\alpha - 1} \left(\int_0^{\gamma} (\kappa_2(\gamma, \zeta) z(\zeta) - P_{(m,m)}[\kappa_2](\eta, \zeta) z_m(\zeta)) dB(\zeta) \right) d\gamma \right|^2 \right].
$$

(5.16)

Since, $|\eta - \gamma| \le 1$, then by using the Cauchy–Schwarz inequality

$$
\mathcal{T}_4 \le 4 \left| \frac{\lambda_2^2}{(\Gamma(\alpha))^2} \right| E\left[\left| \eta \int_0^{\eta} \left(\int_0^{\gamma} (\kappa_2(\gamma, \zeta) z(\zeta) - P_{(m,m)}[\kappa_2](\eta, \zeta) z_m(\zeta)) dB(\zeta) \right)^2 d\gamma \right| \right]
$$

= 4 \left| \frac{\lambda_2^2}{(\Gamma(\alpha))^2} \right| \eta \int_0^{\eta} E\left[\left(\int_0^{\gamma} (\kappa_2(\gamma, \zeta) z(\zeta) - P_{(m,m)}[\kappa_2](\eta, \zeta) z_m(\zeta)) dB(\zeta) \right)^2 \right] d\gamma. (5.17)

Now, by using the Itô isometry property, following is obtained:

$$
T_{4} \leq 4 \left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| \eta \int_{0}^{\eta} E \left[\int_{0}^{\gamma} (\kappa_{2}(\gamma, \zeta)z(\zeta) - P_{(m,m)}[\kappa_{2}](\eta, \zeta)z_{m}(\zeta))^{2} d\zeta \right] d\gamma
$$

\n
$$
\leq 4 \left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| \eta \int_{0}^{\eta} E \left[\int_{0}^{\gamma} (2|\kappa_{2}(\gamma, \zeta)(z(\zeta) - z_{m}(\zeta))|^{2} + 2 |(\kappa_{2}(\gamma, \zeta) - P_{(m,m)}[\kappa_{2}](\eta, \zeta)) (z_{m}(\zeta)) - z(\zeta) + z(\zeta))|^{2} d\zeta \right] d\gamma
$$

\n
$$
\leq 4 \left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| \eta \int_{0}^{\eta} E \left[\int_{0}^{\gamma} (2K_{2}^{2}|e(\zeta)|^{2} + 4S_{2}^{2}(m)|e(\zeta)|^{2} + 4\mathcal{M}^{2} S_{2}^{2}(m)) d\zeta \right] d\gamma
$$

\n
$$
= 4 \left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| \eta \left[(2K_{2}^{2} + 4S_{2}^{2}(m)) E \left[\int_{0}^{\eta} \int_{0}^{\gamma} |e(\zeta)|^{2} d\zeta d\gamma \right] + 4\mathcal{M}^{2} S_{2}^{2}(m) \frac{\eta^{2}}{2} \right].
$$

\n(5.18)

By changing the order of integration

$$
\mathcal{T}_4 \le 4 \left| \frac{\lambda_2^2}{(\Gamma(\alpha))^2} \right| \eta(2\mathcal{K}_2^2 + 4\mathcal{S}_2^2(m)) ||e(\zeta)||^2 + 8 \left| \frac{\lambda_2^2}{(\Gamma(\alpha))^2} \right| \eta^3 \mathcal{M}^2 \mathcal{S}_2^2(m). \tag{5.19}
$$

Now, by substituting Eqs. [\(5.12\)](#page-10-1), [\(5.15\)](#page-11-0), [\(5.19\)](#page-11-1) into Eq. [\(5.10\)](#page-9-4)

$$
||e(\eta)||^2 \le \frac{4}{(\Gamma(\alpha))^2} \mathcal{P}^2(m)
$$

+4 $\left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) ||e(\zeta)||^2 + 8 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| \mathcal{M}^2 \mathcal{S}_1^2(m)$

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$$
+4\left|\frac{\lambda_2^2}{(\Gamma(\alpha))^2}\right|(2\mathcal{K}_2^2+4\mathcal{S}_2^2(m))||e(\zeta)||^2+8\left|\frac{\lambda_2^2}{(\Gamma(\alpha))^2}\right|\mathcal{M}^2\mathcal{S}_2^2(m). \quad (5.20)
$$

Then,

$$
||e(\eta)|| \leq \sqrt{\frac{\frac{4}{(\Gamma(\alpha))^2} \mathcal{P}^2(m) + 8\mathcal{M}^2 \frac{(\lambda_1^2 \mathcal{S}_1^2(m) + \lambda_2^2 \mathcal{S}_2^2(m))}{\Gamma(\alpha))^2}}{1 - \frac{4}{\Gamma(\alpha))^2} [\lambda_1^2 (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) + \lambda_2^2 (2\mathcal{K}_2^2 + 4\mathcal{S}_2^2(m))]}.
$$
(5.21)

5.2 Convergence analysis

Theorem 5.2.1 *Let* $z(\eta)$ *and* $z_m(\eta)$ *be the exact and approximate solutions of Eq.* [\(1.1\)](#page-1-0) *respectively. And*

1. $|z(\eta)| \leq M, \forall \eta \in [0, 1],$ 2. $|\kappa_i(\eta, \zeta)| \leq \mathcal{K}_i, i = 1, 2, \forall (\eta, \zeta) \in [0, 1] \times [0, 1],$ 3. $\frac{4}{\Gamma(\alpha))^2} [\lambda_1^2 (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) + \lambda_2^2 (2\mathcal{K}_2^2 + 4\mathcal{S}_2^2(m))] < 1.$

Then $z_m(\eta) \to z(\eta)$ *as* $m \to \infty$ *in* L^2 *.*

Proof Consider the SFIDE as follows:

$$
{}^{C}D^{\alpha}z(\eta) = g(\eta) + \lambda_1 \int_0^{\eta} \kappa_1(\eta, \zeta) z(\zeta) d\zeta + \lambda_2 \int_0^{\eta} \kappa_2(\eta, \zeta) z(\zeta) dB(\zeta), \quad \eta \in [0, 1].
$$
\n(5.22)

Using the same explanation as that used to prove the previous theorem, we can get to the following:

By using Eqs. [\(5.12\)](#page-10-1), [\(5.15\)](#page-11-0) and[\(5.18\)](#page-11-2)

$$
||e(\eta)||^{2} \leq \frac{4\eta}{(\Gamma(\alpha))^{2}} \mathcal{P}^{2}(m)
$$

+4 $\left| \frac{\lambda_{1}^{2}}{(\Gamma(\alpha))^{2}} \right| E\left[\eta^{2}(2\mathcal{K}_{1}^{2} + 4\mathcal{S}_{1}^{2}(m)) \int_{0}^{\eta} |e(\zeta)|^{2} d\zeta + 2\eta^{4} \mathcal{M}^{2} \mathcal{S}_{1}^{2}(m) \right]$
+4 $\left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| \eta \left[(2\mathcal{K}_{2}^{2} + 4\mathcal{S}_{2}^{2}(m)) E\left[\int_{0}^{\eta} \int_{0}^{\gamma} |e(\zeta)|^{2} d\zeta d\gamma \right] + 4\mathcal{M}^{2} \mathcal{S}_{2}^{2}(m) \frac{\eta^{2}}{2} \right].$ (5.23)

Since $\eta \leq 1$, then

$$
||e(\eta)||^{2} \leq \frac{4}{(\Gamma(\alpha))^{2}} \mathcal{P}^{2}(m)
$$

+4
$$
\left| \frac{\lambda_{1}^{2}}{(\Gamma(\alpha))^{2}} \right| (2\mathcal{K}_{1}^{2} + 4\mathcal{S}_{1}^{2}(m)) \int_{0}^{\eta} ||e(\zeta)||^{2} d\zeta + 8 \left| \frac{\lambda_{1}^{2}}{(\Gamma(\alpha))^{2}} \right| \mathcal{M}^{2} \mathcal{S}_{1}^{2}(m)
$$

+4
$$
\left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| (2\mathcal{K}_{2}^{2} + 4\mathcal{S}_{2}^{2}(m)) \int_{0}^{\eta} ||e(\zeta)||^{2} d\zeta + 8 \left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| \mathcal{M}^{2} \mathcal{S}_{2}^{2}(m).
$$
(5.24)

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Now,

$$
||e(\eta)||^{2} \leq \frac{4}{(\Gamma(\alpha))^{2}} \mathcal{P}^{2}(m) + 8\mathcal{M}^{2} \left(\left| \frac{\lambda_{1}^{2} \mathcal{S}_{1}^{2}(m)}{(\Gamma(\alpha))^{2}} \right| + \left| \frac{\lambda_{2}^{2} \mathcal{S}_{2}^{2}(m)}{(\Gamma(\alpha))^{2}} \right| \right) + \left(4 \left| \frac{\lambda_{1}^{2}}{(\Gamma(\alpha))^{2}} \right| (2\mathcal{K}_{1}^{2} + 4\mathcal{S}_{1}^{2}(m)) + 4 \left| \frac{\lambda_{2}^{2}}{(\Gamma(\alpha))^{2}} \right| (2\mathcal{K}_{2}^{2} + 4\mathcal{S}_{2}^{2}(m)) \right) \int_{0}^{\eta} ||e(\zeta)||^{2} d\zeta.
$$
\n(5.25)

Let.

$$
\delta(m) = \frac{4}{(\Gamma(\alpha))^2} \mathcal{P}^2(m) + 8\mathcal{M}^2 \left(\left| \frac{\lambda_1^2 \mathcal{S}_1^2(m)}{(\Gamma(\alpha))^2} \right| + \left| \frac{\lambda_2^2 \mathcal{S}_2^2(m)}{(\Gamma(\alpha))^2} \right| \right),
$$

$$
L_1 = \left(4 \left| \frac{\lambda_1^2}{(\Gamma(\alpha))^2} \right| (2\mathcal{K}_1^2 + 4\mathcal{S}_1^2(m)) + 4 \left| \frac{\lambda_2^2}{(\Gamma(\alpha))^2} \right| (2\mathcal{K}_2^2 + 4\mathcal{S}_2^2(m)) \right).
$$

Therefore,

$$
||e(\eta)||_2^2 \le \delta(m) + L_1 \int_0^{\zeta} ||e(\zeta)||^2 d\zeta.
$$
 (5.26)

Applying Grönwall inequality, we obtain

$$
||e(\eta)||_2^2 \le \delta(m)(1 + L_1 \int_0^{\zeta} e^{L_1(\eta - \zeta)} d\zeta).
$$
 (5.27)

It implies

$$
||e(\eta)||_2^2 \to 0 \text{as} m \to \infty \text{in} L^2.
$$

So, $z_m(\eta)$ converges to $z(\eta)$ as $m \to \infty$ in L^2 .

6 Applications of the proposed method

Three examples are solved in this section using the proposed numerical approach that was described in the previous section.

Example 1 Consider the following fractional order stochastic integro-differential equation:

$$
{}^{C}D_{\eta}^{\alpha}z(\eta) = \frac{-\eta^{5}e^{\eta}}{5} + \frac{6\eta^{2.25}}{\Gamma(3.25)} + \int_{0}^{\eta} e^{\eta}\zeta z(\zeta)d\zeta
$$

$$
+ \lambda_{2} \int_{0}^{\eta} e^{\eta}\zeta z(\zeta)dB\zeta, \quad \zeta, \eta \in [0, 1], \tag{6.1}
$$

with the initial condition $z(0) = 0$. The exact solution of Eq. [\(6.1\)](#page-13-1) is not available. If $\alpha = 0.75$ and $\lambda_2 = 0$, the exact solution is $z(\eta) = \eta^3$ and the approximate solution is obtained by the proposed SVFPs method. Table [1](#page-14-0) represents the absolute error comparison between two methods based on the orthonormal Chelyshkov polynomials (OCPs) and SVFPs. For the numerical solution of Eq. [\(6.1\)](#page-13-1) for various values of α with $m = 4$ and $m = 6$, the proposed operational matrix collocation approach is employed. Newton cotes nodes have

been selected from the collocation points. Tables [2](#page-15-0) and [3](#page-16-0) provide the comparison between numerical solutions obtained by the operational matrix method based on OCPs and SVFPs for the above problem for different values of *m*. The plot of SVFPs solutions for different values of α with $m = 4$ and $m = 6$ are shown in Figs. [1](#page-19-1) and [2,](#page-19-2) respectively.

Example 2 Consider the following fractional order stochastic integro-differential equation:

$$
{}^{C}D_{\eta}^{\alpha}z(\eta) = \frac{7}{12}\eta^{4} - \frac{5}{6}\eta^{3} + \frac{2\eta^{2-\alpha}}{\Gamma(3-\alpha)} + \frac{\eta^{1-\alpha}}{\Gamma(2-\alpha)} + \int_{0}^{\eta} (\eta + \zeta)z(\zeta)d\zeta
$$

$$
+ \lambda \int_{0}^{\eta} \zeta z(\zeta)dB\zeta, \quad \zeta, \eta \in [0, 1], \tag{6.2}
$$

with the initial condition $z(0) = 0$. The exact solution of Eq. [\(6.2\)](#page-14-1) is not known. For the numerical solution of Eq. [\(6.2\)](#page-14-1) for different values of α , the proposed operational matrix collocation method is utilised. Newton cotes nodes have been selected from the collocation points. Tables [4](#page-17-0) and [5](#page-18-0) provide the numerical solutions comparison obtained by the operational matrix method based on OCPs and SVFPs for the above problem for different values of *m*. The plot of SVFPs solutions, for different values of α with $m = 4$ and $m = 6$ are shown in Figs. [3](#page-20-0) and [4,](#page-20-1) respectively.

Example 3 Consider the following fractional order stochastic integro-differential equation:

$$
{}^{C}D_{\eta}^{\alpha}z(\eta) = -\frac{\eta^{3}}{3} + \frac{\Gamma(2)\eta^{1-\alpha}}{\Gamma(2-\alpha)} + \int_{0}^{\eta}\zeta z(\zeta)d\zeta + \lambda \int_{0}^{\eta}z(\zeta)dB\zeta, \quad \zeta, \eta \in [0,1], \tag{6.3}
$$

with the initial condition $z(0) = 0$. The exact solution of Eq. [\(6.3\)](#page-14-2) is unknown. The proposed operational matrix collocation technique is used to solve Eq. [\(6.3\)](#page-14-2) numerically for different values of α with $m = 4$ and $m = 6$. From the collocation points, Newton cotes nodes have been chosen. With respect to the above-mentioned problem, Tables [6](#page-21-0) and [7](#page-22-0) compare the numerical solutions derived using the operational matrix technique based on OCPs and SVFPs for different values of *m*. The plot of SVFPs solutions for different values of α with $m = 4$ and $m = 6$ are shown in Figs. [5](#page-23-2) and [6,](#page-23-3) respectively.

Table 4 Comparison of approximate solutions by using the OCOM and SVFOM methods for *m* $=$ 4 in Example 2

Fig. 1 The approximate solution graph by the SVFOM method for $m = 4$ (Example [1\)](#page-13-2)

Fig. 2 The approximate solution graph by the SVFOM method for $m = 6$ (Example [1\)](#page-13-2)

7 Conclusion

The main objective of this work is to apply the operational matrix method to solve SFIDE. In this study, a novel stochastic operational matrix has been generated for the first time using shifted Vieta–Fibonacci polynomials. The SFIDE has been transformed into a system of algebraic equations. With the use of these operational matrices, the resultant algebraic system of equations is numerically solved by applying the collocation technique. The error bound and the convergence analysis of the proposed numerical technique have also been described. The precision and effectiveness of the suggested numerical technique are demonstrated using three

Fig. 3 The approximate solution graph by the SVFOM method for $m = 4$ (Example [2\)](#page-14-3)

Fig. 4 The approximate solution graph by the SVFOM method for $m = 6$ (Example [2\)](#page-14-3)

different examples. The numerical experiments reveal that the proposed numerical scheme based on SVFPs and OCPs based numerical method have the best agreement of results. As a consequence, it is clear from the numerical experiment results that the proposed numerical approach is extremely effective, accurate, and reliable. In future, we have a plan to work on fractional stochastic integro-differential equations with the ABC fractional derivative.

1 1.02111 1.0212 1.0212 1.0212 1.0212 1.0212 1.0212 1.0256 1.0256 1.0256 1.0256 1.0256

Table 6 Comparison of approximate solutions by using the OCOM and SVFOM methods for

m

 $=$ 4 in Example [3](#page-14-4)

Table 7 Comparison of approximate solutions by using the OCOM and SVFOM methods for *m* $= 6$ in Example 3

Fig. 5 The approximate solution graph by the SVFOM method for $m = 4$ (Example [3\)](#page-14-4)

Fig. 6 The approximate solution graph by the SVFOM method for $m = 6$ (Example [3\)](#page-14-4)

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Data Availability This article includes all the data that were generated or analyzed during this research.

References

Agarwal P, El-Sayed AA, Tariboon J (2021) Vieta–Fibonacci operational matrices for spectral solutions of variable-order fractional integro-differential equations. J Comput Appl Math 382:113063

Ahmad B, Nieto JJ (2011) Riemann-Liouville fractional integro-differential equations with fractional nonlocal integral boundary conditions. Bound Value Probl 1:1–9

- Arikoglu A, Ozkol I (2009) Solution of fractional integro-differential equations by using fractional differential transform method. Chaos Solitons Fractals 40(2):521–529
- Atman KG, Şirin H (2020) Nonlocal phenomena in quantum mechanics with fractional calculus. Rep Math Phys 86(2):263–270
- Behera S, Saha Ray S (2022) An efficient numerical method based on Euler wavelets for solving fractional order pantograph Volterra delay-integro-differential equations. J. Comput. Appl. Math. 406:113825
- El-Misiery AEM, Ahmed E (2006) On a fractional model for earthquakes. Appl Math Comput 178(2):207–211 Kamrani M (2016) Convergence of Galerkin method for the solution of stochastic fractional integro differential
- equations. Optik 127(20):10049–10057 Lodge AS, McLeod JB, Nohel JA (1978) A nonlinear singularly perturbed Volterra integrodifferential equation occurring in polymer rheology. Proc R Soc Edinb Sect A Math 80(1–2):99–137
- Machado JT, Lopes AM (2015) Analysis of natural and artificial phenomena using signal processing and fractional calculus. Fract Calc Appl Anal 18:459–478
- Meral FC, Royston TJ, Magin R (2010) Fractional calculus in viscoelasticity: an experimental study. Commun Nonlinear Sci Numer Simul 15(4):939–945
- Mirzaee F, Alipour S (2020) Cubic B-spline approximation for linear stochastic integro-differential equation of fractional order. J Comput Appl Math 366:112440
- Mirzaee F, Samadyar N (2017) Application of orthonormal Bernstein polynomials to construct a efficient scheme for solving fractional stochastic integro-differential equation. Optik 132:262–273
- Mirzaee F, Samadyar N (2019) On the numerical solution of fractional stochastic integro-differential equations via meshless discrete collocation method based on radial basis functions. Eng Anal Bound Elem 100:246– 255
- Mirzaee F, Alipour S, Samadyar N (2019) Numerical solution based on hybrid of block-pulse and parabolic functions for solving a system of nonlinear stochastic Itô-Volterra integral equations of fractional order. J Comput Appl Math 349:157–171
- Momani S, Odibat Z (2006) Analytical approach to linear fractional partial differential equations arising in fluid mechanics. Phys Lett A 355(4–5):271–279
- Øksendal B (2003) Stochastic differential equations. Springer, Berlin, pp 65–84
- Sadri K, Hosseini K, Baleanu D, Salahshour S, Park C (2022) Designing a matrix collocation method for fractional delay integro-differential equations with weakly singular kernels based on Vieta–Fibonacci polynomials. Fractal Fract 6(1):2
- Saha Ray S (2015) Fractional calculus with applications for nuclear reactor dynamics. CRC Press, Boca Raton
- Saha Ray S, Patra A (2013) Numerical solution of fractional stochastic neutron point kinetic equation for nuclear reactor dynamics. Ann Nucl Energy 54:154–161
- Saha Ray S, Singh P (2021) Numerical solution of stochastic Itô-Volterra integral equation by using Shifted Jacobi operational matrix method. Appl Math Comput 410:126440
- Sahu PK, Saha Ray S (2018) Comparison on wavelets techniques for solving fractional optimal control problems. J Vib Control 24(6):1185–1201
- Scalas E, Gorenflo R, Mainardi F (2000) Fractional calculus and continuous-time finance. Physica A Stat Mech Appl 284(1–4):376–384
- Sun H, Zhang Y, Baleanu D, Chen W, Chen Y (2018) A new collection of real world applications of fractional calculus in science and engineering. Commun Nonlinear Sci Numer Simul 64:213–231
- Taheri Z, Javadi S, Babolian E (2017) Numerical solution of stochastic fractional integro-differential equation by the spectral collocation method. J Comput Appl Math 321:336–347
- Yzbaı S, Sezer M, Kemancı B (2013) Numerical solutions of integro-differential equations and application of a population model with an improved Legendre method. Appl Math Model 37(4):2086–2101

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