#### **ORIGINAL PAPER**



# Effect of SiO<sub>2</sub> Addition on Chromium Transitions in Borate Glasses

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### Abstract

Melt-quenching technique was used to prepare borosilicate glasses of composition  $xSiO_2$ - (75-x)  $B_2O_3$ -24.7Li<sub>2</sub>O-0.3Cr<sub>2</sub>O<sub>3</sub> (x = 0, 10, 20, 30, 40 and 50 mol%). With increasing the former content, both  $Cr^{6+}$  and  $Cr^{3+}$  optical transitions undergo strong intensity variations. The crystal field strength (10Dq) is increased, while Racah parameters (C,B) are decreased with increasing SiO<sub>2</sub> content. The ratio estimated Dq/B values confirmed a weak crystal field environment for  $Cr^{3+}$ ions, and further suggest a more covalent bond character. The ESR results revealed that Si-contained samples exhibit a strong microwave absorption signal, and justified the presence of both  $Cr^{3+}$  and  $Cr^{6+}$ (charge transfer into  $3d^0 2p^{6+} = 3d^1Cr^{5+}$ ) oxidation states. The fingerprints of borate and silicate relevant structural groups were clearly identified from FTIR spectroscopy.

Keywords Borosilicate glasses · Chromium hexa/trivalent · Ligand field theory · ESR spectra · FTIR

# **1** Introduction

Glasses doped with transition metal (TM) oxides are interesting materials for several potential applications, due to their technologically unique properties and commercial availability. For example, these glasses could be integrated as active components in diverse optical, electrical, thermal, and mechanical devices [1–7]. Furthermore, TM ions play a central role nowadays in probing glass and/or crystalline structures, since they own outer *d*-orbitals of broad radial distribution responsible for their superior sensitivity to minute changes in the surrounding actions [8–11]. Additionally, materials having multivalence states, such as Cr for instance, are convenient as cathode elements in rechargeable batteries, since they also offer high energy density and capacitance [12, 13].

Among several transition metals, Cr<sup>3+</sup> ions are commonly used as activators in different luminescence materials and to color glasses, particularly when introduced in small quantities

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in the glass matrix, beside their large influence over other relevant physical properties such as the insulating strength of the glass samples [12-17]. Furthermore,  $Cr^{3+}$ ions in  $3d^{3}$  configurations have excellent VIS-NIR broadband luminescence, laser performance (suitability for solid state lasers) and applications in solar collectors [1, 2, 6, 18, 19]. The special glass characteristics, such as high glass forming ability, lower glass transition temperature, high thermal stability, and transmission, could be utilized by using alkali ions in varying types and quantities. In this context, Li<sub>2</sub>O ions turn out to be of particular interest given their smallest ionic radius compared to other alkali ions, which have shown to enhance photoluminescence in colored glasses [20]. Moreover, the addition of SiO<sub>2</sub> to alkali borate glasses is known to enhance its durability, hardness, and the transmission to UV light [21]. Hence the mixed former alkali-borosilicate glasses offer promising host to accommodate different modifying oxides and ions from transition metal oxides (TMO), such as Cr<sub>2</sub>O<sub>3</sub>. Although alkali borate glasses are highly transparent to visible light, their strong optical absorption in ultraviolet UV region significantly mask relevant high energy optical transitions [17]. Furthermore, the availability of boron in both BO3 and BO4 structural units and the strong B-O covalent bonds render borates glasses highly stable [22]. Chromium ions normally participate in glass networks with CrO<sub>4</sub>  $(Cr^{6+})$ ,  $CrO_6$   $(Cr^{3+})$  and to lesser extent with  $CrO_3$  structural units [21, 23]. Among varies oxidation states (e.g., Cr<sup>2+</sup>, Cr<sup>3+</sup>, Cr<sup>5+</sup>or Cr<sup>6+</sup>) chromium ions are most frequently found in two

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valence states. These are  $Cr^{3+}$ , which together with  $CrO_6$ structural units act as modifiers, and  $Cr^{6+}$  acting as network formers with  $CrO_4$ . Glasses doped with  $Cr^{3+}$  and  $Cr^{6+}$  ions revealed distinct colors which can be observed by UV-VIS spectroscopy [18, 24]. In the present study, lithium borosilicate glasses doped with  $Cr_2O_3$  were prepared. These glasses were characterized by FTIR, UV–VIS Optical absorption. The latter allowed the extraction of ligand filed parameters (Dq, Racah interelectronic repulsion parameters, B and C). Additionally, the electron spin resonance (ESR) was used to shed the light into the different Cr oxidation states.

# **2 Experimental Setup**

Following the melt quenching technique, the glass compositions  $xSiO_2$ - (75-x) B<sub>2</sub>O<sub>3</sub>-24.7Li<sub>2</sub>O-0.3Cr<sub>2</sub>O<sub>3</sub>(x = 0, 10, 20, 30, 40and 50 mol%) were prepared. The mixtures were ground together in an agate mortar and transferred into a porcelain crucible. The batches were melted in an electrical furnace under ordinary atmospheric conditions, in the temperature between 1100 and 1200 °C for one hour until a bubble free liquid was formed. The molten were quenched rapidly to room temperature between two polished copper plates. Green colored glasses with high optical quality and transparency were obtained. The amorphous nature of the prepared glasses was ensured from X-ray diffraction (XRD) measurements, performed using XRD-6000 Shimadzu X-ray diffractometer, with a Cu K $\alpha$  radiation of wavelength  $\lambda = 1.54$  Å for an angular range between 5 and 80°. Optical absorption spectra of the samples were recorded using JenWay-6405-type UV-VIS spectrophotometer in 190 to 1000 nm spectral range. Electron spin resonance (ESR) data were collected using EMX Brukertype spectrometer, operating at the X-band frequency with 100 kHz field modulation and 10 mW microwave power. The magnetic field was varied between 75 to 6000 G. For reasonable quantitative spin analysis, the weight of the powder is kept fixed for all samples. Infrared absorption spectra were recorded in the range 400-1600 cm<sup>-1</sup> using FTIR Nicolet 6700 spectrometer, on powdered samples pressed into disks after mixing with high purity KBr.

## 3 Results and Discussion

## 3.1 X-Ray Diffraction

Figure 1 shows X-ray diffraction (XRD) patterns for Si = 0, 30 and 50 mol% samples, taken as representative examples. The spectra exhibit no traces of crystalline phases, i.e., no continuous or discrete sharp peaks, but rather the amorphous characteristic broad hump at  $2\theta \sim 23^{\circ}$ , which is a typical feature of borate glasses. The absence of well-defined diffraction peaks

rules out the occurrence of long-range atomic arrangement and, therefore, clearly confirms the glassy nature of the samples [25–30]. The second weak halo around  $2\theta \sim 45^{\circ}$ , which is better defined for the SiO<sub>2</sub>-free sample, is commonly associated to the phase separation character of borate glasses [30].

The intensity of XRD hump at  $2\theta \sim 23^{\circ}$  was found to decrease with the amount of Silicon oxide, with an incremental broadening of the hump, reflecting an enhanced amorphous character of the prepared samples and substantial loss of the short-range order [31, 32]. This behavior is quantified in the inset of Fig. 1, by depicting the variation of the area under the hump, obtained from a standard Gaussian fitting, with the SiO<sub>2</sub> content.

### 3.2 Optical Properties

Figure 1 shows the optical absorption spectra of all glass samples taken at room temperature. An overall shift of the absorption edge of ~0.1 eV towards lower wavelengths (i.e., blue shift) was observed while increasing the SiO<sub>2</sub> content. Additional spectral features related to discrete Cr transitions were observed, the identification of which calls for a detailed deconvolution process for all glass samples, as exemplified at the inset of Fig. 2. Specifically, the spectra exhibited strong absorption bands at ~338 and ~370 nm which are characteristic for chromium ions in Cr<sup>6+</sup> oxidation state, and are attributed to  ${}^{4}A_{2g} \rightarrow {}^{4}T_{1g}$  and  ${}^{4}A_{2g} \rightarrow {}^{2}A_{1g}$  optical transitions, respectively [1-5, 24, 33]. These bands undergo spectral sharpening and intensity enhancement with SiO<sub>2</sub> contents, as shown Fig. 2. Additionally, they were found to shift to higher wavelength (338-348 nm and 370-390 nm) with increasing the silicate oxide content. The origin of the two UV bands assigned to Cr<sup>6+</sup> oxidation state, stems from charge transfer with p levels [1-3, 34, 35], which in the presence of crystal field splits into two states that show up as two optical transitions in the visible and UV spectral regimes [1-3, 34-36].

Indeed, optical transitions in the visible regime within the transparent glass host is responsible for the glass color. Two such bands are located at ~421 and ~619 nm as shown in Fig. 2, which are resulting from d-d transitions of  $Cr^{3+}$ ions (d<sup>3</sup>) in octahedral symmetry glass environment (strength of crystal field and covalency), and are assigned to  ${}^{4}A_{2g} \rightarrow {}^{4}T_{1g}$ and  ${}^{4}A_{2g} \rightarrow {}^{4}T_{2g}$  transitions, respectively [1-3, 34-37]. Additionally, the spectra contain a weak fine structure in the form of spectral dips at ~770 nm. These faint Fano antiresonances dips are present because of the of spin-orbit interaction and the level mixing between  ${}^{2}E$  and  ${}^{4}T$  and  ${}^{2}T_{1}$  and  ${}^{4}T_{2}$  [1–3, 34–38]. Contrary, the band located at ~421 nm predominantly defines the yellowish color of the glass, while the ~619 nm band contributes to the greenish color [1-3, 18]. The intensity of the absorption bands at ~421 and 619 nm increases with increasing SiO<sub>2</sub> content. This trend is fully supported by the observed glass color transformation from dark-





20°



**Fig. 2** Optical absorption spectra for all investigated samples of composition xSiO<sub>2</sub>- (75-x) B<sub>2</sub>O<sub>3</sub>-24.7Li<sub>2</sub>O-0.3Cr<sub>2</sub>O<sub>3</sub> and the inset is Si = 40 and Cr-free Si = 40 samples. Brown, green and orange curves are the deconvolution results from the Si = 40 sample

green towards yellow-green with  $SiO_2$  content. The two absorption bands at 421 and 619 nm are shifted to higher wavelength (i.e., red shift) with increasing  $SiO_2$  content.

The optical absorption spectra presented in Fig. 2 provide additional information regarding the interactions within the host, such as the crystal field strength (10Dq). The band located at ~619 ( $v_1$ ) defines the magnitude of 10Dq [1–3, 16, 25, 37];

$$10Dq = v_1 \tag{1}$$

The Racah parameters B &C defining the inter-electron Coulomb repulsion between 3d electrons of  $Cr^{3+}$ ions, are obtained following the equations [1–3, 15, 33]:

$$B = \frac{(2\nu_1 - \nu_2)(\nu_2 - \nu_1)}{(27\nu_1 - 15\nu_2)} \tag{2}$$

$$C = \frac{v_3 - 4B - v_1}{3}$$
(3)

where  $v_2$  refers to the band positioned at ~421 nm and  $v_3$  is the position of Cr<sup>6+</sup> band at the average sum of 337 and 369 nm bands. By substituting  $v_1$ ,  $v_2$  and  $v_3$ , in (cm<sup>-1</sup>) energy units, obtained from the deconvolution process, into Eqs. (1), (2), (3), the 10Dq, B and C ligand field parameters were obtained and presented in Fig. 3 (a,b). The 10Dq and (B, C) parameters clearly follow an opposite behavior with SiO<sub>2</sub> content, reflecting a relatively weaker inter-electronic repulsion within d-shell. Consequently, the chemical bonds between the ligands and  $Cr^{3+}$  ions acquire more ionic character [1–3, 19, 38]. This may be illustrated by the slightly longer Cr-O distances in glass materials than oxide crystals [6]. The values of B are much lower than  $B = 1030 \text{ cm}^{-1}$ , i.e., the value of the free  $Cr^{3+}$  ion [7]. It is well-established that the relation between Racah parameters is  $C \approx 3.6$  B, which is reasonably satisfied for the present system as listed in Table 1. Similarly, the ratio Dq/B defines the strength of the ligand field, where Dq/B less or greater than 2.3 indicates the weak or strong crystal field, respectively. In the present studied

**Table 1**The estimated band gap energy, ligand field parameters (C/B,<br/>Dq/B), and the bond formation parameter(h) as a function of SiO2 content

$SiO_2$ content (mol. %)	Band gap energy	Ligand field parameters		
	E <sub>g</sub> (eV)	C/B	Dq/B	h
0	3.63	3.62	1.99	0.54
10	3.63	3.59	2.05	0.65
20	3.73	3.57	2.10	0.74
30	3.70	3.57	2.13	0.80
40	3.74	3.53	2.15	0.83
50	3.72	3.51	2.17	0.87

glass, the ratio of Dq/B changes from 1.99 to 2.17 with increasing of SiO<sub>2</sub> content [1–3, 10, 39, 40]. This indicates that  $Cr^{3+}$ ions are in the weak crystal field regime, and further suggests more covalent bond character for this glass system [17, 33]. This can be supported by estimating the bonding formation (the nephelauxetic, h) of the ligand encoded within Racah parameter B as [1–3, 8, 25]:

$$h = \frac{\left[ \left( B_{free} - B \right) / B_{free} \right]}{K_{cr^{3+}}} \tag{4}$$

where,  $B_{free}$  is Racah parameter for the free (gaseous)  $Cr^{3+}$  ion and k is the central  $Cr^{3+}$ ion, which take the values  $B_{free} =$ 918 cm<sup>-1</sup> and  $K_{Cr3+} = 0:21$  [1–3, 36, 38]. Therefore, for all glass samples, B is lower than  $B_{free}$ , as given in Table 1. The calculated nephelauxetic h parameter is increased with SiO<sub>2</sub>content, indicating a gradual increase in the covalent environment of  $Cr^{3+}$  as well as increased d-electrons localization [24, 26, 38].

The measurement of the overall absorption edge is a useful average method for the investigation of the momentumintegrated electronic structure in crystalline and amorphous/ glassy materials [39]. The absorption coefficient ( $\alpha$ ) stands as the main factor in estimating the ( $\alpha$ hv)<sup>0.5</sup> and ln( $\alpha$ ) parameters for large and low absorption, respectively. The optical band gap (E<sub>g</sub>) can be estimated according to the equation [1–3, 41]:

$$\alpha h \nu = A \left( h \nu - E_g \right)^n \tag{5}$$

The best fit to the optical absorption edge was obtained for n = 2, which indicates allowed indirect band transition for the glass samples. The values of Eg as obtained from this fitting for all investigated samples are given in the Table 1. With increasing SiO<sub>2</sub> contents, the optical band gap is enlarged consistent with the reduced number of NBO's suggested later from FTIR results. On the other hand, the reduced number of non-bridging oxygens and the rigidity reduce the degree of localized states within the glass network resulting in wider band gap and shifting the absorption edge towards lower wavelength [21, 25]. Fluctuations of the gap energy within ~0.1 eV was obtained for this glass system with average value of  $3.68 \pm 0.05$  eV. In fact, the higher amorphous degree reported from XRD data for SiO<sub>2</sub> contained samples, which signals reduced band gap size, compensates the large gap induced by NBO's, so that the final gap size is only partially  $(\pm 0.05 \text{ eV})$  changed with SiO<sub>2</sub> additives.

#### 3.3 ESR Spectral Studies

The room temperature ESR spectra for x = 0, 20 and 40 mol% samples doped with fixed 0.3 mol%  $Cr_2O_3$  are shown in Fig. 4. It is observed that all Cr characteristic resonances lose intensity with increasing SiO<sub>2</sub> content. The ESR spectra

**Fig. 3** Ligand field parameters composition dependence; (**a**) crystal field strength (10Dq), and (**b**) Racah parameters (B.C)



exhibit resonances with effective g values at 4.82 (1430 G), 4.08 (1688 G), 2.25 (3061 G) and 1.93 (3571 G), where G refers to the magnetic field in Gauss units. The first resonance at ~1430 G confirms the presence of isolated Cr<sup>3+</sup> ions subject to strong crystal field in distorted octahedral configuration [1, 2, 9, 13, 17, 24, 42, 43]. The second strong absorption at ~1688 G has its origin in  $Cr^{3+}$ ions [1, 2, 44, 45]. The magnetic resonance of these glasses arises from the 3d<sup>3</sup> electron in paramagnetic Cr<sup>3+</sup>ions (The Cr<sup>3+</sup> free ions have 3d<sup>3</sup> electrons and  ${}^{4}F_{3/2}$  ground state). Similar to optical absorption spectra, the Cr<sup>3+</sup>ions ground state in an octahedral crystal field is <sup>4</sup>A<sub>2g</sub> level [14]. On the other hand, higher field absorption located at ~3061 G is due to exchange coupled pairs or large  $Cr^{3+}$ clusters [14, 17, 23]. The last resonance at  $\sim$ 3571 G (g = 1.93) confirms the presence of  $Cr^{5+}$  ions, through the identification of its charge transfer state (both naturally existing and the ones produced after the charge transfer between  $Cr^{6+}$  and  $O^{2-}$ ) [1, 2, 34, 35, 39]. Notice that  $Cr^{5+}$  signal at g = 1.93 follows an unexpected behavior. Specifically, its intensity decreases with increasing the SiO<sub>2</sub> content in clear contrast with optical absorption data. In order to resolve this apparent conflict, the ESR spectra of Cr-free samples with SiO<sub>2</sub> content x = 0 and x = 40 mol% are depicted in the inset of Fig. 4. It is clear that the Si-contained sample exhibits a strong microwave absorption signal centered at ~3000 G. Therefore, the superposition of this signal with the intrinsic  $Cr^{5+}$  resonance is responsible for the inverted behaviour.

Finally, in the presence of lower symmetry field component and spin-orbit coupling, the degeneracy of the spin states is lifted up into two Kramer's doublets with the observed resonances of g values at 2.25 and 1.93. The signal obtained at g = 4.82 refers to isolated  $Cr^{3+}$  ions within the glass

Fig. 4 ESR spectra for selected glass samples (x = 0, 20 and 40 mol%), the inset figure shows the ESR signal of the Cr-free of Si = 0 and Si = 40 samples



network, while the intense signal at g = 4.08 arises due to magnetic exchange coupling between  $Cr^{3+}$  and  $Cr^{3+}$  ion pairs [14, 22, 25, 36].

# **3.4 FTIR Results**

The FTIR absorption spectra of the prepared samples are shown in Fig. 5. The  $SiO_2$ -free sample is characterized by broad bands with complex features which are mainly

consisting of vibrational bands from combined  $BO_3$  and  $BO_4$  structural units, in addition to a weak low energy resonance originating from the modifier and TM-dopant (Li and Cr ions) [46–49]. The main absorption bands are centered at 690 cm<sup>-1</sup> (is generated by the B-O-B bending vibrations of BO<sub>3</sub> units), 960 cm<sup>-1</sup> (could be attributed to B-O asymmetric stretching of tetrahedral BO<sub>4</sub> groups) and 1380 cm<sup>-1</sup> (is assigned to the stretching vibrations of the BO<sub>3</sub> triangles) [46–49]. While the modifier band is found at wavenumbers

Fig. 5 FTIR spectra of all prepared samples



less than 600 cm<sup>-1</sup>, for the lithium containing glasses the weak band at (439-470) cm<sup>-1</sup> is due to Li ionic vibrational units [2, 50], and rather likely overlaps with  $Cr^{6+}$  band [1, 2, 12]. When doped with SiO<sub>2</sub>, a number of drastic changes in intensity and band position take place for all those vibrational modes. Specifically, the modifier band observed below 600 cm<sup>-1</sup> gains intensity, after SiO<sub>2</sub> doping, resulting from Si-O-Si bending modes of bridging oxygen's [49, 51, 52]. Likewise, the BO<sub>4</sub> band shows higher intensity due to the formation of Si-O-Si and B-O-Si bonds which contribute vibrational modes at 950–1065 and  $\sim 1050$  cm<sup>-1</sup>, respectively. These linkages were combined with the band B-O-B tetrahedral structural with increasing SiO<sub>2</sub> in the glass network. All FTIR bands are shifted towards higher wavenumber with increasing the concentration of  $SiO_2$  [2, 53]. In contrast to  $BO_4$  and modifier bands, the BO<sub>3</sub> bands at 690 and 1380 cm<sup>-1</sup> are significantly reduced. This is a clear indication of the fact that BO<sub>3</sub> borate groups and NBOs are decreasing in the glass structure [49, 52]. Of particular importance, is the behaviour of the shoulder around 1240 cm<sup>-1</sup> which is due to B-O asymmetric band and reflects the number of non-bridging oxygens (NBOs) of BO<sub>3</sub>triangles [1, 54, 55]. This shoulder exhibits dramatic intensity reduction until it eventually vanishes for the highest SiO<sub>2</sub> concentration here explored. Such a less NBO's structure might be attributed to the decreases in the degree of localization of electrons and donor centers lead to large energy gaps confirming the optical absorption results [42, 49, 56].

# 4 Conclusion

The optical transitions of Cr ions in borosilicate glass system of composition  $xSiO_2$ - (75-x)  $B_2O_3$ -24.7Li<sub>2</sub>O-0.3Cr<sub>2</sub>O<sub>3</sub> (x =0, 10, 20, 30, 40 and 50 mol%) were investigated. The addition of silicon oxide on borate glasses enhanced their amorphous nature and simultaneously produced suitable modification in Cr<sup>3+</sup> and Cr<sup>6+</sup> optical transitions (i.e., promotes the formation of Cr<sup>6+</sup> in the glassy samples). It was noted that the replacement of Boron oxide by Silicon oxide influences the ESR spectra rather clearly. Indeed, SiO<sub>2</sub> absorbs significant portion of the microwave responsible for the spin transition mechanism, leading to an opposite behavior for the intensity of the transformed Cr<sup>6+</sup> band (i.e. Cr<sup>5+</sup>). The present compositions stand as good candidates for multifunction optical applications such as optical switches and UV-filters.

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