GENERAL AND APPLIED PHYSICS

Graphene‑Au‑Coated Plasmonic Sensor Based on D‑Shaped Bezier Polygonal Hollow Core Photonic Crystal Fiber

J. Divya1 · S. Selvendran1 · A. Sivanantha Raja² · A. Sivasubramanian1

Received: 16 May 2021 / Accepted: 28 July 2021 / Published online: 4 August 2021 © Sociedade Brasileira de Física 2021

Abstract

A D-shaped photonic crystal fber (PCF)-based plasmonic sensor is proposed for detecting refractive index (RI) variations within the range of 1.33–1.38. The center hollow core has formed using Bezier curves, which confnes light within the core by the principle of photonic band gaps. An analyte is injected into a core to adjust its refractive index. Light propagating in the air core is afected by the analyte flled in the core, causing plasmon waves to be produced on the metal surface. On top of this structure, a chemically stable gold (Au) metal layer is placed as a plasmonic material. Numerical studies are carried out using the fnite element method (FEM). The gold layer thickness is optimized and obtained an average sensitivity of 5600 nm/RIU. The sensor's sensitivity is enhanced by sandwiching graphene between an analyte and a gold flm. The average sensitivity of the sensor increases by 2000 nm/RIU after graphene is incorporated, having reached 7600 nm/RIU according to the results obtained. The sensor can efficiently detect slight variations in the analyte indices and, thus, it is used in various chemical and biosensing applications.

Keywords Photonic crystal fiber · Surface plasmon resonance · Plasmonic sensor · Gold · Graphene

1 Introduction

Photonic crystal fiber (PCF)-based devices have been developed in recent years due to their fexibility and controllability. Unlike optical fber, PCF has a fexible structure and a number of desirable properties, such as endless single mode, high nonlinearity, high birefringence, large mode feld size, ease of flling material, low transmission loss, and controllable dispersion [[1](#page-8-0)[–3](#page-8-1)]. Hollow-core PCF (HCPCF) and non-hollow-core PCF (NHCPCF) are the two most common types of PCF [\[4](#page-8-2), [5\]](#page-8-3). In HCPCF, array of air holes in center act as a core and the silica material is flled in cladding regions with air holes. In NHCPCF, instead of air holes, it has a silica core that is surrounded by air holes [[6](#page-8-4)]. HCPCF use the photonic bandgap (scattering) method for signal propagation, whereas NCHCF uses the index guiding method (total internal refection). In both

 \boxtimes J. Divya divyaselva0605@gmail.com

¹ SENSE, Vellore Institute of Technology, Chennai, Tamil Nadu 600127, India

² Alagappa Chettiar Government College of Engineering and Technology, Tamil Nadu, Karaikudi 630 003, India cases, the difference in refractive index causes the PCF to act as a waveguide, allowing a particular band of light to pass through it [[7](#page-8-5)]. The bands of light propagated can be controlled by adjusting the size and position of the air holes [[8](#page-8-6)].

PCFs are most often used in communications, but their uses have since grown to include sensors [[9](#page-8-7)], polarization flters [[10\]](#page-8-8), multiplexer and demultiplexer [[11\]](#page-8-9), and other felds. A PCF can be used as a sensor by using the surface plasmon resonance (SPR) concept. Metals such as gold (Au), silver (Ag), copper (Cu), aluminum (Al), indium (In), and sodium (Na) are capable of producing surface plasmons (SPs) at the metal and dielectric boundary. These SPs have an oscillating nature and, under certain resonance conditions, have a tendency to propagate at metal surfaces [[12](#page-8-10)]. Because of its higher sensitivity to the refractive index (RI) of the substrate in contact with plasmonic material, SPR techniques are commonly used in many sensing applications. The accuracy and resolution of sensor detection are afected by sensitivity, which is one of the most critical performance parameters of sensors. The PCF-surface plasmon resonance (SPR) sensor's sensitivity can be increased by coating plasmonic material on the inner wall of air holes, which is a difficult task. To address these problems, D-shaped PCF sensors have been proposed; the top circular portion of the fber with air holes has been fully removed and then polished. The polished surface is coated with plasmonic material and a sensing layer. Air holes are placed at regular intervals to guide the light propagation in a specifc direction [[13](#page-8-11)]. A dual-core D-shaped PCF sensor was proposed by Md. Nazmus et al. The gold thickness was optimized, and a wavelength sensitivity of 8000 nm/RIU was obtained [[14](#page-8-12)]. Emranul Haque. et.al., designed a SPR-based PCF sensor. A gold (Au) layer was coated in an open ring placed in top of the D-shaped PCF structure for simple detection. RIs of analyte ranging from 1.18 to 1.36 and 20,000 nm/RIU wavelength sensitivity and 1054 RIU^{-1} maximum amplitude sensitivity were obtained [[15](#page-8-13)].

Rather than increasing the thickness of the metal flm, nanomaterials, such as graphene, are incorporated with the metal surface and decrease the cost or narrow down the spectrum. When in contact with an analyte, efficient charge transfer can cause strong coupling and increased sensitivity at the metal-graphene interface, which can lead to an increase in SPR [[16\]](#page-8-14). As a result, a variety of PCF structures have been proposed. For example, Bin Li et al. proposed an SPR-based liquid refractive index sensor by depositing a controllable thickness gold flm and then coating it with graphene. After the addition of graphene, the sensor's sensitivity increases by 390 nm/RIU, resulting in a fnal sensitivity of 2290 nm/RIU [[17](#page-8-15)]. A graphene-Au-coated

photonic crystal fber (PCF) sensor had been developed by Hongyan Yang et al. Meanwhile, graphene on gold can stably adsorb biomolecules and increase SPR; it can improve sensor sensitivity achieving wavelength sensitivity of 4200 nm/RIU and amplitude sensitivity of 450 RIU−1 [\[18](#page-8-16)].

We propose a plasmonic sensor based on a D-shaped HCPCF with a center core built using Bezier curves that can detect small RI variations in analyte over a range of 1.33–1.38 in this paper. This sensor is used for biosensing applications due to its wide range of RI variations. The fnite element (FEM) approach is used to conduct numerical investigations. Simultaneously, various optimization strategies, such as gold and graphene layer thickness, are calculated. An optimal plasmonic sensor structure is obtained. Coating graphene between an analyte and the gold flm improves the sensor's sensitivity. The average sensitivity of the sensor increased by 2000 nm/ RIU and reached 7600 nm/RIU. Following the introduction chapter, Sect. [2](#page-1-0) summarizes the D-shaped fber design. The results of various optimization techniques are discussing in Sect. [3.](#page-3-0) Finally, in Sect. [4,](#page-8-17) we conclude the outcome and possible progress based on the results.

2 Modeling of D‑Shaped HCPCF Sensor

Figure [1](#page-1-1) shows a schematic diagram of the proposed D-shaped HCPCF sensor. A hexagonal structure with 0.6-µm diameter air holes (d1) arranged in a triangular

Fig. 1 Schematic diagram of the proposed D-shaped hollow core photonic crystal fber sensor

lattice, and the lattice constant (Λ) (distance between the centers of successive air holes) of 0.8 µm. Instead of traditional circles, the middle hollow core is a complicated Bezier polygon structure. Bezier polygon structures are derived using a circular lattice structure with air holes of 0.6-μm diameter, and the lattice constant ($Λ$) of 0.8 μm by keeping an angle of separation of 45° each between them.

The fiber's core is a Bezier polygon structure that works in Dirac mode, a new light trapping technique. The Dirac mode's unique power-law profile allows for loose lateral containment, allowing for long-range interaction between multiple cores. This property leads to the novel coupling and sensing properties for use in optical couplers, laser arrays, and fiber sensors [[19](#page-8-18)]. By filling the air core with the analyte, the hollow core fiber can be used as a sensor.

To construct a D-shape, the top circular portion of the fber with air holes has been fully removed. Gold is used as a plasmonic material because it is chemically stable. The gold layer that has a thickness of 30 nm is placed on the top of the structure. Because of its high surface-to-volume

ratio, broadband optical, and plasmonic properties [[20](#page-8-19)], graphene is coated between the analyte and the gold layer to improve sensing efficiency. Graphene has a thickness of 15 nm. The indices of the analyte put in the core afect light propagation through the air core. The formation of surface plasmon polaritons (SPP) causes light coupled from the core mode to the gold surface. At this phase matching point, the wave vectors of both the light wave and the plasmon wave are the same, so the optimum coupling is observed. The sensor's smaller central air hole of 0.4-µm diameter $(d2)$ is allows the passing wave to penetrate and enter the metal–dielectric interface, which can excite more free electrons. Scattering boundary conditions are established to minimize scattering losses to achieve higher accuracy. The COMSOL Multiphysics software, which is based on the fnite element method (FEM), is used for simulation and performance analysis. The proposed D-shaped PCF sensor can be directly fabricated either by halfway stacking [[21](#page-8-20)] or by fabricated a regular PCF followed by mechanical side polishing methods [[22](#page-8-21)]. The desired gold layer and graphene layer can be deposited using the chemical vapor deposition method [\[23](#page-9-0)]. The analyte can be flled in the center air hole by vacuum pumping technique $[24]$ or infiltration method $[25]$.

The refractive index values for base silica and gold are obtained using the Sellmeier equation [\[26\]](#page-9-3) and Lorentz–Drude model [[27\]](#page-9-4). The complex refractive index equation of graphene is [[28](#page-9-5)]

$$
n_g = \frac{3 + iC_1\lambda}{3}.
$$

where n_g denotes the RI of graphene, λ (μ m) is the wavelength, and $C_1 \approx 5.446 \text{ }\mu\text{m}^{-1}$.

3 Results and Discussion

3.1 Mode Coupling

A change in the RI of the analyte causes a shift in the phase matching point of plasmon mode and core mode. As a

consequence, the resonance wavelength varies depending on the analyte RI, determining the PCF's sensing operation [\[29](#page-9-6)]. The electric feld distributions of x-polarized core mode and y-polarized core mode, respectively, are shown in Fig. [2a](#page-2-0), b, where the light is restricted to the core region. In y-polarization, the coupling efficiency between the SPP mode and the core mode is higher than in x-polarization. Therefore, we consider a y-polarized core mode for study. Figure [2c](#page-2-0) depicts the coupled modes. At the moment, the core mode is phase-matched with the plasmon mode, a resonance peak will appear, and the energy in the core mode will be transferred to the plasmon mode. There exists a frst-order SPP mode, a second-order SPP mode, and a third-order SPP mode in the PCF sensor that are as shown in Fig. [2d](#page-2-0)-f.

Figure [2g](#page-2-0) represents the dispersion relation between the fundamental mode and SPP mode. The real and imaginary values of the effective mode index (n_{eff}) and as a function of wavelength for analyte RI of 1.33 are represented. We found that as the wavelength increases, the real part of the n_{eff} of the core guided mode decreases. At a wavelength of 1160 nm, a slight tilt (S curve) in the real value of n_{eff} indicates that the fber's core mode is in resonance with the plasmon mode. As a result, a portion of the light is coupled into the SPP mode, resulting from loss in fundamental mode. At this point, the imaginary value of n_{eff} reaches its maximum value, which is the resonance wavelength (1160 nm).

Confnement loss [[30\]](#page-9-7) is an important parameter to consider when evaluating the sensor's performance, and is defned as

$$
\text{Confinement Loss} = \frac{40 \times \Pi \times \text{Im}(n_{\text{eff}}) \times 10^4}{\lambda \times \text{ln}(10)} (\text{dB/cm})
$$

where Im (n_{eff}) is the imaginary part of effective mode index.

3.2 Numerical Examination on the Varying Thicknesses of the Gold Layer

The thickness of metal has a significant effect on the efficiency of plasmonic sensors. The gold flm thickness afects the sensor's confnement loss and resonance wavelength. Table [1](#page-3-1) shows the analysis of the diferent thicknesses of

Table 1 Analysis of the sensor for the diferent thicknesses of gold layer for refractive index 1.33 and 1.34

Fig. 3 (**a**) Loss spectrum curve of gold layer thickness (20 nm, 30 nm, and 40 nm) which changes with wavelength. (**b**) Confnement loss factors with respect to wavelength for diferent RI of analyte ranging from 1.33 to 1.38

Fig. 4 (**a**) Loss spectrum curve of Graphene thickness (5 nm, 10 nm, and 15 nm) which changes with wavelength. (**b**) Real and imaginary parts of n_{eff} as a function of wavelength for analyte RI of 1.33 (gold layer along with graphene $(d_{\text{gold}}=30 \text{ nm}$ and $d_{\text{graphene}} = 15 \text{ nm})$. (**c**) Confinement loss factors with respect to wavelength for diferent RI of analyte ranging from 1.33 to 1.38

the gold layer. Figure [3](#page-4-0)a depicts the comparison of the different thicknesses $(10, 20, 30, 40, \text{ and } 50 \text{ nm})$ of gold films. It shows that the confnement loss of PCF decreases as the gold flm thickness increases.

The sensitivity of sensor $[31]$ $[31]$ can be calculated by

$$
S_{\lambda} = \frac{\Delta \lambda_{\text{Peak}}}{\Delta n_a} \text{ (nm/RIU)}.
$$

 $\Delta\lambda_{\text{peak}}$ is the shift in resonance wavelength and Δn_a is the shift in analyte RI value.

From Table [1](#page-3-1), the maximum sensitivity of 6000 nm/RIU is observed for 30-nm gold layer thickness. So, $d_{\text{gold}} = 30$ nm is used for further analysis to improve sensing performance.

Figure [3](#page-4-0)b represents the loss spectrum of the proposed sensor coated with gold for different analyte RIs ranging from 1.33 to 1.38 with a step value of 0.01. As the RI increases, the resonance wavelength shifts towards shorter wavelengths. For the analyte RI value of about 1.34, the observed minimum confnement loss is 502.33 dB/cm at a resonance wavelength of 1100 nm. And the maximum confnement loss of 1006.16 dB/cm has been noticed at a peak resonance wavelength of 880 nm for an RI value of 1.38. These results describe that for the analyte RI value of 1.38, the phase-matching condition has been realized between the core mode to SPP mode at resonance wavelength 880 nm. And the maximum energy transfer takes place between core and SPP mode.

3.3 Numerical Examination on the Varying Thickness of Graphene Layer

As mentioned previously, graphene is layered on the gold film; efficient charge transfer can take place in the goldgraphene interface. The resonance wavelength can be modifed by increasing SPR. And the sensor's sensitivity has increased. The thickness of the graphene with gold flm $(d_{\text{gold}} = 30 \text{ nm})$ affects the sensor's confinement loss and resonance wavelength. Table [2](#page-6-0) represents the analysis of the diferent thickness of graphene layer along with the gold layer (d_{gold} =30 nm). Figure [4](#page-5-0)a shows the comparison of the thickness of 5, 10, 15, 20, and 25 nm graphene. It shows that as the graphene thickness increases, the sensor's confnement loss also increases.

From Table [2,](#page-6-0) the maximum sensitivity of 8000 nm/ RIU is observed for graphene thickness of 5 and 15 nm, but the peak gets broaden in 5-nm thickness. Broadening of loss peak may reduce the detection accuracy. Therefore, d_{graphene} =15 nm is using for better sensing performance.

Figure [4](#page-5-0)b shows the dispersion relation between the fundamental mode and SPP mode. The real and imaginary values of the effective mode index (n_{eff}) and as a function of wavelength for analyte RI of 1.33 are represented. We observed that as the wavelength increases, the real part of the n_{eff} of the core guided mode decreases. At a wavelength of 1340 nm, a deep tilt (S curve) in the real value of n_{eff} indicates that the fber's core mode is in resonance with the plasmon mode. As a result, a portion of the light is coupled into the SPP mode, resulting from loss in fundamental mode. At this point, the imaginary value of n_{eff} reaches its maximum value, which is the resonance wavelength (1340 nm).

Figure [4](#page-5-0)c represents the loss spectrum of the proposed sensor coated with gold and graphene for diferent analyte RIs ranging from 1.33 to 1.38 with a step value of 0.01. As the RI increases, the resonance wavelength shifts towards shorter wavelengths. For the analyte RI value of about 1.38, the observed minimum confnement loss is 421.70 dB/cm at a resonance wavelength of 960 nm. And the maximum confnement loss of 2695.36 dB/cm has been noticed at a peak resonance wavelength of 1340 nm for an RI value of 1.33. These results describe that for the analyte RI value of 1.33, the phase-matching condition has been realized between the core mode to SPP mode at resonance wavelength 1340 nm. And the maximum energy transfer takes place between core and SPP mode.

3.4 Comparative Analysis of With and Without Graphene Layer

The sensor structure was compared with and without graphene coating using the same structural parameters. Figures [2](#page-2-0)g and [4b](#page-5-0) show the real and imaginary values of

Table 2 Analysis of the sensor for the diferent thicknesses of graphene layer along with gold layer thickness of 30 nm for refractive index 1.33 and 1.34

Fig. 5 Resonance wavelength vs analyte RI for sensor sensitivity

measurement

effective mode index (n_{eff}) as a function of wavelength for analyte RI of 1.33 when the sensor is coated with gold $(d_{\text{gold}} = 30 \text{ nm})$ and gold along with graphene $(d_{\text{gold}} = 30 \text{ nm})$ and $d_{\text{graphene}} = 15 \text{ nm}$). A tilt (S curve) in the real value of n_{eff} shows that the core mode of the fber is in resonance with the plasmon mode. A portion of the light is coupled into the SPP mode, and hence, there is a loss in fundamental mode. In Fig. [4b](#page-5-0), the real value of n_{eff} has a long drift than a n_{eff} value in Fig. [2g](#page-2-0), which shows that gold along with graphene has a high resonance wavelength of SPR than the gold alone structure; this leads to enhanced sensitivity.

For further analysis, Figs. [3](#page-4-0)b and [4](#page-5-0)c represent the loss spectrum for diferent analyte RIs ranging from 1.33 to 1.38

with a step value of 0.01 for both structures (with and without graphene). It shows when the sensor has coated with a gold layer, the resonance wavelength ranging from 880 to 1160 nm. Its maximum sensitivity is 6000 nm /RIU, and its average sensitivity is about 5600 nm/RIU. After graphene has layered, the resonance wavelength ranging from 960 to 1340 nm. The maximum sensitivity is 8000 nm/RIU, and its average sensitivity is about 7600 nm/RIU. Compared to the results of sensitivity, the presence of graphene enhances the sensitivity and improves the performance of detection.

Figure [5](#page-7-0) illustrates the variation in sensitivity for the diferent wavelengths. It shows that sensitivity reaches the maximum value of 8000 nm/RIU for the gold layer thickness

Authors	Year	Structure used	Material used	RI detection range (n)	Sensitivity (nm/RIU)
Hongyen Yang et al. [18]	2021	D-shaped PCF sensor	Gold-graphene	$1.32 - 1.41$	4200
Xiangtai Xi et al. [29]	2020	D-shaped SPR biosensor	Gold-graphene	$1.33 - 1.36$	1223
Bin Li et al. $[17]$	2019	D-shaped RI sensor	Gold-graphene	$1.33 - 1.36$	2290
Singh et al. $[27]$	2019	D-shaped sensor	Gold-graphene	$1.32 - 1.35$	5666
Junxia Sun et al. [30]	2019	D-shaped SPR sensor	Gold-graphene	$1.33 - 1.36$	1539
Guowen An et al [28]	2018	D-shaped RI sensor	Gold-graphene	$1.33 - 1.39$	4391
Our proposed work		D-shaped Hollow core PCF sensor	Gold-graphene	$1.33 - 1.38$	For Gold: 5600 For gold and graphene: 7600

Table 3 Performance comparison with previously reported D-shaped plasmonic sensors

of 30 nm coated with a graphene layer thickness of 15 nm. It is due to the enhancement of the confned electric feld onto the top surface of the sensor.

 Performance comparison with previously reported D-shaped plasmonic sensors is detailed in Table [3.](#page-7-1) Our proposed sensor structure yields a sensitivity value of 7600 nm/ RIU, which is much better than the previously reported sensors [\[17](#page-8-15), [18](#page-8-16), [28](#page-9-5), [32](#page-9-9)–[34\]](#page-9-10).

4 Conclusion

In conclusion, we proposed a D-shaped hollow-core PCFbased plasmonic sensor, with the analyte flled in the central core. By changing the analyte's RI, the resonance wavelength has changed. SPR is initially generated by the gold flm, but with an optimized gold layer thickness, sensitivity has improved. In the refractive index range of 1.33 to 1.38, an average sensitivity of 5600 nm/RIU has obtained. After that, we coated a graphene layer between the analyte and the gold flm. The sensor's average wavelength sensitivity increased by 2000 nm/RIU and reaching 7600 nm/ RIU. It is proven that graphene can be used to not only prevent gold from oxidation but also, to improve sensing efficiency. Since the sensor can detect minor changes in analyte indices, it can be used in various chemical and biosensing applications.

References

- 1. J. Yang, L. Zhou, X. Che, J. Huang, X. Li, W. Chen, Photonic crystal fber methane sensor based on modal interference with an ultraviolet curable fuoro-siloxane nano-flm incorporating cryptophane A. Sens. Actuators B Chem. **235**, 717–722 (2016). [https://](https://doi.org/10.1016/j.snb.2016.05.125) doi.org/10.1016/j.snb.2016.05.125
- 2. T. Cheng, Z. Duan, M. Liao, W. Gao, D. Deng, T. Suzuki, Y. Ohishi, A simple all-solid tellurite microstructure optical fber. Opt. Express **21**, 3318–3323 (2013). [https://doi.org/10.1364/](https://doi.org/10.1364/OE.21.003318) [OE.21.003318](https://doi.org/10.1364/OE.21.003318)
- 3. W.C. Wong, C.C. Chan, L.H. Chen, T. Li, K.X. Lee, K.C. Leong, Polyvinyl alcohol coated photonic crystal optical fber sensor for humidity measurement. Sens. Actuators B Chem. **174**, 563–569 (2012).<https://doi.org/10.1016/j.snb.2012.07.032>
- 4. H. Chen et al., A novel polarization splitter based on dual-core photonic crystal fber with a liquid crystal modulation core. IEEE Photon J. **6**(4), 2201109 (2014). [https://doi.org/10.1109/](https://doi.org/10.1109/JPHOT.2014.2337874) [JPHOT.2014.2337874](https://doi.org/10.1109/JPHOT.2014.2337874)
- 5. B. Eggleton, C. Kerbage, P. Westbrook, R. Windeler, A. Hale, Microstructure optical fber devices. Opt. Express **9**, 698–713 (2001).<https://doi.org/10.1364/OE.9.000698>
- 6. X. Zhang, R. Wang, F. Cox, B. Kuhlmey, M. Large, Selective coating of holes in microstructure optical fber and its application to in-fber absorptive polarizers. Opt. Express **15**, 16270– 16278 (2007).<https://doi.org/10.1364/OE.15.016270>
- 7. A. Hassani, M. Skorobogatiy, Design of the microstructure optical fber-based surface plasmon resonance sensors with

D Springer ⊉

enhanced microfuidics. Opt. Express **14**, 11616–11621 (2006). <https://doi.org/10.1364/OE.14.011616>

- 8. H.W. Lee, M.A. Schmidt, R.F. Russell, N.Y. Joly, H.K. Tyagi, P. Uebel, P.S.J. Russell, Pressure assisted melt-flling and optical characterization of Au nano-wires in microstructure fbers. Opt. Express **19**, 12180–12189 (2011). [https://doi.org/10.1364/OE.](https://doi.org/10.1364/OE.19.012180) [19.012180](https://doi.org/10.1364/OE.19.012180)
- 9. S. Selvendran, A. Sivanantharaja, S. Yogalakshmi, A highly sensitive Bezier polygonal hollow core photonic crystal fber biosensor based on surface plasmon resonance. Optik – Int. J. Light and Electron Opt **171**, 109–113 (2018). [https://doi.org/10.](https://doi.org/10.1016/j.ijleo.2018.06.039) [1016/j.ijleo.2018.06.039](https://doi.org/10.1016/j.ijleo.2018.06.039)
- 10. S. Yogalakshmi, S Selvendran and A Sivanantha Raja, Design and analysis of a photonic crystal fber-based polarization flter using surface plasmon resonance. Laser Phys. **26**(2016)
- 11. X. Qiang, K. Li, N. Copner, S. Lin, An ultrashort wavelength multi/demultiplexer via rectangular liquid-infltrated dual-core polymer optical fber. Materials **12**, 1709 (2019). [https://doi.org/](https://doi.org/10.3390/ma12101709) [10.3390/ma12101709](https://doi.org/10.3390/ma12101709)
- 12. U. Fano, The theory of anomalous difraction gratings and of quasi-stationary waves on metallic surfaces. J. Opt. Soc. **31** (3) 213–222.<https://doi.org/10.1364/JOSA.31.000213>
- 13. M. Chen, T. Lang, B. Cao, Y. Yu, C. Shen, D-type optical fber immunoglobulin g sensor based on surface plasmon resonance. Opt. Laser Technol. **131**, 106445 (2020). [https://doi.org/10.](https://doi.org/10.1016/j.optlastec.2020.106445) [1016/j.optlastec.2020.106445](https://doi.org/10.1016/j.optlastec.2020.106445)
- 14. M.N. Sakib, M.B. Hossain, K.F. Al-tabatabaie, I.M. Mehedi, M.T. Hasan, M.A. Hossain, I.S. Amiri, High performance dual core D-shape PCF-SPR sensor modeling employing gold coat. Res. Phys. S2211–3797(19)32981-X, [https://doi.org/10.1016/j.](https://doi.org/10.1016/j.rinp.2019.102788) [rinp.2019.102788](https://doi.org/10.1016/j.rinp.2019.102788)
- 15. E. Haque, M.A. Hossain, F. Ahmed, Y. Namihira, Surface plasmon resonance sensor based on modifed d-shaped photonic crystal fber for wider range of refractive index detection. IEEE Sensors J. <https://doi.org/10.1109/JSEN.2018.2865514>
- 16. B. Song, D. Li, W. Qi, M. Elstner, C. Fan, H. Fang, Graphene on Au (111): A highly conductive material with excellent adsorption properties for high-resolution bio/nano detection and identifcation. Chem. Phys. Chem **11**, 585–589 (2010). [https://doi.](https://doi.org/10.1002/cphc.200900743) [org/10.1002/cphc.200900743](https://doi.org/10.1002/cphc.200900743)
- 17. B. Li, T. Cheng, J. Chen, X. Yan, Graphene-enhanced surface plasmon resonance liquid refractive index sensor based on photonic crystal fber. Sensors **19**, 3666 (2019). [https://doi.org/10.](https://doi.org/10.3390/s19173666) [3390/s19173666](https://doi.org/10.3390/s19173666)
- 18. H. Yang, M. Liu, Y. Chen, L. Guo, G. Xiao, H. Liu, J. Li, L. Yuan, Highly sensitive graphene-Au coated plasmon resonance PCF sensor. Sensors **21**, 818 (2021). [https://doi.org/10.3390/](https://doi.org/10.3390/s21030818) [s21030818](https://doi.org/10.3390/s21030818)
- 19. K. Xie, W. Zhang, A.D. Boardman, H. Jiang, Z. Hu, Y. Liu, M. Xie, Q. Mao, L. Hu, Q. Li, T. Yang, F. Wen, E. Wang, Fiber guiding at the Dirac frequency beyond photonic bandgaps. Light: Sci. Appl. **4** (2015) e304, [https://doi.org/10.1038/lsa.](https://doi.org/10.1038/lsa.2015.77) [2015.77](https://doi.org/10.1038/lsa.2015.77)
- 20. A.B. Kuzmenko, E. van Heumen, F. Carbone, D. van der Marel, Universal optical conductance of graphite. Phys. Rev. Lett. **100**, 117401 (2008). [https://doi.org/10.1103/PhysRevLett.100.](https://doi.org/10.1103/PhysRevLett.100.117401) [117401](https://doi.org/10.1103/PhysRevLett.100.117401)
- 21. N. Zhang, K. Li, Y. Cui, Z. Wu, P.P. Shum, J.L. Auguste, X.Q. Dinh, G. Humbert, L. Wei, Ultra-sensitive chemical and biological analysis via specialty fbers with built-in micro structured optofuidic channels. Lab. Chip **18**(4), 655–661 (2018). [https://](https://doi.org/10.1039/C7LC01247K) doi.org/10.1039/C7LC01247K
- 22. T. Huang, Highly sensitive SPR sensor based on D-shaped photonic crystal fber coated with indium tin oxide at near-infrared wavelength. Plasmonics **12**(3), 583–588 (2017). [https://doi.org/](https://doi.org/10.1007/s11468-016-0301-7) [10.1007/s11468-016-0301-7](https://doi.org/10.1007/s11468-016-0301-7)
- 23. J. Boehm, A. François, H. Ebendorff-Heidepriem, T.M. Monro, Chemical deposition of silver for the fabrication of surface plasmon micro structured optical fbre sensors. Plasmonics **6**(1), 133–136 (2011).<https://doi.org/10.1007/s11468-010-9178-z>
- 24. New design for all-normal near zero dispersion photonic crystal fber with selective liquid infltration for broadband supercontinuum generation at 1.55 μm, Hindawi Publishing Corporation Journal of Photonics, Vol. 2014, Article ID 728592, [https://doi.](https://doi.org/10.1155/2014/728592) [org/10.1155/2014/728592](https://doi.org/10.1155/2014/728592)
- 25. T. Han, Y. Liu, Z. Wang, J. Guo, Z. Wu, S. Wang, Z. Li, W. Zhou, Unique characteristics of a selective-flling photonic crystal fber Sagnac interferometer and its application as high sensitivity sensor. Opt. Express **21**(1), (2013) [https://doi.org/](https://doi.org/10.1364/OE.21.000122) [10.1364/OE.21.000122](https://doi.org/10.1364/OE.21.000122)
- 26. W. Sellmeier, To explain the abnormal color sequence in the spectrum of some substances. Ann. Der. Phys. **219**, 272–282 (1870)
- 27. F. Hao, P. Nordlander, Efficient dielectric function for FDTD simulation of the optical properties of silver and gold nanoparticles. Chem. Phys. Lett. **446**, 115–118 (2007)
- 28. S. Singh, Y.K. Prajapati, Highly sensitive refractive index sensor based on D-shaped PCF with gold-graphene layers on the polished surface. Appl. Phys. A **125**, 437 (2019). [https://doi.org/](https://doi.org/10.1007/s00339-019-2731-5) [10.1007/s00339-019-2731-5](https://doi.org/10.1007/s00339-019-2731-5)
- 29. S. Selvendran, A. Sivanantha Raja, S. Yogalakshmi, A highly sensitive surface plasmon resonance biosensor using photonic crystal fber flled with gold nanowire encircled by silicon lining. Optik S0030–4026(17)31370–0. [https://doi.org/10.1016/j.](https://doi.org/10.1016/j.ijleo.2017.10.157) [ijleo.2017.10.157](https://doi.org/10.1016/j.ijleo.2017.10.157)
- 30. N. Chen, M. Chang, X.D. Zhang, J. Zhou, X.L. Lu, S.L. Zhuang, Highly sensitive plasmonic sensor based on a dual-side polished photonic crystal fber for component content sensing applications. Nanomaterials **9**, 1587 (2019). [https://doi.org/10.3390/](https://doi.org/10.3390/nano9111587) [nano9111587](https://doi.org/10.3390/nano9111587)
- 31. C. Liu, W. Su, Q. Liu, X. Lu, F. Wang, T. Sun, P.K. Chu, Symmetrical dual D-shape photonic crystal fbers for surface plasmon resonance sensing. Opt. Express **26**, 9039–9049 (2018). <https://doi.org/10.1364/OE.26.009039>
- 32. G. An, S. Li, T. Cheng, X. Yan, X. Zhang, X. Zhou, Z. Yuan, Ultra-stable D-shaped optical fber refractive index sensor with graphene-gold deposited platform. Plasmonics (2018). [https://](https://doi.org/10.1007/s11468-018-0788-1) doi.org/10.1007/s11468-018-0788-1
- 33. X. Xi, Xu. Jihua, S. Li, J. Song, W. Yang, Y. Sun, S. Jiang, Y. Han, X. Fan, An Au nanoflm-graphene/D-type fber surface plasmon resonance sensor for highly sensitive specificity bioanalysis. Sensors **20**, 991 (2020). [https://doi.org/10.3390/](https://doi.org/10.3390/s20040991) [s20040991](https://doi.org/10.3390/s20040991)
- 34. J. Sun, S. Jiang, Xu. Jihua, Z. Li, Yu. Chonghui Li, X.Z. Jing, J. Pan, C. Zhang, B. Man, Sensitive and selective SPR sensor employing gold-supported graphene composite flm/D shaped fber for dopamine detection. J. Phys. D: Appl. Phys (2019). <https://doi.org/10.1088/1361-6463/ab08c1>

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.