RESEARCH ARTICLE

Check for updates

Dynamic analysis of microparticle behavior in quad-beam opticfiber optical tweezers

Sheng Hu^{1,2}^(b) · Hai-peng Li¹ · Cheng-liang Zhu^{1,2} · Yong Zhao^{1,2}

Received: 24 March 2021/Accepted: 13 July 2021/Published online: 25 July 2021 © The Optical Society of India 2021

Abstract In recent year, optic-fiber optical tweezers were considered as a valuable manipulated tool, allowing for trapping and sorting numerous interest microparticles in biological and microfluidic application. Especially for dual optic-fiber array with misalignment, vortex optical field can be utilized to cause particle realizing periodic orbital rotation, contributing to measure unknown parameters such as torque, viscosity, and angular momentum. Yet, this optic-fiber array with several misalignments creating more vortex fields remains an interesting work for a large range of microparticle trapping and rotation. Therefore, motion profiles of particle exerting by four optic-fibers forming quad-beam were studied by changing offset distance in our numerical model. Compared to dual beam, the array of such four optic-fibers has much more controllable variables, reflecting the asymmetric radial distance relevant to particle. Herein, these motion characteristics relevant to microparticle have been simulated based on ray optics approximation. Finally, dynamic simulations imply that quad-beam can generate two or more vortexes inducing particle to move in orbital rotation, which have revealed different rotational radius, direction, and velocity. As a results, this optic-fiber array may be readily implemented in cellular trapping and manipulation, exploring biophysical properties of different types of cell.

Sheng Hu husheng@neuq.edu.cn

¹ College of Information Science and Engineering, Northeastern University, Shenyang 110819, China

² Hebei Key Laboratory of Micro-Nano Precision Optical Sensing and Measurement Technology, Qinhuangdao 066004, China **Keywords** Optical tweezers · Optic-fiber optical tweezers · Particle rotation · Applied optics

Introduction

Optic-fiber optical tweezers (OFOTs), as a prominent manipulation approach, enable optofluidic technology to address micro- and nanoparticles of interest in the fields, such as separation, purification, transport, and analysis in recent years [1, 2]. The waveguide of optical fiber not only migrates the majority of particles moving toward desired position [3], but stably traps smaller particles using a low light intensity other than objective lens with high numeric aperture value [4]. Hence, both precise manipulation and strong trapping can be further exploited to understand both the characterization and classification of single cell with free label in life science [5]. Meanwhile, Chen et al. [6, 7] have reported dual-beam optic-fibers with transverse offset realizing the rotational motion of polystyrene bead or cell, which could be used to assist the prediction and diagnosis of disease. In this case, the structure of misalignment related to dual-beam optic-fibers, generating spin and orbital rotation, is more convenient than special optical field, for example circularly polarized light [8, 9], Hermite–Gaussian [10], and Bessel beam [11]. Afterward, Pei and Chen [12] continued to utilize this approach in the measurement of viscosity coefficient. In addition, Li et al. [13] pointed out that different optic-fiber misalignment, optical power, and optic-fiber separation could exhibit three dynamic behaviors. The existence of both optical tweezers and gravity leads to an "8-shaped" trajectory for such a borosilicate particle suspending in air.

In addition, an array of optic-fiber, providing multi-path light source, has already been used in point-of-care

diagnostic platform for medical endoscopes and measurement [14, 15]. To best of our knowledge, it is seldom that optic-fiber array is used to trap and manipulate microparticles. Therefore, we attempt to research on optic-fiber array leading to particle rotation or enrichment through optical radiation pressure. Moreover, this approach based on dynamic simulation offers a detailed theoretical perspective supporting experimental configuration and observation [7, 12, 16]. In this way, the numerical model has been built to study the optic-fiber array with a certain offset distance. Making this quad-beam misalignment, the adjustment referred to both axial and radial distance has been studied to generate various motion behaviors. This provides a technical feasibility serving the characteristic identification of different biological cell due to much more rotational optical fields.

Theory and model

In our research, quad-beam consists of four same opticfibers, and then two pairs are brought face to face in array form. A particle suspending in the solution is exerted by optical pressure from these four optic-fibers, as shown in Fig. 1a. Compared to only one dual beam, quad-beams would be more flexible due to more controllable parameters. In this case, the position of these four optic-fibers can be illustrated in Fig. 1b. To simplify the research work, the numerical model is proposed to study the dynamic behavior of microparticle in two-dimensional zx plane. The coordinate (z_i , x_i) in zx plane corresponds directly to the focal spot position of individual optic-fiber, in which the subscript *i* stands for the corresponding optic-fiber label. The coordinate of a particle should be noted (z_p , x_p). The variable S_0 is a simulated width in square area.

Optical radiation pressure, so-called optical tweezers, is attributed to the interaction between dielectric particle and light to produce an observable dynamic behavior. Owing to wave-particle duality of light, a particle suspending in 657

media acts by optical radiation pressure interpreted with either ray optics (RO) [17, 18] or electromagnetic (EM) [11] model. In comparison between particle size and light wavelength, it is thought as Mie regime in RO approximation to predict optical trapping force if the particle diameter is far greater than wavelength [19]. In contrast, electromagnetic approximation based on Maxwell equation attempted to acquire the strength of the force relying on Maxwell stress tensor or volume-integration method. The calculation of optical radiation pressure relevant to opticfiber has been described by Sidick and co-workers [20]. Since then, many research groups have carried out outstanding work with optic-fiber for particle manipulation by means of their presented theoretical model. Assuming that light with fixed polarization orientation is orthogonal to light propagated from other optic-fiber. Without regard to interference effect, the optical force generated by *i*th opticfiber can be divided into x- and z-axial components, as expressed in Eqs. (1) and (2).

$$F_{zi}(z_{p}, x_{p}) = \frac{n_{1}P_{0}}{c} \frac{2r_{0}^{2}}{\pi} \int_{0}^{\pi} d\varphi \int_{0}^{\theta_{max}} d\theta \sin 2\theta \frac{\exp(-2r^{2}/\omega^{2})}{\omega^{2}R_{c}}$$
$$\times \left\{ q_{s}R_{z} + \frac{q_{g}}{\tan\gamma} \left[R_{z} - R_{c} \frac{R_{z} + r_{0}\cos\theta}{a\cos\gamma} \right] \right\}$$
(1)

$$F_{\rm xi}(z_{\rm p}, x_{\rm p}) = \frac{n_1 P_0}{c} \frac{2r_0^2}{\pi} \int_0^{\pi} d\varphi \int_0^{\theta_{\rm max}} d\theta \sin 2\theta \frac{\exp(-2r^2/\omega^2)}{\omega^2 R_{\rm c}}$$
$$\times \left\{ q_s(r_0 \sin \theta \cos \varphi - d_{\rm ip}) + \frac{q_g}{\tan \gamma} \right.$$
$$\left. \times \left[r_0 \sin \theta \cos \varphi - d_{\rm ip} (1 - \frac{R_{\rm c}}{a \cos \gamma}) \right] \right\}$$
(2)

where n_1 refractive index of medium, c light velocity in vacuum, P_0 power of incident light source, r_0 radius of microparticle, d_{ip} the transverse distance between the

Fig. 1 Schematic structure of quad-beam: **a** Microparticle experienced by optical force from four optic-fibers, **b** Coordinate system for both quad-beam optic-fibers and microparticle in *zx* plane



particle and the optical axis of *i*th optic-fiber. The integral form is assigned to spherical coordinate. Thus, the variables φ and θ are azimuthal and colatitude angle, respectively. The parameters q_s and q_g are related to the fractions of the momentum transferred to the microparticle, as expressed in Eqs. (3) and (4) [19].

$$q_{\rm s} = 1 + R\cos 2\alpha_{\rm i} - T^2 \frac{\cos(2\alpha_{\rm i} - 2\alpha_{\rm t}) + R\cos 2\alpha_{\rm i}}{1 + R^2 + 2R\cos 2\alpha_{\rm t}}$$
(3)

$$q_{\rm g} = -R\sin 2\alpha_{\rm i} + T^2 \frac{\sin(2\alpha_{\rm i} - 2\alpha_{\rm t}) + R\sin 2\alpha_{\rm i}}{1 + R^2 + 2R\cos 2\alpha_{\rm t}} \tag{4}$$

Herein, these intermediate parameters in Eqs. (1) \sim (4) could be calculated by Eq. (5). The detailed reduction of these formulas was cited in Ref. [20]. The sign of R_z is determined by R_c .

$$R = \frac{1}{2} \left[\frac{\sin^{2}(\alpha_{i} - \alpha_{t})}{\sin^{2}(\alpha_{i} + \alpha_{t})} + \frac{\tan^{2}(\alpha_{i} - \alpha_{t})}{\tan^{2}(\alpha_{i} + \alpha_{t})} \right], \qquad T = 1 - R$$

$$R_{c} = z_{p} \left[1 + \left(\frac{\pi n_{1} \omega_{0}^{2}}{\lambda_{0} z_{p}} \right) \right], \qquad R_{z} = \pm \sqrt{R_{c}^{2} - r^{2}}$$

$$\gamma = \cos^{-1} \left(\frac{R_{c}^{2} + a^{2} - r_{0}^{2}}{2aR_{c}} \right), \qquad a = \sqrt{d_{ip}^{2} + (R_{z} + r_{0} \cos\theta)^{2}}$$

$$\alpha_{i} = \cos^{-1} \left[\frac{d_{ip}^{2} + (R_{z} + r_{0} \cos\theta)^{2} - r_{0}^{2} - R_{c}^{2}}{2r_{0}R_{c}} \right], \qquad n_{1} \sin\alpha_{i} = n_{2} \sin\alpha_{t}$$

$$r = \sqrt{d_{ip}^{2} + (r_{0} \sin\theta)^{2} - 2d_{ip}r_{0} \sin\theta \cos\varphi}, \qquad \omega = \omega_{0} \sqrt{1 + \left(\frac{\lambda_{0} z_{p}}{\pi n_{1} \omega_{0}^{2}}\right)}$$
(5)

Noting that value θ_{max} can be obtained by solution for θ in $d_{\text{ip}}^2 + (R_z + r_0 \cos \theta)^2 - r_0^2 - R_c^2 = 0$. Apart from optical radiation pressure, Stokes drag force F_s is used as friction term in Eq. (6).

$$F_{\rm s} = -6\pi r_0 \eta v \tag{6}$$

where η is viscosity of medium and ν is velocity of particle. However, inertial forces are by orders of magnitude smaller than viscous friction effects at microscales and below. Neglecting Brownian stochastic force, Newton's dynamics can be used to calculate the particle behavior in Eq. (7),

$$m\dot{v}(t) = F_{\rm s} + \hat{x}F_{\rm xi} + \hat{z}F_{\rm zi}$$
 (i = 1,...,4) (7)

where m is mass of particle. In our simulation, the differential equation in Eq. (7) could be studied by using Runge– Kutta method in a small temporal discretization.

Results and discussion

Considered the medium containing microparticle, its refractive index n_1 is 1.33, and the viscosity η is 1 mPa·s. The refractive index n_2 for microparticle is 1.59, and its radius r_0 is 5.5 µm. As for incident light source, light wavelength 980 nm with power P_0 equal to 100mW is coupled into the four optic-fibers; beam waist radius ω_0 is

 $3 \mu m$. Herein, the offset distance of these optic-fibers could be studied to simulate the motion behaviors of particles as follows.

Same axial distance

Firstly, the same radial distance ($z_i = z_j$, *i*, *j* = 1,..., 4) regarding four optic-fibers is studied, which have two offset distances at *x*-axial direction. In addition, the radial distance *z* of such optic-fibers is 85 µm concerning original point. Beyond the misalignment at vertical level, the field of optical pressure force can be plotted in Fig. 2a. There are two trapping positions pointed to \pm 10 µm at *x*-axial direction. Relying on dynamic simulation, six particles at different random position could be calculated to describe their own trajectories, as shown in Fig. 2b. These particles below the horizontal level move toward the bottom trapping location, and vice versa.

When the misalignment of four optic-fiber is assumed, there are two vortex fields to induce these particle to perform an orbital rotation within their nearby regions. The position of optic-fibers is (- 85 µm, 15 µm), (- 85 µm, - $6 \,\mu\text{m}$), (85 μm , $-15 \,\mu\text{m}$), (85 μm , $6 \,\mu\text{m}$), respectively. Analyzing one of vortex field, the gradient forces from bottom and top optic-fibers always trap particles moving toward the center position in z-axial direction. Due to the misalignment structure of optic-fiber, the scattering forces from the bottom and top optic-fibers alternately repel the particle away from optic-fiber facet. Thus, the orbital rotation of particle can occur in the dual beam with offset distance, as shown in Fig. 3. It is illustrated that one opticfiber array or more could form more vortexes to cause particle orbital rotation. This have benefited the generation of so many vortical effects to expand the dynamic observation and contributed to compare different types of cell with its rotary velocity, which identifies the biophysical characteristics of biological cells. Furthermore, both the offset distance and the optic-fiber setout can discretionarily be adjusted to form special vortex pattern. Decreasing the distance between four optic-fibers, two vortexes have merged together. Finally, only one vortex leads to the orbital motion for these nearby particles, as shown in Fig. 4. In comparison with the structure in Fig. 3, the distance of left-hand optic-fibers can be lower than the right-hand side. As a result, the top and bottom spins have clockwise and counterclockwise direction, as shown in Fig. 5. The top particles are clockwise motion, but the bottom particles exhibit counterclockwise. Especially, this adjustable function could be expected to utilize in the control of molecular gear-driven motion with co-rotation or anti-rotation [21].

In Chen's work, particle behaviors regarding four different types have been identified including capture, spiral



Fig. 2 Simulation without misalignment of four optic-fibers at symmetric structure: a Optical trapping force for particle with radius of 5.5 µm. Four red arrows on the both sides imply the vertical position of optic-fibers. Meanwhile, the colors and directions of

(a)20

(mm)

(mn)×

arrows represent the magnitudes and directions of optical trapping forces, respectively. Equipotential curve describes strength of optical force. b These motion trajectories of six particles



Fig. 4 Coordinate of opticfibers is (- 85 µm, 12.5 µm), $(-85 \ \mu m, -4 \ \mu m), (85 \ \mu m,$ - 12.5 μm), (85 μm, 4 μm), respectively: a Only one steady vortex field, b Particles sustain same circular motion closing to original point

1: 15 10 10 x (mm)(-10μm, 5μm) (8μm, 9μm) 0 (-2μm, -1μm) (10μm, -15μm) -10 -1(-15 -15 -20 -20 $-25 - 20 - 15 - 10 - 5 0 / z (\mu m)$ -20 -10 0 z(μm) 10 20 15 20 25 5 10 $x 10^{-11}$ (b) 20 (a)20 4.5 15 1(10 (mm) (-12µm, 1µm) (10µm, 9µm) -10 $(-3\mu m, -1\mu m)$ -15 (12µm, -11µm) -2 -20 -20 -10 10 20 0 -25 -20 -15 -10 -5 0 5 10 15 20 25 z(µm)

 $x 10^{-11}$ (b) 20

motion, orbital rotation, and escape. As far as their final position is concerned, both orbital rotation and capture are of great importance which are in favor of study on momentum and rotational inertia. Moreover, the offset distances have been decreased so that vortex field can be restrained. These particles attempt to move toward the region with minimum optical radiation area, as shown in Figs. 6 and 7. Different from Fig. 2, these particles

maintain the motion with tortuous forms instead of straights due to the effect of vortex field. Similar to the position of orbital rotation, the trapping region can occur in the original point, as shown in Fig. 7. This assembly of optic-fiber array is better to capture the particle to satisfy the demand of particle enrichment.

Fig. 5 Coordinate of opticfibers is $(-85 \ \mu\text{m}, 6 \ \mu\text{m})$, $(-85 \ \mu\text{m}, -6 \ \mu\text{m})$, $(85 \ \mu\text{m}, -15 \ \mu\text{m})$, $(85 \ \mu\text{m}, 15 \ \mu\text{m})$, respectively: **a** Two vortex fields with different spin direction, **b** Particles move in either clockwise or anticlockwise motion

Fig. 6 Coordinate of opticfibers is $(-85 \ \mu m, 12 \ \mu m)$, $(-85 \ \mu m, -6 \ \mu m)$, $(85 \ \mu m, -12 \ \mu m)$, $(85 \ \mu m, 6 \ \mu m)$, respectively: **a** Weak vortex field disables the orbital rotation, **b** Each of particles move toward the vortex center

Fig. 7 Coordinate of opticfibers is $(-85 \ \mu m, 12 \ \mu m)$, $(-85 \ \mu m, -3 \ \mu m)$, $(85 \ \mu m, -12 \ \mu m)$, $(85 \ \mu m, 3 \ \mu m)$, respectively: a There is an irregular vortex field due to the superposition of three vortex fields, b Particles could only move toward the original point



Different axial distance

Apart from changing the offset distance relevant to *x*-axial direction, the axial distance in *z*-axial direction can be considered in our numerical model. Due to various possibility for the radial distance of four optic-fiber in *z*-direction, we mainly consider two situations which refer to two combinations, namely fibers 1 and 4, as well as fibers 2 and 3. When the radial distance z_i (i = 1, 4) reduces to 71 µm, but the fibers 2 and 3 are still 85 µm, the vortex field in the asymmetry is inevitable to cause the different particle behaviors on the top and bottom. In Fig. 8a, the vortex field

turns to non-axisymmetric spiral. In the horizontal position, the vortex field in the top and bottom is rather different from Fig. 4a. Compared with each other, the top vortex has a bulge closing to left. It implies that two particles are inclined to move to left-hand side on the top level. Conversely, the bottom particles are close to the right-hand side. According to dynamic simulation, the particle motion can be plotted in Fig. 8b, which have demonstrated the abovementioned speculation. However, the top particles simultaneously move in an approximated ellipsoid. The bottom particles no experienced by strong torques become spiral motion into the trap center. Thus, we also speculate



that the radial distance for fibers 1 and 4 can be controlled to change the rotational radius on the top and bottom region. Cooperating with the adjustment in vertical offset distance, one larger vortex field on the top region than the bottom region can be generated, as shown in Fig. 9a. Therefore, it is obvious to find that the rotational area regarding the top particles is larger than the bottom particles, as shown in Fig. 9b. In Fig. 9b, the position of four particles is similar to Fig. 8b.

Considered the particle enrichment on the center, the misalignment distance can be decreased to play against the top and bottom vortex having a relative balance field. Actually, the quasi-balance position is closed to right-hand side compared to Fig. 7a. In this case, these particles would move toward the horizontal area closing to right-hand side, as shown in Fig. 10. According to this simulated results, this manipulation can be used in the particle microscale migration without the aid of motorized stages by changing the optic-fiber radial distance.

Conversely, the middle of two optic-fibers, namely fiber 2 and 3, has been adjusted for their radial distance z_i . Keeping these particles on same position, the same offset distance illustrates different motion. In Fig. 11, the patterns of the top and bottom vortex occur to exchange so that the rotational radius with respect to the bottom particles could

be greater than the top particles. This verifies that changing the radial distance could realize two different motion results in optic-fiber array. Compared to the change of optic-fibers 1 and 4, the particles are inclined to move toward the left-handed side, as shown in Fig. 12. Decreasing the radial distance between two optic-fibers, the force in the z-axial direction becomes weak, which results in failure to the horizontal motion. Finally, these particles exerted on vertical force only stop on the weakest optical trapping area.

Researching on particle rotational period, the particle A with respect to two kinds of different radial distance, reflecting optic-fibers 1–4 and 2–3, could be compared and plotted in Fig. 13. By applying method for seeking peak, the period in *z*-axial direction is 0.30 s and 0.26 s for particle A in Figs. 9 and 11, respectively. Considered in x-axial direction, the period 0.29 s in Fig. 9 is quite close to 0.27 s in Fig. 11. This implies that shorter major axis of ellipse is higher rotational frequency when particle moves in elliptic orbit. Meanwhile, the movement rule of these bottom particles is consistent with the above result. A particle motion with controllable period behavior based on quad-beam is expected to continue in our future work.

Fig. 9 Coordinate of opticfibers is $(-79 \ \mu m, 14 \ \mu m)$, $(-85 \ \mu m, -5 \ \mu m)$, $(85 \ \mu m, -14 \ \mu m)$, $(79 \ \mu m, 5 \ \mu m)$, respectively: **a** The top vortex field is different from the bottom, **b** The top and bottom particles simultaneously move in periodic movement, but the rotational area on the top is larger than the bottom



Fig. 10 Coordinate of opticfibers is $(-91 \ \mu m, 13 \ \mu m)$, $(-85 \ \mu m, -4 \ \mu m)$, $(85 \ \mu m, -13 \ \mu m)$, $(91 \ \mu m, 4 \ \mu m)$, respectively: a Decentered weak vortex field, b These particles are closed to right-hand side related to original point

Fig. 11 Coordinate of opticfibers is $(-85 \ \mu m, 14 \ \mu m)$, $(-79 \ \mu m, -5 \ \mu m)$, $(79 \ \mu m, -14 \ \mu m)$, $(85 \ \mu m, 5 \ \mu m)$, respectively: The top vortex field is different from the bottom, **b** The top and bottom particles simultaneously move in periodic movement, but the rotational area on the bottom is larger than the top

Fig. 12 Coordinate of opticfibers is $(-85 \ \mu m, 13 \ \mu m)$, $(-91 \ \mu m, -4 \ \mu m)$, $(91 \ \mu m, -13 \ \mu m)$, $(85 \ \mu m, 4 \ \mu m)$, respectively: a Decentered weak vortex field, b These particles are closed to left-hand side related to original point





Conclusion

Attempt to manipulate the motion of microparticle is fascinating work by using optic-fiber tweezers. The particle trapping and rotation behaviors have opened up exciting possibilities for analyzing particle property, size, shape in biomedical application. Thus, an array of optic-fibers having quad-beam has been presented to generate two kinds of effect on the particle motion. In same axial distance, this symmetric structure of optical field could realize two same vortex fields inducing neighboring particle in rotational motion with elliptic orbit. As soon as the offset distance of four optic-fibers change, the trapping location could be either original point or the center of a pair of two optic-fibers. In addition, the axial distance of two different individual optic-fiber could change the top and bottom vortex range, which has benefits for realizing particle motion in different rotational radius. In this end, this quadbeam by using four same optic-fibers, compared to dual beam, could offer an opportunity to generate more vortex or trapping field to manipulate particle motion, which is Fig. 13 The displacement of particle *A* changes with time in Figs. 9 and 11: a Displacement of *z*-axial direction,b Displacement of *x*-axial direction



expected to explore the biophysical properties of various single cells synchronously.

Acknowledgements The work was supported by the National Natural Science Foundation of China (61903069, 62003081), the Scientific and Technical Research Project in Colleges and Universities of Hebei Province (BJ2021101), and the Hebei Natural Science Foundation (F2020501040).

Declarations

Conflict of interest There are no conflicts to declare.

References

- D. Gao, W. Ding, M. Nieto-Vesperinas, X. Ding, M. Rahman, T. Zhang, C. Lim, C.W. Qiu, Optical manipulation from the microscale to the nanoscale: fundamentals, advances and prospects. Light-Sci. Appl. 6, e17039 (2017)
- C.L. Zhang, B.J. Xu, C.Y. Gong, J.T. Luo, Q.M. Zhang, Y. Gong, Fiber optofluidic technology based on optical force and photothermal effects. Micromachines 10, 499 (2019)
- Y. Zhang, H.X. Lei, B.J. Li, Refractive-index-based sorting of colloidal particles using a subwavelength optical fiber in a static fluid. Appl. Phys. Expr. 6, 072001 (2013)
- J.B. Decombe, S. Huant, J. Fick, Single and dual fiber nano-tip optical tweezers: trapping and analysis. Opt. Expr. 21, 30521–30531 (2013)
- T.R. Arey, K.L. Cotner, B. Li, L.L. Sohn, Developments in labelfree microfluidic methods for single-cell analysis and sorting. Wiley Interdisciplinary Reviews-Nanomed. Nanobiotechnol. 11, e1529 (2019)
- X.L. Chen, G.Z. Xiao, K.Y. Yang, W. Xiong, H. Luo, Characteristics of the orbital rotation in dual-beam fiber-optic trap with transverse offset. Opt. Expr. 24, 16952–16960 (2016)
- X. Chen, G. Xiao, X. Han, W. Xiong, H. Luo, B. Yao, Observation of spin and orbital rotation of red blood cell in dual-beam fibre-optic trap with transverse offset. J. Opt. 19, 055612 (2017)
- J.W. Liaw, Y.S. Chen, M.K. Kuo, Spinning gold nanoparticles driven by circularly polarized light. J. Quant. Spectrosc. Radiat. Transfer 175, 46–53 (2016)

- F. Monteiro, S. Ghosh, E.C. van Assendelft, D.C. Moore, Optical rotation of levitated spheres in high vacuum. Phys. Rev. A 97, 051802 (2018)
- F.G. Mitri, Cylindrical particle manipulation and negative spinning using a nonparaxial Hermite-Gaussian light-sheet beam. J. Opt. 18, 105402 (2016)
- F.G. Mitri, Reverse orbiting and spinning of a Rayleigh dielectric spheroid in a J(0) Bessel optical beam. J. Opt. Soc. America B-Opt. Phys. 34, 2169–2178 (2017)
- M.R. Pei, S.L. Jin, X.L. Chen, T.F. Kuang, W. Xiong, X. Han, G.Z. Xiao, H. Luo, Microscopic viscosity measurement using orbital rotation in dual-beam fiber-optic trap with transverse offset. Opt. Eng. 59, 126106 (2020)
- W. Li, N. Li, Y. Shen, Z. Fu, H. Su, H. Hu, Dynamic analysis and rotation experiment of an optical-trapped microsphere in air. Appl. Opt. 57, 823–828 (2018)
- E. Scharf, J. Dremel, R. Kuschmierz, J. Czarske, Video-rate lensless endoscope with self-calibration using wavefront shaping. Opt. Lett. 45, 3629–3632 (2020)
- C. Ozkan Loch, M.A. Eichenberger, M. Togno, S.P. Zinsli, M. Egloff, A. Papa, R. Ischebeck, A.J. Lomax, P. Peier, S. Safai, Characterization of a low-cost plastic fiber array detector for proton beam dosimetry. Sensors 20, 5727 (2020)
- X. Chen, G. Xiao, H. Luo, W. Xiong, K. Yang, Dynamics analysis of microsphere in a dual-beam fiber-optic trap with transverse offset. Opt. Expr. 24, 7575–7584 (2016)
- S.H. Zhang, M. Shao, X. Yang, J.H. Zhou, Calculating optical forces with skew line ray model for Gaussian beam. Optik 203, 163935 (2020)
- A. Devi, A.K. De, Alternate analytic formulation of optical force on a dielectric sphere in the ray optics limit. J. Opt. Soc. America B-Opt.Phys. 35, 244–250 (2018)
- E. Sidick, S.D. Collins, A. Knoesen, Trapping forces in a multiple-beam fiber-optic trap. Appl. Opt. 36, 6423–6433 (1997)
- A. Callegari, M. Mijalkov, A.B. Gokoz, G. Volpe, Computational toolbox for optical tweezers in geometrical optics. J. Opt. Soc. America B-Opt. Phys. 32, B11–B19 (2015)
- Z.S. Jia, Q.P. Guan, H.L. Wang, X.Y. Wang, Theoretical investigation on cis-trans isomerisation of azaphosphatriptycenebased molecular gear. Supramol. Chem. 32, 569–577 (2020)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.