# Microwave absorption properties of  $\text{SiC}(\widehat{a}\text{SiO}_2(\widehat{a})\text{Fe}_3\text{O}_4$  hybrids **in the 2–18 GHz range**

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Abstract: To enhance the microwave absorption performance of silicon carbide nanowires (SiCNWs), SiO<sub>2</sub> nanoshells with a thickness of approximately 2 nm and Fe<sub>3</sub>O<sub>4</sub> nanoparticles were grown on the surface of SiCNWs to form SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids. The microwave absorption performance of the  $\text{SiCaSiO}_2$ @Fe<sub>3</sub>O<sub>4</sub> hybrids with different thicknesses was investigated in the frequency range from 2 to 18 GHz using a free-space antenna-based system. The results indicate that  $\text{SiC}\& \text{SiO}_2\& \text{Fe}_3\text{O}_4$  hybrids exhibit improved microwave absorption. In particular, in the case of an SiC@SiO<sub>2</sub> to iron(III) acetylacetonate mass ratio of 1:3, the microwave absorption with an absorber of 2-mm thickness exhibited a minimum reflection loss of −39.58 dB at 12.24 GHz. With respect to the enhanced microwave absorption mechanism, the Fe<sub>3</sub>O<sub>4</sub> nanoparticles coated on SiC@SiO<sub>2</sub> nanowires are proposed to balance the permeability and permittivity of the materials, contributing to the microwave attenuation.

**Keywords:** silicon carbide nanowires; hybrids; microwave absorption; mechanism; impedance match

## **1. Introduction**

Electromagnetic (EM) pollution by electronic devices and communications equipment has increased interest in investigations of microwave-absorbing materials [1−2]. However, existing absorbing materials often have shortcomings such as high density, a large matching thickness and a narrow absorption band. Efficient microwave absorbers with low weight, good thermal stability, a wide absorption frequency range, and high absorption capability are highly desirable and necessary [3−5]. Silicon carbide (SiC) nanostructures are highly desirable as EM absorbers because of their excellent strength, low density, chemical inertness and good dielectric loss to microwave radiation [6−7]. Compared with other carbon materials, SiC can be used under high-temperature conditions and its loss factor increases with temperature [8−9]. One-dimensional SiC nanowires have been found to possess substantially higher permittivity than SiC nanoparticles [6]. Furthermore, the criteria of being lightweight, thin and absorbing microwaves over a broad frequency bandwidth can be realized by SiC. However, the EM parameters of pure SiC do not satisfy the requirement of impedance matching, which results in strong reflection and weak absorption.

Current microwave absorption research tends to be heavily focused on ferrite-based materials, metal powders, related alloys, and various newly developed nanosized materials because they produce large electric, magnetic, and/or dielectric losses [10]. However, ferrites' intrinsic disadvantages of high density and low strength restrict their widespread application. In addition, their absorption bandwidth is relatively narrow and their magnetic characteristics are lost above the Curie temperature [11].

Fabricating composites by combining low-density dielectric loss materials with strong absorption magnetic loss absorbing materials has been demonstrated as an effective approach to improving the EM-wave absorption ability. Kong *et al.* [12] synthesized  $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>/reduced graphene



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oxide (RGO) hybrids via a solvothermal method. γ-Fe<sub>2</sub>O<sub>3</sub> colloidal nanoparticle clusters were assembled on RGO sheets, which led to effective optimization of EM parameters. This special construction can ensure impedance matching; thus, the material exhibits strong absorbing capability [13]. Fe<sub>3</sub>O<sub>4</sub> is well known to exhibit high chemical stability and a high Curie temperature [14], which suggests that a composite containing  $Fe<sub>3</sub>O<sub>4</sub>$  might exhibit enhanced microwave absorption.

With respect to the fabrication of hybrid materials, the interface characteristics are also important. SiC is well known to exhibit weak hydrophilicity, with a water contact angle of approximately 75°. Such poor wettability limits the excellent performance of SiC when it is used in composites [15]. In our recent work,  $SiO<sub>2</sub>$ -coating-modified SiC nanowires displayed good hydrophilicity because of a large number of silanol groups formed on the surface of the  $SiO<sub>2</sub>$ shells as well as in the interspace between nanowires [16]. In this work,  $\text{SiC}(a)\text{SiO}_2(a)\text{Fe}_3\text{O}_4$  hybrids were produced using a convenient polyol approach. The hybrids show improved microwave absorption compared to pure SiC nanowires. The mechanism of this enhancement is discussed.

## **2. Experimental**

#### **2.1. Materials**

Gangue (SiO<sub>2</sub>, purity > 99%), carbon black, iron(III) acetylacetonate (Fe(acac)<sub>3</sub>), ethyl acetate, and sodium hydroxide were supplied by Sinopharm Chemical Reagent Beijing Co., Ltd. Argon (99.99%) was supplied by Haipu Gas Co., Ltd. Deionized water and ethanol were used in all of the experiments.

#### 2.2. Synthesis of SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids

 $SiC@SiO<sub>2</sub>$  nanowires were produced on a large scale through carbothermal reduction methods using gangue and carbon black as the raw materials [17]. The prepared gray nanowires were treated in an alkaline solution of 1 mol/L NaOH for 1 h to form an ultra-thin  $SiO<sub>2</sub>$  shell on the surface; NaOH solution facilitates the formation of Si–O–Fe bonds and prevents the aggregation of  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles [18]. The preparation procedure of  $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>$  hybrids was as follows:  $100 \text{ mg SiC}(\partial S)$  nanowires were added to 25 mL triethylene glycol solution in an ultrasonic bath to obtain a uniform suspension. The iron precursor,  $Fe (acac)$ <sub>3</sub>, was added to the resulting suspension, and the resulting mixture was sonicated for approximately 30 min to form a well-distributed solution. The mixture was subsequently heated to 280 $^{\circ}$ C at a rate of 3 $^{\circ}$ C·min<sup>-1</sup> under vigorous stirring and argon protection. After cooling to room temperature, 30 mL of ethyl acetate was added to dilute the solution. Finally, the products were magnetically separated, washed with ethanol and deionized water several times and dried in a vacuum oven. To investigate the effect of the  $Fe<sub>3</sub>O<sub>4</sub>$  content on the microwave absorption performance, different mass ratios of  $SiC(\partial S_1O_2)$  to Fe(acac)<sub>3</sub>, i.e., 1:1, 1:2, 1:3, and 1:4, were prepared.

#### **2.3. Characterization**

The phases were identified by X-ray diffraction (XRD) using Cu K<sub>a</sub> radiation (40 kV, 20 mA,  $\lambda = 0.15406$  nm) in the angular range from  $10^{\circ}$  to  $80^{\circ}$  at a scan speed of  $2^{\circ}/\text{min}$ . Fourier transform infrared (FT-IR) spectra were recorded in transmission mode in the spectral scanning range from 2000 to 500 cm<sup>−</sup><sup>1</sup> using a Nicolet-Nexus 470 infrared spectrophotometer. Raman spectra were recorded on an ALMEGA confocal laser micro-Raman spectrometer using Ar-ion laser excitation with a wavelength of 325 nm. X-ray photoelectron spectroscopy (XPS) was performed on a VG Multilab 2009 system. The morphologies of the products were examined by field-emission scanning electron microscopy (SEM; FEI-SIRION, operated at 5 kV). Transmission electron microscopy (TEM) with high-resolution transmission electron microscopy (HRTEM) and selected-area electron diffraction (SAED) were performed on a Tecnai G2 F30 S-TWIN. Magnetic measurements were carried out using a physical property measurement system (PPMS; Quantum Design, PPMS-9).

The relative complex permittivity and permeability as functions of frequency were obtained by the coaxial reflection/transmission method using an Agilent E8363A vector network analyzer in the frequency range from 2 to 18 GHz. The samples for measurement were prepared by milling 50wt% of the samples in a paraffin wax matrix and the experiment was carried out at 50°C. For the test, the obtained mixture was molded into a ring with a 3-mm interior diameter, 7-mm exterior diameter, and 2-mm height. The electromagnetic parameters of the paraffin matrixes filled with samples were investigated by measuring relative complex permittivity,  $\varepsilon_r = \varepsilon' + \varepsilon''$ , and permeability,  $\mu_r = \mu' + \mu''$ , in the frequency range from 2 to 18 GHz. Here the real relative complex permittivity  $(\varepsilon')$  and the imaginary part  $(\varepsilon'')$  of the relative complex permittivity represent the energy storage ability and loss ability, respectively.  $\mu'$  and  $\mu''$  are defined as the real and imaginary parts of the relative complex permeability, respectively. The tangent of dielectric loss  $(\delta)$  and the tangent of magnetic loss  $(\delta_m)$  can be expressed as  $\tan \delta = \varepsilon'' / \varepsilon'$  and

 $\tan \delta_{\rm m} = \mu'' / \mu'$ , respectively [19]. Through a combination of network analysis and transmission line theory [20−21], electromagnetic parameters such as complex permittivity and permeability can be measured.

## **3. Results and discussion**

#### **3.1. Phase and morphology characterization**

The phase of  $SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub>$  hybrids was characterized by XRD, as shown in Fig. 1. For comparison,  $SiC@SiO<sub>2</sub>$  nanowires synthesized in our previous work were also investigated [22]. The major peaks at  $2\theta = 35.7^{\circ}$ , 41.5°, 60.1°, 71.9°, and 75.7° are indexed as the (111), (200), (220), (311), and (222) reflections of SiC (JCPDS card no. 96-101-1032), respectively. A weak broad peak at approximately  $22^{\circ}$ , which corresponds to the amorphous  $SiO<sub>2</sub>$ shells, was observed at high magnification after the NaOH treatment  $[23]$ . The SiO<sub>2</sub> shell plays a key role in the formation of  $SiC(\partial SiO_2(\partial Fe_3O_4$  nanohybrids [24]. After surface modification by Fe<sub>3</sub>O<sub>4</sub>, new peaks at  $2\theta = 30.1^{\circ}$ , 35.5°, 43.1°, 53.5°, 57.1°, and 62.7° appear; these peaks correspond to the (220), (311), (400), (422), (511), and (440) planes of the cubic spinel structure of  $Fe<sub>3</sub>O<sub>4</sub>$  (JCPDS card no. 96-210-1927), respectively. The relative intensity of  $Fe<sub>3</sub>O<sub>4</sub>$ became more distinct with increasing  $Fe<sub>3</sub>O<sub>4</sub>$  content [25]. No other peaks were observed, indicating that  $SiC(\partial S)$  Si $O_2(\partial F)$  nanohybrids were pure.

To obtain detailed information about the structure of the synthesized products, typical SEM and TEM images were collected. Figs. 2(a) and 2(b) represent the SEM images of  $SiC@SiO<sub>2</sub>$  nanowires at low and high magnifications, respectively. The nanowires can be characterized as long and straight filaments possessing smooth surface with a length up to several millimeters. The TEM image (Fig.  $2(c)$ ) shows that the nanowires have a core–shell structure. The corresponding SAED pattern in the inset of Fig. 2(c) indicates that the silica layer is amorphous and that the thickness is approximately 2 nm. After the  $Fe<sub>3</sub>O<sub>4</sub>$  modification on the surface, if we take the  $SiC(\partial_{\Omega}SiO_{2}(\partial_{\Omega}Fe_{3}O_{4} h)$  hybrids with an  $SiC@SiO<sub>2</sub>$ -to-Fe(acac)<sub>3</sub> mass ratio of 1:3 as an example, the surface becomes extremely rough (Figs. 2(d) and  $2(e)$ ). HRTEM analysis (Fig.  $2(f)$ ) indicates that two types of lattice fringes corresponding to the (111) plane of SiC and the (311) plane of  $Fe<sub>3</sub>O<sub>4</sub>$  are observed. This result



Fig. 1. XRD patterns of SiC@SiO<sub>2</sub> nanowires (a) and SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> with SiC@SiO<sub>2</sub> to Fe(acac)<sub>3</sub> mass ratios of **1:1 (b), 1:2 (c), 1:3 (d), and 1:4 (e). The arrow inset shows the magnification part at about 22°.**



**Fig. 2.** SEM (a,b) and TEM (c) images of SiC@SiO<sub>2</sub> nanowires; SEM (d) and TEM (e,f) images of SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids with an  $SiC@SiO<sub>2</sub>$ -to-Fe(acac)<sub>3</sub> mass ratio of 1:3.

is in agreement with the XRD analysis (Fig. 1). The thickness of  $Fe<sub>3</sub>O<sub>4</sub>$  is approximately 5 nm and it is uniformly distributed on the surface. The silica layer is also present, as shown in Fig. 2(f). Even after sonication for an extended period,  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles remain firmly and uniformly anchored onto the surface of  $SiC@SiO<sub>2</sub>$ nanowires, indicating the strong interaction between the two substances.

To further investigate the phase structure of  $SiC@SiO<sub>2</sub>$ nanowires and  $SiCaSiO<sub>2</sub>(aFe<sub>3</sub>O<sub>4</sub>)$  hybrids, FT-IR was carried out, as shown in Fig. 3. The strong bands at 800.99 and 926.50 cm<sup>-1</sup> correspond to stretching vibrations of the Si–C bonds [26]. The peak at  $1107.19 \text{ cm}^{-1}$  corresponds to an Si–O antisymmetric stretching frequency and is attributed to the amorphous silica on the surface [27]. After the Fe<sub>3</sub>O<sub>4</sub> coating was applied, a peak at 583.09 cm<sup>-1</sup> corresponding to an Fe–O vibration appears in the spectrum; this vibration likely originates from  $Fe<sub>3</sub>O<sub>4</sub>$ nanoparticles. Collectively, these spectral results indicate that  $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>$  hybrids were successfully synthesized. The absorption band at  $1622.4 \text{ cm}^{-1}$  is attributed to Si−OH groups in the sample [28].



Fig. 3. FT-IR spectra of SiC@SiO<sub>2</sub> nanowires and SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids.

Fig. 4 shows the Raman spectra of the as-prepared  $SiC(\partial S)$  Si $C_2(\partial S)$  hybrids. In the case of small amounts of coated  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles, two intense peaks are observed at 790 and 950  $\text{cm}^{-1}$ ; these peaks correspond to the transverse phonon (TO) mode and the longitudinal optical (LO) mode of SiC, respectively. The peaks at 1353 and 1580  $cm^{-1}$  are both related to  $sp<sup>3</sup>$  C and  $sp<sup>2</sup>$  C, which may arise from defects in the SiC structure. With increasing amount of  $Fe<sub>3</sub>O<sub>4</sub>$ , the LO mode of SiC at  $958 \text{ cm}^{-1}$  gradually weakens and almost disappears. A similar phenomenon occurs in the cases of the peaks at 1353 and 1580 cm<sup>-1</sup>. By comparison, the intensity

of the peak at  $680 \text{ cm}^{-1}$ , which is the typical characteristic peak of  $Fe<sub>3</sub>O<sub>4</sub>$ , gradually increases, further confirming that  $SiC@SiO<sub>2</sub>(QFe<sub>3</sub>O<sub>4</sub>$  hybrids were successfully synthesized [29−30]. The increase in intensity of  $Fe<sub>3</sub>O<sub>4</sub>$  peaks and decrease in intensity of SiC peaks suggest an increase in the amount of  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles.



Fig. 4. Raman spectra of SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids with SiC@SiO<sub>2</sub>-to-Fe(acac)<sub>3</sub> mass ratios of 1:1, 1:2, 1:3, and 1:4.

Structural information about the synthesized  $\text{SiC}(\partial S)$ and  $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>$  was analyzed by XPS; the corresponding results are shown in Fig. 5. Compared with the XPS spectrum of  $SiC@SiO<sub>2</sub>$  (Fig. 5(a)), that of the  $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>$  hybrid shows an emergent signal corresponding to the characteristic peak of Fe. The two Fe peaks with binding energies of 710.8 and 724.4 eV in the high-resolution Fe 2p spectrum (Fig. 5(b)) are attributed to Fe 2p<sub>3/2</sub> and Fe 2p<sub>1/2</sub> of trivalent Fe in the Fe<sub>3</sub>O<sub>4</sub> [31–32]. In addition, the peak at 530.6 eV is attributed to O 1s; its intensity increases significantly after modification, indicating that Fe<sub>3</sub>O<sub>4</sub> was successfully formed on the SiC $@SiO<sub>2</sub>$  surface. The peak at 284.5 eV is attributed to C 1s. The peaks at 151.7 and 100.8 eV are assigned to Si 2s and Si 2p of  $SiC@SiO<sub>2</sub>$  [33]. Moreover, the intensities of the peaks of C and Si decrease substantially with increasing amount of Fe<sub>3</sub>O<sub>4</sub>, which is attributable to the fact that  $SiC@SiO<sub>2</sub>$  was almost fully covered by  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles.

The field-dependent magnetization curves of  $SiC(\partial_{\Omega}SiO_{2}\partial_{\Omega}Fe_{3}O_{4}$  nanohybrids were also measured at room temperature, as shown in Fig. 6. The  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles endow the  $\text{SiC}(\partial S) = \partial_4 R$  hybrids with ferromagnetism. The saturation magnetization and coercivity were measured to be 13.42, 23.20, 32.64, 43.97 emu $\cdot$ g<sup>-1</sup> and 31.34, 26.21, 32.45, 34.64 Oe, corresponding to  $SiCaSiO<sub>2</sub>$ -to-Fe(acac)<sub>3</sub> mass ratios of 1:1, 1:2, 1:3, and 1:4, respectively. Magnetic  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles contribute to magnetic loss to an EM wave, which results in effective wave-absorbing perfor-



Fig. 5. Full XPS spectra of SiC@SiO<sub>2</sub> and SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids (a) and Fe 2p of SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids (b).



Fig. 6. Magnetic hysteresis loops of SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids with SiC@SiO<sub>2</sub>-to-Fe(acac), mass ratios of 1:1, 1:2, 1:3, **and 1:4.** 

## **3.2. Microwave absorption performance**

Fig. 7 presents the real part  $\varepsilon'$  and imaginary part  $\varepsilon''$ of the complex relative permittivity and the dielectric tangent loss tan $\delta_e$  of SiC $\omega$ SiO<sub>2</sub> nanowires and  $SiC@SiO<sub>2</sub>(QFe<sub>3</sub>O<sub>4</sub>$  hybrids in the frequency range from 2 to

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18 GHz. Compared with the microwave absorption performance of  $SiCaSiO<sub>2</sub>$  nanowires, the values of the complex relative permittivity ( $\varepsilon'$ ,  $\varepsilon''$ ) and dielectric tangent loss tan $\delta_e$  of SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids increase. In addition, the value of tan $\delta_e$  for the sample with an  $SiC@SiO<sub>2</sub>$ -to-Fe(acac)<sub>3</sub> mass ratio of 1:3 is the highest. Notably, in the case of the sample with an  $SiC@SiO<sub>2</sub>$ -to-Fe(acac)<sub>3</sub> mass ratio of 1:1, a sharp shoulder peak appears at the frequency of 15 GHz. This peak is attributed to the presence of fewer coated  $Fe<sub>3</sub>O<sub>4</sub>$ nanoparticles because they well balance the permeability and permittivity of the materials by making their EM impedances match. Thus, most of the EM waves are able to transmit into the materials. The microwaves in the materials are intensively attenuated, which leads to a lower EM reflection. The minimum reflection loss of an EM wave reaches −16.54 dB at 15 GHz, and the absorption bandwidth reaches 3 GHz in the case of an absorber with a thickness in the 1.5-mm range. In the case of an  $SiC@SiO_2$ -to-Fe(acac)<sub>3</sub> mass ratio of 1:4, no dual-loss peaks are observed, which is attributed mainly to the fact that  $\text{SiC}(\partial S)$  is almost fully covered by  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles. The microwave absorption performance of the hybrids mainly arises from the dielectric loss of  $Fe<sub>3</sub>O<sub>4</sub>$ , and attaining an impedance match is difficult. Therefore, the hybrids with a thickness in the 1.5-mm range have no effective absorption band. Given that SiC is a typical dielectric material,  $Fe<sub>3</sub>O<sub>4</sub>$  exhibits dual-loss characteristics. The dipole polarization presented by  $Fe<sub>3</sub>O<sub>4</sub>$ nanoparticles, which mainly results from the natural physical properties of  $Fe<sub>3</sub>O<sub>4</sub>$ , can contribute to the dielectric loss. Dipoles are present in  $Fe<sub>3</sub>O<sub>4</sub>$ , especially when the particle size is nanoscale. The number of surface atoms with unsaturated bonds greatly increases as the size decreases, resulting in an increase of the number of dipoles [35]. In this work, the size of  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles was approximately 5 nm, as characterized in Fig. 2; thus, a large number of dipoles exist. A greater number of dipoles results in a higher dielectric constant. Under the external alternating electric field, the reorientation of dipoles results in polarization and energy dissipation. In this case, a high relative dielectric constant means high dielectric loss because the reorientation must be completed by the migration of ions, which dissipates energy [36]. What is more, the dielectric loss is related to the natural structure of  $Fe<sub>3</sub>O<sub>4</sub>$ . Its chemical formula can be written as  $\text{Fe}_{A}^{3+}[\text{Fe}^{2+}, \text{Fe}^{3+}]_{B}O_{4}^{2-}$ , where Fe cations occupy interstices of a face-centered cubic closed-packed frame of oxygen ions, and A and B denote the tetrahedral sites and octahedral sites, respectively. Electrons can transfer between  $Fe^{2+}$  and  $Fe^{3+}$ ions in B sites as the field (EM wave) is applied. Therefore,

the electronic spin and charge polarization strongly affect the loss [37−38]. Based on the aforementioned analyzes, the higher dielectric loss is related to the dipole polarization

presented by  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles and its natural structures. In addition to these, the interfacial polarization and the related relaxation may also improve the dielectric loss [39].



tangential loss (c) of the SiC@SiO<sub>2</sub> nanowires and SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids with SiC@SiO<sub>2</sub>-to-Fe(acac)<sub>3</sub> mass ratios of 1:0, 1:1, 1:2, **1:3, and 1:4.** 

Fig. 8 presents the real part  $\mu'$  and imaginary part  $\mu''$ of the complex relative permeability and the magnetic tangent loss tan $\delta_{\rm m}$  of the SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids in the frequency range from 2 to 18 GHz. The values of imaginary part  $\mu''$  and magnetic tangent loss tan $\delta_{\rm m}$ increased after  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles were coated in the

low-frequency region from 2 to 10 GHz, indicating that the magnetic loss mainly occurred in the low-frequency range [40]. By comparison, the value of the magnetic loss  $\tan\delta_{\rm m}$  was the highest when the SiC@SiO<sub>2</sub>-to-Fe(acac)<sub>3</sub> mass ratio was 1:4. This result is attributed to the increasing content of  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles.



**Fig. 8. Real part of the complex permeability spectra (a), imaginary part of the complex permeability spectra (b), and the magnetic**  tangential loss (c) of the  $\text{SiC@SiO}_2 \text{@Fe}_3\text{O}_4$  hybrids with  $\text{SiC@SiO}_2$ -to-Fe(acac)<sub>3</sub> mass ratios of 1:1, 1:2, 1:3, and 1:4.

Based on the transmission line theory [41], the reflection loss (RL, dB) of a sample is determined according to the following equations:

$$
Z_{\rm in} = Z_0 (\mu_{\rm r} / \varepsilon_{\rm r})^{1/2} \tanh\left[j(2\pi f d / c)(\mu_{\rm r} / \varepsilon_{\rm r})^{1/2}\right]
$$
 (1)

$$
RL = 20lg \left| \frac{Z_{in} - Z_0}{Z_{in} + Z_0} \right|
$$
 (2)

where  $Z_{\text{in}}$  is the input impedance of the absorber;  $Z_0$  is the impedance of air;  $\mu_r$  and  $\varepsilon_r$  are the relative complex permeability and permittivity of the absorber, respectively; *f* is the frequency of the EM waves; *d* is the thickness of the absorber, *c* is the velocity of EM waves in free space, and *j* is the imaginary unit. When RL is smaller than −10 dB, only 10% of the microwave energy is reflected, whereas its 90% is absorbed. The corresponding frequency range, in which the

RL is smaller than −10 dB, is defined as the effective absorption bandwidth.

The EM absorption performance of  $SiC@SiO<sub>2</sub>$  nanowires and  $SiC(\partial S)O_2(\partial F \partial G_4)$  hybrids is compared in Fig. 9. The EM absorption performance was substantially enhanced after  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles were attached to the  $SiC@SiO<sub>2</sub>$  nanowires. The values of the reflection losses of  $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>$  hybrids with different  $SiC@SiO<sub>2</sub>$ -to-Fe(acac)<sub>3</sub> mass ratios were obtained to be  $-16.54$  (1:1),  $-21.04$  (1:2),  $-39.58$  (1:3), and  $-26.54$  dB (1:4) at 15.44, 6.8, 12.24, and 4.56 GHz, respectively. The optimum state for  $SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub>$  hybrids to attenuate microwaves was obtained when the mass ratio between  $SiC@SiO<sub>2</sub>$  nanowires and Fe(acac)<sub>3</sub> was 1:3. This ratio may well result in balanced permeability and

permittivity of the material by making the microwave impedance match. Most of the microwaves are thus able to be transmitted into the material. The absorption bandwidth with the reflection loss below −10 dB is as high as 3.68 GHz (from 10.8 to 14.48 GHz) for the absorber with a thickness in the 2.0-mm range. Table 1 summarizes the

EM absorption properties for typical related materials [42−46], which reveals that the effective absorption bandwidth in this work is wider while the thickness is much smaller, indicating that the  $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>)$ hybrids synthesized in this work are very promising EM-wave absorptive candidates.



Fig. 9. Calculated reflection loss of SiC@SiO<sub>2</sub> nanowires (a) and SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids synthesized with SiC@SiO<sub>2</sub> **nanowires-to-Fe(acac)**<sub>3</sub> mass ratios of 1:1 (b), 1:2 (c), 1:3 (d), and 1:4 (e).

Sample	Minimum $RL / dB$	Frequency range / GHz $(RL < -10 dB)$	$d_{\rm m}$ / mm	Percentage of absorber / wt%	Ref.
SiC.	$-17.4$	$9.7 - 12.2$	3.0	30	[42]
NiO@SiC	$-46.9$	$8.2 - 12.4$	2.0	50	$[43]$
CNT/SiC	$-37.6$	$12.4 - 17.5$	2.0	20	$[44]$
Fe <sub>3</sub> O <sub>4</sub> /Ppy/CNT	$-25.9$	$7.8 - 12.5$	3.0	20	$[45]$
$Fe3O4(a)ZnO$	$-22.7$	$10.1 - 16.0$	3.5	50	[46]
$SiC@SiO2(@Fe3O4)$ hy- brids in this work	$-39.6$	$10.8 - 14.5$	2.0	50	This work

**Table 1. Comparison of the microwave absorption performance of typical related materials** 

Note: The  $d<sub>m</sub>$  is the effective absorption bandwidth when RL is smaller than −10 dB.

In view of the improved EM wave attenuation performance of  $SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub>$  hybrids, such excellent performance resulted from the natural physical properties of  $SiC@SiO<sub>2</sub>$  and Fe<sub>3</sub>O<sub>4</sub> and the special core/shell structures. There are usually two factors for achieving a low RL value: the incident microwave should transmit into the materials with minimum reflection, and the materials should effectively attenuate the incident microwave [47]. The impedance match between the material and free space is important. This requires that the permittivity and permeability are equal in both the imaginary and the real parts. SiC attenuates microwaves by dielectric loss, which mainly originates from electronic polarization, atomic polarization and intrinsic electric dipole polarization.  $Fe<sub>3</sub>O<sub>4</sub>$  exhibits dual-loss characteristics: it attenuates microwaves by both dielectric and magnetic losses, where the magnetic loss mainly originates from eddy current loss and hysteresis loss. The dipole polarization exhibited by  $Fe<sub>3</sub>O<sub>4</sub>$  nanopar-

ticles contributes to the dielectric loss. Especially under an external alternating electric field, the reorientation of dipoles results in polarization and energy dissipation because the reorientation must be completed by the migration of ions, which dissipates energy [36]. In addition, the special core/shell structure of  $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>$  hybrids plays a vital role in absorbing EM. The interfaces at the core and the shell are clearly observed in the HRTEM image in Fig. 2(c). The interfacial polarization and associated relaxation also contribute to the microwave attenuation [39]. Such improved dielectric loss induced by the interfacial polarization has been observed in other core/shell systems. For example, Ag–Ni core–shell materials showed additional absorption compared to the naked Ni nanoparticles, where the additional absorption was induced by the interfacial polarization [48]. Furthermore, the defects formed at the interfaces due to the crystal mismatch between SiC and  $SiO<sub>2</sub>$  can also serve as polarized centers, leading to an additional dielectric loss. A schematic of the possible EM absorption mechanism for hybrids is shown in Fig. 10. The  $Fe<sub>3</sub>O<sub>4</sub>$  nanoparticles increase the permeability of the complex. Hybrids combine the dielectric loss of  $SiC@SiO<sub>2</sub>$  nanowires with the magnetic loss of Fe<sub>3</sub>O<sub>4</sub>, which may well balance the permeability and permittivity of the materials, making the EM impedance match. Therefore, most of the EM waves can be transmitted into the materials and the microwaves in the materials are intensively attenuated.



 $\sim$  SiC@SiO<sub>2</sub> nanowires  $\bullet$  Fe<sub>3</sub>O<sub>4</sub> nanoparticles  $\sim$  Heat

Fig. 10. Schematic of a possible EM absorption mechanism for the SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids.

# **4. Conclusions**

 $SiC@SiO<sub>2</sub>(@Fe<sub>3</sub>O<sub>4</sub>$  hybrids were successfully produced using a convenient polyol approach. Microstructural analysis revealed that  $Fe<sub>3</sub>O<sub>4</sub>$  nanocrystals were immobilized onto the  $SiC@SiO<sub>2</sub>$  nanowires surface by a strong interaction. The microwave absorption performance was enhanced with the minimum reflection loss of −39.58 dB at 12.24 GHz. In particular, the SiC@SiO<sub>2</sub>@Fe<sub>3</sub>O<sub>4</sub> hybrids with an SiC@SiO<sub>2</sub> nanowires-to-Fe(acac)<sub>3</sub> mass ratio of 1:3 achieved an absorption bandwidth as high as 3.68 GHz (from 10.8 to 14.48 GHz) with a reflection loss below −10 dB for the absorber with a thickness in the 2.0-mm range. The main mechanism was that nanohybrids combined the dielectric loss of  $SiC@SiO<sub>2</sub>$  nanowires with the magnetic loss of Fe<sub>3</sub>O<sub>4</sub>, which may have well balanced the permeability and permittivity of the materials, making the EM impedance match.

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