



# Design of a High Efficiency p-Si Based Heterojunction Solar Cell with $\text{In}_2\text{S}_3$ Window and NiO BSF Layers

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## Abstract

The numerical evaluation performed on the design of n- $\text{In}_2\text{S}_3$ /p-Si/p<sup>+</sup>-NiO solar cell reveals that it can come up with a high efficiency gain along with substantial values in other photovoltaic parameters. The pristine n- $\text{In}_2\text{S}_3$ /p-Si structure imparts a power conversion efficiency, PCE of 23.24%. The selection of NiO in back surface field (BSF) layer makes an improvement of ~0.1 V in open circuit voltage,  $V_{\text{OC}}$  and a slight improvement in short circuit current density,  $J_{\text{SC}}$ . Under single sun and AM 1.5G spectrum, the optimum thickness (Window = 0.2  $\mu\text{m}$ , Base = 350  $\mu\text{m}$ , BSF = 0.2  $\mu\text{m}$ ), doping concentration (Window<sub>donor</sub> =  $1.0 \times 10^{18} \text{ cm}^{-3}$ , Base<sub>acceptor</sub> =  $1.0 \times 10^{17} \text{ cm}^{-3}$ , BSF<sub>acceptor</sub> =  $1.0 \times 10^{20} \text{ cm}^{-3}$ ) and defect density (Window =  $1.0 \times 10^{14} \text{ cm}^{-3}$ , Base =  $1.0 \times 10^{12} \text{ cm}^{-3}$ , BSF =  $1.0 \times 10^{14} \text{ cm}^{-3}$ ) provide an enhanced PCE of 26.74% with the application of NiO as BSF in the pristine structure. The other photovoltaic parameters results with  $V_{\text{OC}} = 0.79 \text{ V}$ ,  $J_{\text{SC}} = 40.37 \text{ mA/cm}^2$  and FF = 83.85%. Insertion of a thin and optimized NiO reduces recombination at the back surface by a potential barrier that enhances  $V_{\text{OC}}$  and current in the device. Such a cost-effective solar cell exhibits enough possibility of fabricating a highly efficient, reliable and promising performance device.

**Keywords** p-Silicon ·  $\text{In}_2\text{S}_3$  window · NiO BSF · High efficiency · SCAPS-1D

## 1 Introduction

P-type (p-Si) based solar cells have received a lot of interest due to their to substantial accessibility, lower cost [1]. In application field they are frequent in space technologies and have proven to be resilient to radiation [2]. Their cell operation mainly involves: (i) absorption of sunlight and corresponding generation of electron–hole pair, (ii) separation of charge carriers at the terminals and (iii) external circuit capable of collecting the carriers [2, 3]. Such a cell considering an infinite number of connections and extremely focused sunlight can achieve a limiting efficiency of 86.8% [4, 5]. In today's marketplace, a variety of tandem, two-layer or multilayered cells are frequently offered with ~30% efficiency under one-sun illumination. For instance, a recent work with the application of methyl ammonium tin mixed halide based materials in different layers of a Si-based tandem solar cell

displays 30.7% PCE with high  $V_{\text{OC}}$  of 2.14 V [6]. Such performance can be improved to nearly 40% under centralized rays, but it may rise the overall complexity and the manufacturing costs [7, 8].

In an effort to find a fantastic answer to the problem at hand, various combinations of window, base and back surface materials have been sought in some recent works and the end result is always a lightweight and uncomplicated construction approach [9–11]. In the similar pattern, we propose a solar cell with three different materials namely silicon, indium sulphide ( $\text{In}_2\text{S}_3$ ) and nickel oxide (NiO) which can form two heterojunctions and theoretically comes up with high standard of exemplary cell performance.

The bandgap value of Si ( $E_g = 1.12 \text{ eV}$  and corresponding wavelength for light absorption  $\lambda_{\text{cut-off}} = 1107 \text{ nm}$ ) are extremely close to the ideal values for single absorber solar-to-electric energy conversion. The amount of light that is coupled into the silicon absorber can be maximized by reducing losses due to parasitic absorption and surface reflection regardless of the stacking or deposition method [12].

With thin film technology, the bandgap of  $\text{In}_2\text{S}_3$  has been found to vary between 2.0–3.7 eV providing high

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transmittance between 380 to 700 nm. In comparison with popular CdS in window,  $\text{In}_2\text{S}_3$  is nontoxic and comes up with high stability and high carrier mobility [13]. Moreover, good transparency and photoconductive ability of  $\text{In}_2\text{S}_3$  has made it an alternative candidate to replace CdS window [14]. Some diverse in-line processing techniques for deposition of  $\text{In}_2\text{S}_3$  thin films such as thermal evaporation, chemical bath deposition, physical vapor deposition, atomic layer deposition, successive physical layer adsorption and reaction, pyrolysis etc. are commonly recognized [13, 15, 16]. Likewise, NiO thin film can also be deposited using varieties of methods like pulsed laser deposition, e-beam technique, chemical vapour deposition, sputtering or even spin coating and inkjet printing [17]. It is possible to tailor some properties of the NiO thin films with annealing or plasma treatment. Additionally, a wide variety of tunings of work function from 3.7 to 6.7 eV, can be achieved by varying the surface dipole, composition, or defect density [18].

Effects like surface recombination will become more prominent as the p-type monocrystalline wafer gets thinner. Hence, measures to control or lessen this effect are ought to be taken. One way to reduce back surface recombination in a p-type wafer is to apply BSF layer with the base structure [19, 20].

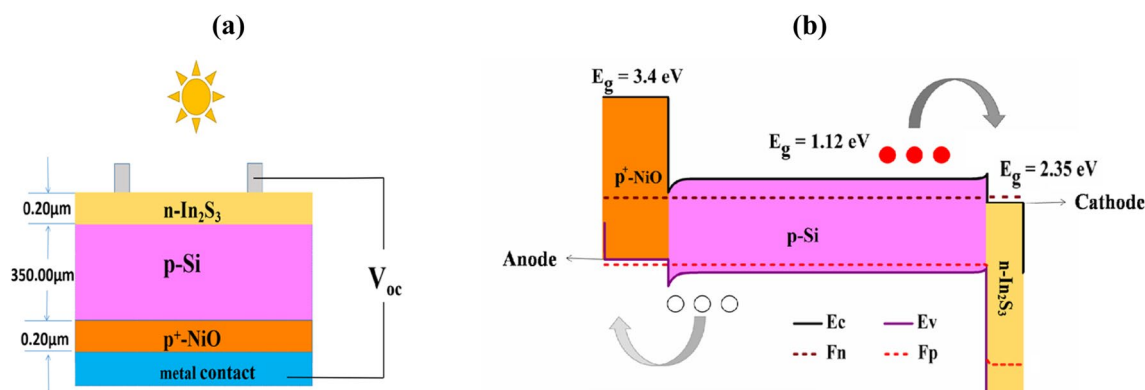
Previously, NiO has demonstrated its capacity to significantly reduce carrier recombination at the contacts when used as a BSF with CdTe based cells, leading to improved  $V_{OC}$  [21]. By several research groups, NiO has been successfully found to be employed as hole transport layer in CdTe and SnS based cells along with some perovskite structures [22–24]. Non-stoichiometric oxygen-rich NiO films develop into reasonably excellent p-semiconductors with wide band gap of 3.7 eV. Properties like exceptional chemical stability, low cost, and magnetic, electric, and optical properties of NiO films have made them a desirable p-type oxide semiconductor for use in optoelectronic

applications [25]. The recombination at the back of solar cells becomes an issue and can directly lower efficiency when charge carriers are slightly out of the field. Studies reveal that the addition of NiO causes increment of potential barrier at the back contact along with reduction of recombination with higher  $V_{OC}$ . Therefore, we suggest the formation of heavily doped  $p^+$ -NiO at the back surface of p-Si centric bulk Si, enabling efficient separation of hole and electron at the back which in turn can reduce their recombination and can increase total quantum efficiency. Thick NiO can lead to a dominant potential barrier causing deterioration in the cell performance [26]. To overcome such complexities, the thickness as well as other parameters like doping concentration and defect density have been optimized here as suggested by the previous studies [27]. To the best of our knowledge, such possibility of  $\text{In}_2\text{S}_3$  and NiO has never been studied with silicon before.

In this work, we propose n- $\text{In}_2\text{S}_3$ /p-Si/ $p^+$ -NiO-based double-junction structure where we perform numerical simulations for obtaining band diagram, J-V characteristics, quantum Efficiency, optimization of associated layers and observing other features related to solar cell. The standard AM 1.5G global spectrum has been utilized throughout the simulations to identify overall performance parameters.

## 2 Device Architecture and Calculation

Figure 1 represents the structured layout of the propounded Si solar cell along with its illuminated energy band diagram. Incident photons from sun first fall on to n-type thinner window and reaches to the n/p window-absorber junction as shown in Fig. 1(a). It is evident that p-Si makes a satisfactory band alignment with n- $\text{In}_2\text{S}_3$  window as well as  $p^+$ -NiO BSF layer as illustrated in Fig. 1(b). The changes in photovoltaic characteristics i.e.  $V_{OC}$ ,  $J_{SC}$ , FF and PCE with optimized values have been discussed later with this work.



**Fig. 1** **a** Proposed p-Si centric solar cell showing layer optimized thickness and **b** its energy band diagram (not to scale) when lit with AM 1.5G solar spectrum

The simulation software SCAPS-1D (version 3.3.07) has been exerted to calculate solar device outcomes [11, 23]. The absorption interpolation model for  $\text{In}_2\text{S}_3$  and NiO have been chosen from traditional model of the simulator. The absorption profile for p-Si has been selected from SCAPS default files. Rest of the parameters required for simulation have been selected from previous reports and are included in Table 1 [24, 28, 29]. With such parameters, the bending of the energy band at junctions have been happened in such a way that an easy path has been created for electron to be collected by Ti metal grid above the window layer. On the contrary, the proper collection of holes has been ensured by the BSF with  $E_V$  values above the Si layer.  $E_V$  indicates the energy of valence band while  $F_n$  and  $F_p$  denote the Fermi energy levels for electron and hole, in turn. The energy level of  $F_p$  lies below the  $E_V$  of NiO and goes inside the energy band gap of  $\text{In}_2\text{S}_3$ . Likewise,  $F_n$  energy level goes over by a small amount of conduction level energy  $E_C$  of  $\text{In}_2\text{S}_3$  and passes between NiO energy bandgap. Note that, the  $E_C$  and  $E_V$  values are at 1.37 and -0.97 eV, respectively at  $\text{In}_2\text{S}_3$  layer that contains a bandgap of 2.35 eV and electron affinity of 4.40 eV [30]. In the rear field, NiO has the electron affinity of 1.45 eV with a bandgap of 3.4 eV creating  $E_C$  and  $E_V$  values at 2.28 and 5.37 eV, respectively [31]. Thus, they create proper dual heterojunctions with Si containing the bandgap of 1.12 eV and electron affinity of 4.05 eV.

Moreover, the following values in Table 2 have been also considered in each layer. In all levels, Gaussian energetic distribution with reference for defect energy above  $E_V (< 2.7)$ , energy level with respect to reference 0.35 eV with characteristic energy of 0.10 eV have been considered.

**Table 1** Properties of different layers in the p-Si solar cell structure

Parameters	n- $\text{In}_2\text{S}_3$	p-Si	$\text{p}^+$ -NiO
Thickness ( $\mu\text{m}$ )	0.20	350.00	0.20
Bandgap, $E_g$ (eV)	2.35	1.12	3.40
Electron affinity (eV)	4.40	4.05	1.45
Dielectric permittivity (relative)	13.50	11.90	10.07
CB effective density of states ( $1/\text{cm}^3$ )	$2.20 \times 10^{17}$	$1.05 \times 10^{19}$	$2.80 \times 10^{19}$
VB effective density of states ( $1/\text{cm}^3$ )	$1.80 \times 10^{19}$	$3.92 \times 10^{18}$	$1.00 \times 10^{19}$
Electron mobility ( $\text{cm}^2/\text{Vs}$ )	$1.0 \times 10^2$	$3.9 \times 10^3$	$1.2 \times 10^1$
Hole mobility ( $\text{cm}^2/\text{Vs}$ )	$2.5 \times 10^1$	$1.9 \times 10^3$	$2.8 \times 10^0$
Shallow uniform donor density, $N_D$ ( $1/\text{cm}^3$ )	$1.0 \times 10^{18}$	0.00	0.00
Shallow uniform acceptor density, $N_A$ ( $1/\text{cm}^3$ )	0.00	$1.0 \times 10^{17}$	$1.0 \times 10^{20}$
Density of defects ( $\text{cm}^{-3}$ ) (above $E_V$ w.r.t. $E_{\text{ref}}$ (eV))	$1.0 \times 10^{14}$	$1.0 \times 10^{12}$	$1.0 \times 10^{14}$
Type of defects	Single acceptor	Single donor	Single donor
Energetic distribution	Gaussian	Gaussian	Gaussian
Energy level with respect to Reference (eV)	1.91	0.35	0.65

**Table 2** Constant properties considered in proposed p-Si solar cell

Parameters	Value
Electron and hole thermal velocity (cm/s)	$1.0 \times 10^7$
Electron capture cross section ( $\text{cm}^2$ )	$1.0 \times 10^{-15}$
Hole capture cross section ( $\text{cm}^2$ )	$1.0 \times 10^{-17}$
Temperature (K)	300
Incident( or bias) light power ( $\text{W}/\text{m}^2$ )	1000
Transmission(%)	100

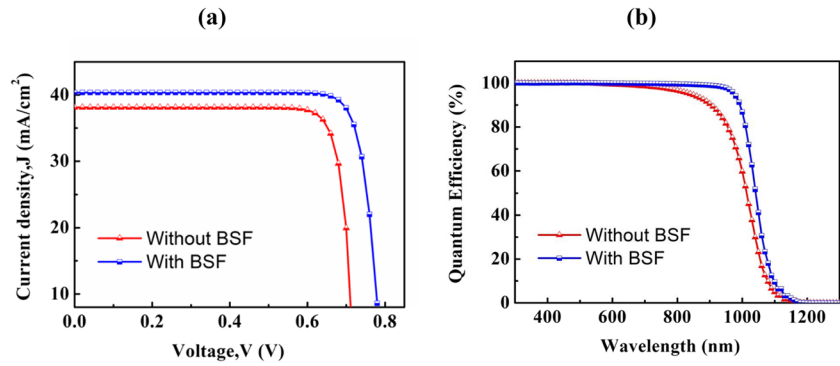
### 3 Results and Discussion

#### 3.1 J-V Characteristics and QE Performance of the Proposed Model

Figure 2(a) signifies the J-V characteristics of the proposed cell along with its photon absorption performance corresponds to wavelength between 300–1400 nm. The pristine structure with single n<sup>+</sup>/p junction cell exhibits 0.72 V of  $V_{OC}$  and  $38.12 \text{ mA}/\text{cm}^2$  of  $J_{SC}$  with corresponding FF of 84.75% and resultant PCE of 23.24%. On the contrary, the addition of only 0.2  $\mu\text{m}$  thick NiO in optimized condition provides an improvement in  $V_{OC}$  as well as a marginal rise in  $J_{SC}$ . The addition of NiO ( $E_g = 3.4 \text{ eV}$ ) creates a proper heterojunction with Si ( $E_g = 1.12 \text{ eV}$ ) and thereby the double heterojunction of n/p/p<sup>+</sup> in the proposed structure provides an enhanced  $V_{OC}$  of 0.79 V. The suppression of carrier recombinations at the surface owing to the diminishing of speed of surface recombination by the electric field of Si/NiO interface escalates both the  $V_{OC}$  and  $J_{SC}$  of the silicon device and exalted the PCE to 26.74% [10, 11, 24].

The QE graph illustrated in Fig. 2(b) further explains corresponding absorption of photons due to addition of NiO

**Fig. 2** The (a) J-V characteristic curves and (b) Quantum efficiency in the proposed optimized Si solar cell



with this regard. According to the QE, all photons with comparable wavelengths less than 1200 nm have been absorbed. It indicates that the entire wavelength has been absorbed by the p-Si base. Figure 2(b) makes evident how the NiO facilitates the extra electron–hole (e–h) separation generated by the photons of slightly longer wavelengths in Si absorber. For comparison, the QE for the pristine structure is ~72% in the wavelength of 970 nm while the proposed model still displays QE of ~96%. In 1050 nm, the QE of pristine structure falls to ~23% and the addition of BSF provides the advantage of having QE of ~46%. In light of this, it is evident that a NiO BSF even with a small thickness can prevent minority carriers from entering the recombination surface, which in turn contributes to the upsurge and persistence of the  $V_{OC}$  and  $J_{SC}$  [20]. Some previous studies with dual heterojunctions also come up with an intensification of this voltage [28, 32–35]. As the bandgap of NiO is a bit higher than that of p-Si, it creates well prevention to the minority carrier progression which in turn reduces surface recombination velocity [10, 11, 24]. As a result,  $J_{SC}$  gets enhanced nearly about 2.25 mA/cm<sup>2</sup> and displays a  $J_{SC}$  of 40.37 mA/cm<sup>2</sup> in the proposed structure [36]. Thus, the amelioration of both  $V_{OC}$  and  $J_{SC}$  with the establishment of n-In<sub>2</sub>S<sub>3</sub>/p-Si/p<sup>+</sup>-NiO causes overall PCE to improve about 3.5% and culminate in 26.74%.

**3.2 Consequences of Changes in Operating Temperature**

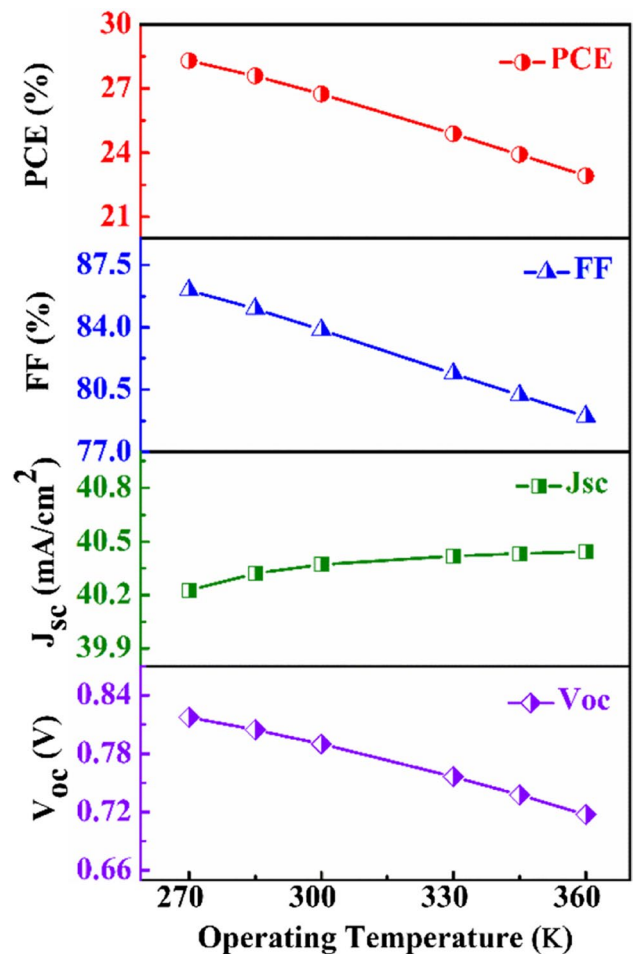
Like the most of the semiconductor devices, the proposed Si solar cell has been modeled at an operating temperature of 300 K. The changes in the temperature follows the conventional cell output pattern as shown in Fig. 3. From previous reports, it is usual for Si based cell that the rise in temperature decreases  $V_{OC}$  and slightly elevates  $J_{SC}$  [37–39]. From the general expression of  $V_{OC}$ , we know,

$$V_{OC} = V_{th} \cdot \ln\left(\frac{J_{SC}}{J_o} + 1\right) \tag{1}$$

The reliance of  $V_{OC}$  on temperature can be explained from the following equation for the temperature range of 273–523 K in p-Si [38].

$$\frac{dV_{OC}}{dT} = \frac{V_{OC}}{T} + V_{th} \frac{1}{J_{SC}} \frac{dJ_{SC}}{dT} - \left( \frac{E_g(0)}{T} + \frac{\alpha T}{(T + \beta)^2} \right) \tag{2}$$

Here,  $J_o$  represents reverse saturation current density,  $E_g(0) = E_g$  indicates bandgap at  $T = 0K$  with T being the absolute



**Fig. 3** The effect of variation in operating temperature on proposed cell parameters

temperature and  $\alpha, \beta$  are two empirical parameters.  $V_{th} = kT/q$  defines the thermal voltage, where  $k$  is the Boltzmann’s constant and  $q$  represents the electron charge. Thus, the variation of temperature ranging from 270–360 K causes a decrement of  $V_{OC}$  from 0.81 V to 0.71 V in the proposed structure. In the same range, we also see a small enhancement of current from 40.22 to 40.44 mA/cm<sup>2</sup> attributed to the fact of capturing more photons leading by a shrunk in bandgap [38]. Moreover, the Fermi–Dirac distribution can represent how carriers populate energy levels [40]. As the bandgap decreases with thermal energy, the  $J_{SC}$  slightly increases as the photon cut-off wavelength shifts to a longer wavelength of the spectrum enabling creation of more e–h pairs. However, such a change for  $J_{SC}$  is very small compared to the change in  $V_{OC}$  [41]. As the change in  $V_{OC}$  dominates, FF and corresponding PCE follows the similar pattern [38]. Thus, FF shows a decrement from 86.07 to 79% and corresponding PCE suffers a fall of 28.3 to 22.92% for the proposed structure.

### 3.3 Series and Shunt Resistance’s Impact on the Proposed Model

Though the optimized structure has been experimented without any parasitic parameters namely series and shunt resistances, the effect of their addition has been delineated in Fig. 4. In reality, it is natural for a cell to show certain series resistance originating from the lump and contacts. In contrast, leakage current forming fabrication defects drive the shunt resistance [10]. In general, low values of series and higher values of shunt resistances of the devices must be achieved in order to design and construct efficient PV

devices [42]. From Fig. 4(a), we hardly find any effect on  $V_{OC}$  and  $J_{SC}$  due to change in series resistance from 0 to 5  $\Omega\cdot\text{cm}^2$ . However, like previous findings, FF drastically falls from 83.35 to 60.8% due to change in 5  $\Omega\cdot\text{cm}^2$  which in turn takes the value of PCE down to 19.4% [10, 42].

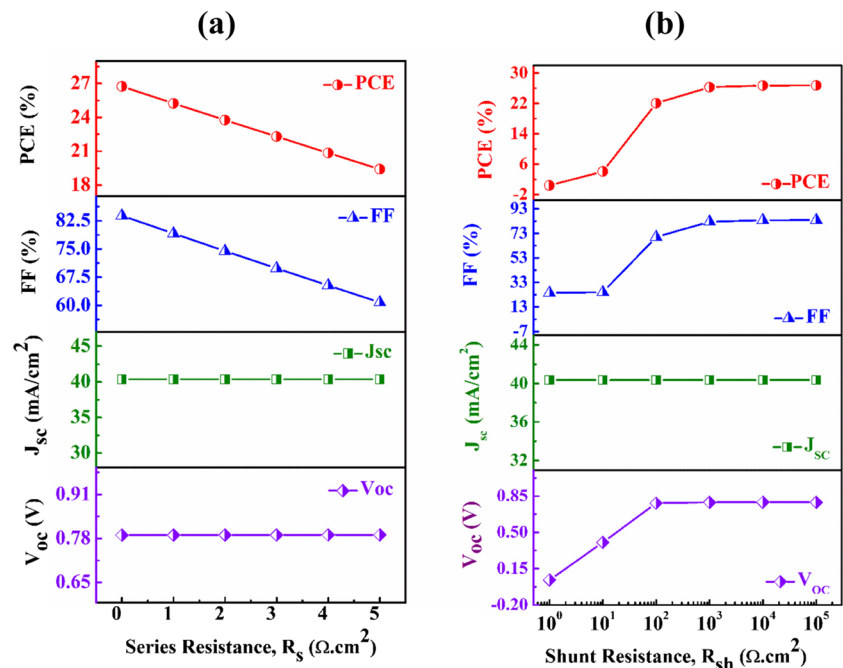
Shunt resistance also has a considerable effect on the observed output as depicted in Fig. 4(b). With shunt resistance starting from the lower order of ten,  $V_{OC}$  advances until it touches its maximum value at 0.79 V. Considering zero shunt the device results the  $J_{SC}$  of 40.37 mA/cm<sup>2</sup>.  $J_{SC}$  remains on the same value as expected from the general theoretical expressions of Si solar cell. Consequently, FF follows the changes of  $V_{OC}$ . At the order of 10<sup>2</sup>  $\Omega\cdot\text{cm}^2$ ,  $V_{OC}$  gets raised at 0.78 V with FF of 70.02%. In the higher order of shunts, we find gradual rise of FF to 83.84% as  $V_{OC}$  reaches at 0.79 V and the corresponding PCE displays the optimized value of 26.74%.

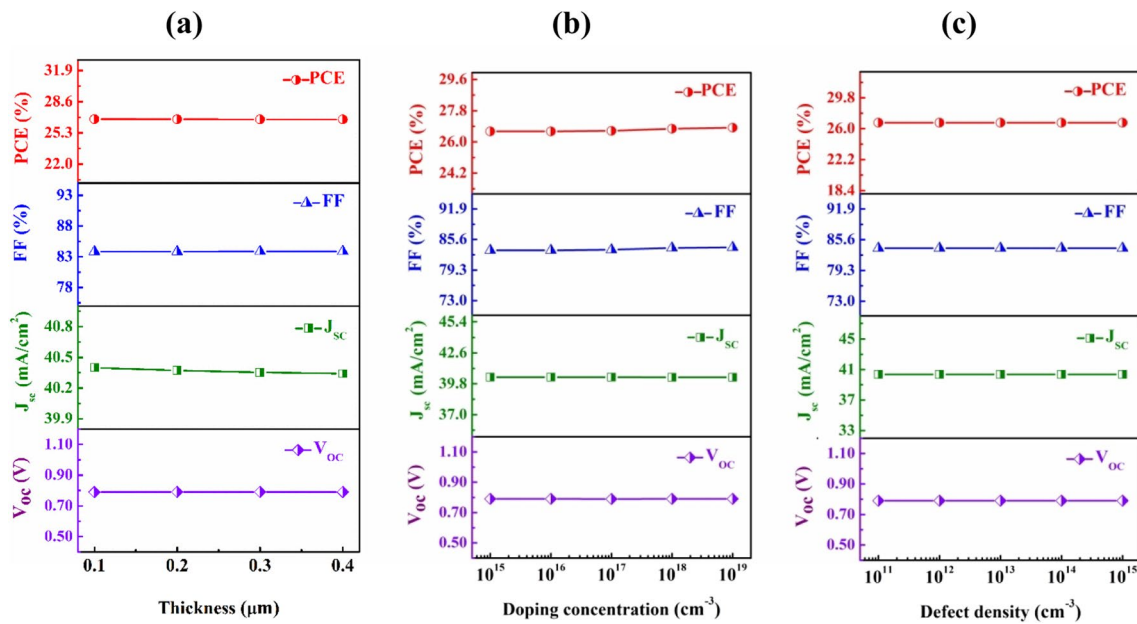
### 3.4 Optimization of Window, Absorber and BSF Layer in the Proposed Model

The simulation has been run with variation of three parameters namely thickness, carrier and defect density in each part to find the highest PCE with considerable amount of  $V_{OC}$ ,  $J_{SC}$  and FF. The simulations result three figures in each layer while keeping the values of corresponding layers in optimum condition. The outputs give an indication on how the device performance may vary during fabrication of real time cells. However, the ultimate target is to find the theoretical optimized condition that gives the highest PCE.

Figure 5(a), (b) and (c) display the consequences of changes in the parameters of In<sub>2</sub>S<sub>3</sub> window layer. The

**Fig. 4** a The ramifications of changing series and (b) shunt resistance to the proposed structure



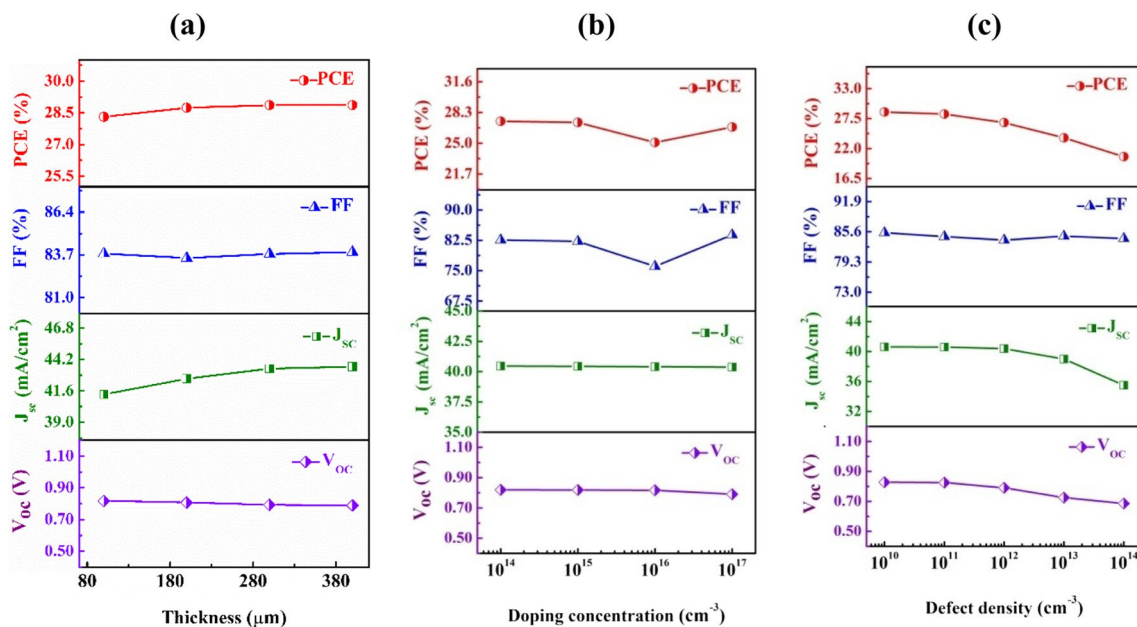


**Fig. 5** a Thickness, b carrier and c defect density variation in In<sub>2</sub>S<sub>3</sub> window layer of proposed Si solar cell

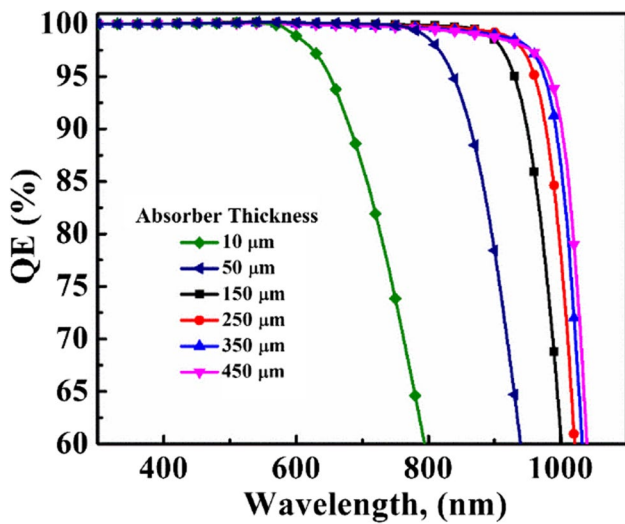
disparity in thickness, doping concentration or defect density in the window does not seem to make a significant effect in the operation of the proposed structure. Here, the breadth of window layer has been deviated between 0.1 and 0.4 μm while the doping concentration and defect densities have been varied between the ranges from 10<sup>15</sup> to 10<sup>19</sup> cm<sup>-3</sup> and 10<sup>11</sup> to 10<sup>15</sup> cm<sup>-3</sup>, respectively. As In<sub>2</sub>S<sub>3</sub> possess wide bandgap of 2.35 eV, a trivial fluctuation occurs in the investigated range of thickness. Even the alteration in other parameters does

not seem to affect the cell performance at all [34, 43]. Thus, the optimum width and carrier concentration have been chosen at 0.2 μm and 10<sup>18</sup> cm<sup>-3</sup>, respectively with considerable defect density of 10<sup>14</sup> cm<sup>-3</sup>.

In Fig. 6, we observe the cell performance due to the similar variation of parameters in *p*-Si absorber. The addition of 300 μm absorber width does not create a notable change in output as illustrated in Fig. 6(a). However, the addendum of thickness provides the opportunity of greater absorption



**Fig. 6** a Thickness, b doping and c defects variation in *p*-Si absorber of proposed solar cell



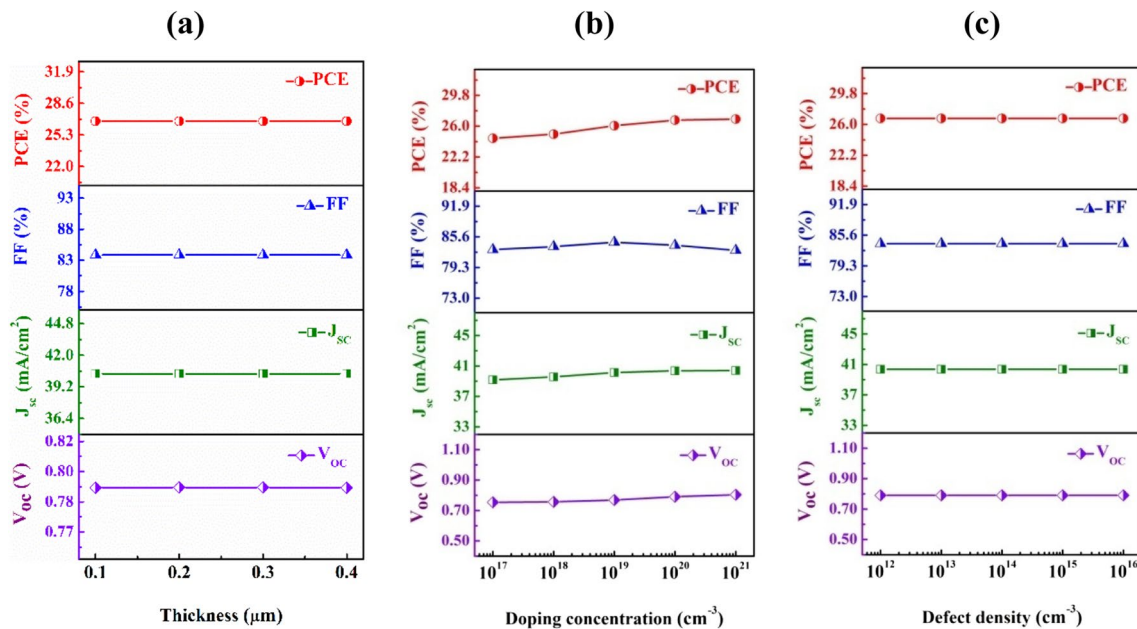
**Fig. 7** Variation of QE for different p-Si thickness in the proposed structure

of photons than before [34, 44]. As a consequence, we see a gradual increment in the values of  $J_{SC}$ . At 350  $\mu\text{m}$  of thickness the  $V_{OC}$  results with 0.79 V and  $J_{SC}$  comes with 43.43  $\text{mA}/\text{cm}^2$ . These values provide an acceptable FF of 83.76% and PCE with 28.86% which tend to rise with further addition.

There are no significant variations in  $V_{OC}$  and  $J_{SC}$  with addendum of higher order of carriers from  $10^{14} \text{ cm}^{-3}$  to  $10^{17} \text{ cm}^{-3}$  as delimitates in Fig. 6(b). FF drops down to

76.06% at  $10^{16} \text{ cm}^{-3}$  which is 82.59% at  $10^{14} \text{ cm}^{-3}$ . This reduction can be explained by a declining ratio between the rates of charge carrier production and recombination in the p-Si layer [45]. At higher acceptor densities, in the order of  $10^{18} \text{ cm}^{-3}$ , the FF again rises to 83.85%. This is attributed to a rise in the electric field (Schottky barrier) close to NiO, as well as an increase in the creation and recombination rates of free charges [45]. Furthermore, FF also gets improved with the small value of series resistance and a drop of diode ideality factor at higher acceptor densities [11, 46]. The efficiency seems to follow the FF pattern and at  $10^{17} \text{ cm}^{-3}$  the optimum efficiency is found to be 26.74% with FF of 83.85%.

It seems obvious that, the cell performance gets degraded with the addition of defect density in the absorber layer. It is seen from the Fig. 6(c) that  $V_{OC}$  faces a detrimental performance with the incremental order of defect densities. At the order of  $10^{10} \text{ cm}^{-3}$ , the PCE comes with 28.72% with voltage of 0.82 V and current of 40.61  $\text{mA}/\text{cm}^2$ . These values get lowered as the order is changed to  $10^{14} \text{ cm}^{-3}$  and result the voltage of 0.68 V with photocurrent of 35.51  $\text{mA}/\text{cm}^2$  and PCE of 20.52%. The SRH carrier recombination mechanism plays the principal role at higher order of defect densities. It causes recombination of carriers which in turn stimulates the dark current and lowers corresponding  $V_{OC}$  as well as the  $J_{SC}$  [34, 43, 47]. With the above findings we choose 350  $\mu\text{m}$  thick p-Si absorber to be doped in the concentration order of  $10^{17} \text{ cm}^{-3}$  while maintaining the defect in the order of  $10^{12} \text{ cm}^{-3}$  to achieve the highest PCE.



**Fig. 8** a Thickness, b doping and c defect density variation in NiO BSF layer of proposed solar cell

Figure 7 illustrates the QE variation over the p-Si thickness ranging from 10 to 450  $\mu\text{m}$ . It is inferred that, with the addition of absorber thickness, the longer wavelength photons take part in generation of additional e–h pairs [6, 45]. For instance, the QE related to 10  $\mu\text{m}$  of Si thickness drops to 60% correspond to incident photon wavelength of 790 nm. Increasing the thickness to 50  $\mu\text{m}$  achieves QE of 99% for same wavelength of photons and similar drop occurs at wavelength of around 940 nm. This is happened as at least a width of 30  $\mu\text{m}$  for silicon is needed to absorb photons of wavelength upto 900 nm [48]. When the layer is too thick, photo generated carriers must travel a larger distance to window and BSF for collection of carriers and a significant loss can occur [45]. QE does not show any major enhancement after the increment of absorber thickness over 400  $\mu\text{m}$ .

The width of NiO BSF layer, doping concentration and defect densities have been varied and the corresponding outcomes are delineated in Fig. 8. All the figures show that the alteration in thickness, doping and defects of NiO has negligible effect on the operation of the Si solar cell. The results are similar to the results of some previous works [10, 29]. The  $V_{\text{OC}}$  is almost constant at 0.79 V with thickness but has got a slight increment along with doping. The  $J_{\text{SC}}$  also remains constant at 40.37  $\text{mA}/\text{cm}^2$  which might have faced a slight improvement with increment in doping orders. Such increment may come with the possible enhancement in the carrier dependent lifetime along with the surface recombination speed and the reduction in excess minority carriers with relative high potential in  $n^+/p/p^+$  interface [49]. The defects in the studied range do not tend to cause any declination in cell performance as the carrier diffusion length is high enough compared to the width of the BSF that can rise recombination current in the device [9, 10]. Therefore, we have chosen an optimum thickness of NiO of 0.2  $\mu\text{m}$  with a doping of  $10^{20} \text{ cm}^{-3}$  with defects of  $10^{14} \text{ cm}^{-3}$ .

## 4 Conclusions

In this work, the performance of a p-Si based cell with  $n\text{-In}_2\text{S}_3$  window has been analyzed with and without the addition of NiO BSF layer. The band diagram of  $n\text{-In}_2\text{S}_3/p\text{-Si}/p^+\text{-NiO}$  solar cell reveals the perfect alignment of each layer. The proposed design comes with a PCE of 26.74% with  $V_{\text{OC}}=0.79 \text{ V}$ ,  $J_{\text{SC}}=40.370 \text{ mA}/\text{cm}^2$  and  $\text{FF}=83.85\%$ . The optimized thickness and doping density of Si are 350  $\mu\text{m}$  and  $1.0 \times 10^{17} \text{ cm}^{-3}$ , in turn. The addition of 0.2  $\mu\text{m}$  of NiO BSF can effectively enhance  $V_{\text{OC}}$  to  $\sim 0.1 \text{ V}$ . The optimization of p-Si base thickness results a  $\text{QE} \sim 86.83\%$  for the photons wavelength of 1000 nm. The changes in other parameters like working temperature, series and shunt resistances have notable impact on the operations of the device.

Therefore, a p-Si wafer based heterojunction arrangement like this could lead to an enhanced efficiency in future solar market.

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**Data Availability** No datasets were generated or analysed during the current study.

## Declarations

**Ethics Approval** Not applicable.

**Consent to Participate** All authors agreed to this submission.

**Consent for Publication** All authors have read and agreed to the published version of the manuscript.

**Competing Interests** The authors declare no competing interests.

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