

Finite Boundary Regularity for Conformally Compact Einstein Manifolds of Dimension 4

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Abstract

We prove that a 4-dimensional C^2 conformally compact Einstein manifold with Hölder continuous scalar curvature and with $C^{m,\alpha}$ boundary metric has a $C^{m,\alpha}$ compactification. We also study the regularity of the new structure and the new defining function. This is a supplementary proof of Anderson's work and an improvement of Helliwell's result in dimension 4.

Keywords Regularity · Einstein manifold · Asymptotically hyperbolic · Harmonic coordinates · Schauder estimate

Mathematics Subject Classification 35J25 · 53A35 · 53C21 · 53C25 · 58J05

1 Introduction

In 1985, Charles Fefferman and Robin Graham [\[10](#page-19-0)] introduced a new method to study the local conformal invariants of manifolds. Similar to *n*-sphere embedded into *n* + 2-dimensional Minkowski space, they tried to embed an arbitrary conformal *n*manifold into an $n + 2$ -dimensional Ricci-flat Lorentz manifold, which they called the ambient space. The ambient spaces were used to produce local scalar conformal invariants. An important part of the ambient space construction is the introduction of conformally compact Einstein metrics for a conformal manifold. The study of conformally compact Einstein metrics could tell us some relationship between the Riemannian structure in the interior and the conformal structure on the boundary. Much progress has been made since then. In recent years, the physics community has also

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become interested in conformally compact Einstein metrics because the introduction of AdS/CFT correspondence in the quantum theory of gravity in theoretic physics by Maldacena [\[24](#page-19-1)].

Let *M* be the interior of a compact $(n+1)$ -dimensional manifold \overline{M} with non-empty boundary ∂M . We call a complete metric g^+ on *M* is $C^{m,\alpha}$ (or $W^{k,p}$) conformally compact if there exits a defining function ρ on \overline{M} such that the conformally equivalent metric

$$
g = \rho^2 g_+
$$

can extend to a $C^{m,\alpha}$ (or $W^{k,p}$) Riemannian metric on \overline{M} . The defining function is smooth on \overline{M} and satisfies

$$
\begin{cases}\n\rho > 0 \text{ in } M \\
\rho = 0 \text{ on } \partial M \\
d\rho \neq 0 \text{ on } \partial M\n\end{cases}
$$
\n(1.1)

Here $C^{m,\alpha}$ and $W^{k,p}$ are usual Hölder space and the Sobolev space. We call the induced metric $h = g|_{\partial M}$ the boundary metric associated to the compactification g. It is easy to see that different defining function induces different boundary metric and every two of the boundary metrics are conformal equivalent. Then the conformal class [*h*] is uniquely determined by $(M, g₊)$. We call $[h]$ the conformal infinity of $g₊$. If in addition, g^+ is Einstein, i.e.

$$
Ric_{g_+} + ng_+ = 0,\t\t(1.2)
$$

then we say $(M, g₊)$ is a conformally compact Einstein manifold.

There are some interesting problems concerning conformally compact Einstein metric. Such as the existence problem, see [\[2](#page-19-2)[,14](#page-19-3)[–16](#page-19-4)[,19](#page-19-5)[,21](#page-19-6)] etc. The unique problem, see $[1,7]$ $[1,7]$. The compactness problem, see $[2,6,7]$ $[2,6,7]$ $[2,6,7]$ $[2,6,7]$.

In this paper, we deal with the boundary regularity problem. Given a conformally compact Einstein manifold (*M*, g^+) and a compactification $g = \rho^2 g^+$, if the boundary metric *h* is $C^{m,\alpha}$, is there a $C^{m,\alpha}$ compactification of g^+ ? This problem was first raised by Fefferman and Graham in 1985 in [\[10\]](#page-19-0) and they observed that if $dim M = n + 1$ is odd, the boundary regularity in general breaks down at the order *n*. If $dim M = n + 1$ is even, the $C^{m,\alpha}$ compactification may exist.

In $[8]$, Chrusciel et al. used the harmonic diffeomorphism at infinity to construct a good structure near boundary where Einstein equation could be written as an elliptic PDE of second order uniformly degenerating at the boundary. That is so-called 'gauge-broken Einstein equation'. Then they use polyhomogeneity result of some specific degenerate equation to obtain a good result of the boundary regularity. We suggest the readers to see [\[3\]](#page-19-11) for more details about these equations. They proved that if the boundary metrics are smooth, the $C²$ conformally compact Einstein metrics have conformal compactifications that are smooth up to the boundary in the sense of $C^{1,\lambda}$ diffeomorphism in dimension 3 and all even dimensions, and polyhomogeneous smooth in odd dimensions greater than 3. This is certainly a very good result in the

sense that they made good use of Einstein equation and gave us a suitable coordinate in infinity to study conformally compact Einstein metrics. I think their method is more geometrical. The condition of that the initial compactification is C^2 in all dimension should be sharp. However, their result only hold for smooth case. It is believed that their method could also be used to prove the finite regularity although we may loss half regularity in this situation.

In [\[1](#page-19-7)[,2](#page-19-2)], Anderson considered the Bach tensor in dimension 4, and proved the finite regularity result. He only assume that the initial compactification *g* is $W^{2,p}$ where $p > 4$. I am not sure whether the $W^{2,p}$ condition is good enough to prove his result. As a supplementary proof, we use Anderson's method to prove his conclusion where we assume that the initial compactification *g* is C^2 and the scalar curvature is C^{σ} for any $\sigma \in (0, 1)$.

In [\[17](#page-19-12)], Helliwell solved the issue in all even dimensions by following Anderson's method. He considered the Fefferman–Graham ambient obstruction tensor instead of Bach tensor in higher dimensions. It is conformally invariant and vanishes for Einstein metrics. Helliwell assumed the initial compactification *g* is at least in $C^{n,\alpha}$ for a $(n+1)$ smooth manifold. It means the original compactification is $C^{3,\alpha}$ for a smooth manifold of dimension 4. Now we reduce the condition $C^{3,\alpha}$ to $C^{2,\sigma}$ to improve his result.

This is the main result:

Theorem 1.1 *Let* (M, g^+) *be a conformally compact Einstein manifold of dimension* 4 *with a C*² *compactification g* = $\rho^2 g^+$ *. If the scalar curvature* $S \in C^{\sigma}(\overline{M})$ *for some* $\sigma > 0$, *the boundary metric h* = $g|_{\partial M} \in C^{m,\alpha}(\partial M)$ *with* $m \geq 2, \alpha \in (0,1)$, *then under a* $C^{2,\lambda}$ *coordinates change,* g^+ *has a* $C^{m,\alpha}$ *conformally compactification* $\tilde{g} = \tilde{\rho}^2 g^+$ *with the boundary metric* $\tilde{g}|_{\partial M} = h$.

Remark 1.2 The new coordinates form $C^{m+1,\alpha}$ differential structure of \overline{M} . $\tilde{\rho}$ is a $C^{m+1,\alpha}$ defining function.

If $g = \rho^2 g^+$ is $C^{2,\sigma}$, then the condition of *S* in Theorem [1.1](#page-2-0) holds automatically. Hence the conclusion is also true.

If the boundary metric h is smooth, then g^+ has a smoothly conformally compactification \tilde{g} with the boundary $\tilde{g}|_{\partial M} = h$.

The condition that "the scalar curvature $S \in C^{\sigma}(\overline{M})$ " seems unnatural and this is because we choose the Yamabe compactification for the new \tilde{g} . This condition is used to improve the regularity of the new defining function and new compactification for the Yamabe equation with Dirichlet data. I think the condition may be removed if we choose another "good" compactification.

It is well known that (see [\[10\]](#page-19-0)) if (M, g^+) is a 4-dimensional conformally compact Einstein metric with boundary metric *h* and $g = r^2 g^+$ is the geodesic compactification associated with *h*, then according to the Gauss lemma, $g_{+} = r^{-2}(dr^{2} + g_{r})$.

$$
g_r = h + g^{(2)}r^2 + g^{(3)}r^3 + \cdots
$$

where $g^{(2)}$ is the Schouten tensor and is determined by *h*. $g^{(3)}$ is determined by g^{+} and *h* and hence it is a non-local term. The rest of power series is determined by $g^{(3)}$ and *h*. This property is also true for higher dimension. From this point of view, Helliwell's condition of $C^{3,\alpha}$ initial compactification seems very natural. That we improve it to $C^{2,\sigma}$ is a big step as we don't need any information about the non-local term.

The outline of this paper is as follows. In Sect. [2,](#page-3-0) we introduce some basic facts about conformally compact Einstein metrics. We show that the Yamabe compactification exists. The conditions in Theorem [1.1](#page-2-0) are unchanged under this compactification. We also consider the Bach equation in dimension 4 and it is an elliptic PDE of second order about Ricci tensor if the scalar curvature is constant. At last, we introduce the harmonic coordinates.

In Sect. [3,](#page-6-0) we deduce some boundary conditions. Including the Dirichlet condition of metric and Ricci curvature, the Neumann condition of Ricci curvature and the oblique derivative condition of metric. We prove that these conditions are true even if the compactification *g* is only C^2 .

In Sect. [4,](#page-12-0) we attempt to prove the main theorem. The first difficulty is C^{α} and $C^{1,\alpha}$ estimate of Ricci curvature. So we present the intermediate Schauder theory to solve the problem. Then we finish our proof with the classical Schauder theory. In the end, with the help of Bach equation, we prove the regularity of defining function in the new coordinates.

2 Preliminaries

Let $(M, g₊)$ be a $n + 1$ -dimensional conformally compact Einstein manifold and $g = \rho^2 g^+$ is a compactification. Then

$$
K_{ab} = \frac{K_{+ab} + |\nabla \rho|^2}{\rho^2} - \frac{1}{\rho} [D^2 \rho(e_a, e_a) + D^2 \rho(e_b, e_b)],\tag{2.1}
$$

$$
Ric = -(n-1)\frac{D^2\rho}{\rho} + \left[\frac{n(|\nabla\rho|^2 - 1)}{\rho^2} - \frac{\Delta\rho}{\rho}\right]g,
$$
\n(2.2)

$$
S = -2n\frac{\Delta\rho}{\rho} + n(n+1)\frac{|\nabla\rho|^2 - 1}{\rho^2}.
$$
 (2.3)

Here *Kab*, *Ric*, *S* are the sectional curvature, Ricci curvature and scalar curvature of *g* and D^2 denote the Hessian. Readers can see [\[4\]](#page-19-13) for the conformal transformation law of curvatures.

If *g* is a C^2 compactification, then from $(2.3), |\nabla \rho| = 1$ $(2.3), |\nabla \rho| = 1$ on ∂M . Then by [\(2.1\)](#page-3-1) K_{+ab} tends to −1 as $\rho \to 0$. Hence a C^2 conformally compact Einstein manifold is asymptotically hyperbolic. Let $D^2 \rho|_{\partial M} = A$ denote the second fundamental form of ∂M in (\overline{M}, g) . The equation [\(2.2\)](#page-3-1) further implies that ∂M is umbilic.

2.1 Constant Scalar Curvature Compactification

Lemma 2.1 *Let* (M, g_+) *be a conformally compact n-manifold with a* $W^{2,p}$ *conformal compactification* $g = \rho^2 g_+$ *where* $p > n/2$ *. Suppose that* $h = g|_{\partial M}$ *is the boundary metric. Then there exits a* $W^{2,p}$ *constant scalar curvature compactification* $\hat{g} = \hat{\rho}^2 g_+$ *with boundary metric h*.

Proof We only need to solve a Yamabe problem with Dirichlet data. Let $\hat{g} = u^{\frac{4}{n-2}}g$, then we consider the equation

$$
\begin{cases} \Delta_{g}u - \frac{n-2}{4(n-1)}Su + \frac{n-2}{4(n-1)}\lambda u^{\frac{n+2}{n-2}} = 0\\ u > 0 \text{ in } \overline{M} \\ u \equiv 1 \text{ on } \partial M \end{cases}
$$
(2.4)

In [\[23\]](#page-19-14), Ma proved that the equation has a $C^{2,\alpha}$ solution if the metric *g* is $C^{2,\alpha}$ when $\lambda = -1$. Now we extend his conclusion in the case that $g \in W^{2,p}$ for some $p > n/2$. Let $\lambda = -1$, and we consider the following functional

$$
I(u) = \frac{1}{2} \int_M \left(|\nabla u|^2 + \frac{n-2}{4(n-1)} S u^2 \right) dv + \frac{n-2}{2n} \int_M \frac{n-2}{4(n-1)} |u|^{\frac{2n}{n-2}} dv
$$

on the set

$$
A = \{u \in H^1(M) : u|_{\partial M} = 1\}.
$$

It is coercive and weakly lower semi-continuous. Then *I* attains its infimum in *A*, which means that $\exists u \in A$, $I(u) = \inf_{v \in A} I(v)$. Since for any $\eta \in H_0^1(M)$, $t \in \mathbb{R}$, $u + t\eta \in A$, we have that

$$
\frac{\mathrm{d}}{\mathrm{d}t}I(u+t\eta)|_{t=0}=0.
$$

Then *u* is a H^1 weak solution. By the Sobolev embedding theorem it follows that $u \in L^{\frac{2n}{n-2}}$. Now let

$$
f(x) = \frac{n-2}{4(n-1)} \left(S + u^{\frac{4}{n-2}} \right),
$$

then *u* is a weak solution of $-\Delta_g u + fu = 0$, $u|_{\partial M} = 1$ and $f \in L^{\frac{n}{2}}$. By a standard method in PDE we can infer that $u \in L^q$ for any $q \ge 2$. (One can see more details in [\[5,](#page-19-15) Theorem 2.3]) So $fu \in L^{p'}$ for any $p' < p$ and it implies that $u \in W^{2,p'}$. If choose $p' > n/2$, then *u* is Hölder continuous. Then $f \in L^p$, and finally we get that $u \in W^{2,p}$. The strong maximum principle tells us that *u* is positive in M. Ч

If $g \in C^2$ and $S_g \in C^{\sigma}$ for some $\sigma > 0$, we know that Eq. [2.4](#page-4-0) has a $C^{2,\sigma}$ solution *u*. Then $\hat{g} = u^{\frac{4}{n-2}}g$ is still C^2 , and the new defining function $\hat{\rho} = u\rho \in C^{2,\sigma}$. In the following of this section, we don't distinguish g with \hat{g} . When we refer to the compactification *g*, we mean the scalar curvature of *g* is −1 near the boundary and the defining function is $C^{2,\sigma}$.

2.2 The Bach Equation

For a 4-dimensional manifold, the Bach tensor is a conformal invariant and vanishes for Einstein metric, see [\[4\]](#page-19-13). In local coordinates,

$$
B_{ij} = P_{ij,k}^{\ \ k} - P_{ik,j}^{\ \ k} - P^{kl} W_{kijl} \tag{2.5}
$$

where $P_{ij} = \frac{1}{2}R_{ij} - \frac{S}{12}g_{ij}$ is the Schouten tensor.

Let $\{y^{\beta}\}_{\beta=0}^{3}$ be the smooth structure on \overline{M} and when restricted on ∂M , $\{y^{i}\}_{i=1}^{3}$ is smooth structure of ∂M . From above we can assume that $g \in C^{\infty}(M) \cap C^2(\overline{M})$, $S_g \equiv$ -1 . Then the fact that g_{+} is Einstein and [\(2.5\)](#page-5-0) imply that

$$
\Delta Ric_{\alpha\beta} = \Gamma * \partial Ric + \mathcal{Q}
$$
 (2.6)

in y-coordinates. Here $\Delta = g^{\alpha\beta}\partial_{\alpha}\partial_{\beta}$, Γ is the Christoffel symbol of *g*, $\Gamma * \partial Ric$ denote the bilinear form of Γ and∂ *Ric* and Q denotes a quadratic curvature term.

2.3 The Harmonic Coordinates Near Boundary

In the rest of the paper, if there are no special instructions, any use of indices will follow the convention that Roman indices will range from 1 to *n*, while Greek indices range from 0 to *n*.

We call the coordinates $\{x^{\beta}\}_{\beta=0}^{n}$ harmonic coordinates with respect to *g* if

$$
\Delta_g x^\beta = 0
$$

for $0 \leq \beta \leq n$. We are now going to construct harmonic coordinates in a neighbourhood of ∂ *M* if *g* is smooth.

In fact, if $g \in C^{1,\alpha}, \alpha \in (0,1)$ for any point $p \in \partial M$, there is a neighbourhood *V* and smooth structure $\{y^{\beta}\}_{\beta=0}^{n}$ where $y^{0}|_{\partial M} = 0$. Then by solving a local Dirichlet problem:

$$
\begin{cases} \Delta_g x^{\beta} = 0 \text{ in } V \\ x^{\beta} |_{V \cap \partial M} = y^{\beta} |_{V \cap \partial M} \end{cases}
$$
 (2.7)

there is a $C^{2,\alpha}$ solution by [\[13\]](#page-19-16) and we have the Schauder estimate:

$$
\| x^{\beta} - y^{\beta} \|_{C^{2,\alpha}(V)} \leq C (\|\Delta(x^{\beta} - y^{\beta})\|)_{C^{\alpha}(V)} + \| x^{\beta} - y^{\beta} \|_{C^{2,\alpha}(\partial V)} = C \|\Delta y \|_{C^{\alpha}(V)}
$$

We can assume that the y-coordinates is the normal coordinates at *p*, then $\Delta y(p) = 0$. Hence if *V* is small enough, $x^{\beta} - y^{\beta} \|_{C^{2,\alpha}(V)}$ tends to 0. $\{x^{\beta}\}_{\beta=0}^{n}$, $0 \le \beta \le n$ is a coordinate around *p*.

In particular, if $g \in C^2$, then the solution $x \in C^{2,\alpha}(y)$ for any $\alpha \in (0, 1)$. Hence

$$
g_{\alpha\beta} = g\left(\frac{\partial}{\partial x^{\alpha}}, \frac{\partial}{\partial x^{\beta}}\right) \in C^{1,\alpha}(\overline{M})
$$

In harmonic coordinates $\{x^{\beta}\}_{\beta=0}^n$, the Ricci tensor could be written as:

$$
\Delta g_{ij} = -2R_{ij} + Q(g, \partial g)
$$

where $Q(g, \partial g)$ is a polynomial of *g* and ∂g . For more details, one can see [\[9](#page-19-17)].

Here we refer to the special coordinates constructed in Sect. [4](#page-12-0) in [\[11](#page-19-18)]. Instead of the harmonic coordinates above, those coordinates may also be useful in our situation, and may also help us to deal with it in higher dimension of even number. That's an interesting problem.

3 The Boundary Conditions

In this section, we derive a boundary problem for *g* and Ricci curvature of a conformal compact Einstein manifold in the harmonic coordinates as defined in Sect. [2.](#page-3-0) We do it locally, that is, for any $p \in \partial M$, there is a neighbourhood *V* contains p and a local harmonic chart $\{x^{\beta}\}\$. Let $D = V \cap \partial M$ be the boundary portion and let $g \in C^2(V)$ be the Yamabe compactification. We will give the Dirichlet and Neumann boundary conditions of *g* and $Ric(g)$ on D. Here we state that the boundary conditions in this section hold for all dimension.

In fact, as it is showed in [\[17](#page-19-12)[,18](#page-19-19)] that, if *g* is $C^{3,\alpha}$ compact, we have following boundary conditions:

Proposition 3.1 *Let* $(M, g⁺)$ *be a n* + 1*-dimensional conformally compact Einstein manifold with a* $C^{3,\alpha}$ *Yamabe compactification* $g = \rho^2 g^+$. $g|_{\partial M} = h$ *is the boundary metric. Suppose that* $\{x^{\beta}\}_{\beta=0}^{n}$ *are any coordinates near the boundary such that* x_0 *is defining function and* {*xⁱ* } *n ⁱ*=⁰ *are coordinates of* [∂] *^M*. *We have:*

$$
g_{ij} = h_{ij},\tag{3.1}
$$

$$
R_{ij} = \frac{n-1}{n-2} \left(Ric_h \right)_{ij} + \left(\frac{1}{2n} S - \frac{1}{2(n-2)} S_h \right) h_{ij} + \frac{n-1}{2n^2} H^2 h_{ij}, \tag{3.2}
$$

$$
R_{0i} = -\left(g^{00}\right)^{-\frac{1}{2}} \frac{n-1}{n} \frac{\partial H}{\partial x_i} - \frac{g^{0j}}{g^{00}} R_{ij},\tag{3.3}
$$

$$
R_{00} = \frac{1}{(g^{00})^2} \left(g^{0i} g^{0j} R_{ij} + g^{00} \left(\frac{1}{2} \left(S - S_h \right) - \frac{n-1}{2n} H^2 \right) \right), \tag{3.4}
$$

$$
N\left(R_{0i}\right) = \left(g^{00}\right)^{-\frac{1}{2}}\left(-g^{j\beta}\partial_{\beta}R_{ji} + g^{\eta\beta}\Gamma_{i\beta}^{\tau}R_{\eta\tau}\right),\tag{3.5}
$$

 W *khere* $N = \frac{\nabla x_0}{|\nabla x_0|} = (g^{00})^{-\frac{1}{2}} g^{0\beta} \partial_\beta$ *be the unit norm vector on* ∂*M* and $R_{\alpha\beta}$, *S*, *H* are *Ricci curvature, scalar curvature and mean curvature respect to g*.

The formula (3.1) is trivial and (3.5) is deduced by the second Bianchi identity and the fact that the scalar curvature is constant near the boundary. Here we briefly recall the proof of The formula [\(3.2\)](#page-6-1), [\(3.3\)](#page-6-1) and [\(3.4\)](#page-6-1). For a $C^{3,\alpha}$ conformally compact Einstein metric, there is a unique $C^{2,\alpha}$ geodesic compactification with the same boundary metric $[20, \text{Lemma 5.1}]$ $[20, \text{Lemma 5.1}]$. Then for such a C^2 geodesic compactification, we have a good formula for Ricci curvature and scalar curvature on the boundary. At last, we use the Ricci formula under conformal change to get (3.2) , (3.3) and (3.4) .

In this section, we will show that the formula (3.2) , (3.3) and (3.4) still hold for C^2 conformally compact Einstein metric.

In fact, if *g* is C^2 conformally compact, then there exists a sequence of $C^{3,\alpha}(\overline{M})$ metrics g_k which converge to g in C^2 norm in smooth structure of \overline{M} . However, g_k are not conformal Einstein in general. In the following, we omit the index *k* and assume that *g* is a $C^{3,\alpha}$ metric on \overline{M} . By choosing a defining function ρ satisfying $|\nabla \rho|_g = 1$ on ∂M , we make $g^+ = \rho^{-2}g$. Then with Taylor theorem, there is a $C^{2,\alpha}$ function *b* such that $|\nabla \rho|^2 = 1 + b\rho$ near the boundary.

$$
Ric = -(n-1)\frac{D^2\rho}{\rho} + \left[\frac{n(|\nabla\rho|^2 - 1)}{\rho^2} - \frac{\Delta\rho}{\rho}\right]g + \frac{F}{\rho},\tag{3.6}
$$

$$
S = -2n\frac{\Delta\rho}{\rho} + n(n+1)\frac{|\nabla\rho|^2 - 1}{\rho^2} + \frac{trF}{\rho},
$$
\n(3.7)

where $F = \rho(Ric_{g+} + ng_+) = \rho Ric_g + (n-1)D^2 \rho - (nb - \Delta \rho)g \in C^{1,\alpha}(\overline{M}).$ Now we prove the following formulas:

$$
R_{0i} = -(g^{00})^{-\frac{1}{2}} \frac{n-1}{n} \frac{\partial H}{\partial x_i} - \frac{g^{0j}}{g^{00}} R_{ij} + Q(F, DF, h, Dg, H),
$$

\n
$$
R_{00} = \frac{1}{(g^{00})^2} \left(g^{0i} g^{0j} R_{ij} + g^{00} \left(\frac{1}{2} (S - S_h) - \frac{n-1}{2n} H^2 \right) \right) + Q(F, DF, h, Dg, H),
$$

\n
$$
R_{ij} = \frac{n-1}{n-2} (Ric_h)_{ij} + \left(\frac{1}{2n} S - \frac{1}{2(n-2)} S_h \right) h_{ij} + \frac{n-1}{2n^2} H^2 h_{ij} + Q(F, DF, h, Dg, H).
$$

\n(3.8)

Here $h = g|_{\partial M}$, *H* is the mean curvature, *Q* is a polynomial and $Q(F, DF, h, Dg, H)$ = 0 if $F = DF = 0$ on ∂M . We will use three lemmas to prove [\(3.8\)](#page-7-0).

First, there is a unique $C^{2,\alpha}$ geodesic compactification of g^+ with boundary metric *h* and denote it by $\bar{g} = r^2 g^+$. Let $\bar{g} = u^2 g$ where $u = \frac{r}{\rho}$ satisfying that $\equiv 1$ on the boundary and $u \in C^{2,\alpha}$. Then $\overline{F} = r(Ric_{g+} + ng_+) = uF$ is still $C^{1,\alpha}(\overline{M})$. We will calculate the boundary curvature of \bar{g} and notice that the second fundamental form of \overline{g} at ∂M is not 0, but determined by the tensor *F*.

Lemma 3.2 *Suppose that* $\bar{g} = r^2 g_+$ *is a* C^2 *conformally compactification of manifold* $(M, g₊)$ *with boundary metric h. Then on the boundary* ∂M ,

$$
\bar{S} = \frac{n}{n-1}(S_h) + Q(\bar{F}, D\bar{F}),
$$
\n(3.9)

$$
\bar{R}_{ij} = \frac{n-1}{n-2} (Ric_h)_{ij} - \frac{1}{2(n-1)(n-2)} S_h h_{ij} + Q(\bar{F}, D\bar{F}, h, D\bar{g}).
$$
 (3.10)

Here S and R_{ij} are the scalar curvature and Ricci curvature of \bar{g} . Q is a polynomial *satisfying* $Q(F, DF, h, D\overline{g}) = 0$ *if* $F = DF = 0$.

Proof Let us choose the coordinates (r, y^1, \ldots, y^n) , near ∂ *M* such that $\bar{g} = dr^2 + g_r$, i.e.

$$
g_{ri} = g^{ri} = 0, g_{rr} = g^{rr} = 1.
$$

According to Gauss Codazzi equation,

$$
\begin{split} \bar{R}_{ij} &= \bar{g}^{\alpha\beta} \bar{R}_{i\alpha\beta j} \\ &= \bar{g}^{kl} ((R_h)_{iklj} + \bar{A}_{il} \bar{A}_{kj} - \bar{A}_{ij} \bar{A}_{kl}) + \bar{R}_{irrj} \\ &= (R_h)_{ij} + \bar{g}^{kl} \bar{A}_{il} \bar{A}_{kj} + \bar{H} \bar{A}_{ij} + \bar{R}_{irrj} . \end{split} \tag{3.11}
$$

Taking trace with respect to *i* and *j*,

$$
\bar{R}_{rr} = \frac{1}{2} (\bar{S} - S_h + \bar{H}^2 - \bar{g}^{ij} \bar{g}^{kl} \bar{A}_{il} \bar{A}_{kj}).
$$
\n(3.12)

Then

$$
\begin{split} \bar{R}_{irrj} &= \bar{g}(\bar{\nabla}_{\partial_i} \bar{\nabla}_{\partial_r} \partial_r, \partial_j) - \bar{g}(\bar{\nabla}_{\partial_r} \bar{\nabla}_{\partial_i} \partial_r, \partial_j) - \bar{g}(\bar{\nabla}_{[\partial_r, \partial_i]} \partial_r, \partial_j) \\ &= -\partial_r \bar{g}(\bar{\nabla}_{\partial_i} \partial_r, \partial_j) + \bar{g}(\bar{\nabla}_{\partial_i} \partial_r, \bar{\nabla}_{\partial_r} \partial_j) \\ &= -\partial_r \bar{A}_{ij} + \bar{A}^2(\partial_i, \partial_j). \end{split} \tag{3.13}
$$

From (3.6) and (3.7) , we have:

$$
\begin{aligned}\n\bar{R}_{ij} &= -(n-1)\frac{\bar{A}_{ij}}{r} - \frac{\bar{\Delta}r}{r}\bar{g}_{ij} + \frac{\bar{F}_{ij}}{r},\\ \n\bar{R}_{ri} &= \frac{\bar{F}_{ri}}{r},\\ \n\bar{R}_{rr} &= -\frac{\bar{\Delta}r}{r} + \frac{\bar{F}_{rr}}{r},\\ \n\bar{S} &= -2n\frac{\bar{\Delta}r}{r} + \frac{tr\bar{F}}{r}.\n\end{aligned} \tag{3.14}
$$

 \overline{R} ic and \overline{S} is continuous on \overline{M} , so on ∂M (*r* = 0) we have:

$$
\bar{A}_{ij} = \frac{1}{n-1} (\bar{F}_{rr} h_{ij} - \bar{F}_{ij}), \n\bar{H} = \bar{\Delta}r = \bar{F}_{rr} = \frac{1}{2n} tr \bar{F}.
$$
\n(3.15)

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Hence

$$
\begin{aligned}\n\bar{R}_{ij} &= -(n-1)\partial_r \bar{A}_{ij} - \partial_r \bar{\Delta}r \bar{g}_{ij} - \bar{\Delta}r \partial_r \bar{g}_{ij} + \partial_r \bar{F}_{ij}, \\
\bar{R}_{ri} &= \partial_r \bar{F}_{ri}, \\
\bar{R}_{rr} &= -\partial_r \bar{\Delta}r + \partial_r \bar{F}_{rr}, \\
\bar{S} &= -2n\partial_r \bar{\Delta}r + \partial_r tr \bar{F}.\n\end{aligned}
$$
\n(3.16)

Combining all the formulas above, we get that

$$
\bar{S} = \frac{n}{n-1} \left(S_h - \bar{H}^2 + |\bar{A}|_h^2 - \frac{1}{n} \partial_r tr \bar{F} \right)
$$

=
$$
\frac{n}{n-1} \left(S_h - \bar{F}_{rr}^2 + \frac{1}{(n-1)^2} \left(n \bar{F}_{rr}^2 + \bar{F}_{rr} tr_h \bar{F} + |\bar{F}|_h^2 \right) - \frac{1}{n} \partial_r tr \bar{F} \right)
$$
^(3.17)

which is (3.9) .

$$
\bar{R}_{ij} = \frac{n-1}{n-2}((R_h)_{ij} + \bar{g}^{kl}\bar{A}_{il}\bar{A}_{kj} + \bar{H}\bar{A}_{ij}) - \bar{F}_{rr}\partial_r\bar{g}_{ij} + \partial_r\bar{F}_{ij} \n+ \frac{1}{n-2}(\bar{A}_{ij}^2 - \frac{1}{2}(\bar{S} - S_h + \bar{H}^2 - |A|^2) + \partial_r\bar{F}_{rr})h_{ij}
$$
\n(3.18)

Noticing that \bar{A}_{ij} is totally determined by \bar{F} and *h*, hence [\(3.10\)](#page-7-2) holds.

Lemma 3.3 *Let* $g = \rho^2 g_+$ *be a* $C^{3,\alpha}$ *conformally compact metric of* (M, g_+) *and* $g = r^2 g_+$ *be* $C^{2,\alpha}$ *geodesic compactification with the same boundary metric* $g|_{\partial M} =$ $g|_{\partial M} = h$. *Let* $r = u\rho$, $A = D^2 \rho$, *then* $A|_{\partial M} = \overline{A} - u_r h$.

Proof In the local coordinates $(r, y^1, y^2, \ldots, y^n)$ near ∂M , $\overline{A}_{ij} = -\overline{\Gamma}^r_{ij}$. Then the relationship between the connection ∇ of *g* and $\overline{\nabla}$ of \overline{g} is:

$$
\Gamma_{ij}^r = \overline{\Gamma}_{ij}^r - \frac{1}{u} (\delta_j^r u_i + \delta_i^r u_j - g_{ij} u_r) = \frac{1}{u} u_r h_{ij}.
$$

 $g = u^{-2}\bar{g}$, *grad_g* = u^2 *grad_g*, then

$$
A_{ij} = D^2 \rho(\partial_i, \partial_j) = g(\nabla_{\partial_i} \nabla \rho, \partial_j) = -g(\nabla \rho, \nabla_{\partial_i} \partial_j)
$$

\n
$$
= -\Gamma^r_{ij} g(\nabla \rho, \partial_r) = -\Gamma^r_{ij} \bar{g}(\bar{\nabla} \rho, \partial_r)
$$

\n
$$
= -\Gamma^r_{ij} \bar{g}(\bar{\nabla}(\frac{r}{u}), \partial_r) = -\Gamma^r_{ij} \bar{g}(\frac{u\bar{\nabla}r - r\bar{\nabla}u}{u^2}, \partial_r)
$$

\n
$$
= -\Gamma^r_{ij} \bar{g}(\bar{\nabla}r, \bar{\nabla}r) = \bar{A}_{ij} - u_r h_{ij}
$$
\n(3.19)

Ч

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$$
\Box
$$

Lemma [3.3](#page-9-0) tells us that $u_r = \frac{H - H}{n}$. Using the fact that $u|_{\partial M} \equiv 1$,

$$
\bar{\nabla}u = \frac{H - H}{n} \bar{\nabla}r.
$$

Lemma 3.4 *Suppose that g*, *g are defined as in Lemma* ¯ *[3.3,](#page-9-0) then on the boundary* ∂ *M*,

$$
R_{ri} = \frac{n-1}{n} \frac{\partial (H - H)}{\partial x_i} + Q(\bar{F}, D\bar{F}, H),
$$

\n
$$
R_{rr} = \frac{1}{2} (S - S_h) - \frac{n-1}{2n} H^2 + Q(F, DF, H),
$$

\n
$$
R_{ij} = \bar{R}_{ij} + \left(\frac{1}{2n} (S - \bar{S})\right) h_{ij} + \frac{n-1}{2n^2} H^2 h_{ij} + Q(\bar{F}, D\bar{F}, H).
$$
 (3.20)

Here $Q(\bar{F}, D\bar{F}, H) = 0$ *if* $F = DF = 0$. *Proof* Let $g = u^{-2}\overline{g}$, then

$$
Ric = \overline{R}ic + (n-1)\frac{\overline{D}^2u}{u} + \left(\frac{\overline{\Delta}u}{u} + \frac{n|\overline{\nabla}u|_{\overline{g}}^2}{u^2}\right)\overline{g}.
$$

We also know that

$$
\bar{\Delta}u = div \bar{\nabla}u = div \left(\frac{\bar{H} - H}{n} \bar{\nabla}r\right) = \frac{\bar{H} - H}{n} \bar{\Delta}r + \frac{\partial_r(\bar{H} - H)}{n},
$$

$$
\bar{D}^2 u(\partial_i, \partial_j) = \bar{g}(\bar{\nabla}_{\partial_i} \bar{\nabla}u, \partial_j) = \frac{\bar{H} - H}{n} \bar{A}_{ij},
$$

$$
\bar{D}^2 u(\partial_i, \partial_r) = u_{ir} = \frac{1}{n} \frac{\partial(\bar{H} - H)}{\partial x_i},
$$

$$
\bar{D}^2 u(\partial_r, \partial_r) = \frac{\partial_r(\bar{H} - H)}{n} = \bar{\Delta}u - \frac{\bar{H} - H}{n} \bar{\Delta}r.
$$

Then on ∂*M*, we conclude that

$$
R_{ri} = \bar{R}_{ri} + \frac{n-1}{n} \frac{\partial(\bar{H} - H)}{\partial x_i} = \partial_r \bar{F}_{ri} + \frac{n-1}{n} \frac{\partial(\bar{H} - H)}{\partial x_i},
$$

\n
$$
R_{rr} = \bar{R}_{rr} + n\bar{\Delta}u - \frac{(n-1)(\bar{H} - H)}{n} \bar{\Delta}r + \frac{(\bar{H} - H)^2}{n},
$$

\n
$$
R_{ij} = \bar{R}_{ij} + \left(\bar{\Delta}u + \frac{(\bar{H} - H)^2}{n}\right)h_{ij} + (n-1)\frac{\bar{H} - H}{n} \bar{A}_{ij}.
$$
 (3.21)

Taking trace with respect to *i* and *j*,

$$
S = \bar{S} + 2n\bar{\Delta}u + \frac{n+1}{n}(\bar{H} - H)^2 - \frac{(n-1)(H - H)}{n}\bar{\Delta}r + \frac{n-1}{n}(\bar{H} - H)\bar{H}.
$$

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Then

$$
\bar{\Delta}u = \frac{1}{2n} \left(S - \bar{S} - \frac{n+1}{n} (\bar{H} - H)^2 + \frac{(n-1)(\bar{H} - H)}{n} \bar{\Delta}r - \frac{n-1}{n} (\bar{H} - H)\bar{H} \right).
$$
\n(3.22)

The result follows from (3.21) and (3.22) .

In the end, Lemmas 3.2 and 3.4 imply (3.8) .

With the preparation above, let's consider a C^2 conformally compact Einstein metric $g = \rho^2 g^+$ on (\overline{M}, y) . We can choose a sequence of $C^{3,\alpha}$ metric g_k which converge to *g* in $C^2(\overline{M})$ norm. Let $\rho_k = \frac{\rho}{|\nabla^{g_k} \rho|_{g_k}}$, so $\rho_k \in C^{3,\alpha}(\overline{M})$ and $|\nabla^{g_k} \rho_k|_{g_k} \equiv 1$ on ∂M . Let $g_k^+ = (\rho_k)^{-2} g_k$, then g_k is a $C^{3,\alpha}$ conformally compactification of (M, g_k^+) with defining function ρ_k . Defining $F_k = \rho_k (Ric_{g_k^+} + ng_k^+)$ as above, then the formula of *Ric_{gk}* on ∂*M* is like [\(3.10\)](#page-7-2) and F_k converge to 0 in $C^1(\partial M)$ norm.

Finally, as the Ricci curvature of g_k converge to that of g uniformly, we conclude that [\(3.7\)](#page-7-1), [\(3.19\)](#page-9-1) and [\(3.20\)](#page-10-2) hold.

3.1 Other Boundary Conditions

We see that if the metric *g* in Lemma [3.3](#page-9-0) is conformally Einstein, then $A = 0$ on ∂M and the boundary is umbilic. This conclusion is also true even if g is C^2 compact and in this case the geodesic compactification is at least $C¹$. Then we have

$$
A_{ij} = \frac{H}{n} h_{ij}.
$$

Taking the derivative of the equation above along ∂ *M*,

$$
\partial_k A_{ij} = \frac{\partial_k H}{n} h_{ij} + \frac{H}{n} \partial_k h_{ij}.
$$

Combining it with (3.3) , we get that

$$
\partial_k A_{ij} = -\frac{1}{n-1} (g^{00})^{\frac{1}{2}} \left(R_{0k} + \frac{g^{0j}}{g^{00}} R_{ij} \right) \tag{3.23}
$$

Technically, this is not a boundary condition because both sides are of the second derivative of *g*. However, this plays an important role in proving the regularity and we will use the condition later.

If we choose harmonic coordinates, we also have the following boundary condition:

$$
g^{\eta\beta}\partial_{\eta}\left(g_{\alpha\beta}-\frac{1}{2}\partial_{\alpha}g_{\eta\beta}\right)=0
$$
\n(3.24)

This is just the local expression of $\Delta_{g} x^{\alpha} = 0$.

$$
\Box
$$

4 Proof of the Main Theorem

We prove the main theorem in this section with the Bach equation in harmonic coordinates and some boundary conditions in last section. Firstly, let's recall some intermediate Schauder theory of elliptic PDE in [\[12](#page-19-21)[,22\]](#page-19-22), i.e. C^{α} and $C^{1,\alpha}$ estimate.

4.1 Intermediate Schauder Estimate

Suppose Ω is a bounded convex domain in \mathbb{R}^n and a is a positive number satisfying $a = k + \beta$ ($k \in \mathbb{N}, \beta \in (0, 1]$) Defining

$$
|u|_a = \sum_{|\alpha| \le k} |D^{\alpha} u|_0 + \sum_{|\alpha|=k} \sup_{x, y \in \Omega} \frac{|D^{\alpha} u(x) - D^{\alpha} u(u)|}{|x - y|^{\beta}}.
$$

Let $H_a(\Omega)$ denote the Hölder space of functions with finite norm $|u|_a$ on Ω , i.e. $H_a(\Omega) = C^{k,\beta}(\overline{\Omega})$. Setting

$$
\Omega_{\delta} = \{x \in \Omega | dist(x, \partial \Omega) > \delta\}
$$

Let *b* be a number satisfying $a + b \ge 0$ and define

$$
|u|_{a,\Omega}^{(b)} = \sup_{\delta>0} \delta^{a+b} |u|_{a,\Omega_{\delta}}
$$

Let $H_a^{(b)}(\Omega)$ denote the space of functions *u* in $H_a(\Omega_\delta)$, $(\forall \delta > 0)$ such that $|u|_{a,\Omega}^{(b)}$ is finite. Let $H_a^{(b-0)}(\Omega)$ be the space of functions *u* in $H_a^{(b)}(\Omega)$ such that if $\delta \to 0$, then $\delta^{a+b}|u|_{a,\,\Omega} \to 0.$

Basic properties: (the following constant *C* depends on a, b, Ω .)

- 1. $H_a^{(-a)}(\Omega) = H_a(\Omega) = C^{k, \beta}(\overline{\Omega})$. Noticing that if *a* is positive integer, $H_a(\Omega) =$ $C^{a-1,1}(\overline{\Omega})$;
- 2. If $b \ge b'$, then $|u|_{a,\Omega}^{(b)} \le C|u|_{a,\Omega}^{(b')}$;
- 3. If $0 \le a' \le a, a' + b \ge 0$ and *b* is not a non-positive integer, then $|u|_{a'}^{(b)}$ a', Ω ^{\leq} $C|u|_{a,\Omega}^{(b)};$
- 4. If $0 \le c_j \le a + b, a \ge 0, j = 1, 2$, then

$$
|uv|_a^{(b)} \le C(|u|_a^{(b-c_1)}|v|_0^{(c_1)} + |u|_0^{(c_2)}|v|_a^{(b-c_2)})
$$

Specially, if *u* and *v* are continuous functions (bounded), then $|uv|_a^{(b)} \leq C(|u|_a^{(b)} +$ $|v|_a^{(b)}$).

With the preparation above, we could state the intermediate Schauder theory. Assuming that Ω is a bounded C^{γ} domain where $\gamma \geq 1$ and a, b are not integer satisfying

$$
0 < b \le a, \ a > 2, \ b \le \gamma
$$

Let

$$
P = \sum_{|\alpha| \le 2} p_{\alpha}(x) D^{\alpha}
$$

be the elliptic differential operator of second order on $\overline{\Omega}$ where

$$
p_{\alpha} \in H_{a-2}^{(2-b)}(\Omega), \text{ if } |\alpha| \le 2
$$

\n
$$
p_{\alpha} \in H_0(\Omega), \text{ if } |\alpha| = 2
$$

\n
$$
p_{\alpha} \in H_{a-2}^{(2-|\alpha|-0)}(\Omega), \text{ if } b < |\alpha|.
$$

Then we have:

Lemma 4.1 [\[12,](#page-19-21) Theorem 6.1] *Let P, a, b be defined as above. If* $p_0 \le 0$ *and the principal part of P is positive, then the Dirichlet problem*

$$
Pu = f \ in \ \Omega, \ \ u = u_0 \ on \ \partial \Omega
$$

has a unique solution $u \in H_a^{(-b)}(\Omega)$ *for every* $f \in H_{a-2}^{(2-b)}(\Omega)$ *and* $u_0 \in H_b(\partial \Omega)$ *, and we have*

$$
u_a^{(-b)}(\Omega) \le C(|u|_{b,\partial\Omega} + |f|_{a-2}^{(2-b)}(\Omega))
$$

We also have the following regularity result:

Lemma 4.2 [\[12,](#page-19-21) Theorem 6.3] *Let* Ω , *P*, *a*, *b satisfy the hypotheses in Lemma [4.1,](#page-13-0) and let* $u \in C^0(\overline{\Omega}) \cap C^2(\Omega)$, $u|_{\partial\Omega} \in H_b(\partial\Omega)$, $Pu \in H_{a-2}^{(2-b)}(\Omega)$. *Then it follows that* $u \in H_a^{(-b)}(\Omega)$.

For the boundary oblique derivative conditions, we have the following lemma:

Lemma 4.3 [\[22](#page-19-22), Theorem 3] *Let a, b be non-integer and* $1 < b \le a$, $a > 2$ *and* $Ω ⊂ ℝⁿ$ *be bounded domain with* H_b *boundary. Let*

$$
P = \sum_{|\alpha| \le 2} p_{\alpha}(x) D^{\alpha} \text{ in } \Omega, \quad M = \sum_{|\alpha| \le 1} m_{\alpha}(x) D^{\alpha} \text{ on } \partial \Omega.
$$

Here

$$
\sum_{|\alpha|=2} p_{\alpha} \xi^{\alpha} \ge c |\xi|^2 \,\forall \xi \in \mathbb{R}^n, \quad \sum_{|\alpha|=1} m_{\alpha} v^{\alpha} > 0
$$

where c is a positive number. We also let

$$
p_{\alpha} \in H_{a-2}^{(2-b)}(\Omega), \text{ if } |\alpha| \le 2; \quad m_{\alpha} \in H_{b-1}(\partial \Omega) \text{ if } |\alpha| \le 1,
$$

$$
p_{\alpha} \in H_{a-2}^{(0-0)} \text{ if } |\alpha| = 2 \text{ and } b < 2.
$$

(a) If $p_0 \leq 0$, $m_0 < 0$, then the oblique derivative problem

$$
Pu = f \ in \ \Omega, \quad Mu = g \ on \ \partial\Omega \tag{4.1}
$$

has a unique solution $u \in H_a^{(-b)}(\Omega)$ *for every* $f \in H_{a-2}^{(2-b)}(\Omega)$ *and* $g \in H_a^{(2)}(\Omega)$ *<i>and a H_{b−1}*(∂Ω). *Moreover*,

$$
u_a^{(-b)}(\Omega) \le C(|g|_{b-1,\partial\Omega} + |f|_{a-2}^{(2-b)}(\Omega)).
$$

 $\mathcal{L}(b)$ *If* $u \in C^0(\overline{\Omega}) \cap C^2(\Omega)$ *is a solution of* (4.1) *with* $f \in H_{a-2}^{(2-b)}(\Omega), g \in H_{b-1}(\partial \Omega)$ *and the directional derivative* $\sum_{|\alpha|=1} m_{\alpha}$ *exists at each point of* ∂Ω, *then u* ∈ $H_a^{(-b)}(\Omega)$.

4.2 The *C***1***,-* **Regularity of Ricci Curvature**

For a C^2 conformally compact Einstein metric $g = \rho^2 g_+, \rho \in C^{2,\sigma}$, we know that $Ric \in C^0(\overline{M})$ in the initial smooth y-coordinates. We observe that from [\(2.2\)](#page-3-1)

$$
\rho Ric = -(n-1)D^2 \rho + \left[\frac{n(|\nabla \rho|^2 - 1)}{\rho} - \Delta \rho\right]g = Q(\partial g, \partial^2 \rho) \in C^{\sigma}(\overline{M}, \{y\}).
$$

Now we compute the metric and curvature in harmonic coordinates $\{x^{\beta}\}_{\beta=0}^{3}$. As *g* is C^2 , $x \in C^{2,\lambda}(y)$, $\forall \lambda \in (0, 1)$. Then in x-coordinates, we have that

$$
Ric\left(\frac{\partial}{\partial x^{\alpha}},\frac{\partial}{\partial x^{\beta}}\right) = \frac{\partial y^{\gamma}}{\partial x^{\alpha}}\frac{\partial y^{\tau}}{\partial x^{\alpha}} Ric\left(\frac{\partial}{\partial y^{\gamma}},\frac{\partial}{\partial y^{\tau}}\right) \in C^{0}(\overline{M},\{x\})
$$
(4.2)

$$
\rho Ric\left(\frac{\partial}{\partial x^{\alpha}},\frac{\partial}{\partial x^{\beta}}\right) = \rho \frac{\partial y^{\gamma}}{\partial x^{\alpha}} \frac{\partial y^{\tau}}{\partial x^{\alpha}} Ric\left(\frac{\partial}{\partial y^{\gamma}},\frac{\partial}{\partial y^{\tau}}\right) \in C^{\sigma}(\overline{M},\{x\})\tag{4.3}
$$

By lemma [4.4](#page-14-1) below, we conclude that $Ric \in H_{\sigma}^{(1-\sigma)}(\overline{M})$.

Lemma 4.4 *Suppose that f is a continuous function on* \overline{M} *and* $\rho f \in C^{\sigma}(\overline{M})$, *then* $f \in H_{\sigma}^{(1-\sigma)}(\overline{M}).$

Proof As $|\nabla \rho| \equiv 1$ on ∂M , we can assume that $\frac{1}{2} \leq |\nabla \rho| \leq 2$ on $\partial M \times [0, \epsilon)$ for a small $\epsilon > 0$. Let

$$
\Omega_{\delta} = \{x \in M | dist(x, \partial M > \delta), M_{\delta} = \{x \in M | \rho(x) > \delta\}.
$$

A direct calculation shows that

$$
\Omega_\delta\subset M_{\frac{\delta}{2}}\subset \Omega_{\frac{\delta}{4}}
$$

So we don't distinguish Ω_{δ} and M_{δ} when studying the definition of $|u|_{a,\Omega}^{(b)}$. Since $\rho f \in C^{\sigma}(\overline{M})$, for any $x, y \in M_{\delta}$,

$$
\frac{|\rho(x)f(x) - \rho(y)f(y)|}{d^{\sigma}(x, y)} \leq C.
$$

Then

$$
C \ge \frac{|\rho(x)f(x) - \rho(x)f(y) + \rho(x)f(y) - \rho(y)f(y)|}{d^{\sigma}(x, y)}
$$

$$
\ge \rho(x)\frac{|f(x) - f(y)|}{d^{\sigma}(x, y)} - |f(y)|\frac{|\rho(x) - \rho(y)|}{d^{\sigma}(x, y)},
$$

which means

$$
\rho(x)\frac{|f(x)-f(y)|}{d^{\sigma}(x,y)} \leq C + |f|_{0,M_\delta}|\rho|_{\delta}.
$$

By assumption, *f* is continuous, in particular, *f* is bounded. As a consequence, $\delta |f|_{\sigma, M_{\delta}} < C'$ for any $\delta > 0$. This proves the lemma. \Box

Lemma 4.5 *In harmonic coordinates,* $g \in H_{2+\sigma}^{(-1-\sigma)}(\overline{M})$.

Proof In harmonic charts,

$$
\Delta g_{\alpha\beta} = -2R_{\alpha\beta} + Q(g, \partial g)
$$

Let $a = 2 + \sigma$, $b = 1 + \sigma$, then according to Lemma [4.2,](#page-13-1) $g_{\alpha\beta} \in H_{2+\sigma}^{(-1-\sigma)}(\overline{M})$. \square \Box

Now we have that $g \in H_{2+\sigma}^{(-1-\sigma)}(\overline{M})$, so the curvature $Rm \in H_{\sigma}^{(1-\sigma)}(\overline{M})$. By linear transformation of tensor in coordinate system [similar to [\(4.2\)](#page-14-2)], *Rm* is still continuous in x-coordinates. Recall that Q in [\(2.6\)](#page-5-1) is the quadratic term of curvature, then $Q \in$ $H_{\sigma}^{(1-\sigma)}(\overline{M})$ from the basic property 4 in Sect. [4.1.](#page-12-1)

As $g \in C^2(\overline{M}, y)$, [3.2\)](#page-6-1), [\(3.3\)](#page-6-1) hold in y-coordinates on ∂M . In the harmonic coordinates $\{x^{\beta}\}_{\beta=0}^3$,

$$
\frac{\partial x^{\alpha}}{\partial y^{i}}|_{\partial M} = \frac{\partial x^{\alpha}|_{\partial M}}{\partial y^{i}} = \delta^{\alpha}_{i}.
$$

Then on ∂ *M*,

$$
Ric\left(\frac{\partial}{\partial y^i}, \frac{\partial}{\partial y^j}\right) = \frac{\partial x^{\gamma}}{\partial y^i} \frac{\partial x^{\tau}}{\partial y^j} Ric\left(\frac{\partial}{\partial x^{\gamma}}, \frac{\partial}{\partial x^{\tau}}\right) = \delta_i^{\gamma} \delta_j^{\tau}
$$

$$
Ric\left(\frac{\partial}{\partial x^{\gamma}}, \frac{\partial}{\partial x^{\tau}}\right) = Ric\left(\frac{\partial}{\partial x^i}, \frac{\partial}{\partial x^j}\right),
$$

$$
Ric\left(\frac{\partial}{\partial x^0},\frac{\partial}{\partial x^\alpha}\right)=\frac{\partial y^\gamma}{\partial x^0}\frac{\partial y^\tau}{\partial x^\alpha}Ric\left(\frac{\partial}{\partial y^\gamma},\frac{\partial}{\partial y^\tau}\right).
$$

For any $p \in \partial M$, consider the $C^{2,\lambda}$ harmonic chart $(V, \{x^{\theta}\}_{\theta=0}^3)$ around p. Let *D* = *V* ∩∂*M* be the boundary portion. Then the Bach equation [\(2.6\)](#page-5-1) could be written as

$$
\Delta Ric_{\alpha\beta} = \frac{\partial}{\partial y^{\theta}} f^{\theta}_{\alpha\beta} + \mathcal{Q}.
$$
 (4.4)

Here $f_{\alpha\beta}^{\theta} = \Gamma * Ric \in H_{\sigma}^{(1-\sigma)}(\overline{M})$, $\theta = 0, 1, 2, 3$. We will firstly deal with the R_{ij} term where $1 \leq i, j \leq 3$. Consider the following equations:

$$
\begin{cases}\n\Delta u_{ij}^0 = f_{ij}^0 \text{ in } V \\
\frac{\partial}{\partial y^0} u_{ij}^0 = 0 \text{ on } D \\
\Delta u_{ij}^k = f_{ij}^k \text{ in } V \\
u_{ij}^k = 0 \text{ on } D\n\end{cases}
$$
\n(4.5)

where $k = 1, 2, 3$. By Lemmas [4.1](#page-13-0) and [4.3,](#page-13-2) the 4 equations above have solutions in $H_{2+\sigma}^{(-1-\sigma)}(V)$. Let $\tilde{R}_{ij} = R_{ij} - \partial_{\theta} u_{ij}^{\theta}$, then

$$
\Delta \tilde{R}ic_{ij} = \mathcal{Q} + \mathcal{Q}(g, \partial g, \partial^2 g, \partial u, \partial^2 u) \in H_{\sigma}^{(1-\sigma)}(V)
$$

From Lemma [4.2](#page-13-1) and the boundary conditions of R_{ij} , we have that $\tilde{R}_{ij} \in H_{2+\sigma}^{(-1-\sigma)}(V)$, which means that $\tilde{R}_{ij} \in C^{1,\sigma}(V)$ and $R_{ij} \in C^{\sigma}(V)$. We could also prove that $\tilde{R}_{00} \in$ $C^{1,\sigma}(V)$ and $R_{00} \in C^{\sigma}(V)$ in the same way.

To study the regularity of R_{0i} , $i = 1, 2, 3$, we need to consider the following 12 equations:

$$
\begin{cases} \Delta u_{0i}^{\theta} = f_{0i}^{\theta} \text{ in } V\\ N(u_{0i}^{\theta}) = -(g^{00})^{-\frac{1}{2}}g^{j\theta}R_{ji} + P_{i}^{\theta}(\partial g) \text{ on } D \end{cases}
$$
(4.6)

Here $\theta = 0, 1, 2, 3$ and $P_i^{\theta}(\partial g)$ is a polynomial of *g* and g^{-1} to be determined, hence in $C^{\sigma}(D)$. Lemma [4.3](#page-13-2) tells us that these 12 equations have solutions $u_{0i}^{\theta} \in H_{2+\sigma}^{(-1-\sigma)}(\overline{M})$. Now let $\tilde{R}_{0i} = R_{0i} - \partial_{\theta} u_{0i}^{\theta}$, then

$$
\Delta \tilde{R}ic_{0i} = \mathcal{Q} + \mathcal{Q}(g, \partial g, \partial^2 g, \partial u, \partial^2 u) \in H_{\sigma}^{(1-\sigma)}(V).
$$

We recall the Neumann boundary condition (3.5) :

$$
N(R_{0i}) = (g^{00})^{-\frac{1}{2}} (-g^{j\beta} \partial_{\beta} R_{ji} + g^{\eta\beta} \Gamma_{i\beta}^{\tau} R_{\eta\tau}).
$$

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Then

$$
N(\tilde{R}_{0i}) = N(R_{0i} - \partial_{\theta}u_{0i}^{\theta}) = N(R_{0i}) - N(\partial_{\theta}u_{0i}^{\theta})
$$

= $(g^{00})^{-\frac{1}{2}}(-g^{j\beta}\partial_{\beta}R_{ji} + g^{\eta\beta}\Gamma_{i\beta}^{\tau}R_{\eta\tau}) + \partial_{\theta}((g^{00})^{-\frac{1}{2}}g^{j\theta}R_{ji}$
 $- P_{i}^{\theta}(\partial g)) + Q(\partial g, \partial u)$
= $(g^{00})^{-\frac{1}{2}}g^{\eta\beta}\Gamma_{i\beta}^{\tau}R_{\eta\tau} - \partial_{\theta}P_{i}^{\theta}(\partial g) + Q(\partial g, \partial u, R_{ij}).$

So if we select some good polynomial $P_i^{\theta}(\partial g)$, we could make that there is no second derivative of metric *g* in $(g^{00})^{-\frac{1}{2}} g^{\eta\beta} \Gamma_{i\beta}^{\tau} R_{\eta\tau} - \partial_{\theta} P_i^{\theta}(\partial g)$. In other words,

$$
N(\tilde{R}_{0i}) = Q(\partial g, \partial u, R_{ij}) \in C^{\sigma}(D)
$$

We again use Lemma [4.3](#page-13-2) to conclude that $\tilde{R}_{0i} \in H_{2+\sigma}^{(-1-\sigma)}(V)$. So $\tilde{R}_{0i} \in C^{1,\sigma}(V)$ and $R_{0i} \in C^{\sigma}(V)$.

Now we have proved that $R_{\alpha\beta} \in C^{\sigma}(V)$ for all $0 \le \alpha, \beta \le 3$, so $f^{\theta}_{\alpha\beta} = \Gamma * Ric \in$ $C^{\sigma}(V)$. Then the solutions of Eqs. [4.5](#page-16-0) and [4.6](#page-16-1) $u^{\theta}_{\alpha\beta}$ are in $C^{2,\sigma}(V)$. Finally, we get that $R_{\alpha\beta} \in C^{1,\sigma}(V)$ by the same method above.

Thus we have finished the first step of the proof, i.e. $Ric \in C^{1,\sigma}(\overline{M})$ in harmonic charts.

4.3 The *Cm,˛* **Regularity of Metric in Harmonic Charts**

We have already shown that $g_{\alpha\beta} \in C^{1,\lambda}$ for any $\lambda \in (0, 1)$ in harmonic charts, then

$$
\Delta g_{\alpha\beta} = -2R_{\alpha\beta} + Q(g, \partial g) \tag{4.7}
$$

If $1 \leq i, j \leq 3$, we have the boundary conditions

$$
g_{ij}=h_{ij},
$$

So $g_{ij} \in C^{2,\lambda}$.

Let A_{ij} be the second fundamental form,

$$
A_{ij} = \frac{1}{2} (g^{00})^{\frac{1}{2}} g^{0\beta} (\partial_\beta g_{ij} - \partial_i g_{\beta j} - \partial_j g_{\beta i}).
$$

Since $Ric \in C^{1,\sigma}(\overline{M})$, according to [\(3.23\)](#page-11-1), $A_{ij} \in C^{2,\sigma}(\partial M)$. Combining it with that $g_{ii} \in C^{2,\lambda}(\overline{M}),$

$$
\partial_j g_{i0} + \partial_i g_{j0} \in C^{2,\sigma}(\partial M)
$$
 (4.8)

Recall the boundary condition [\(3.24\)](#page-11-2)

$$
g^{\eta\beta}\partial_{\eta}\left(g_{\alpha\beta}-\frac{1}{2}\partial_{\alpha}g_{\eta\beta}\right)=0.
$$

Let $\alpha = 0$, and with [\(4.8\)](#page-17-0) we conclude that

$$
\left(g^{j0}\partial_j + \frac{1}{2}g^{00}\partial_0\right)g_{00} \in C^{2,\sigma}(\partial M)
$$
\n(4.9)

So $g_{00} \in C^{2,\lambda}(\overline{M})$.

Let $\alpha = i$ in [\(3.24\)](#page-11-2), and with [\(4.8\)](#page-17-0) we get that

$$
\left(g^{j0}\partial_j + \frac{1}{2}g^{00}\partial_0\right)g_{i0} \in C^{2,\sigma}(\partial M)
$$
\n(4.10)

So $g_{i0} \in C^{2,\lambda}(\overline{M})$. Now we have proved that *g* is $C^{2,\lambda}$ in harmonic charts. Hence ${x^{\theta}}_{\theta=0}^{3}$ form a $C^{3,\lambda}$ differential structure of \overline{M} . Repeat the steps above, we could improve the regularity of metric *g* gradually, and finally $g \in C^{m,\alpha}(\overline{M}, x)$. Hence ${x^{\theta}}_{\theta=0}^{3}$ form a $C^{m+1,\alpha}$ differential structure of \overline{M} .

4.4 Regularity of the Defining Function

We already show that $\rho \in C^{2,\sigma}(\overline{M})$ and ρ is smooth in interior. Then the only thing is to study the boundary regularity of the defining function. For any $p \in \partial M$, take the harmonic chart (V, x) of *p* and let $D = V \cap \partial M$, We could also assume that $g_{\alpha\alpha} = 1$, $g_{ij} = g_{02} = g_{03} = \cdots = g_{0n} = 0$ (*i* $\neq j$), $g_{01} = -\delta$ at *p* where $\delta \in (0, 1)$ is sufficiently close to 1. according to (2.2) and (2.3)

$$
Ric - \frac{Sg}{n+1} = -(n-1)\frac{D^2\rho}{\rho} + \frac{n-1}{n+1}\frac{\Delta\rho}{\rho}g.
$$

Locally, when acting on $(\frac{\partial}{\partial x^0}, \frac{\partial}{\partial x^1}),$

$$
\Delta \rho - (n+1) \cdot g_{01}^{-1} \cdot D^2 \rho \left(\frac{\partial}{\partial x^0}, \frac{\partial}{\partial x^1}\right) = \frac{n+1}{n-1} \cdot g_{01}^{-1} \cdot \rho \left(Ric_{01} - \frac{Sg_{01}}{n+1}\right)
$$
\n(4.11)

If $1 - \delta$ is small enough, then the left side of the formula above is a elliptic operator around *p*. Since $\rho|_D \equiv 0$, $\rho \in C^{m,\alpha}(x)$.

In order to improve the $C^{m+1,\alpha}$ regularity of ρ , we need that $\rho(Ric_{01})$ in [\(4.11\)](#page-18-0) is at least *Cm*−¹,α. Actually,

$$
\Delta(\rho Ric) = \rho \Delta(Ric) + Ric \Delta \rho + 2g(\nabla \rho, \nabla Ric)
$$

The right side of this formula is $C^{m-3,\alpha}$ with the help of Bach equation. $\rho Ric|_{\partial M} \equiv 0$, so $\rho(Ric_{01}) \in C^{m-1,\alpha}$, and the defining function ρ is $C^{m+1,\alpha}$.

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