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Numerical solutions of the fractal foam drainage equation

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Abstract

The fow of liquid relative to the bubbles is called drainage. This paper presents a study of the numerical solution of a non-linear foam drainage equation with timefractional derivative. We use the two-scale approach which is formulated by combining the fractional complex transform and the homotopy perturbation method (HPM). With the aid of the fractional complex transform, frst, we transform the problem into its diferential partner and then HPM is applied to obtain the He's polynomials which are highly and powerful support for non-linear problems. Further, we put forward the theory of the two-scale approach which reveals the sketch between fractional complex transform and the solution of non-linear foam drainage equation. The signifcant results illustrate that this approach does not require any assumption while it reduces the heavy calculation without any restrictive variable. This approach also sheds a bright light on practical applications of fractal calculus.

Keywords Two-scale method · Fractional derivative · Fractal foam drainage equation · He's polynomials

Mathematics Subject Classifcation 35R11 · 35A22

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1 Introduction

Underground fuid fow in fractured and porous media exhibits a variety of hydromechanical phenomena, which have been intriguing the research community for several decades (Schmidt and Steeb [2019](#page-9-0)). Foam drainage equation is observed one of the most important partial diferential equation which in everyday of activities, both natural and industrial. The study of foam drainage equation is a simple model which leads to the fow of liquid through channels and nodes between the bubbles. Consider the non-linear foam drainage equation with time-fractional derivative such as

$$
D_t^{\alpha} \Omega + \frac{\partial}{\partial x} \left(\Omega^2 - \frac{\sqrt{\Omega}}{2} \frac{\partial \Omega}{\partial x} \right) = 0, \tag{1}
$$

with initial condition

$$
\Omega(x,0) = f(x),\tag{2}
$$

 Ω is the cross-section of a channel formed where three films meet which usually indicated as *Plateau border*, *x* denotes the scaled position and t be the time respectively. α is fractal dimension and $\frac{\partial^{\alpha}}{\partial x^{\alpha}}$ is He's fractional derivative (He [2018](#page-8-0), [2020a,](#page-8-1) [b;](#page-8-2) Ain et al. [2020](#page-9-1); Shen and He 2020).

Recent research in foams and emulsions has focused on three topics often treated individually but are, in fact, interdependent: drainage, coarsening, and rheology; see Fig. [1.](#page-1-0) For $\alpha = 1$, Eq. [\(1](#page-1-1)) can be reduced to classical foam drainage equation which has wide applications in particular care of commodities such as creams, oil, lotions, scrubbing and clothes cleaning (Stone et al. [2002](#page-9-2)), chemical industries, mineral processing, and structural material sciences (Hilgenfeldt et al. [2001](#page-8-3)), aluminum metals (Schultz et al. [2000](#page-9-3)), thin porous layer (Koursari et al. [2019\)](#page-8-4). During the manufacturing of foam, the substance is in the liquid state, and fuid can change round while the bubble arrangement remains relatively unchanged. Alam ([2015](#page-7-1)) applied *G*� ∕*G*-expansion method to obtain the exact

solution of the foam drainage equation. Parand and Delkhosh ([2018\)](#page-9-4) used the collocation method to fnd the semi-analytical solution of the nonlinear foam drainage equation. Islam and Akbar ([2018](#page-8-5)) used *G*� ∕*G*- expansion method with the assistance of the fractional complex transformation to fnd the wave solution of the space–time fractional foam drainage equation. Various approaches have been established to identify the approximate and analytical solution of foam drainage equation such Adomian decomposition method (Helal and Mehanna [2007](#page-8-6)), reduced diferential transform method (Gubes et al. [2015\)](#page-8-7), homotopy analysis method (Singh et al. [2016](#page-9-5)), and Haar wavelets method (Arbabi et al. [2016](#page-7-2)). The main focus of the present work is to formulate the strategy being pursued and to obtain the numerical solution of the fractal foam drainage equation. We also aim to confirm that the two-scale method is powerful, efficient, and promising in handling scientifc and engineering problems (Elias-Zuniga et al. [2021a,](#page-8-8) [b,](#page-8-9) [c](#page-8-10), [d](#page-8-11); He [2020c](#page-8-12); Wang [2021;](#page-9-6) Wang et al. [2019](#page-9-7)).

2 Basic idea of the homotopy perturbation method

To illustrate the basic concept of the homotopy perturbation method, consider the following non-linear functional equation (Li and Nadeem [2019;](#page-8-13) Nadeem and Li [2019](#page-9-8)),

$$
A(u) - f(r) = 0, \quad r \in \Omega,
$$
\n(3)

with boundary conditions

$$
B\left(u, \frac{\partial u}{\partial n}\right) = 0, \quad r \in \Gamma,\tag{4}
$$

where *A* is a general functional operator, *B* is a boundary operator, $f(r)$ is a known analytic function, and Γ is the boundary of the domain Ω . The operator *A* can generally be divided into two operators, *L* and *N*, where *L* is a linear and *N* being a nonlinear operator. Therefore, Eq. ([3\)](#page-2-0) can be written as follows

$$
L(u) + N(u) - f(r) = 0.
$$
 (5)

Using the homotopy technique, we construct a homotopy $v(r, p)$: $\Omega \times [0, 1] \rightarrow \mathbb{R}$ that satisfes

$$
H(v, p) = (1 - p)[L(v) - L(u_0)] + p[L(v) - N(v) - f(r)],
$$
\n(6)

or

$$
H(v, p) = L(v) - L(u0) + pL(u0) + p[N(v) - f(r)] = 0,
$$
\n(7)

where $p \in [0, 1]$, is called homotopy parameter, and u_0 is an initial approximation for the solution of Eq. (3) (3) , which satisfies the boundary conditions. According to HPM, we can use p as a small parameter and assume that the solution of Eq. ([7\)](#page-2-1) can be written as a power series in *p*

$$
v = v_0 + pv_1 + p^2 v_2 + \dotsb. \tag{8}
$$

Considering $p = 1$, the approximate solution of Eq. (3) (3) will be obtained as follows

$$
u = \lim_{p \to 1} v = v_0 + v_1 + v_2 + v_3 + \dotsb. \tag{9}
$$

Additionally, some recent applications of HPM can be viewed in (He et al. [2020a,](#page-8-14) [b](#page-8-15); He and Dib [2020a,](#page-8-16) [b](#page-8-17); Skrzypacz et al. [2020;](#page-9-9) Anjum and He [2020](#page-7-3); He and Jin [2020](#page-8-18)).

3 Fractional complex transform

During the modeling of a problem, the dimension and scale are extremely valuable things since the various scales and dimensions will lead to remarkable results and properties for the same confguration. A fractional complex transform is a scientifc approach that converts a fractional diferential equation into a fractal space in a continuous space and is defned as (He et al. [2012;](#page-8-19) Li and He [2010](#page-8-20); Wang and Yao [2020](#page-9-10))

$$
\Delta S = \frac{\Delta t^{\alpha}}{\Gamma(1+\alpha)},\tag{10}
$$

where ΔS is the smaller scale and Δt is the larger scale. On a smaller scale, the timefractional foam drainage equation behaves discontinuously, especially at the peak of the solitary wave. On the other hand, the larger scale predicts a smooth solitary wave. The transformation is given in Eq. [\(10](#page-3-0)) is an approximate one to convert a fractal space on a small scale to a smooth space with a large scale. To under this, consider an example of a tree that stops growing at night, so when we use a scale of 24 h, it grows continuously, while when we measure it in 12 h, it becomes discontinuous. So, Eq. (10) (10) is also called the two-scale transform (Ain and He [2019;](#page-7-4) He and Ji [2019](#page-8-21); He and Ain [2020\)](#page-8-22), which was geometrically studied using fractal theory. The results of any particular problem depend on the scale. For example, on an observable scale, the liquid is steady; accordingly, Newton's laws can be executed, otherwise invalid in case of molecular scale. For example, on an observable scale, Newton's law is efective only if the liquid is steady, otherwise, it is illegitimate. In other words, if the fow is independent of time, then Newton's law is valid while if the fow is dependent on time, then Newtons' law becomes invalid. So, the two-scale transform is to convert a fractal space on a small scale to an approximate smooth space on a larger scale. Some recent work on fractal calculus can be studied in Wang [\(2020a,](#page-9-11) [b\)](#page-9-12), Wang and Wang ([2020a](#page-9-13), [b](#page-9-14), [2021](#page-9-15)) andd Wang et al. [\(2021](#page-9-16)).

4 Analysis of the method

In this section, we obtain an analytical solution of Eq. [\(1](#page-1-1)). We frst use the transformation

$$
\Omega(x,t) = u^2(x,t). \tag{11}
$$

Thus, Eq. ([1\)](#page-1-1) becomes as

$$
D_t^{\alpha} u^2 + \frac{\partial}{\partial x} \left(u^4 - \frac{u}{2} \cdot \frac{\partial}{\partial x} (u^2) \right) = 0,
$$

\n
$$
D_t^{\alpha} u^2 + \frac{\partial}{\partial x} \left(u^4 - u^2 \frac{\partial u}{\partial x} \right) = 0,
$$

\n
$$
2u D_t^{\alpha} u + 4u^3 \frac{\partial u}{\partial x} - 2u \frac{\partial u}{\partial x} \cdot \frac{\partial u}{\partial x} - u^2 \frac{\partial^2 u}{\partial x^2} = 0,
$$

which may also be written as

$$
D_t^{\alpha} u = \frac{1}{2} u u_{xx} - 2u^2 u_x + u_x^2, \qquad (12)
$$

with initial condition

$$
u(x,0) = -\sqrt{c} \tanh(\sqrt{c}x),\tag{13}
$$

where *c* is the velocity of the wavefront (Alquran [2014\)](#page-7-5). The initial condition (13) (13) is taken in such a way that it satisfes the problem to fnd a particular solution. Initially, we will apply the fractional complex transform to write it in its partner diferential equation such as

$$
S = \frac{t^{\alpha}}{\Gamma(1+\alpha)}.\tag{14}
$$

So, Eq. [\(12\)](#page-4-1) can be written as

$$
\frac{\partial u}{\partial S} + 2u^2 \frac{\partial u}{\partial x} - \left(\frac{\partial u}{\partial x}\right)^2 - \frac{1}{2}u \frac{\partial^2 u}{\partial x^2}.
$$
 (15)

The initial guess will be taken as $u_0 = -\sqrt{c} \tanh(\sqrt{cx})$. The HPM together with He's polynomials will be applied on Eq. ([15](#page-4-2)) as follows

$$
\frac{\partial u_1}{\partial S} + 2u_0^2 \frac{\partial u_0}{\partial x} - \left(\frac{\partial u_0}{\partial x}\right)^2 - \frac{1}{2}u_0 \frac{\partial^2 u_0}{\partial x^2} = 0, \quad u_1(x,0) = 0
$$

$$
\frac{\partial u_2}{\partial S} + 2\left(u_0^2 \frac{\partial u_1}{\partial x} + 2u_0 u_1 \frac{\partial u_0}{\partial x}\right) - 2\frac{\partial u_0}{\partial x} \frac{\partial u_1}{\partial x} - \frac{1}{2}\left(u_0 \frac{\partial^2 u_1}{\partial x^2} + u_1 \frac{\partial^2 u_0}{\partial x^2}\right) = 0, \quad u_2(x, 0) = 0
$$

$$
\frac{\partial u_3}{\partial S} + 2\left(u_0^2 \frac{\partial u_2}{\partial x} + 2u_0 u_1 \frac{\partial u_1}{\partial x} + (u_1^2 + 2u_0 u_2) \frac{\partial u_0}{\partial x}\right) - \left(\left(\frac{\partial u_1}{\partial x}\right)^2 + 2\frac{\partial u_0}{\partial x} \frac{\partial u_2}{\partial x}\right)
$$

$$
- \frac{1}{2}\left(u_0 \frac{\partial^2 u_2}{\partial x^2} + u_1 \frac{\partial^2 u_1}{\partial x^2} + u_2 \frac{\partial^2 u_0}{\partial x^2}\right) = 0, \quad u_3(x, 0) = 0
$$

$$
\frac{\partial u_4}{\partial S} + 2\left(u_0^2 \frac{\partial u_3}{\partial x} + 2u_0 u_1 \frac{\partial u_2}{\partial x} + (u_1^2 + 2u_0 u_2) \frac{\partial u_1}{\partial x} + (2u_1 u_2 + 2u_0 u_3) \frac{\partial u_0}{\partial x}\right)
$$

$$
- \left(2\frac{\partial u_1}{\partial x} \frac{\partial u_2}{\partial x} + 2\frac{\partial u_0}{\partial x} \frac{\partial u_3}{\partial x}\right) - \frac{1}{2}\left(u_0 \frac{\partial^2 u_3}{\partial x^2} + u_1 \frac{\partial^2 u_2}{\partial x^2} + u_2 \frac{\partial^2 u_1}{\partial x^2} + u_3 \frac{\partial^2 u_0}{\partial x^2}\right)
$$

$$
= 0, \quad u_4(x, 0) = 0
$$

Proceeding in the same way, we can obtain the high-order approximations.

$$
u_0(x, S) = -\sqrt{c} \tanh(\sqrt{c}x),
$$

$$
u_1(x, S) = c^2 \operatorname{sech}^2(\sqrt{c}x)S,
$$

$$
u_2(x, S) = c^{\frac{7}{2}} \operatorname{sech}^2(\sqrt{c}x) \tanh(\sqrt{c}x)S^2,
$$

$$
u_3(x, S) = -\frac{1}{12}c^4(-1 - 13c + 10c \cosh(2\sqrt{cx}) - (-1 + c)\cosh(4\sqrt{cx})) \operatorname{sech}^6(\sqrt{cx})S^3,
$$

$$
u_4(x, S) = -\frac{1}{96}c^{\frac{11}{2}}\left(-46 - 582c + (-41 + 289c)\cosh(2\sqrt{cx}) + 6(1 + c)\cosh(4\sqrt{cx}) + \cosh(6\sqrt{cx}) - c\cosh(6\sqrt{cx})\right)\sech^8(\sqrt{cx})\tanh(\sqrt{cx})S^4,
$$

the rest of the components can also be found in the same manner. Thus, the approximate solution according to HPM has the following form

$$
u(x,t) = -\sqrt{c} \tanh(\sqrt{c}x) + c^2 \mathrm{sech}^2(\sqrt{c}x)S + c^{\frac{7}{2}} \mathrm{sech}^2(\sqrt{c}x) \tanh(\sqrt{c}x)S^2.
$$

Thus, by using Eq. (14) (14) , we get

$$
u(x,t) = -\sqrt{c} \tanh(\sqrt{c}x) + c^2 \mathrm{sech}^2(\sqrt{c}x) \left[\frac{t^{\alpha}}{\Gamma(1+\alpha)} \right]
$$

$$
+ c^{\frac{7}{2}} \mathrm{sech}^2(\sqrt{c}x) \tanh(\sqrt{c}x) \left[\frac{t^{\alpha}}{\Gamma(1+\alpha)} \right]^2.
$$

This series leads to the exact solution for Eq. [\(1](#page-1-1)) when $\alpha = 1$.

$$
\Omega(x,t) = -\sqrt{c} \tanh \sqrt{c(x-ct)}. \quad x \le ct \tag{16}
$$

Using the Mathematica package 11.0.1, we illustrate the graphical representations and physical behaviors of the fractal foam drainage equation. For case $\alpha = 1$ $\alpha = 1$, we compare the obtained results of the approximate solution of Eq. (1) and the exact solution of Eq. (16) (16) , which is depicted in Figs. [2](#page-6-0) and [3](#page-6-1). From Fig. [2,](#page-6-0) it is obvious that when $\alpha = 1$, the solution is nearly identical to the exact solution. The solution graphs have declared that the obtained results are almost identical and confrm worthy contact with the exact solution, which aids us to understand the nature of the fractal foam drainage equation. Moreover, Fig. [4](#page-6-2) shows the error distribution of the approximate solution and the exact solution which helps us to capture the behavior of the obtained solution. Some recent developments in fractal vibration models can be studied from somewhere else (Chun-Hui et al. [2021a,](#page-8-23) [b](#page-8-24); He et al. [2021a](#page-8-25), He et al. [2021b\)](#page-8-26).

5 Conclusion

In this paper, we suggested a hybrid scheme such as fractional complex transform is coupled with HPM to obtain the analytic solution of the nonlinear timefractional foam drainage equation, which is a simple model of the fow of liquid through channels and nodes between the bubbles. The fractional complex transform performed a simple and excellent approach to convert a fractional diferential equation into its diferential partner which is suitable for the development of the two-scale method. This method is applied directly without any discretization, linearization, and small parameter assumptions, which in actual ruins the physical nature of the problems. The results show that the proposed method is very efficient and powerful. The graphical representation shows that the solution procedure is simple and might be found in wide applications of engineering. The procedure reveals that the semi-inverse method is highly efficient and powerful, and can be generalized to other nonlinear evolution equations with fractal derivatives in future applications.

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