



# Spectroscopic and Attenuation Shielding Studies on $B_2O_3$ - $SiO_2$ -LiF-ZnO- $TiO_2$ Glasses

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## Abstract

Effects of  $TiO_2$  addition on the spectroscopic and nuclear shielding properties of lithium fluoride zinc titanate borosilicate glasses with the form  $59B_2O_3$ - $29SiO_2$ - $2LiF$ -( $10-x$ ) $ZnO$ - $xTiO_2$   $x = 0, 2, 4, 6, 8, \text{ and } 10$  mol % labelled as G1-G6, respectively were investigated. The amorphous state of this system was experimentally evaluated using X-ray diffraction. The spectroscopic properties of the recent glasses positively correlated to the increase in  $TiO_2$  content. Molar Refractivity ( $R_m$ ), molar Polarizability ( $\alpha_m$ ), reflection loss ( $R_L$ ), metallization criterion ( $M$ ), electron Polarizability ( $\alpha^\circ$ ), ionic concentration ( $Ni$ ), titanium-titanium separation ( $dTi-Ti$ ), inter ionic distance ( $Ri$ ), inter-nuclear distance ( $ri$ ), and polaron radius ( $rp$ ) of the synthesized glasses were determined. Our results show that both the band gap and refractive index increased with  $TiO_2$  content while Urbach energy decreased. The nuclear radiation shielding properties of these glasses were explored by utilizing Phy-X/PSD simulations. The lower value of the (MFP) sample contains higher  $TiO_2$  content, so good glasses for  $\gamma$  radiation attenuation are accessible. Fast cross section neutron removal of these glasses enhanced as the concentration of  $TiO_2$  increased. Gamma-ray shielding properties of these glasses were compared with different conventional shielding materials and commercial glasses. Therefore, the investigated glasses have potential uses in gamma shielding applications.

**Keywords** Borosilicate · Phy-X/PSD · Optical · Attenuation · FNRCS

## 1 Introduction

The specifications of materials that can serve a double function over the last century can be regarded as among the disguised targets of several researchers. Transparent glasses can be used in radiation shielding materials as a dual function. For

the development of optically transparent radiation shielding materials, significant numbers of glass research labs are rising day-by-day. These materials are used in optical communication, modern optical devices, and radiation shielding materials where protection is needed from radiation [1–8].

Due to its remarkable features such as good thermal stability, hardness, chemical stability, and so on, borosilicate glasses have become the best substitute for concrete shielding, and borosilicate glasses have improved physical characteristics like transparency and refractive index. Titanium oxide, including borate, silicate, borosilicate, and phosphate glasses, one of the most significant additives in glass systems can be considered. The addition of  $TiO_2$  has been reported to increase the glass system's network stability and mechanical characteristics. Glasses having higher amount of titanate confirmed the structural unit  $TiO_4$ ,  $TiO_5$  and  $TiO_6$  and their physical features depend on their number of coordinates [9].

To modify the characteristics of the glass, glass modifiers are incorporated to the glass network. Alkali halides [10–12] such as LiF, and some transition metals like ZnO, and  $TiO_2$  are included in modifiers. The melting temperature of glass and its viscosity can be affected by glass modifiers while retaining its

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chemical structure, as well as affecting its optical, mechanical, thermal, and shielding ability. To generate mobile ion species  $\text{Li}^+$ ,  $\text{Na}^+$ , etc., halides such as NaF, LiF, are introduced into the glass matrix. So that, halide glasses are excellent reagents for metal ions. Glasses incorporating halide ions have long been studied due to their unique physical characteristics.

The emergence of  $\text{TiO}_2$  into the glass network improved the glass's optical, mechanical, thermal, and shielding characteristics [12, 13]. Because of the good ionic conductivity of these glasses, it is extremely possible to use them in UV optics, solid-state batteries, and radiation protection. Depending on their concentration in the glass matrix, intermediate oxides such as  $\text{TiO}_2$  can act as either a glass modifier or former.  $\text{TiO}_2$  improves the mechanical strength and radiation protection of the host glass matrices. These glasses possess lower photon energy and greater refractive index than other glasses. The significant development of lithium fluoride zinc titanate borosilicate glasses is very important scientifically and technologically. In this study the  $\text{TiO}_2$  increased at expense of ZnO. With an increasing in  $\text{TiO}_2$  content the composition changes in the glass network and increase the glass network connectivity. Due to this increase will be transform of Si–O–Zn into Si–O–Ti, and Zn–O bond strength is (73 kcal/mol) is much lower than Ti–O (73 kcal/mol) [14]. This is clearly demonstrated on the modification role of a  $\text{TiO}_2$  in the glass system. The introduction of  $\text{TiO}_2$  into the glass network enhanced the optical and physical properties of the glass.

It is the best study for the preparation of lithium fluoride zinc titanate borosilicate glasses and their structural, mechanical, shielding radiation, and optical properties. Thus, it is possible to find the lithium fluoride zinc titanate borosilicate glasses are suitable for use in environments exposed to radiation. The purposes of this research are to identify the attenuation proficiency of lithium fluoride zinc titanate borosilicate glasses using Phy-X/PSD [15] and XCOM software and to identify the mechanical and structure of these glasses to determine their suitability as gamma-ray shielding materials.

## 2 Experimental Processes and Techniques

Glass samples in Table 1 were formulated using the melt-quench method in the chemical formula  $29\text{SiO}_2 - 2\text{LiF} - 59\text{B}_2\text{O}_3 - (10-x)\text{ZnO} - x\text{TiO}_2$ , where  $x=0, 2, 4, 6, 8, \text{ and } 10$  mol %.  $\text{SiO}_2$ , LiF,  $\text{H}_3\text{BO}_4$ , ZnO and  $\text{TiO}_2$  are the initial materials for obtaining these glasses. All the initial materials that have been acquired from Sigma-Aldrich Company. Through grinding the blend repetitively to obtain a fine powder, the starting materials were blended. First, to eliminate  $\text{H}_2\text{O}$  and other impurities, the base materials were heated to  $650^\circ\text{C}$  for 1 h. The heat was increased for 45 min to  $1200^\circ\text{C}$ . To reduce the internal stresses, the samples were annealed at  $450^\circ\text{C}$  for 2 h and left to cool slowly to ambient temperature. The glass samples amorphous state was checked using X-ray diffraction (A Philips X-ray diffractometer

**Table 1** Chemical compositions of the prepared glasses (mol, %)

Sample name	Chemical Composition				
	$\text{B}_2\text{O}_3$	$\text{SiO}_2$	LiF	ZnO	$\text{TiO}_2$
G 1	59	29	2	10	0
G 2	59	29	2	8	2
G 3	59	29	2	6	4
G 4	59	29	2	4	6
G 5	59	29	2	2	8
G 6	59	29	2	0	10

PW/1710). Optical parameters were predictable by using spectrophotometer type JASCO, V-670 (Japan). In the research, Phy-X/PSD software was used to compute all radiation parameters, and these parameters are calculated by following equations: The

mean free path (MFP) was predictable as  $MEP = \left(\frac{1}{\mu}\right)$ ,.

Electron density ( $N_{\text{eff}}$ ) was predictable as:  $N_{\text{eff}} = N \frac{Z_{\text{eff}}}{\sum_i F_i A_i}$ ,

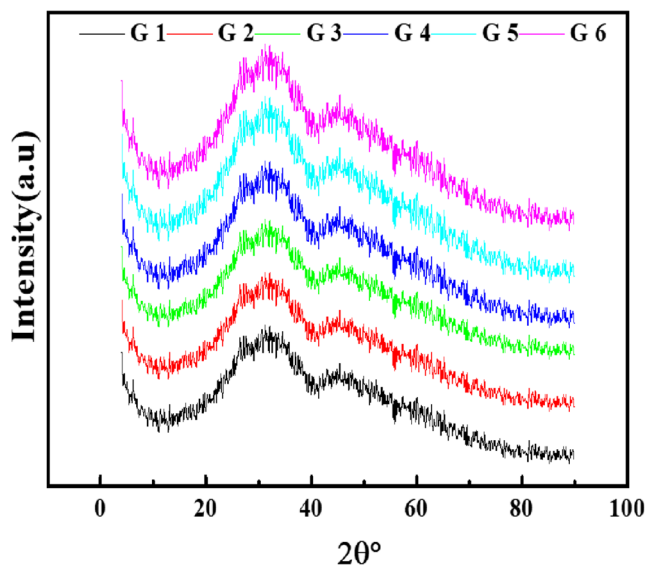
Effective cross-section of removal ( $\Sigma_R$ ) predicted as:

$$\left(\frac{\Sigma_R}{\rho}\right) = \sum_i w_i \left(\frac{\Sigma_R}{\rho}\right)_i \quad \text{and} \quad R = \sum_i \rho_i \left(\frac{R}{\rho}\right)_i$$

## 3 Results and Discussions

### 3.1 Physical Observations

As exemplified in Fig. 1, no discrete lines, no sharp peaks, were confirmed by XRD patterns and signify that the samples have a high amorphous state. The width of the halo varies through one sample to another, but no indication of the crystalline phase has been shown in all the samples.



**Fig. 1** XRD of the studied glasses

Titanium ion concentration computed as,

$$Ti^{+4} = \left( \frac{6.023 \times 10^{23} \times \text{mol fraction of cation} \times \text{valency of Ti}}{V_m} \right), \text{ where}$$

$V_m$  molar volume, titanium ion concentration increased due to the molar volume decrease. Inter-ionic distance  $R_i = \left( \frac{1}{\text{Concentration of Ti}} \right)^{\frac{1}{3}}$ , inter-nuclear distance ( $r_i$ ) and polaron radius ( $r_p$ ) computed as  $rp = \frac{1}{2} \left( \frac{\pi}{6N} \right)^{\frac{1}{3}}$ ,  $ri = \left( \frac{1}{N} \right)^{\frac{1}{3}}$  [16–32]. Titanium – Titanium separation ( $dTi-Ti$ ) computed as  $(dTi-Ti) = \left( \frac{V_m^B}{N} \right)^{\frac{1}{3}}$  and  $V_m^B = \frac{V_m}{2(1-2X_n)}$ , is the borate molar volume in 1 mol and  $x_n$  is considered as the mole fraction of  $B_2O_3$  and  $N$  is Avogadro’s number. It has been confirmed that these perceived values reduce with  $TiO_2$  due to the molar volume decrease.

Molar refractivity ( $R_m$ ), molar polarization ( $\alpha_m$ ), and loss of reflection  $R_L$ , [33–37] of the prepared glasses are computed as:  $R_m = V_m \left( 1 - \sqrt{E_{opt.}/20} \right)$ ,  $\alpha_m = \left( \frac{3}{4\pi N} \right) R_m$ ,  $R_L = \left( \frac{R_m}{V_m} \right)$ , where  $E_{opt.}$  Optical bandgap. The values of these parameters decreased due to the molar volume decrease. The indicator for metallization is estimated as  $M = 1 - \frac{R_m}{V_m}$ . Because of the reduction in molar volume and ( $R_L$ ), metallization is reduced. In Table 2, the above predicted results are shown.

The coordinated average number is a significant criterion for BO or NBO connection confirmation and characterized as  $m = \sum n_{ci} X_i$  where cation coordination is  $n_{ci}$  and  $X_i$  is mole fraction of the glass compoenet. It was noticed that m increases with increase in  $TiO_2$  content. Calculate the number of bonds per unit as  $n_b = \frac{N_A}{V_m} \sum n_{ci} X_i$ . It was discovered that perceived  $n_b$  through  $TiO_2$  content increased.

The bulk module ( $K$ ) and glass transition temperature ( $T_{g(thero.)}$ ) are straight comparative to optical bandgap  $E_{opt}$  and predictable as  $K_{th} = -478.93 + 200.13 E_{opt.}$ ,  $T_{g(thero.)} = -701.87 + 403.33 E_{opt.}$ . It suggested that these increment with  $TiO_2$  because of an increase in bandgap energy.

Cohesive energy values are predicted as  $CE(Kcal/mol) = 18.17 + 4.53 * E_{opt.}$  (e. V) or  $CE(e. V/atom) = 0.792 + 0.198 *$

$E_{opt.}$  (e. V).  $E_{opt.}$  has a direct proportion with cohesive energy. Consequently, with the rise in the  $TiO_2$  content, the cohesive energy value rises.

The theoretical prediction of band-gap energy  $E_{opt}^{Th} = \sum E_{opt} X_i$ ,  $E_{opt}^{Th} = 3.71 * \chi$  and  $E_{opt}^{Th} = 1.08 * \chi^2 - 1.36$ , where  $X_i$  is mole fraction of the glass compoenet and  $\chi$  is electronegativity. The theoretically bandgap rate increases with an increasing of  $TiO_2$ . The whole rising energy bandgap values is probable because of the formation of the BO.

Theoretical optical basicity  $\Lambda_{th}$  can be computed for glass samples  $\Lambda_{th} = \sum \Lambda_{th} X_i$  where  $\Lambda$  is optical basicity of oxides and  $X_i$  fraction of each oxide. This decrease in the values of the  $\Lambda_{th}$  is because of the differences of  $\Lambda_{th}$  between  $TiO_2$  (0.97), and ZnO (1.13) [33–37].

The two-photon absorption factor TPA ( $\beta$ ) cm / GW, expressed as  $\beta = 36.67 - 8.1 E_{opt}$  where  $E_{opt}$  bandgap energy.  $\beta$  value was observed to reduce with the increase in the content of  $TiO_2$  because of increase in  $E_{opt.}$

The ionic and covalent character of glasses could be estimated as the electronegativity difference.  $\Delta X = \sum X_i \Delta X_i$ , where  $\Delta X_i = X_O - X_M$ ,  $I_b = [1 - e^{(-0.25)(X^2)}]$  where  $I_b$  Ionicity. With the increment in  $TiO_2$  content and reduce in covalency  $I_c$ , it was found that ( $I_b$ ) enhanced.

The glass network is influenced by the total number of mechanical constraints and computed as  $N_{con} = N_{bs} + N_{bb}$  where  $N_{bb}$  is bond bending constraints and  $N_{bs}$  is bond stretching,  $N_{bb} = \frac{\sum xi m}{2}$ ,  $N_{bs} = \sum xi(2m - 3)$ . The  $N_{con}$ ,  $N_{bs}$  and  $N_{bb}$ . With an increment in  $TiO_2$ , the overall constraints of  $N_{con}$  are expected to increase. Floppy modes considered as  $M_f = 2 - \frac{5m}{6}$ , cross-linking density  $D_{CL}$  considered as  $D_{cl} = N_{con} - 2$ ,  $CN_{eff} = \frac{2}{5} N_{con} + 3$ . Results are calculated increase with increasing in  $TiO_2$  content. From the data result, it can therefore be suggested that the glass ‘s retained its 2D network with an increment in  $TiO_2$ .

Lone-pair electrons (L) play a major role in the formation of glass and thermal stability. It calculated as  $L = V - m$

**Table 2** Various physical parameters of the studied glasses

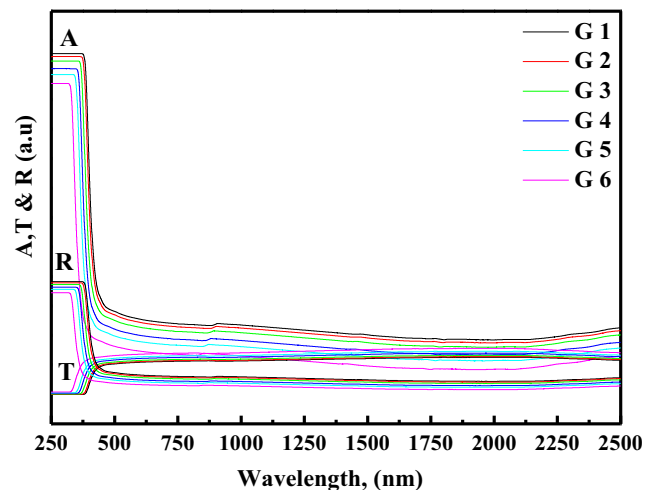
Samples	G 1	G 2	G 3	G 4	G 5	G 6
Molar Refractivity $R_m$ (cm <sup>3</sup> /mol)	15.52	14.745	14.15	13.36	12.7	12.22
Molar Polarizability $\alpha_m$ (Å <sup>3</sup> )	6.18	5.85	5.6	5.3	5.04	4.85
Reflection loss ( $R_2$ )	0.587	0.5865	0.5862	0.586	0.585	0.5848
Metallization criterion (M)	0.413	0.4135	0.4138	0.414	0.4144	0.4147
Electron Polarizability ( $\alpha^\circ$ )	2.68	2.67	2.67	2.67	2.669	2.668
Ion conc. ( $N_i$ ) (10 <sup>21</sup> ions/cm <sup>3</sup> )	–	1.92	3.99	6.35	8.9	11.56
Inter ionic Distance $R_i$ (Å <sup>o</sup> )	–	8.179	6.41	5.49	4.91	4.5
Inter-nuclear distance, $r_i$ (Å)	–	9.474	7.436	6.385	5.71	5.24
Polaron radius, $r_p$ (Å)	–	2.72	2.135	1.8	1.64	1.5
Ti-Ti separation( $d_{Ti-Ti}$ ), nm	0.5	0.492	0.483	0.47	0.46	0.45

Where,  $V$  is valence electron number and  $m$  coordinated average number. The values of  $L&V$  are directly related to the  $\text{TiO}_2$ . This is because of an increase in the rigidity of the glass matrix. The increment in the  $L$  values confirmed that there is a reduction in stress energy. The increment number of coordinates could further signify the above. In Table 3, the above predicted results are shown.

### 3.2 Optical Investigations

Figure 2 exemplifies the absorbance (A) and transmittance (T) and reflectance (R) of glass samples. Optical absorption was quantified to understand the optical characteristics of the prepared glasses [16–32]. Then, with the same approach, we computed the required absorption coefficient for these glasses. There have been indications of decreasing the absorption coefficient as Fig. 3.

$E_{opt}$  is observed as  $\alpha$ .  $h\nu = C(h\nu - E_{opt})^s$  where  $s = 2$ .  $E_{opt}$  values were computed and are illustrated in Fig. 4 and Table 3. The indirect optical band that was collected was

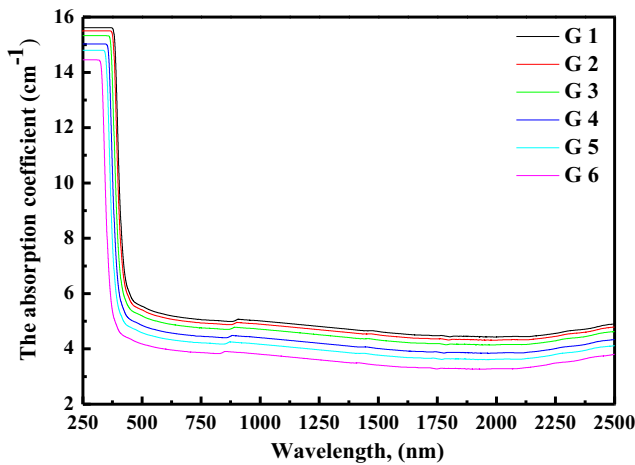


**Fig. 2** Absorbance (A), Transmittance (T), and Reflectance (R) of prepared glasses a function of wavelength and with varying quantities of  $\text{TiO}_2$

identified to be 3.41 for G 1, 3.42 for G 2, 3.425 for G 3, 3.43 for G 4, 3.435 for G 5 and 3.44 for G 6. Particularly,

**Table 3** Various physical and optical parameters of the studied glasses

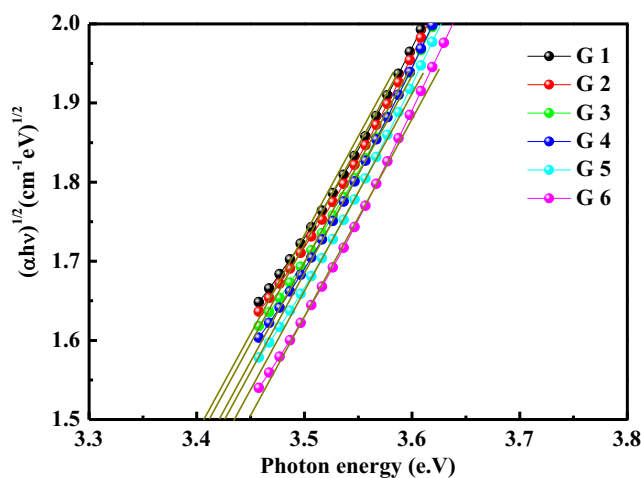
Samples	G 1	G 2	G 3	G 4	G 5	G 6
Average coordination number (m)	3.74	3.78	3.82	3.86	3.9	3.94
Number of bonds per unit volume $n_b$ ( $10^{29} \text{ m}^{-3}$ )	8.52	9.05	9.53	10.2	10.8	11.4
Indirect bandgap (e.V)	3.41	3.42	3.425	3.43	3.435	3.44
Number of oxygen atom	2.45	2.47	2.49	2.51	2.53	2.55
Glass transition temperature (K)	673.5	677.52	679.5	681.55	683.57	685.6
Bulk modulus (Gpa)	203.5	205.5	206.5	207.52	208.52	209.52
Electronegativity ( $\chi$ )	0.9166	0.9193	0.92	0.92	0.92	0.92
Theoretically Electronegativity ( $\chi$ )	1.5116	1.5138	1.516	1.5182	1.5204	1.5226
Optical basicity ( $\chi_{th}$ )	1.2417	1.24035	1.24	1.239	1.23834	1.2376
Theoretically optical basicity ( $\chi_{th}$ )	0.6517	0.6473	0.6429	0.6385	0.6341	0.6341
The two-photon absorption coefficient (TPA)	9.139	9.058	9.0175	8.977	8.9365	8.896
Ionicity ( $I_b$ )	0.435	0.436	0.437	0.438	0.439	0.44
Covalency ( $I_c$ )	0.565	0.564	0.563	0.562	0.561	0.56
bandgap theoretically (e.V)	8.054	8.06	8.066	8.072	8.078	8.078
Duffy bandgap theoretically (e.V)	5.6	5.616	5.62	5.63	5.64	5.65
Di-Qurato bandgap theoretically (e.V)	2.215	2.223	2.244	2.258	2.273	2.288
Cohesive Energy, $CE_{ind}$ (Kcal/mol)	33.62	33.66	33.685	33.71	33.73	33.75
Cohesive Energy, $CE_D$ (Kcal/mol)	1.467	1.469	1.47	1.471	1.472	1.473
Bond-stretching constraints, $N_{bs}$	1.87	1.89	1.91	1.93	1.95	1.97
Bond-bending constraints, $N_{bb}$	2.24	2.28	2.32	2.36	2.4	2.44
Total number of constraints, $N_{con}$	4.11	4.17	4.23	4.29	4.35	4.41
The floppy modes, $M_f$	1.12	1.15	1.18	1.22	1.25	1.28
The cross-linking density, $D_{CL}$	2.11	2.17	2.23	2.29	2.35	2.41
Effective coordination number $CN_{eff}$	4.644	4.668	4.692	4.716	4.74	4.764
The number of the valance electrons, $V$	3.15	3.19	3.23	3.27	3.31	3.31
The lone-pair electrons number, $L_p$	6.89	6.97	7.05	7.13	7.21	7.25
Urbach energy (Eu) (eV)	0.565	0.556	0.544	0.52	0.509	0.48



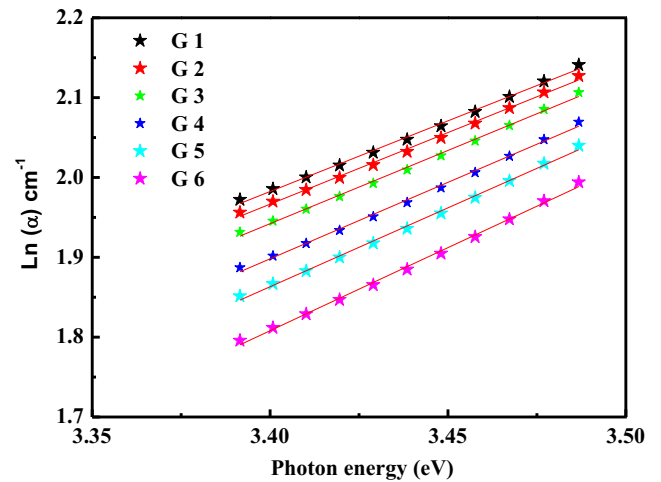
**Fig. 3** The absorption coefficient ( $\text{cm}^{-1}$ ), of prepared glasses a function of wavelength and with varying quantities of  $\text{TiO}_2$

the optical band gap suggests the required energy to absorb light, where it requires more energy to be absorbed than in the bandgap. It has been suggested that the energy gap is increasing with the rise in  $\text{TiO}_2$  content. The raise in the optical bandgap, confirmed that more bridging oxygen (BO) was developed. Hence, (BO) bind energized electrons more strongly than non-bridging oxygen (NBO). As the content of  $\text{TiO}_2$  increased, this increase resulted in a change in the composition of the glass matrix and increased the interconnection. Figure 5 exemplifies the logarithm of absorption coefficient as a function of energy for determining the Urbach energy from the curve slope. In Table 3, the observable Urbach energies are computed. This reduces as the concentration of  $\text{TiO}_2$  raised.

The refractive index of manufacturing glass ( $n$ ) was already computed as:  $n = \frac{(1-R)^2+k^2}{(1+R)^2+k^2}$  where  $k = \alpha\lambda/4\pi$ . ( $n$ ) of manufacturing glasses obtainable in Fig. 6, it was observed



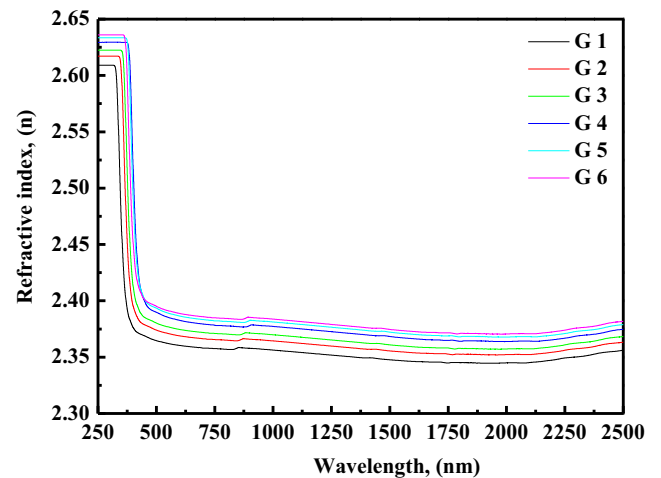
**Fig. 4** Plot of  $(\alpha h\nu)^{1/2}$  against photon energy ( $h\nu$ ) to calculate the direct optical band gap from the intercept of the curves



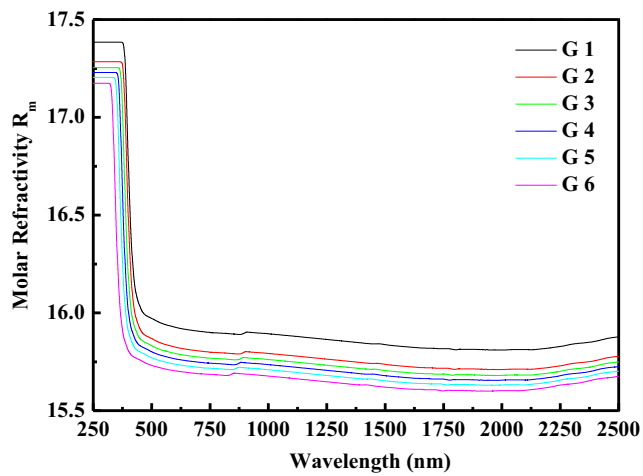
**Fig. 5** Dependence of  $\ln(\alpha)$  upon the photon energy ( $h\nu$ ) for the prepared glasses

that the refractive index of the investigated glasses are increases as density increase. There is a direct correlation between the density and refractive index, i.e., the denser the glass study, the greater the refractive index. So, the refractive index is exactly applicable to reflectance and density, and the molar volume is inversely comparable.

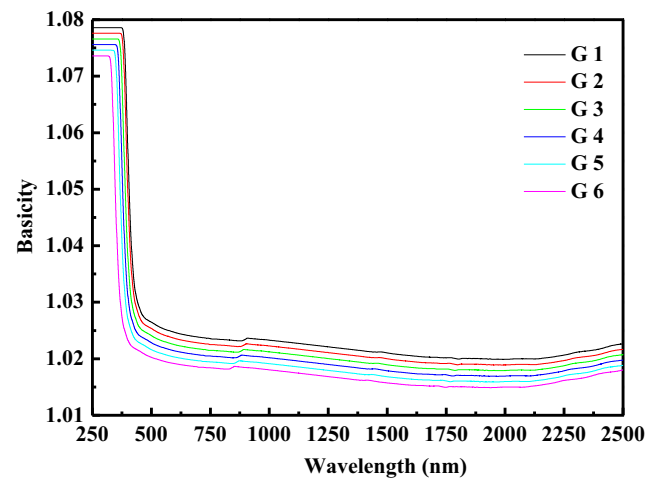
Molar polarization  $R_m$ , and polarizability  $\alpha_m$  of glasses were expected as [33–37]  $R_m = \langle n^2 - 1 | n^2 + 2 \rangle V_m$ ,  $\alpha_m = (3 | 4\pi N) R_m$ , and  $\alpha_0^{2-} = \left[ \frac{V_m (n^2 - 1)}{2.52 (n^2 + 2)} - \sum \alpha_{cat} \right] N_o^{2-}$ . The optical fundamentality of the samples prepared was linked to polarization;  $\Lambda = 1.67 \left( 1 - \frac{1}{\alpha_0} \right)$ . Figs. 7, 8, 9 presented the molar Polarizability  $R_m$ , polarizabilities  $\alpha_m$  and optical basicity  $\Lambda$  respectively of the samples prepared, it was observed that the same trend of refractive index with concentrations of  $\text{TiO}_2$  increased.



**Fig. 6** Refractive index of prepared glasses a function of wavelength and with varying quantities of  $\text{TiO}_2$



**Fig. 7** Molar refractivity of prepared glasses a function of wavelength and with varying quantities of  $\text{TiO}_2$



**Fig. 9** Optical basicity of prepared glasses a function of wavelength and with varying quantities of  $\text{TiO}_2$

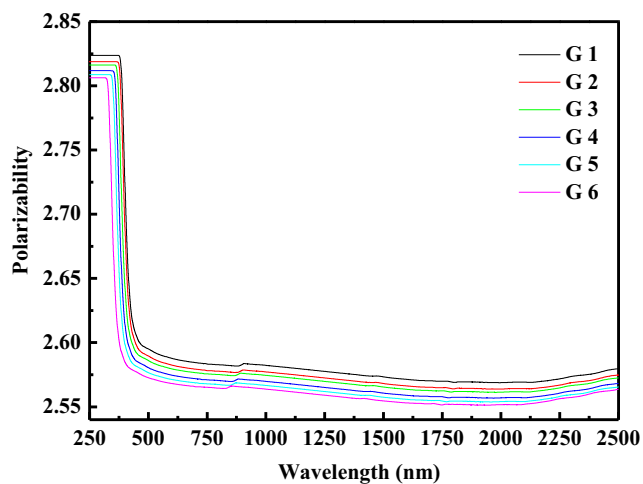
### 3.3 Photon Shielding Features

In this article, the level of protection was evaluated by raising  $\text{TiO}_2$  at the expense of  $\text{ZnO}$  with the composition  $59\text{B}_2\text{O}_3 - 2\text{LiF} - 29\text{SiO}_2 - (10 - x)\text{ZnO} - x\text{TiO}_2$ , ( $0 \leq x \leq 10$ ). In Fig. 10 the mean free path (MFP) was shown. It mentioned that the values of (MFP) are increasing as photon energy increases [38–45]. MFP is an indispensable parameter in assessing a material's gamma shielding competence, and commonly, the lower this value, the higher the shielding performance concerning sample thickness requirements. This inspiration indicated that the growth in energy makes the photon capable of transmitting samples intentionally [38–45]. The lower value of the (MFP) sample contains higher  $\text{TiO}_2$  content, so good glasses for  $\gamma$  radiation attenuation are accessible. The comparison MFP of glasses with standard materials is shown in Fig. 11.

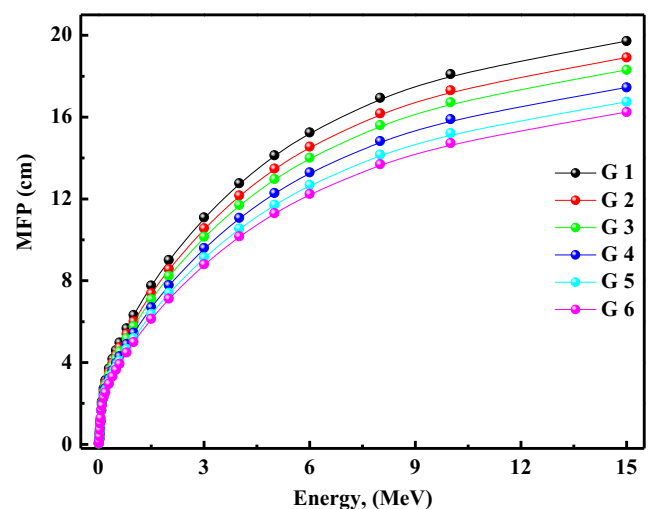
These findings confirmed that glass containing a higher  $\text{TiO}_2$  content had the greatest shielding efficiency.

To evaluate of absorbed doses and dose rates in medical dosimetry for shielding materials, assessment of effective electron density ( $N_{\text{eff}}$ ) is essential. Figure 12 The ( $N_{\text{eff}}$ ) values of samples against energy were illustrated. With the rise of energy, it is shown that ( $N_{\text{eff}}$ ) reduces and then rises. For this reduction, the Compton scattering interaction is accountable. The increase in ( $N_{\text{eff}}$ ) is attributed to the impact of creating pairs at higher energy levels with the increase in  $\text{TiO}_2$  content. Figure 13 demonstrates the pair production of glass system as a versus photon energy. It mentioned that the pair production of these samples in increased with an increase energy.

Figure 14, 15 illustrations the ASC and ESC of glass samples against the energy of photons. It is proposed that the ASC and ESC values will reduce with the increase in energy. Because of the Compton scattering interaction, this reduction



**Fig. 8** Polarizability of prepared glasses a function of wavelength and with varying quantities of  $\text{TiO}_2$



**Fig. 10** The MFP for the prepared glasses as a function of photon

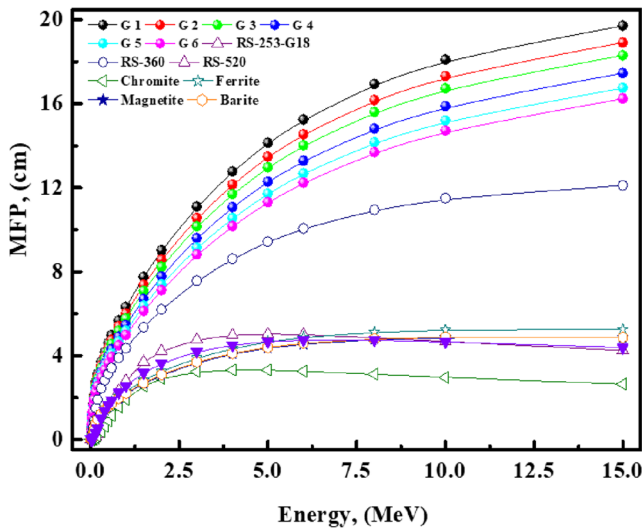


Fig. 11 The comparison MFP of glasses with standard materials

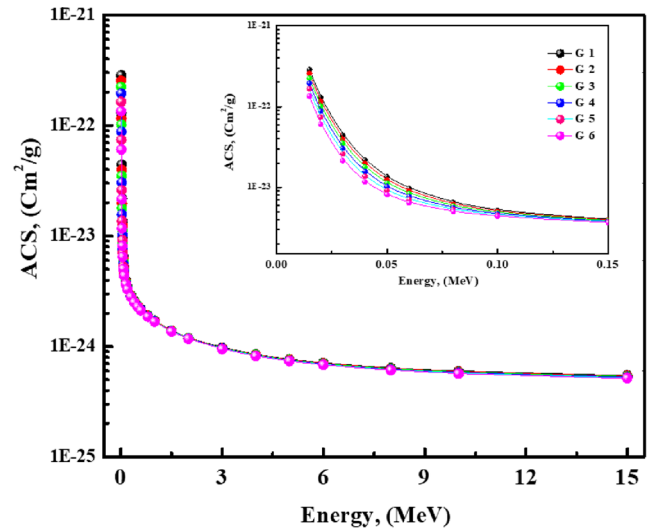


Fig. 14 The ASC for the prepared glasses as a function of photon energy

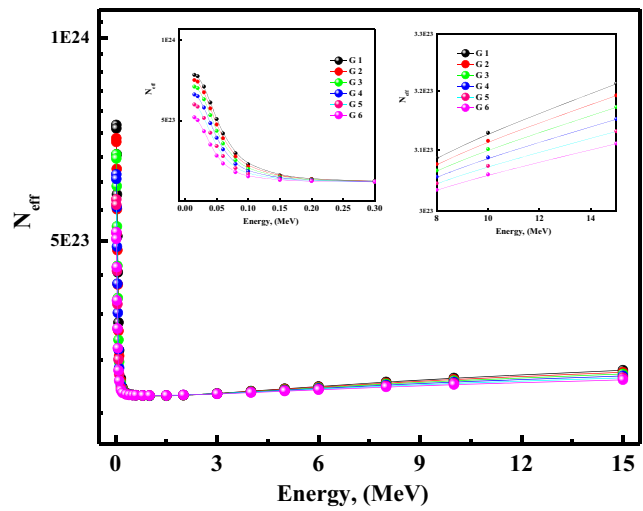


Fig. 12 The ( $N_{eff}$ ) for the prepared glasses as a function of photon energy

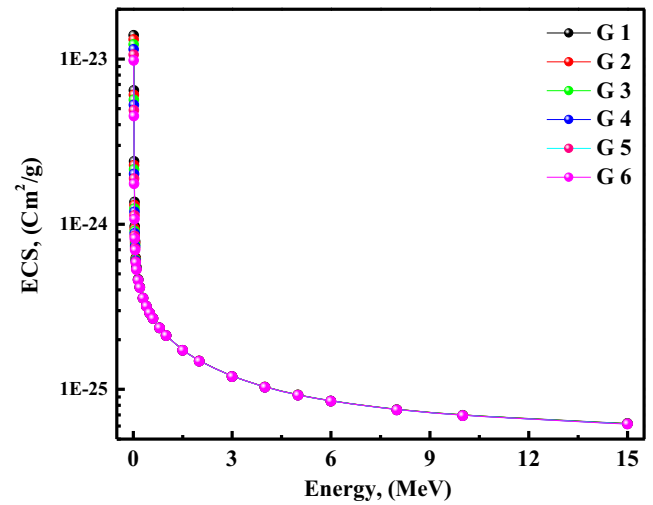


Fig. 15 The ECS for the prepared glasses as a function of photon energy

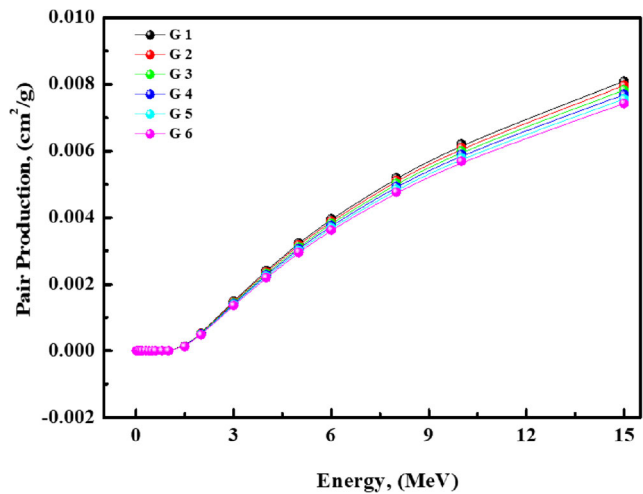


Fig. 13 The pairs production of the prepared glasses as a function of photon energy

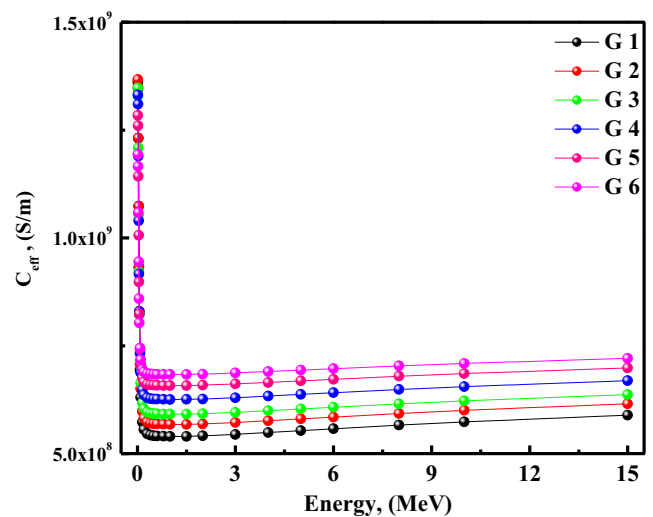
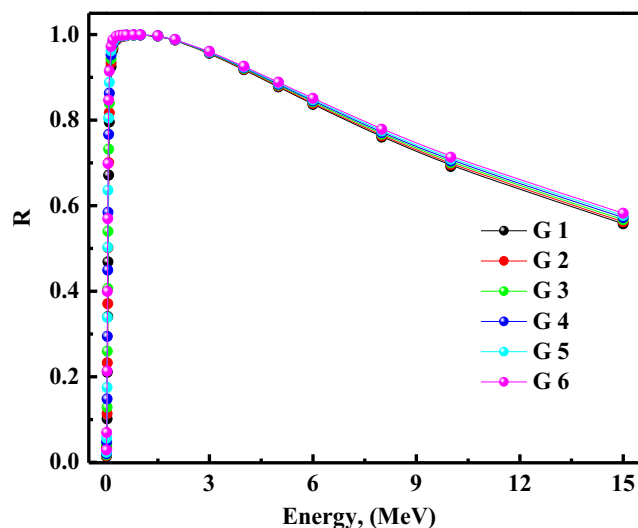


Fig. 16 The  $C_{eff}$  for the prepared glasses as a function of photon energy



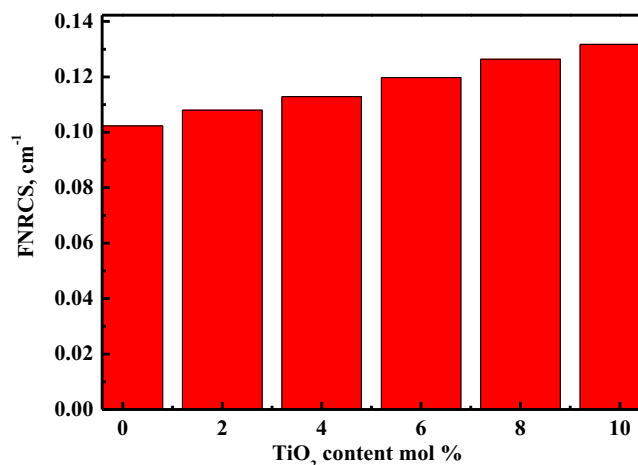
**Fig. 17** Effective removal cross-through ( $\Sigma R$ ) for the prepared glasses as a function of photon energy

occurs. Figure 16 depicted the  $C_{eff}$  against the gamma energy of these glasses. It is proposed that  $C_{eff}$  will reduce with the rise in photon energy. The increase in  $C_{eff}$  corresponded to the impact of pair-creation.

Cross-through effective removal ( $\Sigma R$ ) of examined glasses with energy categorised in Fig. 17. It is indicated that the ( $\Sigma R$ ) increased at lower energy. At greater energy, a reduction in the value of ( $\Sigma R$ ) are detected. Fast cross section neutron removal (FNRCs) is obtainable in Fig. 18. It mentioned that, with  $TiO_2$ , FNRCs enhanced.

## 4 Conclusions

The effect of different  $TiO_2$  concentrations on the spectroscopic and nuclear shielding properties of lithium fluoride



**Fig. 18** FNRCs for the prepared glasses comparison with standard materials

zinc titanate borosilicate glasses with the form  $59B_2O_3-29SiO_2-2LiF-(10-x)ZnO-xTiO_2$ ,  $x=0, 2, 4, 6, 8,$  and  $10$  mol %. The amorphous nature of the glasses was determined by XRD measurements. The coordinated average number, Ionic concentration, bulk module ( $K$ ), glass transition temperature, cohesive energy, theoretically bandgap, ionicity, mechanical constraints, floppy modes, lone-pair electrons, metallization criterion, number of bonds, and electronegativity variation with a composition also increase because of the differences in electronegativity of  $TiO_2$  and  $ZnO$ . Molar volume, covalency, molar refractivity, molar polarizability, reflection loss, electron polarizability, inter ionic distance, inter-nuclear distance, polaron radius, and Ti-Ti separation variation with a composition also decrease because of the molar volume decrease. Our results show that both the band gap and refractive index increased with  $TiO_2$  content while Urbach energy decreased. The nuclear radiation shielding properties of these glasses were explored by utilizing Phy-X/PSD simulations. The lower value of the (MFP) sample contains higher  $TiO_2$  content, so good glasses for  $\gamma$  radiation attenuation are accessible. Fast cross section neutron removal of these glasses enhanced as the concentration of  $TiO_2$  increased. Gamma-ray shielding properties of these glasses were compared with different conventional shielding materials and commercial glasses. Therefore, the investigated glasses have potential uses in gamma shielding applications.

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**Imed Boukhris, Imen Kebaili, and M.S. Al-Buriah:** These authors were listed as they reviewed this paper scientifically and linguistically and helped me in the responses to the reviewers.

**Availability of Data and Material** My manuscript and associated personal data will be shared with Research Square for the delivery of the author dashboard.

## Declarations

The manuscript has not been published elsewhere and that it has not been submitted simultaneously for publication elsewhere.

**Conflict of Interest** The authors declare that they have no conflict of interest. The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

**Consent to Participate** The authors consent to participate.

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