

Improved Analytical Method to Investigate the Dynamic Characteristics of Composite Box Beam with Corrugated Webs

Yulin Feng1,2 · Lizhong Jiang1,2 · Wangbao Zhou1,[2](http://orcid.org/0000-0003-3338-4971)

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Abstract

This study establishes an improved analytical method (IAM) to investigate the dynamic characteristics of composite box beam with corrugated webs (CBBCW), and the IAM has comprehensively considered the efects of several factors, such as the shear lag, interfacial slip, shear deformation and rotational inertia of CBBCW in combination with the characteristics of CBBCW. Further, based on the Hamilton principle, the vibration diferential equation and boundary conditions for CBBCW have been deduced. Finally, an IAM for calculating the dynamic characteristics of CBBCW was proposed. Based on the IAM developed in this study, the natural frequencies of multiple CBBCW cases with diferent spans, shear connection degrees and boundary conditions have been calculated. The results calculated by the IAM have been compared with those calculated by the fnite element method and by the general beam theory. The comparison verifes the efectiveness of the IAM and obtains some conclusions that are meaningful to engineering design, i.e. the shear lag efect of CBBCW increases with increasing shear connection degree and also increases with increasing order of the vibration mode, the shear lag efect of the CBBCW is up to 6.2% in the frst fve orders of the vibration modes and the efect cannot be ignored. In the frst- and second-order vibration modes of the CBBCW cases, the maximum interface slip efect of CBBCW is 28.42% and therefore cannot be ignored. On the other hand, the shear lag efect of CBBCW is usually lower than those of ordinary composite box beam with the same web thickness.

Keywords Composite box beam with corrugated webs · Shear lag · Analytical method · Interface slip · Dynamic characteristic

List of Symbols

- b_2 Width of the concrete cantilever plate
 $2b_3$ Width of the concrete bottom plate
- Width of the concrete bottom plate
- b_4^* Width of the steel beam's top fange
- *b*4 Height of corrugated steel web
- *t*1 Thickness of the concrete roof
- *t*2 Thickness of the concrete cantilever plate
- *t*₃ Thickness of the steel beam bottom plate
- *t*w Thickness of the corrugated steel web

 \boxtimes Wangbao Zhou zhouwangbao@163.com

- School of Civil Engineering, Central South University, 22 Shaoshan Road, Changsha 410075, Hunan Province, People's Republic of China
- ² National Engineering Laboratory for High Speed Railway Construction, Central South University, 22 Shaoshan Road, Changsha 410075, Hunan Province, People's Republic of China
- Thickness of the steel beam's top fange
- t_4^* *t*4 Width of corrugated steel web
- *h_c* Distance between the centroids of the concrete slab and interface
- h_s Distance between the centroids of the steel beam and interface
- b_w Width of horizontal fold
- b_{w1} Width of inclined fold
- *d*w Width of horizontal projection dimension of the inclined fold
- $\theta_{\rm w}$ Inclination angle of the inclined fold
- *G*_{s1} Corrugated steel web's shear modulus
- *E*_s Steel's elastic modulus
- v_s Poisson's ratio of the steel
- $\theta(x,t)$ Rotation angle of the CBBCW cross-section
- *z*t z-Axis coordinates of the centroids of the concrete slab
- *z*_b z-Axis coordinates of the centroids of the steel beam bottom plate

- $\xi(x, t)$ Centroid's longitudinal displacement difference between the concrete slab and the steel beam
- G_c Concrete slab's shear modulus G_s Steel's shear modulus
- **Steel's shear modulus**
- τ_{xvl} Shear stress of the concrete roof
- τ_{xy2} Shear stress of the concrete cantilever plate
- τ_{xy3} Shear stress of the bottom plate
- γ_{xz} Corrugated steel web's shear strain
- τ_{xz} Corrugated steel web's shear stress
 k_{c1} Interfacial slip stiffness
- Interfacial slip stiffness
- k_1 A single stud's slip stiffness
 V_u A single stud's shear capacit
- A single stud's shear capacity
- n_s The number of studs for each row
- *L* Calculated span of the CBBCW
- *r* Shear connection degree
- *f*s Stud's yield strength
- *l*_s Stud's longitudinal spacing
- ρ_c Density of concrete
- ρ _s Density of steel

1 Introduction

Compared with the concrete web, when a corrugated steel web is subjected to an axial pre-pressure, it can be freely compressed because it is of a folded shape along the axial direction. So, there are less deformation constraints on the shrinkage and creep of concrete roof and cantilever plate. Hence, it is more effective in exerting an external pre-stress on the CBBCW. Further, the shear strength of a corrugated steel web is high enough to completely replace the concrete web, thereby reduces the beam's weight signifcantly and eliminating cracking of the concrete webs. Compared with the ordinary steel webs, the corrugated steel webs have a strong out-of-plane stifness, which can efectively avoid the local buckling. Therefore, the CBBCW is a new structure that is popular in application. Recently, CBBCW has been extensively used in the felds of construction, roads, railways and urban rail transit (Kim et al. [2005](#page-11-0); Lho et al. [2014;](#page-11-1) Nguyen et al. [2013;](#page-12-0) Oh et al. [2012;](#page-12-1) Wang et al. [2013](#page-12-2)).

Since the shear stress in a CBBCW is not evenly distributed along the transverse direction of the concrete slab and the bottom plate, when the CBBCW is subjected to shear stress, the longitudinal displacements of the concrete slab and the bottom plate far away from rib plate lag behind those in the vicinity of the rib plate, resulting in a curved transverse distribution of bending normal forces on the concrete slab and the bottom plate. This is known as the shear lag efect, and cannot be ignored if the concrete slab and the bottom plate are wide. Further, since the shear stud between the steel girder and the concrete slab cannot be absolutely rigid as there is the relative slip between the steel girder and the concrete slab even when they are under the complete connection condition. Therefore, the dynamic characteristics of CBBCW are subjected to the combined efect of the shear lag and the slip (Qi and Jiang [2010](#page-12-3); Nie et al. [2007](#page-12-4); Zhou et al. [2012\)](#page-12-5). Studies have also been reported on the method to analyze the dynamic characteristics considering the efects of several factors, such as the shear lag, interfacial slip, shear deformation and moment of inertia of CBBCW, and a series of representative study methods have emerged as well.

The analytical methods to analyze the dynamic characteristics of CBBCWs. Based on the energy variational method, Zhang et al. ([2008\)](#page-12-6) deduced the formula for the natural vibration frequencies of CBBCW and obtained the analytical solution according to the efects of the shear lag and shear deformation. Based on the energy variational method, Li et al. ([2009](#page-11-2)) studied the shear lag effect of the CBBCW under concentrated load and uniform load and deduced the calculation formulas based on the energy variational method. Based on the energy variational method, Chen et al. [\(2016](#page-11-3)) built the lateral displacement distribution pattern of CBBCW considering the infuence of both the shear deformation and shear lag efects. Using a theoretical analysis method to consider both the shear lag and the shear deformation efect, Qiao ([2013\)](#page-12-7) derived formulas which can be used to calculate the defection of a CBBCW.

The fnite element methods to analyze the dynamic characteristics of CBBCWs. Based a fnite element method, Hu and Chen [\(2009\)](#page-11-4) analyzed the shear lag efect of a curved continuous CBBCW and the factors that affect the shear lag. Based on a fnite element method for a three-span continuous CBBCW, Jiang et al. ([2014](#page-11-5)) carried out a comparative study on the shear lag efect under self-weight. Cheng and Yao ([2016](#page-11-6)) developed a simplifed analysis method for predicting the defections of several CBBCW considering the infuence of both shear deformation and shear lag, and the simplifed method was validated with the results of the fnite element analysis. Chen et al. [\(2017\)](#page-11-7) proposed the sandwich beam theory to predict the fexural vibration behavior of CBBCW considering the presence of diaphragms and external prestressing tendons and the interaction between the web shear deformation and the fange local bending. Wu et al. ([2003](#page-12-8)) studied the shear lag efect of a CBBCW through the loaded test and analyzed it using the 3D fnite element theory.

The experimental methods to analyze the dynamic characteristics of CBBCWs. Seven steel-composite I-girders with corrugated webs to were tested to investigate the shear performance, and developed the analytical and numerical models which were verifed by the experimental results (He et al. [2012a](#page-11-8), [b](#page-11-9), [c,](#page-11-10) [2014\)](#page-11-11). Zhou et al. [\(2016a\)](#page-12-9) conducted experimental and theoretical studies on the deformation of a non-prismatic scaled model with several CBBCW in order to quantitatively study the proportional relationship between the bending deformation and the shear deformation. Elamary et al. [\(2017](#page-11-12)) conducted an experimental program to investigate the effect of a top steel flange on the failure mechanism of a CBBCW under bending, and discussed the efects of the top steel fange on the beam stifness, ultimate load, local buckling of the corrugated web, concrete slip, and failure mechanism of the concrete slab.

To sum up, the previous methods for studying the dynamic characteristics of CBBCW mostly suffer from complex deductions, numerous restrictions, and low calculation efficiency. The analytical calculation method that was characterized by less restrictions and higher calculation efficiency had been rarely studied. Therefore, On the basis of the Hamilton principle, this study has comprehensively considered the efects of several factors, such as the shear lag, interfacial slip, shear deformation and moment of inertia of the CBBCW. Then, an IAM for calculating the natural frequencies of the CBBCW has been developed. Using the IAM, the natural frequencies of two simply-supported (SSD) and fxed-supported (FSD) CBBCW with diferent spans have been calculated. To verify the correctness of the IAM developed in this study, the results calculated by the IAM have been compared with those calculated by the GBT and by the FEM. The effects of the shear lag and interfacial slip on the dynamic characteristics of the CBBCW and OCBBW have also been investigated. The IAM is a development of the earlier calculation theory for the dynamic characteristics of CBBCW. Hence, it can be used as a theoretical foundation for future studies on the dynamic characteristics of CBBCW whose conclusions are relevant to engineering design.

2 Theoretical Analyses

2.1 Geometry Shape and Characteristics of CBBCW

The cross-sectional dimensions and the coordinate system of a CBBCW are shown in Fig. [1.](#page-2-0) The geometrical dimensions of the corrugated steel webs is shown in Fig. [2](#page-3-0). As can be seen from Figs. [1](#page-2-0), [2,](#page-3-0) $2b_1$, b_2 , $2b_3$ and b_4 ^{*} are the widths of the concrete roof, cantilever plate, bottom plate and the width of the steel beam's top flange; b_4 is the height of corrugated steel web; t_1 , t_2 , t_3 , t_w , t_4^* are the thicknesses of the concrete roof, cantilever plate, steel beam bottom plate, corrugated steel web and steel beam's top flange, respectively; t_4 is the width of corrugated steel web; h_c is the distance between the centroids of the concrete slab and interface; h_s is the distance between the centroids of the steel beam and interface, and

Fig. 2 Dimensions of corrupted steel web

 $h = h_c + h_s$; b_w , b_{w1} and d_w are widths of horizontal fold, inclined fold, and horizontal projection dimension of the inclined fold; and θ_w is inclination angle of the inclined fold. According to the characteristics of the CBBCW, the following assumptions are reasonable in order to simplify the calculations.

(1) The shear modulus of the corrugated steel web is related to the geometrical shape of its corrugated web (Johnson et al. [1997](#page-11-13), Samanta and Mukhopadhyay [1999](#page-12-10)). The shear modulus G_{s1} can be approximated as:

$$
G_{s1} = \frac{E_s}{2(1 + v_s)} \frac{b_w + d_w}{b_w + d_w \sec \theta_w}
$$
 (1)

where G_{s1} is the shear modulus of the corrugated steel web; E_s is the steel's elastic modulus; v_s is the Poisson's ratio of the steel.

- (2) The corrugated steel webs are like an accordion. Its axial stifness is very small and the apparent elastic modulus E_e is only a few percent of the elastic modulus of the steel plate (He et al. [2014\)](#page-11-11). So, in the calculations, the axial compression stifness of the corrugated steel webs can be ignored, i.e. $E_e = 0$.
- (3) As shown in Fig. [2](#page-3-0), the mass of a longitudinal unit of corrugated steel web is

$$
m_4 = 2(b_{\rm w1}t_{\rm w} + b_{\rm w}t_{\rm w})b_4\rho_{\rm s}
$$
 (2)

Based on the mass-equivalent principle, the equivalent cross-sectional area of the corrugated steel web is

$$
A_{\text{eq}} = t_{\text{weg}} b_4 = \frac{2(b_{\text{wl}} t_{\text{w}} + b_{\text{w}} t_{\text{w}})}{l_{\text{w}}} b_4
$$
 (3)

2.2 GBT and IAM of CBBCW

(1) In the GBT (Zhou et al. [2016b](#page-12-11)), as shown in Fig. [1](#page-2-0), the longitudinal displacement of each point on the concrete slab and the steel beam are:

$$
u_{\text{xpi}}(x, y, z, t) = -(z - z_{t})\theta(x, t) \quad i = 1, 2
$$

\n
$$
u_{\text{xpi}}(x, y, z, t) = -(z - z_{s})\theta(x, t) \quad i = 3, 4
$$
\n(4)

where $\theta(x, t)$ is the cross-section rotation angle function of the CBBCW; z_t , z_b and z_s are the z-axis coordinates of the centroids of the concrete slab, steel beam bottom plate and steel beam, respectively.

(2) Based on the IAM developed in this study, assuming that the longitudinal displacement of each point on the cross-section of the CBBCW is superimposed by three parts: the longitudinal displacement obeying the planesection assumption, the warping displacement caused by the shear lag and the interfacial slip displacement, the longitudinal displacement of each point on the concrete slab and the steel beam can be expressed as:

$$
u_{x i}(x, y, z, t)
$$

=
$$
\begin{cases} -A_s \xi(x, t) / A_0 - (z - z_t) \theta(x, t) + \psi_i(y) U(x, t) & i = 1, 2 \\ A_c \xi(x, t) / nA_0 - (z - z_s) \theta(x, t) + \psi_i(y) U(x, t) & i = 3, 4 \end{cases}
$$
(5)

where $\xi(x, t)$ is the centroid's longitudinal displacement diference between the concrete slab and the steel beam; $A_c = A_1 + A_2$ is the cross-sectional area of the concrete slab; and $A_s = A_3 + A_4^*$ is the effective cross-sectional area of the steel beam; $\psi_i(y)U(x, t)$, $(i = 1, 2, 3, 4)$ are the longitudinal warping displacement function of the concrete roof, cantilever plate, bottom plate and corrugated

steel web, respectively(Lai et al. 2017); $U(x, t)$ is the longitudinal warping amplitude function; $n = E_s/E_c$, E_c is the elastic modulus of the concrete slab; $\psi_i(y) = \alpha_i y^2 / b_i^2 - \alpha_i + 2(\alpha_1 A_1 / n + \alpha_2 A_2 / n + \alpha_3 A_3) / 3A_0$, $(i = 1, 2, 3, 4)$; $\alpha_1 = 1$, $\alpha_2 = b_2^2/b_1^2$, $\alpha_3 = b_3^2 z_b/b_1^2 z_t$, $\alpha_4 = 0$ $; A_1 = 2b_1t_1, A_2 = 2b_2t_2, A_3 = 2b_3t_3, A_4 = 2b_4t_w,$ $A_4^* = 2b_4^*t_4^*, A_0 = (A_1 + A_2)/n + A_3 + A_4^*.$

3 Governing Equation and Solving Process of CBBCW

3.1 Strain and Stress of Cross‑Sectional of CBBCW

According to Eq. [\(5](#page-3-1)), the strain and stress at each point of the cross-section of the CBBCW can be expressed as:

$$
\varepsilon_{xi} = \begin{cases} k_c \frac{\partial \xi}{\partial x} - (z - z_t) \frac{\partial \theta}{\partial x} + \psi_i \frac{\partial U}{\partial x} & i = 1, 2\\ k_s \frac{\partial \xi}{\partial x} - (z - z_s) \frac{\partial \theta}{\partial x} + \psi_i \frac{\partial U}{\partial x} & i = 3, 4 \end{cases}
$$
(6)

$$
\sigma_{xi} = \begin{cases} E_c \left[k_c \frac{\partial \xi}{\partial x} - (z - z_t) \frac{\partial \theta}{\partial x} + \psi_i \frac{\partial U}{\partial x} \right] & i = 1, 2 \\ E_s \left[k_s \frac{\partial \xi}{\partial x} - (z - z_s) \frac{\partial \theta}{\partial x} + \psi_i \frac{\partial U}{\partial x} \right] & i = 3 \end{cases} \tag{7}
$$

$$
\sigma_{xi} = E_e \left[k_s \frac{\partial \xi}{\partial x} - (z - z_s) \frac{\partial \theta}{\partial x} \right] \quad i = 4
$$
\n(8)

$$
\gamma_{xyi} = \frac{\partial \psi_i}{\partial y} U \quad i = 1, 2, 3 \tag{9}
$$

$$
\tau_{xyi} = G_c \frac{\partial \psi_i}{\partial y} U(x) \quad i = 1, 2 \tag{10}
$$

$$
\tau_{xyi} = G_s \frac{\partial \psi_i}{\partial y} U(x) \qquad i = 3 \tag{11}
$$

$$
\gamma_{xz} = \frac{\partial w}{\partial x} - \theta \tag{12}
$$

$$
\tau_{xz} = G_{s1}(\frac{\partial w}{\partial x} - \theta)
$$
\n(13)

where $\varepsilon_{\rm vi}(i=1,2,3,4)$ are the longitudinal normal strains of the roof, cantilever plate, bottom plate and corrugated steel web, respectively; γ_{xvi} ($i = 1, 2, 3$) are the shear strains of the roof, cantilever plate and bottom plate; $\sigma_{\rm vi}(i = 1, 2, 3, 4)$ are the longitudinal normal stresses of the concrete roof, cantilever plate, bottom plate and corrugated steel web, respectively; and $w(x, t)$ is the vertical deflection of the CBBCW; G_c is the shear modulus of the concrete slab; G_s is the steel's shear modulus; $\tau_{\text{xyi}}(i=1,2)$ are the shear stress of the concrete roof and the cantilever plate, respectively; $\tau_{\text{xvi}}(i=3)$ is the shear stress of the bottom plate; γ_{xz} is the shear strain of the corrugated steel web; τ_{xz} is the shear stress of the corrugated steel web.

Since the vertical compression of the concrete slab and the steel beam, and the transverse strains of the plates are very small, they can be ignored (Chen and Wang [2012](#page-11-15); Kashefi et al. [2017](#page-11-16); Ng and Ronagh [2004](#page-12-12)), and the relative interfacial longitudinal displacement $\zeta(x, t)$ can be expressed as:

$$
\zeta(x,t) = \xi + h_c \theta + h_s \theta = \xi + h\theta \tag{14}
$$

According to Eq. (14) (14) (14) , the interfacial shear stress can be expressed as (Nie et al. [2005,](#page-12-13) [2007\)](#page-12-4):

$$
\zeta(x,t) = k_{\rm sl}\zeta(x,t) = k_{\rm sl}(\xi + h\theta)
$$
\n(15)

$$
k_{\rm sl} = k_1 n_{\rm s} / l_{\rm s}, \quad k_1 = 0.66 \, V_{\rm u} \tag{16}
$$

where k_{sl} is the interfacial slip stiffness; k_1 is the slip stiffness of a single stud; $V_u = A_s f_s r l_s / (L n_s)$ is the shear capacity of a single stud; n_s is the number of studs for each row; *L* is the calculated span of the CBBCW; *r* is the shear connection degree; f_s is the stud's yield strength; and l_s is the stud's longitudinal spacing.

3.2 Strain Energy and Kinetic Energy of CBBCW

(1) The strain energy of the CBBCW can be expressed as

$$
V = \frac{1}{2} \int_{L} \left[\sum_{i=1}^{4} \int_{A_{i}} (\sigma_{xi} \varepsilon_{xi} + \tau_{xyi} r_{xyi} + \tau_{xzi} r_{xzi}) dA + \zeta \zeta \right] dx
$$

Substituting Eqs. (6)–(15) into Eq. (17) gives: (17)

$$
V = \frac{1}{2} \int_{L} \left\{ \left(E_{c} k_{c}^{2} A_{c} + E_{s} k_{s}^{2} A_{4*} + E_{s} k_{s}^{2} A_{3} \right) \xi^{\prime 2} - 2 \left[\int_{A_{c}} k_{c} E_{c} (z - z_{t}) dA + \int_{A_{s}} k_{s} E_{s} (z - z_{s}) dA \right] \xi^{\prime} \theta^{\prime} \right.+ \left[\int_{A_{c}} E_{c} (z - z_{t})^{2} dA + \int_{A_{s}} E_{s} (z - z_{s})^{2} dA \right] \theta^{\prime 2} + \left(\int_{A_{c}} E_{c} \psi_{i}^{2} dA + \int_{A_{s}} E_{s} \psi_{i}^{2} dA \right) U^{\prime 2} - 2 \left[\int_{A_{c}} E_{c} (z - z_{t}) \psi_{i} dA + \int_{A_{s}} E_{s} (z - z_{s}) \psi_{i} dA \right] \theta^{\prime} U^{\prime} + 2 \left(\int_{A_{c}} E_{c} k_{c} \psi_{i} dA + \int_{A_{s}} E_{s} k_{s} \psi_{i} dA \right) \xi^{\prime} U^{\prime} \right.+ \left(\int_{A_{c}} G_{c} \psi_{i}^{\prime 2} dA + \int_{A_{3}} G_{s} \psi_{i}^{\prime 2} dA \right) U^{2} + k_{s1} \xi^{2} + G_{s1} A_{4} (w^{\prime} - \theta)^{2} \right\} dx
$$
\n(18)

(2) The kinetic energy of the CBBCW can be expressed as

$T = \frac{1}{2} \int \left(m\dot{w}^2 + \sum \int \rho_i \dot{u}_i^2 dA \right) dx$ (19) $\overline{2}$ \int_L $\sqrt{2}$ $m\dot{w}^2 + \sum_{i=1}^{4}$ $\sum_{i=1}$ \int_{A_i} $\rho_i \dot{u}_i^2 dA$ λ d*x*

where $m = \rho_c A_c + \rho_s (A_s + A_{eq})$; $\rho_1 = \rho_2 = \rho_c$; $\rho_3 = \rho_4 = \rho_4^* = \rho_s$; ρ_c is the density of concrete; and ρ_s is the density of steel.

Substituting Eq. [\(5](#page-3-1)) into Eq. ([19\)](#page-5-0) gives

3.3 Solving the Equations of CBBCW

According to the Hamilton principle,

$$
\delta \int_{t_0}^{t_1} (T - V) dt = 0 \tag{21}
$$

The governing differential equations and the boundary conditions for the bending vibration of a CBBCW are as follows:

$$
T = \frac{1}{2} \int_{L} \left\{ m \dot{w}^{2} + (\rho_{c} k_{c}^{2} A_{c} + \rho_{s} k_{s}^{2} A_{s} + \rho_{s} k_{s}^{2} A_{eq}) \dot{\xi}^{2} + \left[\int_{A_{c}} \rho_{c} (z - z_{t})^{2} dA + \int_{A_{s} + A_{eq}} \rho_{s} (z - z_{s})^{2} dA \right] \dot{\theta}^{2} + \left(\int_{A_{c}} \rho_{c} \psi^{2} dA + \int_{A_{s} + A_{eq}} \rho_{s} \psi^{2} dA \right) \dot{U}^{2} - 2 \left[\int_{A_{c}} \rho_{c} k_{c} (z - z_{t}) dA + \int_{A_{s} + A_{eq}} \rho_{s} k_{s} (z - z_{s}) dA \right] \dot{\xi} \dot{\theta} \tag{20}
$$

$$
-2 \left[\int_{A_{c}} \rho_{c} (z - z_{t}) \psi dA + \int_{A_{s} + A_{eq}} \rho_{s} (z - z_{s}) \psi dA \right] \dot{\theta} \dot{U} + 2 \left(\int_{A_{c}} \rho_{c} k_{c} \psi dA + \int_{A_{s} + A_{eq}} \rho_{s} k_{s} \psi dA \right) \dot{\xi} \dot{U} \right\} dx
$$

$$
\left(\int_{A_c} E_c \psi_i^2 dA + \int_{A_s} E_s \psi_i^2 dA\right) U'' - \left[\int_{A_c} E_c (z - z_t) \psi_i dA + \int_{A_s} E_s (z - z_s) \psi_i dA\right] \theta'' + \left[\int_{A_c} E_c k_c \psi_i dA + \int_{A_s} E_s k_s \psi_i dA\right] \xi''
$$

$$
-\left(\int_{A_c} G_c \psi_i'^2 dA + \int_{A_3} G_s \psi_i'^2 dA\right) U - \left(\int_{A_c} \rho_c \psi^2 dA + \int_{A_s + A_{eq}} \rho_s \psi^2 dA\right) U - \left(\int_{A_c} \rho_c k_c \psi dA + \int_{A_s + A_{eq}} \rho_s k_s \psi dA_1\right) \xi
$$

$$
+\left(\int_{A_c} \rho_c (z - z_t) \psi dA + \int_{A_s + A_{eq}} \rho_s (z - z_s) \psi dA\right) \dot{\theta} = 0
$$
\n(22)

$$
(E_{c}k_{c}^{2}A_{c} + E_{s}k_{s}^{2}A_{4*} + E_{s}k_{s}^{2}A_{3})\xi'' - \left[\int_{A_{c}} k_{c}E_{c}(z - z_{t})dA + \int_{A_{s}} k_{s}E_{s}(z - z_{s})dA\right]\theta'' + \left(\int_{A_{c}} E_{c}k_{c}\psi_{i}dA + \int_{A_{s}} E_{s}k_{s}\psi_{i}dA\right)U'' - k_{s}\xi - \left(\rho_{c}k_{c}^{2}A_{c} + \rho_{s}k_{s}^{2}A_{s} + \rho_{s}k_{s}^{2}A_{eq}\right)\xi + \left[\int_{A_{c}} \rho_{c}k_{c}(z - z_{t})dA + \int_{A_{s}+A_{eq}} \rho_{s}k_{s}(z - z_{s})dA\right]\ddot{\theta}
$$

$$
-\left(\int_{A_{c}} \rho_{c}k_{c}\psi dA + \int_{A_{s}+A_{eq}} \rho_{s}k_{s}\psi dA\right)\ddot{U} = 0
$$
(23)

² Springer

$$
G_{s1}A_{4}(w'-\theta) - \left[\int_{A_{c}} k_{c}E_{c}(z-z_{t})dA + \int_{A_{s}} k_{s}E_{s}(z-z_{s})dA\right] \xi'' + \left[\int_{A_{c}} E_{c}(z-z_{t})^{2}dA + \int_{A_{s}} E_{s}(z-z_{s})^{2}dA\right] \theta''
$$

$$
- \left[\int_{A_{c}} E_{c}(z-z_{t})\psi_{i}dA + \int_{A_{s}} E_{s}(z-z_{s})\psi_{i}dA\right]U'' - k_{s1}\xi h - \left[\int_{A_{c}} \rho_{c}(z-z_{t})^{2}dA + \int_{A_{s}+A_{eq}} \rho_{s}(z-z_{s})^{2}dA\right] \ddot{\theta}
$$

$$
+ \left[\int_{A_{c}} \rho_{c}k_{c}(z-z_{t})dA + \int_{A_{s}+A_{eq}} \rho_{s}k_{s}(z-z_{s})dA\right] \ddot{\xi} + \left[\int_{A_{c}} \rho_{c}(z-z_{t})\psi dA + \int_{A_{s}+A_{eq}} \rho_{s}(z-z_{s})\psi dA\right] \ddot{U} = 0
$$
(24)

$$
G_{s1}A_4(w'' - \theta') - m\ddot{w} = 0\tag{25}
$$

$$
(25)
$$

$$
\left\{\left(\int_{A_c} E_c \psi_i^2 dA + \int_{A_s} E_s \psi_i^2 dA\right) U' \delta U - \left[\int_{A_c} E_c (z - z_t) \psi_i dA + \int_{A_s} E_s (z - z_s) \psi_i dA\right] \theta' \delta U \right\} + \left(\int_{A_c} E_c k_c \psi_i dA + \int_{A_s} E_s k_s \psi_i dA\right) \xi'
$$
\n(26)

$$
\begin{cases}\n\left(E_{c}k_{c}^{2}A_{c} + E_{s}k_{s}^{2}A_{4*} + E_{s}k_{s}^{2}A_{3}\right)\xi' - \left[\int_{A_{c}} k_{c}E_{c}(z - z_{t})dA + \int_{A_{s}} k_{s}E_{s}(z - z_{s})dA\right]\theta'\right] & \left.\int_{\delta\xi} \xi\right|_{0}^{L} = 0 \\
+ \left(\int_{A_{c}} E_{c}k_{c}\psi_{i}dA + \int_{A_{s}} E_{s}k_{s}\psi_{i}dA\right)U'\n\end{cases}
$$
\n(27)

$$
\left\{\left[\int_{A_c} E_c(z-z_t)^2 dA + \int_{A_s} E_s(z-z_s)^2 dA\right] \theta' - \left[\int_{A_c} k_c E_c(z-z_t) dA + \int_{A_s} k_s E_s(z-z_s) dA\right] \xi'\right\} \delta \theta \Big|_{0}^{L} = 0
$$
\n(28)\n
$$
-\left[\int_{A_c} E_c(z-z_t) \psi_i dA + \int_{A_s} E_s(z-z_s) \psi_i dA\right] U'
$$

$$
G_{s1}A_4(w'-\theta)\delta w\Big|_0^L = 0\tag{29}
$$

Let:

$$
U(x,t) = U_1(x)\Gamma(t), \quad \xi(x,t) = \xi_1(x)\Gamma(t), \quad w(x,t) = w_1(x)\Gamma(t), \quad \theta(x,t) = \theta_1(x)\Gamma(t)
$$
\n(30)

where
$$
\Gamma(t) = \sin(\omega t + \varphi)
$$
.
\n
$$
\lambda^k = \partial^k / \partial x^k
$$
\n(31)

Substituting Eqs. (30) (30) – (31) (31) into Eqs. (22) (22) – (25) (25) gives

$$
F_{11}U_1 + F_{12}\xi_1 + F_{13}w_1 + F_{14}\theta_1 = 0
$$
\n(32)

$$
F_{21}U_1 + F_{22}\xi_1 + F_{23} + F_{24}\theta_1 = 0
$$
\n(33)

$$
F_{31}U_1 + F_{32}\xi_1 + F_{33}w_1 - F_{34}\theta_1 = 0
$$
\n(34)

$$
F_{41}U_1 + F_{42}\xi_1 + F_{43}w_1 + F_{44}\theta_1 = 0
$$
\n(35)

The characteristic Eqs. (32) – (35) (35) corresponding to the set of diferential equations is

$$
\begin{vmatrix}\nF_{11} & F_{12} & F_{13} & F_{14} \\
F_{21} & F_{22} & F_{23} & F_{24} \\
F_{31} & F_{32} & F_{33} & F_{34} \\
F_{41} & F_{42} & F_{43} & F_{44}\n\end{vmatrix} = 0
$$
\n(46)

$$
F_{11} = \left(\int_{A_c} E_c \psi_i^2 dA + \int_{A_s} E_s \psi_i^2 dA\right) \lambda^2 - \int_{A_c} G_c \psi_i'^2 dA + \int_{A_3} G_s \psi_i'^2 dA + \left(\int_{A_c} \rho_c \psi^2 dA + \int_{A_s + A_{eq}} \rho_s \psi^2 dA\right) \omega^2\tag{36}
$$

$$
F_{21} = \left(\int_{A_c} E_c k_c \psi_i dA + \int_{A_s} E_s k_s \psi_i dA\right) \lambda^2 + \left(\int_{A_c} \rho_c k_c \psi dA + \int_{A_s + A_{eq}} \rho_s k_s \psi dA\right) \omega^2\tag{37}
$$

$$
F_{41} = -\left(\int_{A_c} E_c(z - z_t) \psi_i dA + \int_{A_s} E_s(z - z_s) \psi_i dA\right) \lambda^2 - \left(\int_{A_c} \rho_c(z - z_t) \psi dA + \int_{A_s + A_{eq}} \rho_s(z - z_s) \psi dA\right) \omega^2\tag{38}
$$

$$
F_{12} = \left(\int_{A_c} E_c k_c \psi_i dA + \int_{A_s} E_s k_s \psi_i dA\right) \lambda^2 + \left(\int_{A_c} \rho_c k_c \psi dA + \int_{A_s + A_{eq}} \rho_s k_s \psi dA\right) \omega^2\tag{39}
$$

$$
F_{22} = (E_{c}k_{c}^{2}A_{c} + E_{s}k_{s}^{2}A_{4*} + E_{s}k_{s}^{2}A_{3})\lambda^{2} - k_{sl} + (\rho_{c}k_{c}^{2}A_{c} + \rho_{s}k_{s}^{2}A_{s} + \rho_{s}k_{s}^{2}A_{eq})\omega^{2}
$$
\n(40)

$$
F_{42} = -\left[\int_{A_c} k_c E_c (z - z_t) dA + \int_{A_s} k_s E_s (z - z_s) dA\right] \lambda^2 - k_{\rm sl} h - \left[\int_{A_c} \rho_c k_c (z - z_t) dA + \int_{A_s + A_{\rm eq}} \rho_s k_s (z - z_s) dA\right] \omega^2 \tag{41}
$$

$$
F_{33} = G_{s1}A_4\lambda^2 + m\omega^2, \quad F_{43} = G_{s1}A_4\lambda, \quad F_{34} = -G_{s1}A_4\lambda, \quad F_{13} = F_{23} = F_{31} = F_{32} = 0 \tag{42}
$$

$$
F_{14} = -\left(\int_{A_c} \rho_c (z - z_t) \psi dA + \int_{A_s + A_{eq}} \rho_s (z - z_s) \psi dA\right) \omega^2 - \left(\int_{A_c} E_c (z - z_t) \psi_i dA + \int_{A_s} E_s (z - z_s) \psi_i dA\right) \lambda^2\tag{43}
$$

$$
F_{24} = -k_{\rm sl}h - \left(\int_{A_{\rm c}} \rho_{\rm c}k_{\rm c}(z-z_{\rm t})\mathrm{d}A + \int_{A_{\rm s}+A_{\rm eq}} \rho_{\rm s}k_{\rm s}(z-z_{\rm s})\mathrm{d}A\right)\omega^2 - \left(\int_{A_{\rm c}} k_{\rm c}E_{\rm c}(z-z_{\rm t})\mathrm{d}A + \int_{A_{\rm s}} k_{\rm s}E_{\rm s}(z-z_{\rm s})\mathrm{d}A\right)\lambda^2\tag{44}
$$

$$
F_{44} = -G_{s1}A_4 + \left[\int_{A_c} E_c (z - z_t)^2 dA + \int_{A_s} E_s (z - z_s)^2 dA \right] \lambda^2 - k_{\rm sl} h^2 + \left[\int_{A_c} \rho_c (z - z_t)^2 dA + \int_{A_s + A_{\rm eq}} \rho_s (z - z_s)^2 dA \right] \omega^2 \tag{45}
$$

Letting λ_i ($i = 1, 2, ..., 8$) be the characteristic root of the characteristic Eq. (46) (46) , the solutions to the set of differential Eqs. (32) (32) – (35) can be expressed as:

 $|B(\omega)| = 0$. Solving $|B(\omega)| = 0$ can obtain the natural vibration frequency of a CBBCW.

$$
U_1 = \sum_{i=1}^{8} a_i C_{1i} \exp(\lambda_i x), \quad \xi_1 = \sum_{i=1}^{8} a_i C_{2i} \exp(\lambda_i x), \quad w_1 = \sum_{i=1}^{8} a_i C_{3i} \exp(\lambda_i x), \quad \theta_1 = \sum_{i=1}^{8} a_i C_{4i} \exp(\lambda_i x)
$$
(47)

where a_i is the integral constant; $C_i = \{C_{1i}, C_{2i}, C_{3i}, C_{4i}\}^T$ are the characteristic vector corresponding to the characteristic root λ_i .

According to Eqs. (26) – (29) (29) (29) , the usual boundary conditions of a CBBCW can be expressed as: The natural boundary conditions of the CBBCW with SSD are $U'_1 = 0$, $\theta'_1 = 0$, $\xi'_1 = 0$, $w_1 = 0$; The natural boundary conditions of the CBBCW with FSD are $U_1 = 0$, $\theta_1 = 0$, $\xi_1 = 0$, $w_1 = 0$. According to the boundary conditions of the SSD and FFD CBBCW, each of the two ends of the CBBCW has four boundary conditions. Substituting Eq. ([30\)](#page-6-0) into the boundary conditions and giving the eight equations with integral constants, the matrix expressions of the eight equations can be expressed as (Xi) and (2017) (2017) (2017) $[B(\omega)]\{a\} = 0$, ${a}$ } = { ${a_1, a_2, ..., a_8}$ }^T, [$B(\omega)$]{ a } = 0 is non-zero only if

4 Verifcation of IAM

In order to verify the efectiveness of the IAM developed in the study, two simple-supported CBBCW and two fxedsupported CBBCW with diferent spans have been used as cases for analyses. The natural frequencies calculated by the GBT, IAM and FEM. The calculation results are shown in Tables [1](#page-8-0), [2](#page-8-1).

The fnite element software ANSYS has been used to simulate the CBBCW cases, in which the concrete slab has been modeled by SOLID65 element, the steel beam has been modeled by SHELL43 element, and the stud has been modeled by COMBIN14 element, where the stud's longitudinal slip stifness, namely the spring stifness of COMBIN14, has been calculated according to Eq. [\(16](#page-4-4)). The vertical interface interactions between the concrete slab and the steel beam

Table 2 Calculation result comparison for the natural frequencies of CBBCW $(L=12 \text{ m})$

Table 1 Calculation result comparison for the natural frequencies of CBBCW

 $(L=10 \text{ m})$

were simulated by coupling the free degrees of the nodes in the vertical direction at the same position, i.e., the vertical separation between the concrete slab and the steel beam was ignored (Deng et al. [2017](#page-11-17); Jiang et al. [2018;](#page-11-18) Xiang et al. [2017](#page-12-14)). The boundary conditions at the ends of the FE models were simulated by constraining the free degrees in both the vertical and transverse directions for the simply supported end, and by constraining the free degrees in the vertical, transverse and longitudinal directions for the fxed supported end (Zhou et al. [2016c](#page-12-15)).

The mechanical parameters and geometric parameters of the CBBCW cases are as follows: $r = 0.25$, $r = 2.00$, $\rho_s = 7.87 \times 10^{-9} \text{ t mm}^{-3}$, $\rho_c = 2.570 \times 10^{-9} \text{ t mm}^{-3}$, $\mu_s = 0.30$, $\mu_c = 0.18$, $b_1 = 800$ mm, $b_3 = 900$ mm, $b₂ = 400$ mm, $b₄ = 400$ mm, $b₅ = 200$ mm, $t₁ = 120$ mm, $t_2 = 120$ mm, $t_3 = 25$ mm, $t_5 = 25$ mm, $t_4 = 12$ mm, $E_s = 2.0 \times 10^5 \text{ MPa}$, $E_c = 4.5 \times 10^4 \text{ MPa}$. E_a denotes the error between the calculation results of the IAM and FEM for natural frequencies of CBBCW.

From Tables [1](#page-8-0) and [2,](#page-8-1) it can be seen that, the IAM and FEM have been used to calculate the first five-order natural vibration frequencies of the four CBBCW cases. The calculation errors between the two methods are less than 4%. This is an indication that with the comprehensive consideration of the shear deformation, interfacial slip and shear lag, the results by IAM developed in this study are in good agreement with those by the FEM. Therefore, the correctness and reasonability of the IAM are validated. However, in the low natural frequency, the results by GBT are basically the same as those by the FEM and IAM, the calculation errors by GBT are signifcantly larger in the high natural frequency. The accuracy of IAM is further verifed.

5 Analysis of Efect Factors

5.1 The Efect of Interface Slip on Natural Frequencies of CBBCW

Table [3](#page-9-0) shows the effect of interface slip on natural frequencies of the CBBCW of SSD and FSD with the *r* = 0.25 and

 $r = 2.00$, respectively. Where, C_s denotes the effect of interface slip.

As shown in Table [3,](#page-9-0) in the frst-order and second-order vibration modes of the CBBCW cases, the maximum interface slip of the CBBCW is 28.42%. This is an indication that the interface slip corresponding to the low-order vibration mode of the CBBCW is signifcant and the increase in the shear connection degree can signifcantly increase the CBBCW's low-order natural frequency. The interface slip corresponding to the high-order vibration mode of the CBBCW decreases signifcantly with the increasing order of the vibration mode. This is an indication that the efect of the shear connection degree on the high-order vibration mode of the CBBCW has been reduced and the proportion of defection caused by the shear deformation is larger on the high-order vibration curve of the CBBCW.

5.2 The Efect of Shear Lag on Natural Frequencies of CBBCW

Figure [3](#page-10-0) shows the effect of shear lag on natural frequencies of the CBBCW of SSD and FSD with $r = 0.25$ and $r = 2.00$ considered shear lag and not considered shear lag, respectively. C_1 denotes the effect of shear lag.

As shown in Fig. [3](#page-10-0), under the condition with diferent shear connection degrees, the shear lag effect of the CBBCW is signifcant. The shear lag efect of the CBBCW increases with increasing shear connection degree. This is an indication that the longitudinal nonlinear deformation efect of the concrete slab and bottom plate of the CBBCW increase signifcantly with increasing shear connection degree. The shear lag efect of the CBBCW increases with the order of the vibration mode. For the four CBBCW cases, the shear lag effect of the CBBCW is up to 6.2% in the first five orders of the vibration modes. So, the shear lag efect has a certain efect on the high-order vibration mode of the CBBCW.

Fig. 3 Relationship between the efect of shear lag and the mode orders of natural frequencies

Fig. 4 Comparison of CBBCW and OCBBW in terms of shear lag

5.3 Comparison of CBBCW and OCBBW with Equal Web Thickness in Terms of Shear Lag

Figure [4](#page-10-1) shows the comparison of CBBCW and OCBBW with equal web thickness in terms of shear lag under $r = 0.25$ and $r = 2.00$, respectively. C_o denotes the effect of shear lag.

From Fig. [4,](#page-10-1) it can be seen that, except the frst order natural vibration mode, the shear lag efect of the OCBBW are larger than those of the CBBCW, and the relative diferences in the shear lag efect between the CBBCW and the OCBBW increases signifcantly with rising natural vibration mode order.

6 Discussion

Compared with the FEM, the IAM can incarnate the key influence factors and their influence laws affecting the dynamic characteristics of CBBCW more intuitively, saves a lot of time in modeling, model debugging and operating, and provide a theoretical basis for deriving a practical formula for engineering calculation and make up the defciency of numerical simulation analysis.

In order to study the calculation efficiency of IAM, take the cases with spans of 10 m and 12 m (FSD, $r = 0.25$) as examples. The calculation efficiency is expressed by $E_c = (t_{FEM} - t_{IAM})/t_{IAM} \times 100\%$. When the calculated span of CBBCW is 10 m, FEM takes 199 s to calculate the dynamic characteristics, IAM takes 159 s, and the calculation efficiency is improved by 25.2% ; When the calculated span of CBBCW is 12 m, FEM takes 284 s to calculate the

dynamic characteristics, while IAM only takes 160 s, the computational efficiency was improved by 78.6% . It can be found that IAM has a higher calculation efficiency, and with the increase of geometric size of CBBCW, the advantages are more signifcant.

7 Conclusions

With the comprehensive consideration of the characteristics of the CBBCW and factors such as shear lag, interfacial slip, shear deformation and moment of inertia, based on the Hamilton principle, an IAM for calculating the natural frequencies of a CBBCW has been developed. Using the IAM, the natural frequencies of multiple CBBCW cases with different spans, shear connection degrees and boundary conditions have been calculated. The results calculated by the IAM have been compared with those calculated by GBT and by FEM. The conclusions are as follows:

- (1) The results calculated by the IAM are in good agreements with those by the FEM, which validates the correctness of the IAM. The IAM provides a theoretical basis for the further study and application of CBBCW dynamic characteristics.
- (2) The interface slip for low-order vibration mode of CBBCW is signifcant, and the increase in the shear connection degree can signifcantly increase the loworder natural frequency of CBBCW.
- (3) The shear connection degree has little efect on the high-order natural vibration frequencies of CBBCW.
- (4) The shear lag efect of CBBCW increases with increasing shear connection degree, and increases with rising natural vibration mode order.
- (5) The relative diferences in the shear lag efect between the CBBCW and the OCBBW increases signifcantly with rising natural vibration mode order.

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