

Source Identification for Parabolic Equations from Integral Observations by the Finite Difference Splitting Method

Nguyen Thi Ngoc Oanh¹

Received: 19 April 2024 / Accepted: 2 May 2024 / Published online: 11 June 2024 © Institute of Mathematics, Vietnam Academy of Science and Technology (VAST) and Springer Nature Singapore Pte Ltd. 2024

Abstract

We study the problem of reconstructing an unknown source term in parabolic equations from integral observations. It is reformulated into a variational problem in combination with Tikhonov regularization and then a formula for the gradient of the objective functional to be minimized is computed via a solution of an adjoint problem. The variational problem is discretized by the splitting method based on finite difference schemes and solved by the conjugate gradient method. A numerical scheme for numerically estimating singular values of the solution operator in the inverse problem is suggested. Some numerical examples are presented to show the efficiency of the method.

Keywords Source identification \cdot Integral observations \cdot Least squares method \cdot Tikhonov regularization \cdot Conjugate gradient method

Mathematics Subject Classification (2010) 35R30 · 65J20 · 65M32 · 65N21

1 Introduction

The problem of determining a source term in parabolic equations from some observations plays an important role in practice [4, 9, 10]. Because of its importance, many researchers devoted their attention to it [1–3, 5, 7, 8, 12, 14, 17, 18, 22, 24]. For more details, let Ω be a bounded domain in \mathbb{R}^n with boundary $\partial \Omega$. Denote the cylinder $Q := \Omega \times (0, T]$, where T > 0 and $S := \partial \Omega \times (0, T]$. Let

$$a_{ij}, i, j \in \{1, 2, \dots, n\}, b \in L^{\infty}(Q),$$
(1)

$$a_{ij} = a_{ji}, \quad i, j \in \{1, 2, \dots, n\},$$
(2)

$$\lambda \|\xi\|_{\mathbb{R}^n}^2 \le \sum_{i,j=1}^n a_{ij}(x,t)\xi_i\xi_j \le \Lambda \|\xi\|_{\mathbb{R}^n}^2, \quad \forall \xi \in \mathbb{R}^n,$$
(3)

$$0 \le b(x,t) \le \mu_1 \quad \text{a.e. in } Q,\tag{4}$$

⊠ Nguyen Thi Ngoc Oanh oanhntn@tnus.edu.vn

¹ Thai Nguyen University of Sciences, Tan Thinh Ward, Thai Nguyen 250000, Vietnam



$$v \in L^2(\Omega), \quad F \in L^2(Q), \tag{5}$$

 λ and Λ be positive constants and $\mu_1 \ge 0$. (6)

Consider the initial boundary value problem

$$\begin{cases} \frac{\partial u}{\partial t} - \sum_{i,j=1}^{n} \frac{\partial}{\partial x_{i}} \left(a_{ij}(x,t) \frac{\partial u}{\partial x_{j}} \right) + b(x,t)u = F(x,t), \quad (x,t) \in Q, \\ u(x,t) = 0, \quad (x,t) \in S, \\ u(x,0) = v(x), \quad x \in \Omega. \end{cases}$$
(7)

Let F have either one of the following forms

$$F(x,t) = f(x,t)\varphi(x,t) + g(x,t),$$
(8)

$$F(x,t) = f(x)\varphi(x,t) + g(x,t),$$
(9)

$$F(x,t) = f(t)\varphi(x,t) + g(x,t)$$
(10)

with $\varphi(x, t) \in L^2(Q)$ and $g(x, t) \in L^2(Q)$ being given.

We consider the problem of determining f from N integral observations of the solution u

$$l_i u = \int_{\Omega} \omega_i(x) u(x, t) dx = z_i(t), \quad t \in (0, T), \ i = 1, \dots, N$$
(11)

with $\omega_i(x) \in L^{\infty}(\Omega)$, nonnegative almost everywhere and $\int_{\Omega} \omega_i(x) dx > 0$, being weighted functions. Suppose that z_i , i = 1, 2, ..., N are approximately given by z_i^{δ} satisfying

$$\|z_i - z_i^{\delta}\|_{L^2(0,T)} \le \delta.$$

These inverse problems may have many solutions, especially in the case f depends on x and t. Indeed, suppose that the coefficients of (7) are sufficiently smooth. If $\varphi(x, t) \neq 0$ and u(x, t) is given for all $(x, t) \in Q = \Omega \times (0, T)$, the inverse problem has a unique solution

$$f(x,t) = \frac{\frac{\partial u}{\partial t} - \sum_{i,j=1}^{n} \frac{\partial}{\partial x_i} \left(a_{ij}(x,t) \frac{\partial u}{\partial x_j} \right) + b(x,t)u - g(x,t)}{\varphi(x,t)}.$$

We show that if there is a *u* satisfying (11), then there are infinitely many $u \in C^{\infty}(Q)$, $u|_{S} = 0$ satisfying (11). Indeed, for $v(x) \in C^{\infty}(\Omega)$ satisfying (11), consider the following equation

$$\langle \omega_i, v \rangle_{L^2(\Omega)} = \int_{\Omega} \omega_i(x) v(x) dx = 0, \ i = 1, 2, \dots, N.$$
(12)

Denote $\mathcal{P} = \operatorname{span}\{\omega_1, \omega_2, \dots, \omega_N\}$. Then \mathcal{P} is a subspace of $L^2(\Omega)$ and dim $\mathcal{P} \leq N$. So $\mathcal{Q} = \mathcal{P}^{\perp}$ is an infinite-dimensional space. Moreover, we have presentation $v = v^1 + v^2$, where $v^1 \in \mathcal{P}, v^2 \in \mathcal{Q}$ and $\int_{\Omega} v^1(x)v^2(x)dx = 0$. It concludes that there are infinite functions $v \in C^{\infty}(\Omega)$ satisfying equation (12). So, there are infinitely many functions $u(\cdot, t) \in C^{\infty}(\Omega)$ satisfying equation

$$\int_{\Omega} \omega_i(x)u(x,t)dx = 0, \ i = 1, 2, \dots, N.$$

Or, there are infinite functions $u(\cdot, t) \in C^{\infty}(\Omega)$ satisfying (11). We conclude that the inverse problem of finding f from (11) has infinite solutions. Therefore, we have to introduce a notion to its solution.



2 Variational Problem

To introduce the concept of weak solution, we use the standard Sobolev spaces $H^1(\Omega)$, $H^{1,0}_0(\Omega)$, $H^{1,0}(Q)$ and $H^{1,1}(Q)$ [11, 21, 23]. Further, for a Banach space *B*, we define

$$L^{2}(0, T; B) = \{u : u(t) \in B \text{ a.e. } t \in (0, T) \text{ and } \|u\|_{L^{2}(0, T; B)} < \infty\},\$$

with the norm

$$\|u\|_{L^2(0,T;B)}^2 = \int_0^T \|u(t)\|_B^2 dt.$$

In the sequel, we shall use the space W(0, T) defined as

$$W(0,T) = \{ u : u \in L^2(0,T; H_0^1(\Omega)), u_t \in L^2(0,T; (H_0^1(\Omega))') \},\$$

equipped with the norm

$$\|u\|_{W(0,T)}^{2} = \|u\|_{L^{2}(0,T;H_{0}^{1}(\Omega))}^{2} + \|u_{t}\|_{L^{2}(0,T;(H_{0}^{1}(\Omega))')}^{2}$$

We note here that $(H_0^1(\Omega))' = H^{-1}(\Omega)$.

The solution of the problem (7) is understood in the weak sense as follows: A weak solution in W(0, T) of the problem (7) is a function $u(x, t) \in W(0, T)$ satisfying the identity

$$\int_{0}^{T} (u_{t}, \eta)_{H^{-1}(\Omega), H_{0}^{1}(\Omega)} dt + \int_{0}^{T} \int_{\Omega} \left(\sum_{i, j=1}^{n} a_{ij}(x, t) \frac{\partial u}{\partial x_{j}} \frac{\partial \eta}{\partial x_{i}} + b(x, t) u \eta \right) dx dt$$

$$= \int_{0}^{T} \int_{\Omega} (f \varphi \eta + g \eta) dx dt, \quad \forall \eta \in L^{2}(0, T; H_{0}^{1}(\Omega))$$

$$(13)$$

and

$$u(x,0) = v(x), \quad x \in \Omega.$$
(14)

Based on the standard hypotheses (1), (2), (3), (4), (5) and 6, the existence and uniqueness of a solution, as well as an a priori estimate to the problem (7), can be established. More precisely, following [23, Chapter IV] and [21, pp. 141–152] there exists a unique solution in W(0, T) of the problem (7). Furthermore, there is a positive constant c_d independent of a_i, b, f, φ, g and v such that

$$\|u\|_{W(0,T)} \le c_d \left(\|f\varphi\|_{L^2(Q)} + \|g\|_{L^2(Q)} + \|v\|_{L^2(\Omega)} \right)$$

We denote the solution u(x, t) of the problem (7) by u(x, t; f) or u(f) to emphasize its dependence on f. To identify f from (11), we minimize the misfit functional

$$J_0(f) = \frac{1}{2} \sum_{i=1}^{N} \|l_i u(f) - z_i\|_{L^2(0,T)}^2$$
(15)

🖄 Springer

with respect to f. However, this minimization problem is unstable and there might be many minimizers to it. Therefore, we minimize the Tikhonov functional instead of (15). In fact, we minimize

$$J_{\gamma}(f) = \frac{1}{2} \sum_{i=1}^{N} \|l_{i}u(f) - z_{i}\|_{L^{2}(0,T)}^{2} + \frac{\gamma}{2} \|f - f^{*}\|_{L^{2}(Q)}^{2}, f^{*} \in L_{2}(Q)$$
(16)

for the case F has form (8).

...

$$J_{\gamma}(f) = \frac{1}{2} \sum_{i=1}^{N} \|l_{i}u(f) - z_{i}\|_{L^{2}(0,T)}^{2} + \frac{\gamma}{2} \|f - f^{*}\|_{L^{2}(\Omega)}^{2}, f^{*} \in L_{2}(\Omega)$$
(17)

for the case F has form (9).

$$J_{\gamma}(f) = \frac{1}{2} \sum_{i=1}^{N} \|l_{i}u(f) - z_{i}\|_{L^{2}(0,T)}^{2} + \frac{\gamma}{2} \|f - f^{*}\|_{L^{2}(0,T)}^{2}, f^{*} \in L_{2}(0,T)$$
(18)

for the case *F* has form (10). Here, $\gamma > 0$ is the Tikhonov regularization parameter, f^* is an a priori estimation of *f*. By the standard method, we can prove that J_{γ} is Fréchet differentiable and derive a formula for its gradient. As $l_i u(f)$ is affine, the functional J_{α} is strictly convex. Hence, it attains a unique minimizer which we call f^* - *least square solution* to the inverse problems (7) and (11). As the inverse problem may have many solutions, we will see that the choice of f^* is crucial for selecting which one among these solutions to the inverse problem.

Indeed, introducing the adjoint problem

$$\begin{cases} -\frac{\partial p}{\partial t} - \sum_{i,j=1}^{n} \frac{\partial}{\partial x_j} \left(a_{ij}(x,t) \frac{\partial p}{\partial x_i} \right) + b(x,t)p = \sum_{i=1}^{N} \omega_i(x) \left(l_i u(t) - z_i(t) \right), \quad (x,t) \in Q, \\ p(x,t) = 0, \quad (x,t) \in S, \\ p(x,T) = 0, \quad x \in \Omega, \end{cases}$$
(19)

we can prove the following results [17, 19].

Theorem 1 The functional J_{γ} (8) is Fréchet differentiable and its gradient ∇J_{γ} at f has the form

$$\nabla J_{\gamma}(f) = \varphi(x,t)p(x,t) + \gamma(f(x,t) - f^*(x,t)),$$

where p(x, t) is the solution to the adjoint problem (19).

Remark 1 When J_{γ} has the form in (17) or (18), we have

i)

$$\nabla J_{\gamma}(f) = \int_0^T \varphi(x, t) p(x, t) dt + \gamma(f(x) - f^*(x)) \text{ for the functional (17).}$$

ii)

$$\nabla J_{\gamma}(f) = \int_{\Omega} \varphi(x, t) p(x, t) dx + \gamma(f(t) - f^*(t)) \text{ for the functional (18).}$$

2.1 Conjugate Gradient Method

To find the minimizer of (16), we use the conjugate gradient method (CG). It proceeds as follows: Assume that at the k-th iteration we have f^k . Then the next iteration is

$$f^{k+1} = f^k + \alpha^k d^k,$$



with

$$d^{k} = \begin{cases} -\nabla J_{\gamma}(f^{k}) & \text{if } k = 0, \\ -\nabla J_{\gamma}(f^{k}) + \beta^{k} d^{k-1} & \text{if } k > 0, \end{cases}$$
$$\beta^{k} = \frac{\|\nabla J_{\gamma}(f^{k})\|_{L^{2}(Q)}^{2}}{\|\nabla J_{\gamma}(f^{k-1})\|_{L^{2}(Q)}^{2}},$$

and

$$\alpha^k = \operatorname{argmin}_{\alpha \ge 0} J_{\gamma}(f^k + \alpha d^k).$$

To evaluate α^k we denote by $\bar{u}(v, g)$ the solution to the problem

$$\begin{cases} \frac{\partial u}{\partial t} - \sum_{i,j=1}^{n} \frac{\partial}{\partial x_i} \left(a_{ij}(x,t) \frac{\partial u}{\partial x_j} \right) + b(x,t)u = g(x,t), & (x,t) \in Q, \\ u(x,t) = 0, & (x,t) \in S, \\ u(x,0) = v(x), & x \in \Omega \end{cases}$$

with $\tilde{u}[f]$ being the solution to the linear problem

$$\begin{cases} \frac{\partial u}{\partial t} - \sum_{i,j=1}^{n} \frac{\partial}{\partial x_{i}} \left(a_{ij}(x,t) \frac{\partial u}{\partial x_{j}} \right) + b(x,t)u = f(x,t)\varphi(x,t), & (x,t) \in Q, \\ u(x,t) = 0, & (x,t) \in S, \\ u(x,0) = 0, & x \in \Omega. \end{cases}$$

In this case, the observation operators have the form

$$l_{i}u(f) = l_{i}\tilde{u}[f] + l_{i}\bar{u}(v,g) := A_{i}f + l_{i}\bar{u}(v,g), \ i = 1, \dots, N$$
(20)

with A_i being bounded linear operators from $L^2(Q)$ into $L_2(0, T)$. We have

$$J_{\gamma}(f^{k} + \alpha d^{k}) = \sum_{i=1}^{N} \frac{1}{2} \|l_{i}u(f^{k} + \alpha d^{k}) - z_{i}\|_{L^{2}(0,T)}^{2} + \frac{\gamma}{2} \|f^{k} + \alpha d^{k} - f^{*}\|_{L^{2}(Q)}^{2}$$

$$= \sum_{i=1}^{N} \frac{1}{2} \|\alpha A_{i}d^{k} + A_{i}f^{k} + l_{i}\bar{u}(v,g) - z_{i}\|_{L^{2}(0,T)}^{2} + \frac{\gamma}{2} \|\alpha d^{k} + f^{k} - f^{*}\|_{L^{2}(Q)}^{2}$$

$$= \sum_{i=1}^{N} \frac{1}{2} \|\alpha A_{i}d^{k} + l_{i}u(f^{k}) - z_{i}\|_{L^{2}(0,T)}^{2} + \frac{\gamma}{2} \|\alpha d^{k} + f^{k} - f^{*}\|_{L^{2}(Q)}^{2}.$$

Differentiating $J_{\gamma}(f^k + \alpha d^k)$ with respect to α and putting $\frac{\partial J_{\gamma}(f^k + \alpha d^k)}{\partial \alpha} = 0$, after some elementary calculations, we obtain

$$\alpha^{k} = -\frac{\langle d^{k}, \nabla J_{\gamma}(f^{k}) \rangle_{L^{2}(Q)}}{\sum_{i=1}^{N} \|A_{i}d^{k}\|_{L^{2}(0,T)}^{2} + \gamma \|d^{k}\|_{L^{2}(Q)}^{2}}.$$

Since $d^k = -\nabla_{\gamma}(f^k) + \beta^k d^{k-1}$, $r^k = -\nabla J_{\gamma}(f^k)$ and $\langle r^k, d^{k-1} \rangle_{L^2(Q)} = 0$, we have

$$\alpha^{k} = \frac{\|r^{k}\|_{L^{2}(Q)}^{2}}{\sum_{i=1}^{N} \|A_{i}d^{k}\|_{L^{2}(0,T)}^{2} + \gamma \|d^{k}\|_{L^{2}(Q)}^{2}}, \ k = 0, 1, 2, \dots$$

Springer

Thus, the CG has the form

Step 1: Set k = 0, initiate f^0 . Step 2: Calculate $r^0 = -\nabla J_{\gamma}(f^0)$ and set $d^0 = r^0$. Step 3: Evaluate

$$\alpha^{0} = \frac{\|F^{*}\|_{L^{2}(Q)}^{2}}{\sum_{i=1}^{N} \|A_{i}d^{0}\|_{L^{2}(0,T)}^{2} + \gamma \|d^{0}\|_{L^{2}(Q)}^{2}}$$

Set $f^1 = f^0 + \alpha^0 d^0$. Step 4: For $k = 1, 2, \dots$ Calculate

$$r^k = -\nabla J_{\gamma}(f^k), \qquad d^k = r^k + \beta^k d^{k-1}$$

with

$$\beta^{k} = \frac{\|r^{k}\|_{L^{2}(Q)}^{2}}{\|r^{k-1}\|_{L^{2}(Q)}^{2}}.$$

Step 5: Calculate

$$\alpha^{k} = \frac{\|r^{k}\|_{L^{2}(Q)}^{2}}{\sum_{i=1}^{N} \|A_{i}d^{k}\|_{L^{2}(0,T)}^{2} + \gamma \|d^{k}\|_{L^{2}(Q)}^{2}}.$$

Update

$$f^{k+1} = f^k + \alpha^k d^k.$$

2.2 Singular Values

Set

$$\mathcal{A} = (A_1, A_2, \dots, A_N), \qquad z = (z_1, z_2, \dots, z_n)$$

where A_i is defined in (20). The problem of determining f in (7) (f has form in (8) or (9) or (10)) from (11) can be written in the form Af = z, where

$$\mathcal{A}: L^2(Q)\left(L^2(\Omega) \text{ or } L^2(0,T)\right) \to \left(L^2(0,T)\right)^N$$

To characterize the ill-posedness degree of the inverse source problem, we have to estimate the singular values of A, i.e., the eigenvalues of A^*A . In doing so, we proceed as follows.

We will present for the case f depends on both time and space variable, i.e., $\mathcal{A} : L^2(Q) \to (L^2(0,T))^N$. For the operator A_i , we have $A_i^* \tilde{g} = \varphi(x,t) \tilde{p}(x,t)$, where $\tilde{g} \in L^2(0,T)$ and $\tilde{p}(x,t)$ is the solution to the adjoint problem

$$\begin{cases} -\frac{\partial \tilde{p}}{\partial t} - \sum_{i,j=1}^{n} \frac{\partial}{\partial x_{j}} \left(a_{ij}(x,t) \frac{\partial \tilde{p}}{\partial x_{i}} \right) + b(x,t) \tilde{p} = \omega_{i}(x) \tilde{g}, \quad (x,t) \in Q, \\ \tilde{p}(x,t) = 0, \quad (x,t) \in S, \\ \tilde{p}(x,T) = 0, \quad x \in \Omega. \end{cases}$$

From (20), we have

$$J_0(f) = \frac{1}{2} \sum_{i=1}^N \|l_i u(f) - z_i\|_{L^2(0,T)}^2 = \frac{1}{2} \sum_{i=1}^N \|A_i f - (z_i - l_i \bar{u}(v,g))\|_{L^2(0,T)}^2.$$



Hence,

$$J'_{0}(f) = \sum_{i=1}^{N} A_{i}^{*} \left(A_{i} f - (z_{i} - l_{i} \bar{u}(v, g)) \right).$$

If we take z_i such that $z_i = l_i \bar{u}(v, g)$, then due to Theorem 1, we have $J'_0(f) = \sum_{i=1}^N A_i^* A_i f = \varphi(x, t) p^*(x, t)$, where p^* is the solution of the adjoint problem

$$\begin{cases} -\frac{\partial p^*}{\partial t} - \sum_{i,j=1}^n \frac{\partial}{\partial x_j} \left(a_{ij}(x,t) \frac{\partial p^*}{\partial x_i} \right) + b(x,t) p^* = \sum_{i=1}^N \omega_i(x) l_i \tilde{u}[f], \quad (x,t) \in Q, \\ p^*(x,t) = 0, \quad (x,t) \in S, \\ p^*(x,T) = 0, \quad x \in \Omega. \end{cases}$$

Thus, if $f(x,t) \in L^2(Q)$ is given, we can calculate the value $J'_0(f) = \sum_{i=1}^N A_i^* A_i f = \varphi(x,t)p^*(x,t)$. Although we do not know the explicit form of A_i^*A , we can use the Lanczos algorithm [20] to estimate its eigenvalues when we discretize the problem. The algorithm looks as follows:

Initialization: Let $\beta_0 = 0$, $q_0 = 0$ and an arbitrary vector b, calculate $q_1 = \frac{b}{\|b\|}$. Put $Q = q_1$ and k = 0. *Iteration:* For k = 1, 2, 3, ...

$$p = \mathcal{A}^* \mathcal{A} q_k,$$

$$\alpha_k = q_k^T p,$$

$$p = p - \beta_{n-1} q_{n-1} - \alpha_k q_k,$$

$$\beta_k = \|p\|,$$

$$q_{k+1} = \frac{p}{\|\beta_k\|}.$$

We will present some numerical examples showing the efficiency of this algorithm in Section 4.

3 Variational Method for Discretized Problem

In this section, we have to restrict some conditions on the domain and coefficients. We start with Problem (13)–(14). First, we suppose that Ω is the open parallelepiped $(0, L_1) \times (0, L_2) \times \cdots \times (0, L_n)$ in \mathbb{R}^n . Second, in (7), we suppose that $a_{ij} = 0$, if $i \neq j$, and for simplicity from now on we denote a_{ii} by a_i . Following [15, 16, 25] (see also [6, 19]), we subdivide the domain Ω into small cells by the rectangular uniform grid specified by

$$0 = x_i^0 < x_i^1 = h_i < \dots < x_i^{N_i} = L_i, \ i = 1, \dots, n$$

with $h_i = L_i/N_i$ being the grid size in the x_i -direction, i = 1, ..., n. To simplify the notation, we denote by $x^k := (x_1^{k_1}, ..., x_n^{k_n})$, where $k := (k_1, ..., k_n)$, $0 \le k_i \le N_i$. We also denote by $h := (h_1, ..., h_n)$ the vector of spatial grid sizes and $\Delta h := h_1 \cdots h_n$. Let e_i be the unit vector in the x_i -direction, i = 1, ..., n, i.e., $e_1 = (1, 0, ..., 0)$ and so on. Denote by

$$\omega(k) = \{x \in \Omega : (k_i - 0.5)h_i \le x_i \le (k_i + 0.5)h_i, \forall i = 1, \dots, n\}.$$



In the following, Ω_h denotes the set of the indices of all interior grid points and $\bar{\Omega}_h$ denotes the set of the indices of all grid points belonging to $\bar{\Omega}_h$, i.e.,

$$\Omega_h = \{k = (k_1, \dots, k_n) : 1 \le k_i \le N_i - 1, \forall i = 1, \dots, n\}.$$

We also make use of the following sets

$$\Omega_h^i = \{k = (k_1, \dots, k_n) : 0 \le k_i \le N_i - 1, 1 \le k_j \le N_j - 1, \forall j \ne i\}$$

for i = 1, ..., n. For a function u(x, t) defined in Q_T , we denote by $u^k(t)$ its approximate value at (x^k, t) . We define the following forward finite difference quotient with respect to x_i

$$u_{x_i}^k := \frac{u^{k+e_i} - u^k}{h_i}.$$

Now, taking into account the homogeneous boundary condition, we approximate the integrals in (13) as follows

$$\int_{Q} \frac{\partial u}{\partial t} \eta dx dt \approx \Delta h \int_{0}^{T} \sum_{k \in \Omega_{h}} \frac{du^{k}(t)}{dt} \eta^{k}(t) dt,$$
(21)

$$\int_{Q} a_{i}(x,t) \frac{\partial u}{\partial x_{i}} \frac{\partial \eta}{\partial x_{i}} dx dt \approx \Delta h \int_{0}^{T} \sum_{k \in \Omega_{h}^{i}} a_{i}^{k + \frac{e_{i}}{2}}(t) u_{x_{i}}^{k}(t) \eta_{x_{i}}^{k}(t) dt,$$
(22)

$$\int_{Q} b(x,t)u\eta dxdt \approx \Delta h \int_{0}^{T} \sum_{k \in \Omega_{h}} b^{k}(t)u^{k}(t)\eta^{k}(t)dt, \qquad (23)$$

$$\int_{Q} f(x,t)\varphi(x,t)\eta dxdt \approx \Delta h \int_{0}^{T} \sum_{k \in \Omega_{h}} f^{k}(t)\varphi^{k}(t)\eta^{k}(t)dt,$$
(24)

$$\int_{Q} g(x,t)\eta dx dt \approx \Delta h \int_{0}^{T} \sum_{k \in \Omega_{h}} g^{k}(t)\eta^{k}(t) dt.$$
(25)

Here $b^k(t)$, $f^k(t)$, $\varphi^k(t)$, $g^k(t)$ and $a_i^{k+\frac{e_i}{2}}(t)$ are approximations to the functions b(x, t), f(x, t), $\varphi(x, t)$, g(x, t) and $a_i(x, t)$ at the grid point x^k . More precisely, if these functions are continuous at x^k , we take their approximations by their value at x^k and $a_i^{k+\frac{e_i}{2}}(t) = a_i(x^{k+\frac{e_i}{2}}, t)$. Otherwise, we take

$$b^{k}(t) = \frac{1}{|\omega(k)|} \int_{\omega(k)} b(x, t) dx, \qquad f^{k}(t) = \frac{1}{|\omega(k)|} \int_{\omega(k)} f(x, t) dx,$$
$$\varphi^{k}(t) = \frac{1}{|\omega(k)|} \int_{\omega(k)} \varphi(x, t) dx, \qquad g^{k}(t) = \frac{1}{|\omega(k)|} \int_{\omega(k)} g(x, t) dx.$$

and

$$a_i^{k+\frac{e_i}{2}}(t) = \frac{1}{|\omega(k)|} \int_{\omega(k)} a_i(x,t) dx.$$

With the approximations (21), (22), (23), (24) and (25), we have the following discrete analogue of (13)

$$\int_{0}^{T} \left[\sum_{k \in \Omega_{h}} \left(\frac{du^{k}}{dt} + b^{k} u^{k} - f^{k} \right) \eta^{k} + \sum_{i=1}^{n} \sum_{k \in \Omega_{h}^{i}} a_{i}^{k + \frac{e_{i}}{2}} u_{x_{i}}^{k} \eta_{x_{i}}^{k} \right] dt = 0.$$
 (26)

Springer

We note that, using the discrete analogue of integration by parts with boundary condition $u^0 = \eta^0 = 0$ and $u^{N_i} = \eta^{N_i} = 0$, we obtain

$$\begin{split} \sum_{k \in \Omega_h^i} a_i^{k + \frac{e_i}{2}} u_{x_i}^k \eta_{x_i}^k &= \sum_{k \in \Omega_h^i} a_i^{k + \frac{e_i}{2}} \frac{u^{k + e_i} - u^k}{h_i} \frac{\eta^{k + e_i} - \eta^k}{h_i} \\ &= \sum_{k \in \Omega_h^i} a_i^{k + \frac{e_i}{2}} \frac{u^{k + e_i} - u^k}{h_i^2} \eta^{k + e_i} - \sum_{k \in \Omega_h^i} a_i^{k + \frac{e_i}{2}} \frac{u^{k + e_i} - u^k}{h_i^2} \eta^k \\ &= \sum_{k \in \Omega_h} \left(a_i^{k - \frac{e_i}{2}} \frac{u^k - u^{k - e_i}}{h_i^2} - a_i^{k + \frac{e_i}{2}} \frac{u^{k + e_i} - u^k}{h_i^2} \right) \eta^k. \end{split}$$

Hence, replacing this equality into (26), we obtain the following system which approximates the original problem (7)

$$\begin{cases} \frac{d\bar{u}}{dt} + (\Lambda_1 + \dots + \Lambda_n)\bar{u} - \bar{F} = 0, \\ \bar{u}(0) = \bar{v}, \end{cases}$$
(27)

with $\bar{u} = \{u^k, k \in \Omega_h\}$ being the grid function. The function \bar{v} is the grid function approximating the initial condition v and

$$(\Lambda_{i}\bar{u})^{k} = \frac{b^{k}u^{k}}{n} + \begin{cases} \frac{a_{i}^{k-\frac{e_{i}}{2}}}{h_{i}^{2}}(u^{k}-u^{k-e_{i}}) - \frac{a_{i}^{k+\frac{e_{i}}{2}}}{h_{i}^{2}}(u^{k+e_{i}}-u^{k}), 2 \leq k_{i} \leq N_{i} - 2, \\ \frac{a_{i}^{k-\frac{e_{i}}{2}}}{h_{i}^{2}}u^{k} - \frac{a_{i}^{k+\frac{e_{i}}{2}}}{h_{i}^{2}}(u^{k+e_{i}}-u^{k}), k_{i} = 1, \\ \frac{a_{i}^{k-\frac{e_{i}}{2}}}{h_{i}^{2}}(u^{k}-u^{k-e_{i}}) + \frac{a_{i}^{k+\frac{e_{i}}{2}}}{h_{i}^{2}}u^{k}, k_{i} = N_{i} - 1 \end{cases}$$

for $k \in \Omega_h$ and

$$\bar{F} = \{ f^k \varphi^k + g^k, k \in \Omega_h \}.$$

We note that the coefficient matrices Λ_i are positive semi-definite (see, e.g., [19]). The boundedness of the solution of (27) has shown in the following theorem.

Theorem 2 Let \bar{u} be a solution of the Cauchy problem (27). There exists a constant c independent of h and the coefficients of the equation such that

$$\max_{t \in [0,T]} \sum_{k \in \bar{\mathcal{Q}}_h} |\bar{u}^k(t)|^2 + \int_0^T \sum_{i=1}^n \sum_{k \in \mathcal{Q}_h^i} |\bar{u}_{x_i}^k|^2 dt \le c \left(\int_0^T \sum_{k \in \bar{\mathcal{Q}}_h} |\bar{f}^k|^2 dt + \sum_{k \in \bar{\mathcal{Q}}_h} |\bar{v}^k|^2 \right).$$
(28)

Proof For arbitrary $t^* \in (0, T]$, set

$$\bar{\eta}^{k}(t) = \begin{cases} \bar{u}^{k}(t) & \text{if } t \in [0, t^{*}], \\ 0 & \text{if } t \notin [0, t^{*}]. \end{cases}$$

Since

$$\int_0^{t^*} dt \sum_{k \in \overline{\Omega}_h} \bar{u}_t^k(t) \bar{u}^k(t) = \frac{1}{2} \sum_{k \in \overline{\Omega}_h} |\bar{u}^k(t^*)|^2 - \frac{1}{2} \sum_{k \in \overline{\Omega}_h} |\bar{u}^k(0)|^2,$$



and $\bar{u}^k(0) = \bar{v}$, it follows from (26) that

$$\frac{1}{2} \sum_{k \in \bar{\Omega}_{h}} |\bar{u}^{k}(t^{*})|^{2} + \int_{0}^{t^{*}} \left[\sum_{k \in \bar{\Omega}_{h}} \bar{b}^{k} |u^{k}|^{2} + \sum_{i=1}^{n} \sum_{k \in \Omega_{h}^{i}} \bar{a}^{k}_{i} |\bar{u}^{k}_{x_{i}}|^{2} \right] dt$$

$$= \int_{0}^{t^{*}} \sum_{k \in \bar{\Omega}_{h}} \bar{f}^{k} \bar{u}^{k} dt + \frac{1}{2} \sum_{k \in \bar{\Omega}_{h}} |\bar{v}^{k}|^{2}.$$
(29)

Multiplying the both sides of the equality (29) by 2, applying Cauchy's inequality to the first term in the right hand side, noting that $b^k \ge 0$, we obtain

$$\sum_{k\in\bar{\Omega}_{h}} |\bar{u}^{k}(t^{*})|^{2} + 2\int_{0}^{t^{*}} \sum_{i=1}^{n} \sum_{k\in\Omega_{h}^{i}} \bar{a}_{i}^{k} |\bar{u}_{x_{i}}^{k}|^{2} dt$$

$$\leq \int_{0}^{t^{*}} \sum_{k\in\bar{\Omega}_{h}} |\bar{f}^{k}|^{2} dt + \int_{0}^{t^{*}} \sum_{k\in\bar{\Omega}_{h}} |\bar{u}^{k}|^{2} dt + \sum_{k\in\bar{\Omega}_{h}} |\bar{v}^{k}|^{2}.$$
(30)

Put

$$y(t) = \sum_{k \in \bar{\Omega}_h} |\bar{u}^k(t^*)|^2.$$

From (30) we have

$$y(t^*) \le \int_0^{t^*} y(t) dt + \int_0^{t^*} \sum_{k \in \bar{\Omega}_h} |\bar{f}^k|^2 dt + \sum_{k \in \bar{\Omega}_h} |\bar{v}^k|^2.$$

Applying Gronwall's inequality, we obtain

$$y(t^{*}) \leq \left(\int_{0}^{t^{*}} \sum_{k \in \bar{\Omega}_{h}} |\bar{f}^{k}|^{2} dt + \sum_{k \in \bar{\Omega}_{h}} |\bar{v}^{k}|^{2} \right) e^{t}.$$
 (31)

Hence, we have

$$\max_{t \in [0,T]} \sum_{k \in \bar{\Omega}_h} |\bar{u}^k(t)|^2 \le c \left(\int_0^T \sum_{k \in \bar{\Omega}_h} |\bar{f}^k|^2 dt + \sum_{k \in \bar{\Omega}_h} |\bar{v}^k|^2 \right).$$

From the conditions (1), (2) and (3) about the coefficient a_i , the inequalities (30) and (31) we have

$$\int_0^T \left(\sum_{k \in \bar{\Omega}_h} |\bar{u}^k(t)|^2 + \sum_{i=1}^n \sum_{k \in \Omega_h^i} |\bar{u}_{x_i}^k|^2 \right) dt \le c \left(\int_0^T \sum_{k \in \bar{\Omega}_h} |\bar{f}^k|^2 dt + \sum_{k \in \bar{\Omega}_h} |\bar{v}^k|^2 \right).$$

Combining the two inequalities, we obtain the inequality (28).

3.1 Time Discretization

To obtain the finite difference scheme for (27), we divide the time interval [0, T] into M sub-intervals by the points t_i , i = 0, ..., M, $t_0 = 0$, $t_1 = \Delta t$, ..., $t_M = M\Delta t = T$. For

simplifying the notation, we set $u^{k,m} := u^k(t_m)$. We also denote by $F^{k,m} := F^k(t_m)$ and $\Lambda_i^m = \Lambda_i(t_m), m = 0, ..., M$. In the following, we drop the spatial index for simplifying the notation. The finite difference scheme is written as follows

$$\begin{cases} u^{m+1} = m^m + \Delta t [F^m - (\Lambda_1^m + \dots + \Lambda_n^m) u^m)], \\ u^0 = \bar{v}. \end{cases}$$

3.2 Splitting Method

In order to obtain a splitting scheme for the Cauchy problem (27), we also discrete the time interval in the same with finite difference method. We denote $u^{m+\delta} := \bar{u}(t_m + \delta \Delta t)$, $\Lambda_i^m := \Lambda_i(t_m + \Delta t/2)$. We introduce the following implicit two-circle component-by-component splitting scheme [15]

$$\frac{u^{m+\frac{i}{2n}} - u^{m+\frac{i-1}{2n}}}{\Delta t} + \Lambda_i^m \frac{u^{m+\frac{i}{2n}} + u^{m+\frac{i-1}{2n}}}{4} = 0, \quad i = 1, 2, \dots, n-1,$$

$$\frac{u^{m+\frac{1}{2}} - u^{m+\frac{n-1}{2n}}}{\Delta t} + \Lambda_n^m \frac{u^{m+\frac{1}{2}} + u^{m+\frac{n-1}{2n}}}{4} = \frac{F^m}{2} + \frac{\Delta t}{8} \Lambda_n^m F^m,$$

$$\frac{u^{m+\frac{n+1}{2n}} - u^{m+\frac{1}{2}}}{\Delta t} + \Lambda_n^m \frac{u^{m+\frac{n+1}{2n}} + u^{m+\frac{1}{2}}}{4} = \frac{F^m}{2} - \frac{\Delta t}{8} \Lambda_n^m F^m,$$

$$\frac{u^{m+1-\frac{i-1}{2n}} - u^{m+1-\frac{i}{2n}}}{\Delta t} + \Lambda_i^m \frac{u^{m+1-\frac{i-1}{2n}} + u^{m+1-\frac{i}{2n}}}{4} = 0, \quad i = n-1, n-2, \dots, 1,$$

$$u^0 = \bar{v}.$$
(32)

Equivalently,

$$\left(E_{i} + \frac{\Delta t}{4}\Lambda_{i}^{m}\right)u^{m+\frac{i}{2n}} = \left(E_{i} - \frac{\Delta t}{4}\Lambda_{i}^{m}\right)u^{m+\frac{i-1}{2n}}, \quad i = 1, 2, \dots, n-1,$$

$$\left(E_{n} + \frac{\Delta t}{4}\Lambda_{n}^{m}\right)\left(u^{m+\frac{1}{2}} - \frac{\Delta t}{2}F^{m}\right) = \left(E_{n} - \frac{\Delta t}{4}\Lambda_{n}^{m}\right)u^{m+\frac{n-1}{2n}},$$

$$\left(E_{n} + \frac{\Delta t}{4}\Lambda_{n}^{m}\right)u^{m+\frac{n+1}{2n}} = \left(E_{n} - \frac{\Delta t}{4}\Lambda_{n}^{m}\right)\left(u^{m+\frac{1}{2}} + \frac{\Delta t}{2}F^{m}\right),$$

$$\left(E_{i} + \frac{\Delta t}{4}\Lambda_{i}^{m}\right)u^{m+1-\frac{i-1}{2n}} = \left(E_{i} - \frac{\Delta t}{4}\Lambda_{i}^{m}\right)u^{m+1-\frac{i}{2n}}, \quad i = n-1, n-2, \dots, 1,$$

$$u^{0} = \bar{v},$$

$$(33)$$

where E_i is the identity matrix corresponding to Λ_i , i = 1, ..., n. The splitting scheme (33) can be rewritten in the following compact form

$$\begin{cases} u^{m+1} = B^m u^m + \Delta t C^m (f^m \varphi^m + g^m), & m = 0, \dots, M - 1, \\ u^0 = \bar{v}, \end{cases}$$
(34)

with

ۣ ؊

$$B^m = B_1^m \cdots B_n^m B_n^m \cdots B_1^m, \qquad C^m = C_1^m \cdots C_n^m,$$

where $B_i^m := (E_i + \frac{\Delta t}{4} \Lambda_i^m)^{-1} (E_i - \frac{\Delta t}{4} \Lambda_i^m), i = 1, \dots, n.$

3.3 Discretized Variational Problem

To complete the variational method for multi-dimensional cases, we use the splitting method for the forward problem and take the discretized functional

$$J_{0}^{h,\Delta t}(\bar{f}) := \frac{\Delta t}{2} \sum_{i=1}^{N} \sum_{m=1}^{M} \left[\Delta h \sum_{k \in \Omega_{h}} \omega_{i}^{k} u^{k,m}(\bar{f}) - z_{i}^{m} \right]^{2},$$
(35)

where $u^{k,m}(\bar{f})$ shows its dependence on the right-hand side term \bar{f} and m is the index of grid points on time axis. The notation $\omega_i^k = \omega_i(x^k)$ indicates the approximation of the function $\omega_i(x)$ in Ω_h at points x^k . Normally, we take as its average over the cell where x_k is located.

For minimizing the problem (35) by the conjugate gradient method, we first calculate the gradient of objective function $J_0^{h,\Delta t}(\bar{f})$ and it is shown by the following theorem

Theorem 3 The gradient $\nabla J_0^{h,\Delta t}(\bar{f})$ of the objective function $J_0^{h,\Delta t}$ at \bar{f} is given by

$$\nabla J_0^{h,\Delta t}(\bar{f}) = \Delta t \sum_{m=0}^{M-1} (C^m)^* \varphi^m \eta^m,$$
(36)

where η satisfies the adjoint problem

$$\begin{cases} \eta^m = (B^{m+1})^* \eta^{m+1} + \psi^{m+1}, & m = M - 1, M - 2, \dots, 0, \\ \eta^M = 0, \end{cases}$$
(37)

with

$$\psi^{k,m} = \Delta h \sum_{i=1}^{N} \omega_i^k \left(\sum_{k \in \Omega_h} \omega_i^k u^{k,m} - z_i^m \right), \ k \in \Omega_h, \ m = 0, \dots, M.$$

Here the matrix $(B^m)^*$ is given by

$$(B^{m})^{*} = \left(E_{1} - \frac{\Delta t}{4}\Lambda_{1}^{m}\right)\left(E_{1} + \frac{\Delta t}{4}\Lambda_{1}^{m}\right)^{-1}\dots\left(E_{n} - \frac{\Delta t}{4}\Lambda_{n}^{m}\right)\left(E_{n} + \frac{\Delta t}{4}\Lambda_{n}^{m}\right)^{-1} \times \left(E_{n} - \frac{\Delta t}{4}\Lambda_{n}^{m}\right)\left(E_{n} + \frac{\Delta t}{4}\Lambda_{n}^{m}\right)^{-1}\dots\left(E_{1} - \frac{\Delta t}{4}\Lambda_{1}^{m}\right)\left(E_{1} + \frac{\Delta t}{4}\Lambda_{1}^{m}\right)^{-1}$$

Proof For an infinitesimally small variation δf of f, we have from (35) that

$$\begin{split} J_0^{h,\Delta t}(\bar{f}+\delta\bar{f}) - J_0^{h,\Delta t}(\bar{f}) &= \frac{\Delta t}{2} \sum_{i=1}^N \sum_{m=1}^M \left[\Delta h \sum_{k \in \Omega_h} \omega_i^k u^{k,m}(\bar{f}+\delta\bar{f}) - z_i^m \right]^2 \\ &- \frac{\Delta t}{2} \sum_{i=1}^N \sum_{m=1}^M \left[\Delta h \sum_{k \in \Omega_h} \omega_i^k u^{k,m}(\bar{f}) - z_i^m \right]^2 \\ &= \frac{\Delta t}{2} \sum_{i=1}^N \sum_{m=1}^M \sum_{k \in \Omega_h} (\Delta h \omega_i^k w^{k,m})^2 \end{split}$$



$$+\Delta t \sum_{i=1}^{N} \sum_{m=1}^{M} \Delta h \sum_{k \in \Omega_{h}} \omega_{i}^{k} w^{k,m} \left[\Delta h \sum_{k \in \Omega_{h}} \omega_{i}^{k} u^{k,m} (\bar{f}) - z_{i}^{m} \right]$$
$$= \frac{\Delta t}{2} \sum_{i=1}^{N} \sum_{m=1}^{M} \sum_{k \in \Omega_{h}} (\Delta h \omega_{i}^{k} w^{k,m})^{2} + \Delta t \sum_{i=1}^{N} \sum_{m=1}^{M} \Delta h \sum_{k \in \Omega_{h}} w^{k,m} \psi_{i}^{k,m}$$
$$= \frac{\Delta t}{2} \sum_{i=1}^{N} \sum_{m=1}^{M} \sum_{k \in \Omega_{h}} (\Delta h \omega_{i}^{k} w^{k,m})^{2} + \Delta t \sum_{i=1}^{N} \sum_{m=1}^{M} \langle w^{m}, \psi_{i}^{m} \rangle, \qquad (38)$$

where $w^{k,m} := u^{k,m}(\bar{f} + \delta \bar{f}) - u^{k,m}(\bar{f})$ and $\psi_i^{k,m} = \Delta h \omega_i^k (\sum_{k \in \Omega_h} \omega_i^k u^{k,m} - z_i^m), k \in \Omega_h$. It follows from (34) that w is the solution to the problem

$$\begin{cases} w^{m+1} = A^m w^m + \Delta t C^m \delta \bar{f} \varphi^m, & m = 0, \dots, M - 1, \\ w^0 = 0. \end{cases}$$
(39)

Taking the inner product of both sides of the *m*th equation of (39) with an arbitrary vector $\eta^m \in \mathbb{R}^{N_1 \times \cdots \times N_n}$, summing the results over $m = 0, \dots, M - 1$, we obtain

$$\sum_{m=0}^{M-1} \langle w^{m+1}, \eta^m \rangle = \sum_{m=0}^{M-1} \langle B^m w^m, \eta^m \rangle + \sum_{m=0}^{M-1} \langle \Delta t C^m \delta \bar{f} \varphi^m, \eta^m \rangle$$

$$= \sum_{m=0}^{M-1} \langle w^m, (B^m)^* \eta^m \rangle + \sum_{m=0}^{M-1} \langle \Delta t C^m \delta \bar{f} \varphi^m, \eta^m \rangle.$$
(40)

Here $(B^m)^*$ is the adjoint matrix of B^m .

Taking the inner product of both sides of the first equation of (37) with an arbitrary vector w^{m+1} , summing the results over m = 0, ..., M - 1, we obtain

$$\sum_{m=0}^{M-1} \langle w^{m+1}, \eta^m \rangle = \sum_{m=0}^{M-1} \langle w^{m+1}, (B^{m+1})^* \eta^{m+1} \rangle + \sum_{m=0}^{M-1} \langle w^{m+1}, \psi^{m+1} \rangle$$

$$= \sum_{m=1}^{M} \langle w^m, (B^m)^* \eta^m \rangle + \sum_{m=1}^{M} \langle w^m, \psi^m \rangle.$$
(41)

Note that $w^0 = \eta^M = 0$, from (40) and (41), we have

$$\sum_{m=1}^{M} \langle w^m, \psi^m \rangle = \sum_{m=0}^{M-1} \langle \Delta t C^m \delta \bar{f} \varphi^m, \eta^m \rangle.$$
(42)

On the other hand, it can be proved by induction that $\sum_{i=1}^{N} \sum_{m=1}^{M} \sum_{k \in \Omega_h} (\omega_i^k w^{k,m})^2 = o(\|\delta \bar{f}\|)$. Hence, from (38) and (42), we obtain

$$J_0^{h,\Delta t}(\bar{f} + \delta \bar{f}) - J_0^{h,\Delta t}(\bar{f}) = \sum_{m=0}^{M-1} (\delta \bar{f}, \Delta t(C^m)^* \varphi^m \eta^m) + o(\|\delta \bar{f}\|).$$

Consequently, the gradient of the objective function J_0^h can be written as

$$\frac{\partial J_0^{h,\Delta t}(\bar{f})}{\partial \bar{f}} = \Delta t \sum_{m=0}^{M-1} (C^m)^* \varphi^m \eta^m.$$

Note that, since the coefficient matrices Λ_i^m , i = 1, ..., n, m = 0, ..., M-1 are symmetric, we have

$$(B^{m})^{*} = \left(E_{1} - \frac{\Delta t}{4}\Lambda_{1}^{m}\right)\left(E_{1} + \frac{\Delta t}{4}\Lambda_{1}^{m}\right)^{-1}\dots\left(E_{n} - \frac{\Delta t}{4}\Lambda_{n}^{m}\right)\left(E_{n} + \frac{\Delta t}{4}\Lambda_{n}^{m}\right)^{-1}$$
$$\times \left(E_{n} - \frac{\Delta t}{4}\Lambda_{n}^{m}\right)\left(E_{n} + \frac{\Delta t}{4}\Lambda_{n}^{m}\right)^{-1}\left(E_{1} - \frac{\Delta t}{4}\Lambda_{1}^{m}\right)\left(E_{1} + \frac{\Delta t}{4}\Lambda_{1}^{m}\right)^{-1}$$

and

$$(C^m)^* = \left(E_n - \frac{\Delta t}{4}\Lambda_n^m\right) \left(E_n + \frac{\Delta t}{4}\Lambda_n^m\right)^{-1} \left(E_1 - \frac{\Delta t}{4}\Lambda_1^m\right) \left(E_1 + \frac{\Delta t}{4}\Lambda_1^m\right)^{-1}.$$

e proof is complete.

The proof is complete.

The conjugate gradient method for the discretized function (35) can be written by following steps:

Step 1. Given an initial approximation f^0 and calculate the residual $\hat{r}^0 = \sum_{i=1}^{N} [l_i u(f^0) - z_i]$ by solving the splitting (32) with f being replaced by initial approximation f^0 and set k = 0.

Step 2. Calculate the gradient $r^0 = -\nabla J_{\gamma}(f^0)$ given in (36) by solving the adjoint problem (37). Then we set $d^0 = r^0$.

Step 3. Calculate

$$\alpha^{0} = \frac{\|r^{0}\|^{2}}{\sum\limits_{i=1}^{N} \|l_{i}d^{0}\|^{2} + \gamma \|d^{0}\|},$$

where $l_i d^0$ can be calculated from the splitting scheme (32) with f being replaced by d^0 and g(x, t) = 0, v = 0. Then, we set

$$f^1 = f^0 + \alpha^0 d^0.$$

Step 4. For $k = 1, 2, \ldots$, calculate $r^k = -\nabla J_{\gamma}(f^k), d^k = r^k + \beta^k d^{k-1}$, where

$$\beta^k = \frac{\|r^k\|^2}{\|r^{k-1}\|^2}.$$

Step 5. Calculate α^k

$$\alpha^{k} = \frac{\|r^{k}\|^{2}}{\sum_{i=1}^{N} \|l_{i}d^{k}\|^{2} + \gamma \|d^{k}\|},$$

where $l_i d^k$ can be calculated from the splitting scheme (32) with f being replaced by d^k and g(x, t) = 0, v = 0. Then, set

$$f^{k+1} = f^k + \alpha^k d^k.$$

4 Numerical Example

To illustrate the performance of the proposed algorithm, we present in this section some numerical tests. These algorithms were implemented in Matlab and run on a personal laptop with 11th Gen Intel(R) Core(TM) i5 2.4Mhz 2419 Mhz 4 Core(s) 8 Logical Processors.

☑ Springer

4.1 One-Dimensional Problems

In this subsection, we present some numerical examples to estimate singular values and determine f. Let $\Omega = (0, 1)$ and T = 1. Consider the one-dimensional system

$$u_t - (au_x)_x = f\varphi(x, t) + g(x, t), \ x \in (0, 1), \ 0 \le t \le 1,$$

$$u(0, t) = u(1, t) = 0, \ 0 \le t \le 1,$$

$$u(x, 0) = v, \ x \in (0, 1),$$

where

$$a = 2xt + x^{2}t + 1$$
; $v = \sin(2\pi x)$ and $\varphi(x, t) = (x^{2} + 1)(t^{2} + 1)$.

For discretization, we take the grid size to be 0.02 in x and t. We take 3 observations at $x^{10} = 0.2$, $x^{25} = 0.5$ and $x^{35} = 0.7$. The weighted functions $\omega_i(x)$, i = 1, 2, 3 are chosen as follows

$$\omega_{1}(x) = \begin{cases} \frac{1}{2\varepsilon} & \text{if } x \in (x^{10} - \varepsilon, x^{10} + \varepsilon) \\ 0 & \text{otherwise} \end{cases} \quad \text{with } \varepsilon = 0.01,$$
$$\omega_{2}(x) = \begin{cases} \frac{1}{2\varepsilon} & \text{if } x \in (x^{25} - \varepsilon, x^{25} + \varepsilon) \\ 0 & \text{otherwise} \end{cases} \quad \text{with } \varepsilon = 0.01,$$
$$\omega_{3}(x) = \begin{cases} \frac{1}{2\varepsilon} & \text{if } x \in (x^{35} - \varepsilon, x^{35} + \varepsilon) \\ 0 & \text{otherwise} \end{cases} \quad \text{with } \varepsilon = 0.01.$$

Approximate singular values of A for the case f depends only on time variable t and space variable x are drawn in Fig. 1. From this figure, we see that the singular values for the case when f depends only on x is much smaller than that for the case f depends only on t. Therefore, the problem of reconstructing f = f(x) is much more ill-posed than f = f(t).

Now we present numerical results for reconstructing f(x, t). We test three types of f(x, t): smooth, non-smooth and discontinuous in the following examples.



Fig. 1 Approximation singular values: (a) f depends only on x; (b) f depends only on t



Example 1

$$f(x,t) = \sin(\pi x)\sin(\pi t)$$

Example 2

$$f(x,t) = \begin{cases} 2t \text{ if } t \le 1/2 \text{ and } t \le x \text{ and } x \le 1-t, \\ 2(1-t) \text{ if } t \ge 1/2 \text{ and } t \ge x \text{ and } x \ge 1-t, \\ 2x \text{ if } x \le 1/2 \text{ and } x \le t \text{ and } t \le 1-x, \\ 2(1-x) \text{ otherwise.} \end{cases}$$

Example 3

$$f(x,t) = \begin{cases} 1, & 0.25 \le x, t \le 0.75, \\ 0 & \text{otherwise.} \end{cases}$$

In all of three above examples, the initial guess $f^* = 0$, $02(\operatorname{rand}(N_x, M) - 0, 5) + f$, noisy level $\delta = 0$, 02, $\gamma = 10^{-2}$ and the initial iteration of the conjugate gradient method $f^0 = 0$. Numerical solutions are presented in Figs. 2, 3 and 4.



Fig.2 Example 1. The exact solution in comparison with the numerical solution: (a) Exact function f(x, t); (b) Reconstruction of f; (c) Comparison of the exact and approximation solutions at x = 0, 24; (d) Comparison of the exact and approximation solutions at x = 0, 5





Fig.3 Example 2. The exact solution in comparison with the numerical solution: (a) Exact function f(x, t); (b) Reconstruction of f; (c) Comparison of the exact and approximation solutions at x = 0, 24; (d) Comparison of the exact and approximation solutions at x = 0, 5



Fig.4 Example 3. The exact solution in comparison with the numerical solution: (a) Exact function f(x, t); (b) Reconstruction of f; (c) Comparison of the exact and approximation solutions at x = 0, 24; (d) Comparison of the exact and approximation solutions at x = 0, 5

4.2 Two-Dimensional Problems

We consider the domain $\Omega = (0, 1) \times (0, 1)$, T = 1 and denote the space variable $x = (x_1, x_2)$. We take 4 observation distributed in 4 parts: $(0, 0.5) \times (0, 0.5)$, $(0.5, 1) \times (0, 0.5)$, $(0.5, 1) \times (0.5, 1)$ and $(0, 0.5) \times (0.5, 1)$.

Consider the system

.

.

$$\begin{aligned} u_t - (a_1 u_{x_1})_{x_1} - (a_2 u_{x_2})_{x_2} + a(x, t)u &= f\varphi(x, t) + g(x, t), \ (x, t) \in Q, \\ u(0, x_2, t) &= u(1, x_2, t) = u(x_1, 0, t) = u(x_2, 1, t) = 0, \ 0 < t \le T, \\ u(x, 0) &= v, \ x \in \Omega. \end{aligned}$$

The grid sizes are chosen 0.02 in x and in t. The weighted functions $\omega_i(x)$, i = 1, 2, 3, 4 are chosen as follows

$$\omega_{1}(x) = \begin{cases} \frac{1}{2\varepsilon} & \text{if } x \in (0, 24 - \varepsilon, 0, 24 + \varepsilon) \times (0, 24 - \varepsilon, 0, 24 + \varepsilon) \\ 0 & \text{otherwise} \end{cases} \quad \text{with } \varepsilon = 0, 01,$$

$$\omega_{2}(x) = \begin{cases} \frac{1}{2\varepsilon} & \text{if } x \in (0, 74 - \varepsilon, 0, 74 + \varepsilon) \times (0, 24 - \varepsilon, 0, 24 + \varepsilon) \\ 0 & \text{otherwise} \end{cases} \quad \text{with } \varepsilon = 0, 01,$$

$$\omega_{3}(x) = \begin{cases} \frac{1}{2\varepsilon} & \text{if } x \in (0, 24 - \varepsilon, 0, 24 + \varepsilon) \times (0, 74 - \varepsilon, 0, 74 + \varepsilon) \\ 0 & \text{otherwise} \end{cases} \quad \text{with } \varepsilon = 0, 01,$$

$$\omega_{4}(x) = \begin{cases} \frac{1}{2\varepsilon} & \text{if } x \in (0, 74 - \varepsilon, 0, 74 + \varepsilon) \times (0, 74 - \varepsilon, 0, 74 + \varepsilon) \\ 0 & \text{otherwise} \end{cases} \quad \text{with } \varepsilon = 0, 01.$$

We test our algorithm for three cases f: (1) f = f(t), (2) f = f(x) and (3) f = f(x, t).

Example 4 We choose the a priori estimation $f^* = 0$, regularization parameter $\gamma = 10^{-2}$, $f^0 = 0$, noise level $\delta = 0, 02$ and

$$a_1(x,t) = a_2(x,t) = 0.2(1 - 0.5\cos(3\pi x_1)\cos(3\pi x_2)\cos(3\pi t)),$$

$$a = x_1^2 + x_2^2 + 2x_1t + 1, \ v = \sin(\pi x_1)\sin(\pi x_2),$$

$$\varphi(x,t) = (x_1^2 + 3)(x_2^2 + 3)(t^2 + 3).$$

We suppose that f depends only on the time variable and has the form

1)

$$f(t) = \sin(2\pi t)$$

2)

$$f(t) = \begin{cases} 2t & \text{if } t < 0.5, \\ 2(1-t) & \text{otherwise.} \end{cases}$$

3)

$$f(t) = \begin{cases} 1 & \text{if } 0.25 \le t \le 0.75, \\ 0 & \text{otherwise.} \end{cases}$$

The numerical results of Example 4 are shown in Fig. 5.





Fig. 5 Example 4: the exact solution in comparison with the numerical solution: (a) f is of the form 1); (b) f is of the form 2); (c) f is of the form 3)

Example 5 We choose the a priori estimation $f^* = 0, 02(\operatorname{rand}(N_1, N_2) - 0, 5) + f$, regularization parameter $\gamma = 10^{-2}, f^0 = 0$, noise level $\delta = 0, 02$ and

$$a_1(x,t) = a_2(x,t) = a = 1, \ a = x_1^2 + x_2^2 + 2x_1t + 1$$

$$v = \sin(\pi x_1)\sin(\pi x_2), \ \varphi(x,t) = (x_1^2 + 1)(x_2^2 + 2)(t^2 + 2).$$

We suppose that f depends only on the space variable and has the form

. .

1)

$$f(x_1, x_2) = \sin(\pi x_1) \sin(\pi x_2).$$

. .

2)

$$f(x_1, x_2) = \begin{cases} 2x_2 & \text{if } x_2 \le 0.5 \text{ and } x_2 \le x_1 \le 1 - x_2, \\ 2(1 - x_2) & \text{if } x_2 \ge 0.5 \text{ and } x_2 \ge x_1 \ge 1 - x_2, \\ 2x_1 & \text{if } x_1 \le 0.5 \text{ and } x_1 \le x_2 \le 1 - x_1, \\ 2(1 - x_1) & \text{otherwise.} \end{cases}$$



3)

302

$$f(x_1, x_2) = \begin{cases} 1 & \text{if } 0.25 \le x_1 \le 0.75 \text{ and } 0.25 \le x_2 \le 0.75, \\ 0 & \text{otherwise.} \end{cases}$$

The numerical results of Example 5 are shown in Figs. 6, 7 and 8.

Example 6 We choose the a priori estimation $f^* = 0, 02(\operatorname{rand}(N_1, N_2, M) - 0, 5) + f$, regularization parameter $\gamma = 10^{-2}, f^0 = 0$, noise level $\delta = 0, 02$ and

$$a_1(x,t) = a_2(x,t) = a = 0.5, \ a = x_1^2 + x_2^2 + 2x_1t + 1$$
$$v = \sin(\pi x_1)\sin(\pi x_2), \ \varphi(x,t) = (x_1^2 + 2)(x_2^2 + 2)(t^2 + 2).$$

We suppose that f depends on both the space and time variable as follows

$$f(x_1, x_2, t) = \sin(\pi x_1) \sin(\pi x_2) t$$

The results of Example 6 are shown in Fig. 9.

We now discuss on the role of f^* . We will see that its choice is important in the case the inverse problem has many solutions.

We assume that f depends only on time variable. This guarantee the uniqueness solution to inverse problem. We take some different values for f^* . However, the choice of f^* does not affect much the numerical solution. The information of this test as in the case f depends only on time variable as in Example 4, regularization parameter $\gamma = 10^{-2}$, $f^0 = 0$, noise level $\delta = 0, 02$. The numerical results with $f^* = 0$, $f^* = 2$ and $f^* = 5$ are presented in Fig. 10 and Table 1 are not much different from each other.



Fig. 6 Example 5, form 1): the exact solution in comparison with the numerical solution: (a) Exact function f; (b) Reconstruction of f; (c) Point-wise error; (d) Comparison at $x_1 = 1/2$





Fig. 7 Example 5, form 2): the exact solution in comparison with the numerical solution: (a) Exact function f; (b) Reconstruction of f; (c) Point-wise error; (d) Comparison at $x_1 = 1/2$

In the case when the solution is not unique, the choice of f^* is crucial. As mention above, there may be infinitely many solutions to the inverse problem, the prediction f^* plays a significant role for selecting the solution. We use the system as in the case f depends both on time and space variables as in Example 6, regularization parameter $\gamma = 10^{-2}$, $f^0 = 0$,



Fig. 8 Example 5, form 3): the exact solution in comparison with the numerical solution: (a) Exact function f; (b) Reconstruction of f; (c) Point-wise error; (d) Comparison at $x_1 = 1/2$

ŵ



Fig.9 Example 6. The exact solution in comparison with the numerical solution at t = 1/2: (a) Exact function f; (b) Reconstruction of f; (c) Point-wise error; (d) Comparison at $x_1 = 1/2$ and t = 1/2



Fig. 10 Exact solution and its approximation with $f^* = 0$, $f^* = 2$, $f^* = 5$

Table 1 L^2 -error with prediction $f^* = 0, f^* = 2, f^* = 5$	f^*	0	2	5
	L^2 – error	0.070528	0.077008	0.89275



Fig. 11 The exact solution in comparison with the numerical solution with $f^* = f_1^*$: (a) Exact solution; (b) Reconstruction of f; (c) Point-wise error



Fig. 12 The exact solution in comparison with the numerical solution with $f^* = f_2^*$: (a) Exact solution; (b) Reconstruction of f; (c) Point-wise error



Fig. 13 The exact solution in comparison with the numerical solution with $f^* = f_3^*$: (a) Exact solution; (b) Reconstruction of f; (c) Point-wise error

Table 2 L^2 -error with the prediction f_1^*, f_2^*, f_2^*		f_{1}^{*}	f_{2}^{*}	f_3^*
1 01 02 03	L^2 – error	0,23757	0,26129	0,30358



Fig. 14 The exact solution in comparison with the its approximation with 9 observations: (a) $f = \sin(2\pi t)$; (b) $f = \begin{cases} 2t & \text{if } t < 0.5, \\ 2(1-t) & \text{otherwise} \end{cases}$; (c) $f = \begin{cases} 1 & \text{if } 0.25 \le t \le 0.75, \\ 0 & \text{otherwise.} \end{cases}$

noise level $\delta = 0, 02$. By varying f^* near f, we can see that the conjugate gradient method will reconstruct the approximation which is closest f^* .

In the test, if we choose f^* by

$$f_1^* = 0, 02 \left(\operatorname{rand}(N_1, N_2, M) - 0, 5 \right) + f,$$

$$f_2^* = 0, 1 \left(\operatorname{rand}(N_1, N_2, M) - 0, 5 \right) + f,$$

$$f_3^* = 0, 5 \left(\operatorname{rand}(N_1, N_2, M) - 0, 5 \right) + f.$$

	$f = \sin(2\pi t)$	$f = \begin{cases} 2t \text{ if } t < 0.5, \\ 2(1-t) \text{ otherwise} \end{cases}$	$f = \begin{cases} 1 & \text{if } 0.25 \le t \le 0.75, \\ 0 & \text{otherwise} \end{cases}$
3 observations	0,052077	0,055625	0,074178
9 observations	0,049649	0,050122	0,054525

Table 3 L^2 -error with 3 observations and 9 observations

The numerical results are presented as in Figs. 11, 12, 13 and Table 2. We can see that if f^* is not close to the exact f, the algorithm cannot reconstruct the chosen f, but maybe the other one.

In the last example, we will test in case we have more observations. The priori estimation $f^* = 0$, noise level $\delta = 0, 02$, regularization parameter $\gamma = 10^{-2}$, $f^0 = 0, a_1(x, t), a_2(x, t), a(x, t)$ and the initial condition v are chosen as in Example 4. The grid sizes are chosen 0.02 in x and in t. We choose 9 observations in domains $(0, 0, 34) \times (0, 0, 34), (0, 0, 34) \times (0, 34, 0, 68), (0, 0, 34) \times (0, 68, 1), (0, 34, 0, 68) \times (0, 0, 34), (0, 34, 0, 68) \times (0, 68, 1)$. The results for reconstructing f are shown in Fig. 14. The comparison of the error between 3 observations and 9 observations is presented in Table 3. We can see that the numerical results for the case of 9 observations are better than that for the case of 3 observations.

Acknowledgements This work was supported by Vietnam Ministry of Education and Training under grant number B2024-TDV-12.

References

- 1. Borukhov, V.T., Vabishchevich, P.N.: Numerical solution of an inverse problem of source reconstructions in a parabolic equation. Mat. Model. **10**(11), 93–100 (1998). (**Russian**)
- Erdem, A., Lesnic, D., Hasanov, A.: Identification of a spacewise dependent heat source. Appl. Math. Model. 37, 10231–10244 (2013)
- Hào, D.N.: A noncharacteristic Cauchy problem for linear parabolic equations II: a variational method. Numer. Funct. Anal. Optim. 13, 541–564 (1992)
- Hào, D.N.: Methods for Inverse Heat Conduction Problems. Peter Lang Verlag, Frankfurt/Main, Bern, New York, Paris (1998)
- Hào, D.N., Huong, B.V., Oanh, N.T.N., Thanh, P.X.: Determination of a term in the right-hand side of parabolic equations. J. Comput. Appl. Math. 309, 28–43 (2017)
- Hào, D.N., Thành, N.T., Sahli, H.: Splitting-based gradient method for multi-dimensional inverse conduction problems. J. Comput. Appl. Math. 232, 361–377 (2009)
- Hamdi, A.: Identification of a time-varying point source in a system of two coupled linear diffusionadvection-reaction equations: application to surface water pollution. Inverse Prob. 25, 115009 (2009)
- Hasanov, A., Pektaş, B.: A unified approach to identifying an unknown spacewise dependent source in a variable coefficient parabolic equation from final and integral overdeterminations. Appl. Numer. Math. 78, 49–67 (2014)
- 9. Isakov, V.: Inverse Source Problems. Amer. Math. Soc, Providence, RI (1990)
- 10. Isakov, V.: Inverse Problems for Partial Differential Equations, 2nd edn. Springer, New York (2006)
- Ladyzhenskaya, O.A.: The Boundary Value Problems of Mathematical Physics. Springer-Verlag, New York (1985)
- Lavrent'ev, M.M., Maksimov, V.I.: On the reconstruction of the right-hand side of a parabolic equation. Comput. Math. Math. Phys. 48, 641–647 (2008)
- Ling, L.V., Takeuchi, T.: Point sources identification problems for heat equations. Commun. Comput. Phys. 5, 897–913 (2009)
- Ling, L.V., Yamamoto, M., Hon, Y.C., Takeuchi, T.: Identification of source locations in two-dimensional heat equations. Inverse Prob. 22, 1289–1305 (2006)
- 15. Marchuk, G.I.: Methods of Numerical Mathematics. Springer-Verlag, New York (1975)
- Marchuk, G.I.: Splitting and alternating direction methods. In: Ciaglet, P.G., Lions, J.-L. (eds). Handbook of Numerical Mathematics. Volume 1: Finite Difference Methods. Elsevier Science Publisher B.V., North-Holland, Amsterdam (1990)
- Oanh, N.T.N., Huong, B.V.: Determination of a time-dependent term in the right hand side of linear parabolic equations. Acta Math. Vietnam. 41, 313–335 (2016)
- Prilepko, A.I., Tkachenko, D.S.: The Fredholm property and the well-posedness of the inverse source problem with integral overdetermination. Comput. Math. Math. Phys. 43, 1338–1347 (2003)



- Thành, N.T.: Infrared Thermography for the Detection and Characterization of Buried Objects. PhD thesis, Vrije Universiteit Brussel, Brussels, Belgium (2007)
- 20. Trefethen, L.N., Bau, D. III.: Numerical Linear Algebra. Philadelphia: SIAM (1997)
- Tröltzsh, F.: Optimal Control of Partial Differential Equations: Theory. Methods and Applications. Amer. Math. Soc., Providence, Rhode Island (2010)
- Vabishchevich, P.N.: Numerical solution of the problem of the identification of the right-hand side of a parabolic equation. Russian Math. (Iz. VUZ) 47(1), 27–35 (2003)
- 23. Wloka, J.: Partial Differential Equations. Cambridge Univ. Press, Cambridge (1987)
- Yamamoto, M.: Conditional stability in determination of densities of heat sources in a bounded domain. In: Control and Estimation of Distributed Parameter Systems: Nonlinear Phenomena (Vorau, 1993), 359–370, Internat. Ser. Numer. Math. 118, Birkhäuser, Basel (1994)
- 25. Yanenko, N.N.: The Method of Fractional Steps. Springer-Verlag, Berlin, Heidelberg, New York (1971)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Springer Nature or its licensor (e.g. a society or other partner) holds exclusive rights to this article under a publishing agreement with the author(s) or other rightsholder(s); author self-archiving of the accepted manuscript version of this article is solely governed by the terms of such publishing agreement and applicable law.