#### **RESEARCH ARTICLE**



# Facile Fabrication of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Yolk–Shell Spheres as Anode **Material for Lithium Ion Batteries**

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#### **Abstract**

Transition metal oxides have been actively exploited for application in lithium ion batteries due to their facile synthesis, high specific capacity, and environmental-friendly. In this paper,  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$  yolk–shell (Y–S) spheres, used as anode material for lithium ion batteries, were successfully fabricated by Stöber method. XRD patterns reveal that Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres possess a good crystallinity. But the diffraction peaks' intensity of  $Fe<sub>3</sub>O<sub>4</sub>$  crystals in the composites is much weaker than that of bare Fe<sub>3</sub>O<sub>4</sub> spheres, indicating that the outer anatase TiO<sub>2</sub>@C layer can cover up the diffraction peaks of inner Fe<sub>3</sub>O<sub>4</sub> spheres. The yolk–shell structure of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C spheres is further characterized by TEM, HAADF-STEM, and EDS mapping. The yolk–shell structure is good for improving the cycling stability of the inner  $Fe<sub>3</sub>O<sub>4</sub>$  spheres during lithium ions insertion–extraction processes. When tested at 200 mA/g, the Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres can provide a stable discharge capacity of 450 mAh/g over 100 cycles, which is much better than that of bare Fe<sub>3</sub>O<sub>4</sub> spheres and TiO<sub>2</sub>@C spheres. Furthermore, cyclic voltammetry curves show that the composites have a good cycling stability compared to bare  $Fe<sub>3</sub>O<sub>4</sub>$  spheres.

**Keywords**  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$  yolk–shell spheres  $\cdot$  Cycle performance  $\cdot$  Lithium ion batteries  $\cdot$  Anode material

# **Introduction**

Rechargeable lithium ion batteries (LIBs) have been regarded as the most popular electrochemical energy storage devices used in hybrid electrical vehicles, electrical vehicles, and handheld electronic equipment [[1–](#page-8-0)[3\]](#page-8-1). Lots of transition metal oxides, such as  $SnO<sub>2</sub>$  [[4,](#page-8-2) [5](#page-8-3)],  $Co<sub>3</sub>O<sub>4</sub>$  [[6,](#page-8-4) [7](#page-8-5)], MnO<sub>2</sub> [\[8](#page-8-6), [9\]](#page-8-7), NiO [[10,](#page-8-8) [11](#page-8-9)], and Fe<sub>3</sub>O<sub>4</sub> [[12,](#page-8-10) [13\]](#page-8-11), have been used as anodes in LIBs to date. Iron oxides, as an alternative of graphite, have been recognized as promising candidates due to their high specifc capacities (above 1000 mAh/g), abundant resources, environmental benignity, and low cost. However, the serious volume expansion of magnetic iron oxides during the lithiation–delithiation processes is a fatal defect for application in LIBs. A series of associated problems such as the cracking of active materials, the collapse of

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School of Chemistry and Chemical Engineering, Jinggangshan University, Jiangxi 343009, China electrode structure, and electrode polarization [\[14–](#page-8-12)[18\]](#page-8-13) result in an extremely bad cycle stability and rate performance. Resolving such problems would thus be benefcial.

Nowadays, coating a layer of stable electrochemical materials is considered to be efective for stabilizing the electrochemical properties of iron oxides in the lithium ions insertion–extraction processes, such as  $Fe<sub>3</sub>O<sub>4</sub>$ @porous carbon [\[14\]](#page-8-12), Fe<sub>3</sub>O<sub>4</sub>/carbon composites [\[15\]](#page-8-14), mesoporous  $Fe<sub>3</sub>O<sub>4</sub>@C$  submicrospheres [\[16](#page-8-15)], and porous Fe<sub>3</sub>O<sub>4</sub>@C hollow spheres [\[19\]](#page-9-0). Zhang et al. [[20](#page-9-1)] enhanced the cycling performance by coating a layer of TiO<sub>2</sub> to form  $Fe<sub>2</sub>O<sub>3</sub>$ @  $TiO<sub>2</sub>$  core–shell (C–S) nanorods, and the results showed high and stable delivered capacities of approximately 860 mAh/g at a high current density of 1000 mA/g. Porous  $TiO<sub>2</sub>$  was chosen as the shell, which benefts electrolyte infltration, restricts excessive volume changes, and prevents collapse and dissolution into the electrolyte of  $Fe<sub>2</sub>O<sub>3</sub>$  during the discharge–charge processes [\[20](#page-9-1)]. The design of yolk–shell (Y–S) structure has also proved to be efective for improving the cycling stability of iron oxides [[21](#page-9-2)]. For example, to increase the conductivity of materials, a unique architecture with Y–S Fe<sub>2</sub>O<sub>3</sub> $\odot$ C composites was synthesized by the classic Stöber method, which can prevent the reunion of

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active materials and support an interior space to alleviate the volume expansion and rupture during discharge–charge experiments. The materials could be found to possess a good cycling stability and provide an excellent rate performance [[21\]](#page-9-2). To relieve the capacity fading of bare metal oxides, other metal oxides with Y-S structures, such as  $Cr_2O_3@$ TiO<sub>2</sub> Y–S octahedrons [[22](#page-9-3)], SnO<sub>2</sub>@C Y–S nanospheres [[23\]](#page-9-4), MnO<sub>r</sub>/C Y–S nanorods [[24](#page-9-5)], and  $Co<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$ Y–S spheres [\[25\]](#page-9-6), have also been synthesized using a variety of methods to achieve the good electrochemical properties.

In this work, we reported a simple method to synthesize  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S$  spheres for alleviating the capacity attenuation of bare  $Fe<sub>3</sub>O<sub>4</sub>$  spheres as anodes in LIBs. The materials were physically characterized to investigate the structure, synthesis process, crystal structure, components, and element valence. When tested as anodes of LIBs, the  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S$  spheres show the excellent cycling stability and good rate performance. At 200 mA/g, the reversible capacity of approximately 450 mAh/g is much higher than that of bare  $Fe<sub>3</sub>O<sub>4</sub>$  spheres (160.2 mAh/g) after 100 cycles. Therefore, constructing Y–S structures is an acceptable strategy to promote the electrochemical performance of transition metal oxides when applied in LIBs.

#### **Experimental**

#### Synthesis of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S Spheres

 $Fe<sub>3</sub>O<sub>4</sub>$  spheres were synthesized using a method reported in Ref. [\[12\]](#page-8-10). First, 0.64 g of sodium dodecyl benzene sulfonate was added to 60 mL of ethylene glycol to form a clear solution, and 2.16 g of  $FeCl<sub>3</sub>·6H<sub>2</sub>O$  and 3.72 g of sodium acetate were then mixed in the above solution under vigorous stirring for several hours. The obtained solution was transferred into a 100-mL Tefon-lined autoclave, and the oven temperature was increased to 190 °C for 15 h. Finally, a brown product was collected by centrifugation and washed with ethanol/water five times.

The synthesis method of  $Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>@TiO<sub>2</sub> C-S$ spheres was similar to that reported in Ref. [[26\]](#page-9-7). The assynthesized  $Fe<sub>3</sub>O<sub>4</sub>$  spheres were dispersed in a mixed solution of 150 mL ethanol/water (*V*/*V*=4:1) using ultrasonication and vigorous agitation. 1 mL of ammonia and 1 mL of tetraethylorthosilane (TEOS) were successively and slowly added to the above solution to form  $Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>$ C–S spheres. 0.1 g of  $Fe<sub>3</sub>O<sub>4</sub> @ SiO<sub>2</sub> C–S$  spheres was then dispersed in an ethanol (120 mL)/water (0.5 mL) solution containing hydroxypropyl cellulose (HPC) by repeated ultrasonication and continuous vigorous stirring for 10 h. Finally, tetra-*n*-butyl titanate (TBOT, 0.6 mL) ethanol solution was slowly added into the above mixed solution to form  $Fe<sub>3</sub>O<sub>4</sub>$  @  $SiO<sub>2</sub>@TiO<sub>2</sub> C-S$  spheres.

Through hydrothermal carbonization treatment, a layer of carbon was coated to improve the conductivity of the asprepared samples. First, the  $Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>@TiO<sub>2</sub> C-S$  spheres were dispersed in 60 mL of glucose (300 mg) aqueous solution under ultrasonication. Second, the mixture was transferred into a 100-mL Tefon-lined autoclave and then heated at 180 °C for 6 h in an oven to obtain the  $Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>@TiO<sub>2</sub>@C C-S$ spheres. Third, the sample was subsequently annealed at 500 °C for 4 h to further carbonize in an argon atmosphere at a heating rate of 2 °C/min. Finally, the sacrificial  $SiO<sub>2</sub>$  layer was removed by immersing in 60 mL of LiOH aqueous solution (0.06 mol/L) to obtain the final product (Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres) at 90 °C.

#### **Material Characterization**

Scanning electron microscopy (SEM, S4800, HITACHI, Japan) and transmission electron microscopy (TEM, Tecnai G2 F20, FEI Company, USA) were employed to observe the morphologies and sizes of samples. Energy-dispersive X-ray spectrometer (EDS, HITACHI, Japan) was used to further investigate the elements and their distribution in the materials. X-ray powder difractometer (XRD, Bruker AXS, D8-Focus) was employed to determine the crystal structure of samples by using Cu-K $\alpha$  irradiation in a range of 10 $\degree$  to 80 $\degree$ . The Raman spectra (DXR Microscope, Thermo Electron Corporation, USA) were collected using Raman microscope to characterize the carbon atom and crystal structure of  $Fe_3O_4@TiO_2@C$ Y–S spheres.

#### **Electrochemical Measurements**

CR2032 coin-type cells were assembled in an Ar-flled glove box to measure the electrochemical properties of samples. The working electrodes were composed of active materials, acetylene black, and polyvinylidene fuoride at a mass ratio of 7:2:1. The loading mass weight density of active materials is approximately 1.0–1.1 mg/cm<sup>2</sup>. 1.0 mol/L LiPF<sub>6</sub> in a mixed solvent of ethylene carbonate and diethyl carbonate (DEC)  $(V/V=1:1)$  served as the electrolyte, and circular metal Li foils were used as counter electrodes in the CR2032 coin-type cells. The discharge–charge measurements of cells were conducted on a Land battery test system in a potential range of 3.0 V to 0.01 V versus Li/Li<sup>+</sup>. Cyclic voltammetry (CV) tests were conducted on a CHI660D electrochemical workstation at different scan rates of 0.10 mV/s and 0.50 mV/s between 3.0 and 0.01 V.

### **Results and Discussion**

Figure [1](#page-2-0)a shows a schematic diagram of the stepwise synthesis of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres, as described in the experimental section. First,  $Fe<sub>3</sub>O<sub>4</sub>$  spheres with various sizes were synthesized using a simple hydrothermal method. SEM image shows that the diameters of large  $Fe<sub>3</sub>O<sub>4</sub>$  spheres are about 200–300 nm, while those of smaller  $Fe<sub>3</sub>O<sub>4</sub>$  spheres are only dozens of nanometers (Fig. [1](#page-2-0)b). Second, a thin  $SiO<sub>2</sub>$  layer was coated on the surface of  $Fe_3O_4$  spheres to form  $Fe_3O_4@SiO_2$  C–S spheres with a very smooth surface by the hydrolysis of TEOS in an alkaline aqueous solution (Fig. [1c](#page-2-0)). Third, the outer TiO<sub>2</sub> shell was coated on the surface of the  $SiO<sub>2</sub>$  layer to obtain  $Fe<sub>3</sub>O<sub>4</sub> @ SiO<sub>2</sub> @ TiO<sub>2</sub> C-S spheres by the hydroly$ sis of TBOT. SEM image shows that the obtained mate-rial became much coarser (Fig. [1d](#page-2-0)). Fourth, the  $Fe<sub>3</sub>O<sub>4</sub>@$  $SiO<sub>2</sub>@TiO<sub>2</sub>@C C-S$  spheres were formed by the hydrolysis of glucose, followed by calcination at 500 °C. Finally, the middle  $SiO<sub>2</sub>$  layer as sacrificial template was removed by immersing in a LiOH aqueous solution and the  $Fe<sub>3</sub>O<sub>4</sub>$  @

 $TiO<sub>2</sub>@C Y-S$  spheres were obtained. Figure [1e](#page-2-0) shows that the fnal product has an uneven size distribution ranging from dozens of nanometers to 800 nm. It is also of note that many spheres (Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>, Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>@TiO<sub>2</sub>, and  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C spheres$  were reunited after coating  $SiO<sub>2</sub>$ and  $TiO<sub>2</sub>$  layers.

The sizes, morphologies, and internal structures of samples were characterized using TEM. As evident from the TEM images shown in Fig. [2](#page-3-0)a, the  $Fe<sub>3</sub>O<sub>4</sub>$  spheres have a size distribution ranging from several nanometers to 300 nm. After hydrolysis of TEOS, the  $SiO<sub>2</sub>$  shell with a thickness of 20–30 nm was coated on the surface of  $Fe<sub>3</sub>O<sub>4</sub>$  spheres to form  $Fe<sub>3</sub>O<sub>4</sub> @ SiO<sub>2</sub> C-S spheres$ . Meanwhile, a lot of small  $SiO<sub>2</sub>$  $SiO<sub>2</sub>$  $SiO<sub>2</sub>$  spheres were synthesized in this process (Fig. 2b). After coating the TiO<sub>2</sub> $@C$  shell and then removing the SiO<sub>2</sub> layer, it can be found from TEM images (Fig. [2c](#page-3-0), d) that  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$  Y–S spheres have a very unambiguous internal void space, and the width between the  $Fe<sub>3</sub>O<sub>4</sub>$  sphere and  $TiO<sub>2</sub>$  shell is about 20 nm. Notably, a large number of smaller spheres with sizes ranging from dozens to hundreds of nanometers exist in the fnal product. It can be attributed to three reasons: (1) Small Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub> @C Y–S spheres



<span id="page-2-0"></span>**Fig. 1 a** Schematic diagram showing stepwise synthesis of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres. SEM images of **b** Fe<sub>3</sub>O<sub>4</sub> spheres, **c** Fe<sub>3</sub>O<sub>4</sub> @SiO<sub>2</sub> C–S spheres, **d** Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>@TiO<sub>2</sub> C–S spheres, and **e** Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub> @C Y–S spheres

<span id="page-3-0"></span>



and  $TiO<sub>2</sub>@C$  hollow spheres were synthesized using small  $Fe<sub>3</sub>O<sub>4</sub>@SiO<sub>2</sub>$  C–S spheres and SiO<sub>2</sub> spheres as templates, respectively; (2) the aggregation of  $TiO<sub>2</sub>$  nanoparticles occurred in the hydrolysis process of TBOT; and (3) the collapse of outer layer  $TiO<sub>2</sub>$  shells resulted in the formation of smaller  $TiO<sub>2</sub>$  nanocrystals in the alkaline thermal etching process.

EDS analysis and elemental mapping were applied to determine the components, elements, and their distribution in Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres. EDS analysis clearly shows that the composites contain the elements of titanium, oxide, iron, and silicon (Fig. [3a](#page-3-1)), and the silicon content is very low. Calculation of the elemental contents (Table [1\)](#page-4-0) shows that the mass ratio of  $Fe<sub>3</sub>O<sub>4</sub>$  to TiO<sub>2</sub> is 2.33:10. The



<span id="page-3-1"></span>**Fig. 3 a** EDS analysis image of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres; **b** HAADF image and EDS elemental mapping images of oxygen, titanium, iron, and silicon in the  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S$  spheres

<span id="page-4-0"></span>

high angle annular dark feld (HAADF) image shows the existence of an unambiguous void space/boundary layer between Fe<sub>3</sub>O<sub>4</sub> sphere and TiO<sub>2</sub>@C shell. Figure [3b](#page-3-1) shows that oxygen element is distributed in the entire sphere, while iron element is only distributed in the center of the material, which further proves that the  $Fe<sub>3</sub>O<sub>4</sub>$  sphere is located inside  $Fe<sub>3</sub>O<sub>4</sub> @ TiO<sub>2</sub> @ C Y-S sphere. In contrast, silicon element$ is discretely dispersed in the  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S$  sphere, and this behavior is obviously diferent from that of other elements and is mainly ascribed to the absorption of only a small amount of SiO<sub>2</sub> on the surface of  $Fe_3O_4@TiO_2@C$ Y–S spheres during the alkaline thermal etching process.

The X-ray powder diffraction (XRD) patterns of  $Fe<sub>3</sub>O<sub>4</sub>$ spheres and  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$  Y–S spheres are shown in Fig. [4a](#page-4-1), and curve (i) shows that the difraction peaks of 2*θ* at 18.5°, 30.3°, 35.6°, 43.3°, 53.8°, 57.3°, 62.8°, 71.4°, and 74.4° are indexed to the planes (111), (220), (311), (400), (422), (511), (440), (620), and (533) of magnetic cubic phase  $Fe<sub>3</sub>O<sub>4</sub>$  crystals (JCPDS No. 19-0629), respectively. When the Fe<sub>3</sub>O<sub>4</sub> spheres were parceled by TiO<sub>2</sub>@C to form Fe<sub>3</sub>O<sub>4</sub>@  $TiO<sub>2</sub>@C Y-S$  spheres, the diffraction peaks (marked by the symbol  $\cdot$ ) of Fe<sub>3</sub>O<sub>4</sub> crystals become much weaker, as shown in curve (ii) of Fig. [4](#page-4-1)a. The 2*θ* at 25.36°, 37.96°, 48.16°, 55.02°, 62.26°, and 68.8° corresponds to the planes (110), (004), (200), (211), (213) and (116) of anatase  $TiO<sub>2</sub>$  crystals.

The higher crystallinity of  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S$  spheres can be ascribed to the lower carbon content and the lower alkali thermal temperature, and this is evidently diferent from that of  $Co_3O_4@TiO_2@C$  Y–S spheres in Ref. [[25](#page-9-6)]. The Raman spectrum of  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S$  spheres is shown in Fig. [4b](#page-4-1). The peaks at 198.4 cm<sup>-1</sup> and 614.9 cm<sup>-1</sup> are derived from the anatase  $TiO<sub>2</sub>$  crystals  $[27]$  $[27]$ , and the peak at 415.2  $cm^{-1}$  is ascribed to iron oxides [\[28](#page-9-9)]. Furthermore, the elemental carbon of  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$  Y–S spheres is confirmed from the D-band  $(1320 \text{ cm}^{-1})$  and G-band  $(1620 \text{ cm}^{-1})$  [[29](#page-9-10), [30](#page-9-11)].

The galvanostatic cycling performances and coulombic efficiencies curves of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres, Fe<sub>3</sub>O<sub>4</sub> spheres, and  $TiO<sub>2</sub>@C$  hollow spheres at 200 mA/g are pre-sented in Fig. [5](#page-5-0). The first discharge capacity of  $Fe<sub>3</sub>O<sub>4</sub>$ @ TiO<sub>2</sub> $@C Y-S$  spheres is 778.1 mAh/g with a very low coulombic efficiency of 48.2%. Such a large irreversible capacity loss is mainly derived from the irreversible decomposition of electrolyte, the solid electrolyte interface (SEI) flms formed in the electrode, and the trapping of lithium ions inside TiO<sub>2</sub> and Fe<sub>3</sub>O<sub>4</sub> lattices [\[12](#page-8-10), [14](#page-8-12)[–17](#page-8-16), [19](#page-9-0), [20\]](#page-9-1). The second discharge capacity is reduced to only 523.6 mAh/g, and the subsequently delivered capacities show a slightly declining trend. To the 18th cycle, the discharge capacity reaches the lowest value of 402.8 mAh/g and then rises slowly to around 450 mAh/g over 100 cycles. This occurs because electrolyte infltration gradually increases in the porous materials to improve the electron transfer and Li-ion migration rate with the number of cycles  $[31]$  $[31]$ . The coulombic efficiencies are nearly 100% after fve cycles, which indicates that a layer of stable SEI flms has formed on the surface of the electrode, and the  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y-S$  spheres have a good electrochemical Li-inserted reversibility. In contrast, the bare  $Fe<sub>3</sub>O<sub>4</sub>$  spheres have poor electrochemical properties (Fig. [5](#page-5-0)a). The frst discharge capacity reaches 1082 mAh/g with a high coulombic efficiency of  $70.7\%$  (Fig. [5b](#page-5-0)), but the



<span id="page-4-1"></span>**Fig. 4 a** XRD patterns of Fe<sub>3</sub>O<sub>4</sub> spheres and Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres; **b** Raman spectrum of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres



<span id="page-5-0"></span>Fig. 5 a Galvanostatic cycling tests and b coulombic efficiencies of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres, Fe<sub>3</sub>O<sub>4</sub> spheres, and TiO<sub>2</sub>@C hollow spheres at a current density of 200 mA/g

delivered capacities of the bare  $Fe<sub>3</sub>O<sub>4</sub>$  spheres decay very quickly. After 10 cycles and 100 cycles, the discharge capacities drop to 307.1 mAh/g and 160.2 mAh/g, respectively, which are far lower than those of the composites under the same test conditions. As such, the high and stable reversible capacities of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres are attributed to two aspects: First, the internal void space formed by the removal of  $SiO<sub>2</sub>$  effectively alleviates the volumetric variation in Fe<sub>3</sub>O<sub>4</sub> spheres, and second, the TiO<sub>2</sub> $@C$  shell effectively restrains the pulverized inner  $Fe<sub>3</sub>O<sub>4</sub>$  spheres to relieve the electrode polarization and capacity fading, thus leading to an obvious improvement in cycling stability [[14,](#page-8-12) [16](#page-8-15), [19–](#page-9-0)[21](#page-9-2)]. The first discharge capacity of  $TiO<sub>2</sub>@C$  hollow spheres is 646.5 mAh/g with a very low coulombic efficiency of  $30.1\%$ , but the second discharge capacity is only 343.0 mAh/g (Fig. [5\)](#page-5-0). After 100 cycles, the discharge capacity is maintained at 286.5 mAh/g. The delivered capacities of TiO<sub>2</sub>@C hollow spheres are obviously much lower than those of  $Fe_3O_4@TiO_2@C$  Y–S spheres over 100 cycles. These obtained data clearly show that the capacities of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres are largely related to those of the inner  $Fe<sub>3</sub>O<sub>4</sub>$  spheres. In addition, the first coulombic efficiency of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres (48.2%) is lower than that of  $Fe<sub>3</sub>O<sub>4</sub>$  spheres (70.7%) and higher than that of  $TiO<sub>2</sub>@C$  hollow spheres (30.1%), implying that the irreversible capacity of  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$  Y–S spheres is largely derived from the outer  $TiO<sub>2</sub>@C$  layer in the first cycle.

Figure [6](#page-6-0) illustrates the galvanostatic discharge–charge voltage curves for the 1st, 2nd, 10th, 30th, 50th, and 100th cycles of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres, Fe<sub>3</sub>O<sub>4</sub> spheres, and  $TiO<sub>2</sub>@C$  hollow spheres at 200 mA/g in a voltage range from 3.0 to 0.01 V. As shown in Fig. [6a](#page-6-0), the voltage plateaus of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres at about 1.75 V and 2.0 V are attributed to the insertion–extraction of lithium ions in the anatase  $TiO<sub>2</sub>$  layer with good crystallinity [[20](#page-9-1)], while the voltage plateaus at about 0.75 V and 1.5 V correspond to

the redox reactions between  $Fe<sub>3</sub>O<sub>4</sub>$  and Fe [\[12,](#page-8-10) [14–](#page-8-12)[17](#page-8-16), [19](#page-9-0)]. Notably, the formation of SEI flms occurs in a voltage range below 0.75 V [\[19](#page-9-0), [20](#page-9-1)].

The galvanostatic discharge–charge voltage profles of pure  $Fe<sub>3</sub>O<sub>4</sub>$  spheres are shown in Fig. [6](#page-6-0)b for comparison. It is evident that the same voltage plateaus at approximately 0.75 V and 1.5 V correspond to the redox reactions between  $Fe<sub>3</sub>O<sub>4</sub>$  and Fe. Besides, the capacity fading is fast in the frst fve cycles, and then, it slows from the 6th cycle to the 10th cycle. The gradual increase in the space between the discharge–charge voltage plateaus suggests that the cell polarization caused by the high resistance, low ionic transport, and low electronic conductivity increases sharply with an increase in cycles during the frst 10 cycles. Such poor electrochemical properties are mainly ascribed to the pulverization of active material and the crack of electrode caused by the severe volume expansion of bare  $Fe<sub>3</sub>O<sub>4</sub>$  spheres. The pulverized  $Fe<sub>3</sub>O<sub>4</sub>$  can lead to a loose connection between the particles and the current collector [\[20](#page-9-1)]. Then, the crack of electrode results in the rupture of SEI flms, which forces the new electrode surface to be exposed to the electrolyte, and unstable SEI layers are then repeatedly formed and more electrolyte is consumed [\[12](#page-8-10), [16](#page-8-15)]. Following this, the capacities and voltage plateaus tend to become stable.

The discharge–charge voltage profiles of  $TiO<sub>2</sub>@C$  hollow spheres are shown in Fig. [6](#page-6-0)c, which indicates that the plateaus at about 1.75 V and 2.0 V are ascribed to the redox process of anatase phase  $TiO<sub>2</sub>$  crystals [\[20\]](#page-9-1), and no other voltage plateaus are observed. In summary, the stable capacities and voltage plateaus indicate that the  $TiO<sub>2</sub>@C$ layer has good electrochemical properties for application in LIBs. Therefore, coating the  $TiO<sub>2</sub>@C$  layer is a good way for improving the cycling stability of the inner  $Fe<sub>3</sub>O<sub>4</sub>$  spheres.

Figure [7a](#page-6-1) presents the rate capability of  $Fe_3O_4@TiO_2@C$ Y–S spheres at various current densities for 20 cycles. At a low current density of 50 mA/g, the frst discharge capacity



<span id="page-6-0"></span>**Fig. 6** Discharge–charge voltage profiles of **a** Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres, **b** Fe<sub>3</sub>O<sub>4</sub> spheres, and **c** TiO<sub>2</sub>@C hollow spheres at different cycles



<span id="page-6-1"></span>**Fig. 7 a** Rate performance and **b** discharge–charge voltage profiles of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres at various current densities

of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres is 950 mAh/g, and then, the capacities are stabilized at about 500 mAh/g in subsequent cycles. In addition, when the constant current densities are switched to 100 mA/g, 200 mA/g, 400 mA/g, 800 mA/g, and 1600 mA/g, the discharge capacities of  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$ Y–S spheres remain at about 460 mAh/g, 410 mAh/g, 320 mAh/g, 150 mAh/g, and 62 mAh/g, respectively. When the current density returns to 50 mA/g, the reversible capacities of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres still reach more than 500 mAh/g, indicating the excellent cycle stability and good reversibility.

Figure [7](#page-6-1)b shows the discharge–charge voltage profles of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres at various current densities. It is evident that at low current densities (50–400 mA/g),

the fading of electrochemical reversible capacities, and the space of discharge–charge voltage plateaus increase slowly with the increasing current density. However, at high current densities (400–1600 mA/g), the fading of electrochemical reversible capacities, and the space of discharge–charge voltage plateaus increase sharply with the increasing current density. This can be attributed to three factors: (1) The middle layer  $TiO<sub>2</sub>$  is a kind of semiconductor material that afects the charge transfer and lithium ions difusion in the interior  $Fe<sub>3</sub>O<sub>4</sub>$  spheres; (2) the lower carbon content results in the low electron transfer and poor lithium ions difusion; and (3) the accumulation or depletion of lithium ions on the electrode causes the concentration polarization of batteries due to slow difusion at high current densities.

CV curves of bare Fe<sub>3</sub>O<sub>4</sub> spheres and Fe<sub>3</sub>O<sub>4</sub> @TiO<sub>2</sub> @C Y–S spheres were measured at a lower scan rate of 0.10 mV/s in a potential range from 3.0 to 0.01 V (Fig. [8](#page-7-0)a, b). From Fig. [8](#page-7-0)a, in the first cycle, an obvious larger reduction current in the region of below 0.75 V should be mainly ascribed to the formation of SEI flms and the reduction process of  $Fe<sub>3</sub>O<sub>4</sub>$ , while the oxidation peak at 1.68 V can be attributed to the oxidation process from Fe to  $Fe<sub>3</sub>O<sub>4</sub>$  [\[12,](#page-8-10) [16](#page-8-15)].

The related electrochemical reaction process is presented in Eqs.  $(1)$  $(1)$ – $(3)$  $(3)$ . Subsequently, the reduction peaks firstly shift to about 1.3 V and then shift to the lower potential, while the oxidation peaks constantly shift to the higher potential. The gradually decreased peak area and the increased peak space indicate the continuously increased rupture and volume expansion of  $Fe<sub>3</sub>O<sub>4</sub>$  spheres, and the electrode polarization, resulting in a poor cycling stability and low reversible capacity  $[10-14]$  $[10-14]$  $[10-14]$ . This is well consistent with the galvanostatic discharge–charge tests of pure  $Fe<sub>3</sub>O<sub>4</sub>$  spheres. From Fig. [8](#page-7-0)b, the electrochemical performance of  $Fe<sub>3</sub>O<sub>4</sub>@$ TiO<sub>2</sub>@C Y–S spheres obviously exceeds that of bare Fe<sub>3</sub>O<sub>4</sub> spheres under the same test conditions. In the frst cycle, a large and irreversible reduction current below 1.0 V corresponds to the formation of SEI flms and the irreversible insertion of lithium ions in TiO<sub>2</sub> and Fe<sub>3</sub>O<sub>4</sub> crystals [\[12,](#page-8-10) [14](#page-8-12)[–16](#page-8-15), [20](#page-9-1)]. A pair of symmetrical redox CV peaks situated at 2.0 V and 1.75 V are attributed to the redox potential of TiO<sub>2</sub> [\[20\]](#page-9-1), while the redox CV peaks at 1.2 V and 0.75 V correspond to the redox potential of  $Fe<sub>3</sub>O<sub>4</sub>$  [[12](#page-8-10), [14](#page-8-12)[–16](#page-8-15), [20](#page-9-1)].

The subsequent CV curves of  $Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C$ Y–S spheres were further measured to investigate the



<span id="page-7-0"></span>**Fig. 8** First ten CV curves of **a** Fe3O4 spheres and **b** Fe3O4@TiO2@C Y–S spheres at a scan rate of 0.10 mV/s and **c** subsequent ffty times CV curves of Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres at 0.50 mV/s in a potential range from 3.0 to 0.01 V versus Li/Li<sup>+</sup>

electrochemical kinetics at a high scan rate of 0.50 mV/s (Fig. [8](#page-7-0)c). Compared with the tested CV curves at 0.10 mV/s, the redox peaks potential of  $TiO<sub>2</sub>$  shifts to 2.25 V and 1.75 V, while the anodic peak of  $Fe<sub>3</sub>O<sub>4</sub>$  at 1.2 V (marked using one red circle in Fig. [8](#page-7-0)b) also shifts to 1.5 V (marked using another red circle in Fig. [8c](#page-7-0)). The wider space between the redox peaks indicates the increasing electrode polarization at a higher scan rate, which is in good agreement with the rate performance tests. In addition, there are no obvious changes in the shifts of peak areas and positions in the subsequent 50 cycles, indicating that the Fe<sub>3</sub>O<sub>4</sub>@TiO<sub>2</sub>@C Y–S spheres possess a good cycling stability as LIBs' anodes. The results match well with the above galvanostatic discharge–charge cycle tests. The electrode reactions are described as follows:

$$
Fe_3O_4 + 2Li^+ + 2e^- \rightarrow Li_2Fe_3O_4 \tag{1}
$$

 $Li_2Fe_3O_4 + 6Li^+ + 6e^- \rightarrow 3Fe + 4Li_2O$  (2)

 $3Fe + 4Li_2O \rightarrow Fe_3O_4 + 8Li^+ + 8e^-$  (3)

$$
C + xLi^{+} + xe^{-} \rightarrow Li_{x}C
$$
 (4)

$$
TiO2 + xLi+ + xe- \rightarrow LixTiO2
$$
 (5)

$$
Li_xTiO_2 \rightarrow TiO_2 + xLi^+ + xe^-
$$
 (6)

## **Conclusions**

The Stöber method was employed to successfully synthesize  $Fe_3O_4@TiO_2@C$  yolk–shell (Y–S) spheres for LIB anodes. The  $TiO<sub>2</sub>@C$  shell and void space effectively alleviate the severe volume expansion of  $Fe<sub>3</sub>O<sub>4</sub>$  spheres and restrict the rupture of electrodes, which thus improves the electrochemical performance throughout the cycles. At a constant current density of 200 mA/g, the Fe<sub>3</sub>O<sub>4</sub>@  $TiO<sub>2</sub>@C Y-S$  spheres can release a high reversible capacity of approximately 450 mAh/g after 100 cycles, which is far higher than that of bare  $Fe<sub>3</sub>O<sub>4</sub>$  spheres (160.2 mAh/g) and  $TiO<sub>2</sub>@C$  hollow spheres (286.5 mAh/g). These results also demonstrate that the design of the Y–S structure can be applied in developing new type electrode materials for LIBs.

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