

# Prediction of optimum design variables for maximum heat transfer through a rectangular porous fin using particle swarm optimization†

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(Manuscript Received December 4, 2017; Revised May 13, 2018; Accepted May 22, 2018) --

## **Abstract**

The current study deals with the optimization of significant variables which govern the heat transfer from a porous fin in convective medium using particle swarm optimization (PSO). The temperature distribution of the fin is obtained analytically by a perturbation technique called homotopy perturbation method (HPM). To validate the temperature distribution obtained by HPM, finite difference method is employed. Next a significance analysis has been carried out to identify important variables that play a vital role in transferring heat from the porous fin. The set of variables thus obtained was then optimized by PSO to enhance the heat transfer rate. Reflective boundary condition is incorporated in the PSO to prevent particles from wandering in the infeasible region. The convergence plots of the variables show the effectiveness of PSO in solving non linear problems of this magnitude which are often encountered in the analysis of heat transfer through porous fins.

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*Keywords*: Particle swarm optimization; Porous fin; Homotopy perturbation; Heat transfer

# **1. Introduction**

Fins are an integral part of any device that generates heat during its working. The extensive research in this field has introduced the concept of porous fins for transferring better heat by improving the surface area of convection [1-3]. Advantages like material saving, weight minimization and increased surface area make these fins a good choice for a number of applications where weight, cost and space are major constraints.

Hatami et al. [4] studied analytically the heat transfer through porous fins made of  $Si<sub>3</sub>N<sub>4</sub>$  and Al by considering t temperature dependent heat generation. Again shape of the fin plays a major role in deciding the heat transfer capability of a fin. Thus, Bhanja and Kundu [5] determined analytically the temperature distribution and performance parameters of a constructal T shaped fin. A comparison between porous and solid straight fins of various profiles on the basis of heat transfer rate was done by Kundu et al. [6]. In another work by Bhanja et al. [7], the heat transfer rate through porous pin fins is found to be better than the solid fins of similar dimension and profile. In an analytical study involving the moving porous fins, Bhanja et al. [8] performed an optimization analysis by considering the effects of both convection and radiation.

Recently Hazarika et al. [9] performed an analytical study of a constructal T shaped porous fin under dehumidifying conditions. An effort has been made by Kundu and Bhanja [10] to study the optimum design analysis of porous fins considering a highly non-linear equation. The study was done for three different models. On the other hand, Das [11] proposed an inverse solution of a cylindrical porous fin exposed in both convective and radiative environment.

Recently, focus has been laid on optimizing the important parameters of heat transferring devices such as fins and heat exchangers [12-14]. The non linear equations obtained in these applications demand meta-heuristic algorithms to find the global optimized values of the parameters. Powerful nature inspired global optimization techniques such as genetic algorithm (GA), particle swarm optimization (PSO) and similar bio inspired stochastic algorithms have been used in a number of applications [15, 16] in the past few years. Recently a number of works have been performed on PSO related to the area of heat transfer [17, 18].

To the best of author's knowledge, there has not been any effort to study the multivariable optimization of porous fins using meta-heuristic algorithm. In the current work, particle swarm optimization (PSO), a swarm based meta-heuristic algorithm is selected to obtain the actual heat transfer rate by optimizing the important variables. The faster convergence ability and its derivative free nature make PSO a good choice for optimizing these kinds of problems. Therefore, the main

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<sup>†</sup> Recommended by Associate Editor Ji Hwan Jeong

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Fig. 1. Schematic diagram of a rectangular straight porous fin.

aim behind this work is to optimize the four key parameters in order to obtain the maximum heat transfer rate for a specified fin volume. Heat transfer rate is calculated by solving analytically the governing differential equation using homotopy perturbation method (HPM).

### **2. Problem definition**

This study considers a rectangular porous fin as shown is Fig. 1. Being porous in nature, the fin allows fluid to penetrate through it and thus Darcy model is used to understand and analyze the interaction between the fluid and the porous medium in convective medium. In this work heat is transferred to the surrounding through the pores due to fluid penetration and the convective heat transfer through the solid part of the fin. Heat generation within the fin and radiation heat transfer is neglected and the temperature is considered to vary along one direction (x-axis) only. There is no contact resistance at the fin base and the wall. Porous medium is homogeneous, isotropic and saturated with a single phase fluid and the physical properties of solid as well as fluid are invariable except density of the fluid that may affect the buoyancy term where Boussinesq approximation is employed. The fin tip is considered to transfer heat through convection. rrounding through the ports due to fluid enertration and<br>
exerceive has transfer through the solid part of the fin.<br>
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exerceived not the solid energetance is c surrounding through the pores due to fluid penetration and<br>convective heat transfer through the solid part of the fin.<br>effected and the temperature is considered to vary along one<br>elected and the temperature is considered *The term in the second of the term and the term of the secondary conditions*<br> *Tansfer through the solid part of the fin.*<br> *The boundary conditions*<br> *The boundary conditions*<br> *The boundary conditions* in non-dimension *solution that the finital differential equality conditions in and particular and the finit and reduction heat that the m* saturate with a single phase indet and are physical photons.<br>
Ses of solid as well as fluid are invariable except density of<br>
fluid that may affect the buoyancy term where Boussinesq where  $\theta_i = (T_i - T_a)/(T_b -$ <br>
heat throu

#### *2.1 Governing equations*

Applying energy balance and using Fourier's law in a differential element of the fin, we obtain the governing differential equation for the porous fin as

$$
\frac{d^2T}{dx^2} - \frac{2h(1-\phi)}{k_{\text{eff}}} (T - T_a) - \frac{mC_p}{k_{\text{eff}} W t \Delta x} (T - T_a) = 0.
$$
 (1)

The mass flow rate,  $\dot{m}$  (in kgs<sup>-1</sup>) of the fluid passing through fin pores is defined as

$$
\dot{m} = \rho V W \Delta x \tag{2}
$$

 $\rho$  is the fluid density in kgm<sup>-3</sup> and *W* in the fin width in m.

The fluid velocity  $V$  (in ms<sup>-1</sup>) is estimated by Darcy's law as [2]

$$
V = gK\beta(T - T_a)/\gamma \tag{3}
$$

*e* and *Technology* 32 (9) (2018) 4495~4502<br>
e fluid velocity V (in ms<sup>-1</sup>) is estimated by Darcy's law as<br>  $V = gK \beta (T - T_a)/\gamma$ . (3)<br>
K is the permeability of the porous fin material in m<sup>2</sup>. The<br>
ective thermal conductivit K is the permeability of the porous fin material in  $m<sup>2</sup>$ . The effective thermal conductivity (in  $Wm^{-1}K^{-1}$ ) of the porous fin can be obtained from following expression *e* fluid velocity *V* (in ms<sup>-1</sup>) is estimated by Darcy's law as<br>  $V = gK\beta (T - T_a)/\gamma$ . (3)<br> *K* is the permeability of the porous fin material in m<sup>2</sup>. The<br>
ective thermal conductivity (in Wm<sup>-1</sup>K<sup>-1</sup>) of the porous fin<br> *k x* and *Technology* 32 (9) (2018) 4495-4502<br> **c** fluid velocity V (in ms<sup>-1</sup>) is estimated by Darcy's law as<br>  $V = gK\beta (T - T_a)/\gamma$ . (3)<br>
K is the permeability of the porous fin material in m<sup>2</sup>. The<br>
recive thermal conductiv between  $V$  ( in ms<sup>-1</sup>) is estimated by Darcy's law as<br>  $(\Gamma - T_n)/\gamma$ . (3)<br>
(3)<br>
The male conductivity (in Wm<sup>-1</sup>K<sup>L1</sup>) of the porous fin<br>
male conductivity (in Wm<sup>-1</sup>K<sup>L1</sup>) of the porous fin<br>
and conductivity (in Wm<sup>-1</sup>K<sup>L</sup>  $V = gK\beta(T - T_a)/\gamma$ . (3)<br>
K is the permeability of the porous fin material in m<sup>2</sup>. The<br>
ties the permeability of the porous fin material in m<sup>2</sup>. The<br>
tective thermal conductivity (in Wm<sup>-1</sup>K<sup>-1</sup>) of the porous fin<br>
the obta *x* and Technology 32 (9) (2018) 4495-4502<br> **c** fluid velocity  $V$  (in ms<sup>4</sup>) is estimated by Darcy's law as<br>  $V = gK \beta (T - T_s)/\gamma$ . (3)<br> *K* is the permeability of the porous fin material in m<sup>2</sup>. The<br>
be obtained from follow

$$
k_{\text{eff}} = \phi k_{\text{f}} + (1 - \phi)k_{\text{s}} \tag{4}
$$

where  $k_f$  and  $k_s$  are the thermal conductivities of the fluid and fin material respectively in  $Wm^{-1}K^{-1}$ .

. The following dimensionless parameters are used

$$
(X; \psi; k_{\scriptscriptstyle R}; \theta) = (x/L; t/L; k_{\scriptscriptstyle S}/k_{\scriptscriptstyle f}; (T - T_a)/(T_b - T_a)) \tag{5a}
$$

$$
\Theta = k_{\text{eff}} / k_{\text{f}} = \phi + k_{\text{R}} - \phi k_{\text{R}}
$$
\n(5b)

$$
[m_1; m_2] = [(Ra Da)/(w^2 \Theta); (2Nu(1-\phi)]/(w^2 \Theta)] \qquad (5c)
$$

$$
k_{\text{eff}} = \phi k_{f} + (1 - \phi)k_{s}
$$
\n(4)  
\nhere  $k_{f}$  and  $k_{s}$  are the thermal conductivities of the fluid and  
\nmaterial respectively in Wm<sup>-1</sup>K<sup>-1</sup>.  
\nThe following dimensionless parameters are used  
\n
$$
(X; \psi; k_{s}; \theta) = (x/L; t/L; k_{s}/k_{f}; (T - T_{a})/(T_{b} - T_{a}))
$$
\n(5a)  
\n
$$
\Theta = k_{\text{eff}}/k_{f} = \phi + k_{R} - \phi k_{R}
$$
\n(5b)  
\n
$$
[m_{1}; m_{2}] = [(Ra Da)/(w^{2} Θ); {2Nu(1-φ)}/(w^{2}Θ)]
$$
\n(5c)  
\n
$$
[Ra; Da; Nu] = [{\phi C_{p}g\phi(T_{b} - T_{a})t^{3}}/((\gamma k_{f}); K/t^{2}; ht/k_{f})].
$$
\n(5d)  
\nUsing Eqs. (2)–(5), Eq. (1) can be expressed as  
\n
$$
d^{2}θ/dX^{2} - m_{1}θ^{2} - m_{2}θ = 0.
$$
\n(6)  
\n2 **Boundary conditions**  
\nThe boundary conditions in non-dimensional form are  
\n
$$
θ = 1
$$
 at  $X = 1$  (Fin base)  
\n
$$
[dθ/dX]_{x=0} = H = (Nu/Θ\psi)θ_{t}
$$
 (Convective tip)  
\n(7a)  
\nhere  $θ_{t} = (T_{t} - T_{a})/(T_{b} - T_{a})$  is the dimensionless tip tem-  
\nrature.  
\n3 **Temperature distribution**  
\nThe present problem adopts the homotopy perturbation  
\nrtho the result of the relationship of the *h*-component distribution

$$
d^2\theta/dX^2 - m_1\theta^2 - m_2\theta = 0.
$$
 (6)

#### *2.2 Boundary conditions*

The boundary conditions in non-dimensional form are

$$
\theta = 1 \text{ at } X = 1 \quad \text{(Fin base)} \tag{7a}
$$

$$
\left[d\theta/dX\right]_{X=0} = H = \left(Nu/\Theta\psi\right)\theta_t \quad \text{(Convective tip)} \tag{7b}
$$

where  $\theta_i = (T_i - T_a)/(T_b - T_a)$  is the dimensionless tip temperature.

# *2.3 Temperature distribution*

where  $k_j$  and  $k_i$  are the thermal conductivities of the fluid and<br>fin material respectively in Wm<sup>-1</sup>K<sup>-1</sup>.<br>The following dimensionless parameters are used<br> $(X_i \psi, k_s; \theta) = (x/L_i/L_i k, / k_j \cdot (T - T_s)/(T_s - T_s))$  (5a)<br> $\Theta = k_{\sigma j}/k_j = \phi + k_k$ The present problem adopts the homotopy perturbation method [19, 20] for obtaining the temperature distribution from the non-linear governing differential equations. According to HPM, the homotopy of Eq. (6) is written as re  $\theta_i = (T_i - T_a)/(T_b - T_a)$  is the dimensionless tip tem-<br>ture.<br> **Elemperature distribution**<br>
ne present problem adopts the homotopy perturbation<br>
nod [19, 20] for obtaining the temperature distribution<br>
i the non-linear gover **Boundary conditions**<br>
The boundary conditions in non-dimensional form are<br>  $\theta = 1$  at  $X = 1$  (Fin base) (7a)<br>  $\left[ d\theta / dX \right]_{x=0} = H = \left( Nu / \Theta \psi \right) \theta$ , (Convective tip) (7b)<br>
ere  $\theta_i = \left( T_i - T_a \right) / \left( T_b - T_a \right)$  is the dimensio The present problem alops the homotopy perturbation<br>
ethod [19, 20] for obtaining the temperature distribution<br>
om the non-linear governing differential equations. Accord-<br>
g to HPM, the homotopy of Eq. (6) is written as<br>  $\left[ d\theta/dX \right]_{x=0} = H = \left( Nu/\Theta \psi \right) \theta$ , (Convective tip) (7b)<br>
ere  $\theta_i = \left( T_i - T_a \right) / \left( T_b - T_a \right)$  is the dimensionless tip tem-<br>
ature.<br> **Temperature distribution**<br>
The present problem adopts the homotopy perturbation<br>
the pre rature.<br> **Temperature distribution**<br>
The present problem adopts the homotopy perturbation<br>
the d [19, 20] for obtaining the temperature distribution<br>
m the non-linear governing differential equations. Accord-<br>
to HPM, the

$$
d^{2}v/dX^{2} - d^{2}\theta_{0} /dX^{2} + p d^{2}\theta_{0} /dX^{2} + p \left( -m_{1}v^{2} - m_{2}v \right) = 0.
$$
\n(8)

$$
v = v_0 + pv_1 + p^2 v_2 + p^3 v_3 + \dots \tag{9}
$$

The first approximation is considered as

$$
\theta_0 = \theta_t + HX. \tag{10}
$$

Substituting Eq. (9) into Eq. (8) and equating the terms with T. Deshamukhya et al. / Journal of Mechanica<br>
Substituting Eq. (9) into Eq. (8) and equating the terms with<br>
the identical powers of *p* gives<br>  $p^0: d^2v_0/dX^2 - d^2\theta_0/dX^2 = 0$  (11)<br>  $p^1: d^2v_1/dX^2 + d^2\theta_0/dX^2 - m_1v_0^2 - m_$ 

$$
p^0: d^2v_0/dX^2 - d^2\theta_0/dX^2 = 0 \tag{11}
$$

$$
p^{1}:d^{2}v_{1}/dX^{2}+d^{2}\theta_{0}/dX^{2}-m_{1}v_{0}^{2}-m_{2}v_{0}=0
$$
\n(12)

$$
v_1|_{X=0} = 0, \ \left[ d v_1 / d X \right]_{X=0} = 0 \tag{13}
$$

$$
p \cdot a \quad v_2/aX = 2m_1v_0v_1 - m_2v_1 - 0 \tag{14}
$$
\n
$$
v_2|_{v_1} = 0, \quad [dv_2/dX]_{v_2} = 0 \tag{15}
$$

$$
\begin{bmatrix} 2|_{X=0} & 0, & \lfloor \frac{uv_2}{uv_1} \rfloor_{X=0} & 0 \end{bmatrix}
$$
 (10)

$$
v_1|_{x_1} = 0, \quad [dv_1/dX]_{x_1} = 0 \tag{17}
$$

$$
\begin{array}{c}\n\left[x_{3}\right]_{X=0} - 0, & \left[ u v_{3} \right] u A \left[ u_{3} - 0 \right] \\
\left[ u v_{3} u_{3} + u_{3} u_{3} \right]_{X=0} - 0\n\end{array}
$$

# Solving Eqs. (11)-(17) yields the following

$$
v_0 = \theta_t + HX \tag{18}
$$

$$
v_1 = X^2 (m_2 \theta_t + m_1 \theta_t^2)/2 + X^3 (Hm_2 + 2Hm_1 \theta_t)/6 + X^4 H^2 m_1 / 12
$$

s (i.e., i.e., i.e., i.e., 
$$
d\bar{d}
$$
)  
\n
$$
p^0: d^2v_0/dX^2 - d^2\theta_0/dX^2 = 0
$$
\n
$$
p^1: d^2v_1/dX^2 + d^2\theta_0/dX^2 = m_1v_0^2 - m_2v_0 = 0
$$
\n(11) 
$$
p^1: d^2v_1/dX^2 + d^2\theta_0/dX^2 = m_1v_0^2 - m_2v_0 = 0
$$
\n(12) 
$$
p^2: d^2v_2/dX^2 - 2m_1v_0v_1 - m_2v_1 = 0
$$
\n(13) 
$$
p^2: d^2v_2/dX^2 - 2m_1v_0v_1 - m_2v_1 = 0
$$
\n(14) 
$$
p^2: d^2v_2/dX^2 - 2m_1v_0v_1 - m_2v_1 = 0
$$
\n(15) 
$$
p^2: d^2v_2/dX^2 - m_1v_1^2 - 2m_1v_0v_2 - m_2v_2 = 0
$$
\n(16) 
$$
y^1 = 0, [dv_2/dX]_{u=0} = 0, [dv_2/dX]_{u=0} = 0
$$
\n(17) 
$$
Maximize F(Z) = Q_a = \Theta \psi[a\theta/dX]_{u=0}
$$
\n(18) 
$$
v_0 = \theta, + HX
$$
\n(18) 
$$
v_1 = Z^2(\theta, \psi, Nu, Da)
$$
\n(19) 
$$
U = h^2\overline{U}/k_f^2 = Nu^2/\psi
$$
\n(10) 
$$
V = \pi^2(m_2\theta, + 3m_1\theta, \theta^2 + 10Hm_1\theta, \theta, + H
$$

$$
\theta = \lim_{v \to v_0} v = v_0 + v_1 + v_2 + v_3 + \dots \tag{22}
$$

## *2.4 Actual heat transfer rate and fin efficiency*

The actual heat transfer rate is calculated by applying Fourier's law of heat conduction at fin base. Actual heat transfer rate  $(q_a)$  and ideal heat transfer rate  $(q_i)$  in dimensional form are *a*  $\lim_{p\to 1} v = v_0 + v_1 + v_2 + v_3 + ...$  (22) fitness landscape<br> *computational editional heat transfer rate and fin efficiency*<br> *computational editional edition at fin base.* Actual heat transfer<br> *a* and ideal heat transfer

$$
\begin{bmatrix} q_a \\ qi \end{bmatrix} = \begin{bmatrix} k_{\text{eff}} W \left[ dT/dx \right]_{t=L} \\ 2hL(t+W)(1-\phi)(T_b - T_a) + \dot{m}C_p(T_b - T_a) + h t W(T_b - T_a) \end{bmatrix}
$$
\n(23)

Actual heat transfer rate  $(Q_a)$  and ideal heat transfer rate  $(Q_i)$ per unit width in dimensionless form are as follows:

$$
\begin{bmatrix} Q_a \\ Q_i \end{bmatrix} = \frac{1}{k_f (T_b - T_a)} \begin{bmatrix} q_{\text{act}} \\ q_{\text{ideal}} \end{bmatrix} = \begin{bmatrix} \Theta \psi [d\theta/dX]_{x=1} \\ \{2Nu(1-\phi)\}/\psi + RaDa/\psi + Nu \end{bmatrix}.
$$
 iteration  
(24) inertia,

Thus fin efficiency is given by

$$
\eta = Q_a/Q_i \,. \tag{25}
$$

#### **3. Optimization analysis**

*T. Deshamukhya et al. / Journal of Mechanical Science and Technology 32 (9) (2018) 4495-4502*<br>
Substituting Eq. (9) into Eq. (8) and equating the terms with **3. Optimization analysis**<br>
identical powers of *p* gives Since 1 Deshamukhya et al. / Journal of Mechanical Science and Technology 32 (9) (2018) 4495-4502<br>
3 Aubstituting Eq. (9) into Eq. (8) and equating the terms with<br>
3. **Optimization analysis**<br>
Since transferring heat effectively *X X* 2 2 *M*<sub>*N*</sub> *x*<sup>2</sup> = 0 *X*  $\int f(x^2 - x^2)y \, dx^2 - \int f(x^2 - y^2)y \, dx^2 + \int f(x^2 - y^2)y \, dx^2 + \int f(x^2 - y^2)y \, dx^2 - \int f(x^2 - y^2)y \, dx^2 - \int f(x^2 - y^2)y \, dx$ 2 1 *P*  $d^3v_x/dX^2 - m_yv_y = 0$ <br>
2 1 *p*  $d^3v_y/dX^2 = m_yv_x^2 - 2m_yv_y - m_zv_z = 0$ <br>
2 1 **a** fin, so here the aim is to maximize the p<sup>2</sup> :  $d^2v_x/dX^2 - d^2\theta_y/dX^2 = 0$ <br>
2 **a** *p*<sup>2</sup> :  $d^2v_x/dX^2 + d^2\theta_y/dX^2 = m_yv_y = m_zv_y = 0$ <br>
2 **a** *p*<sup>2</sup> : *<sup>X</sup> <sup>X</sup> v dv dX* <sup>=</sup> <sup>=</sup> <sup>=</sup> é ù <sup>=</sup> ë û (15) 3 2 2 2 7. *Deshamukhya et al. /Journal of Mechanical Science and Technology 32 (9) (2018) 4495-4502*<br>
identical powers of *p* gives<br>
identical powers of *p* gives<br>  $p^0 : d^2v_y/dX^2 - d^2\theta_y/dX^2 = 0$ <br>  $p^1 : d^2v_y/dX^2 + d^2\theta_y/dX^2 - m_1v_0^$ *X X*  $\partial$  *x X x*  $\partial$  *x x x a d x d x d d x a d x d x d x d x d x d x d x d x d x d x d x d x d x d x d x d x d x d x T. Deshamukhya et al. / Journal of Mechanical Science and Technology 32*<br>
Substituting Eq. (9) into Eq. (8) and equating the terms with **3. Optimization**<br>
identical powers of *p* gives<br>  $p^0$ :  $d^2v_0/dX^2 - d^2\theta_0/dX^2 = 0$ T. Deshamulthya et al. / Journal of Mechanical Science and Technology 32 (9) (2018) 44<br>
Substituting Eq. (9) into Eq. (8) and equating the terms with<br>  $p^6: d^2v_0/dX^2 - d^2v_0/dX^2 = 0$ <br>  $p^1: d^2v_1/dX^2 + d^2\theta_0/dX^2 - m v_0^2 - m v_$ Substituting Eq. (9) into Eq. (8) and equating the terms with **3. Optimization analysis**<br>
identical powers of *p* gives<br>  $p^6: d^2v_n/dX^2 - d^2\theta_n/dX^2 = 0$ <br>  $p^1: d^2v_n/dX^2 + d^2\theta_n/dX^2 - m_v v_n^2 - m_v v_n = 0$ <br>
(11) a fin, so here the ai are powers or  $p$  gives<br>  $\lim_{a\to 0} \left| \frac{dx^2 - a^2 \theta_a}{dx^2 + a^2 \theta_a} \right| dx \le 0$ <br>  $\left| \frac{dx}{dx} \right|_{x=0} = 0$ <br>  $\left| \frac{dx_1}{dx} \$ Since transferring heat effectively is the<br>  $X^2 - d^2\theta_0/dX^2 = 0$  (11) a fin, so here the aim is to maximize the<br>  $X^2 - a^2\theta_0/dX^2 - m_1v_0^2 - m_2v_0 = 0$  (12) parameters which have been selected three<br>  $Q_0$  [ $dx_1/dX$ ] $_{x=0} =$ Since transferring heat effectively is the primary function of a fin, so here the aim is to maximize the actual heat transfer rate for a specified fin volume by optimizing four crucial fin parameters which have been selected through significant parameter analysis discussed in subsequent section. In this work the positions of each particle (in PSO) is represented in a four dimensional space by  $\phi$ ,  $\psi$ , *Da* and *Nu*. The objective function of the present work thus can be represented by and Technology 32 (9) (2018) 4495-4502<br> **Optimization analysis**<br>
Since transferring heat effectively is the primary function of<br>
in, so here the aim is to maximize the actual heat transfer<br>
e for a specified fin volume by (9) (2018) 4495-4502<br> **analysis**<br>
g heat effectively is the primary function of<br>
nim is to maximize the actual heat transfer<br>
If in volume by optimizing four crucial fin<br>
have been selected through significant pa-<br>
scusse *moe and Technology 32 (9) (2018) 4495-4502*<br> **3. Optimization analysis**<br>
Since transferring heat effectively is the primary function of<br>
at fin, so here the aim is to maximize the actual heat transfer<br>
parameters of a sp **Optimization analysis**<br>
Since transferring heat effectively is the primary function of<br>
in, so here the aim is to maximize the actual heat transfer<br>
for a specified fin volume by optimizing four crucial fin<br>
frameters wh ameters which have been selected through significant pa-<br>neter analysis discussed in subsequent section. In this work<br>positions of each particle (in PSO) is represented in a four<br>nensional space by  $\phi$ ,  $\psi$ ,  $Da$  and  $Nu$ meter analysis discussed in subsequent section. In this work<br>positions of each particle (in PSO) is represented in a four<br>enensional space by  $\phi$ ,  $\psi$ ,  $Da$  and  $Nu$ . The objective<br>oteion of the present work thus can be r eters which have been selected through significant pa-<br>
r analysis discussed in subsequent section. In this work<br>
sitions of each particle (in PSO) is represented in a four<br>
sional space by  $\phi$ ,  $\psi$ , Da and Nu. The obje aalysis discussed in subsequent section. In this work<br>
ons of each particle (in PSO) is represented in a four<br>
all space by  $\phi$ ,  $\psi$ ,  $Da$  and  $Nu$ . The objective<br>
f the present work thus can be represented by<br>  $ize F(Z) = Q_a = \$ 

$$
Maximize F(Z) = Q_a = \Theta \psi \left[ d\theta / dX \right]_{X=1}
$$
 (26)

The volume per unit width of a rectangular fin, which is the constraint here, in dimensionless form depends on  $Nu$  and  $\psi$ function of the present work thus can be represented by<br>
Maximize  $F(Z) = Q_a = \Theta \psi \left[d\theta/dX\right]_{X=1}$  (26)<br>
where  $Z = [\phi, \psi, Nu, Da]$ .<br>
The volume per unit width of a rectangular fin, which is the<br>
constraint here, in dimensionless f

$$
U = h^2 \overline{U} / k_f^2 = N u^2 / \psi. \tag{27}
$$

$$
0.4 \le \phi \le 0.6; \ 0.01 \le \psi \le 0.10; 10 \le Nu \le 50; \tag{28}
$$

#### **4. Particle swarm optimization (PSO)**

(1 $\ln_{mn} \pi_2^2 \theta_2^2 + m_2^2 \theta_2^2 + 20m_1^2 m_2 \theta_2^2 + 10m_1^2 \theta_2^2 + 120$  (28)<br>
23 (11)  $\pi_1 m_1^2 \theta_1^2 / 23 + 28m_1^2 m_2 \theta_1^2 / 24$  (21)<br>
23 (11)  $\pi_1 m_1^2 \theta_1^2 / 23 + 28m_1^2 m_1 \theta_1^2 / 25 + 28m_1^2 m_1^2 m_1 \theta_1^2 / 25 + 28m_1^$ *b*  $x^* (H^2m_m, +2H^2m_1^2\theta) / 72 + x^2H^2m_1^2/352$ <br>  $b_1 = x^* (1m_m\theta_1^2\theta_1^2 + m_2^2\theta_1 + 20m_1^2m_1\theta_1^2) / 720$ <br>  $b_2 = \frac{x^* (1m_1m_2\theta_1^2 + m_2\theta_1 + 20m_1^2m_1\theta_1^2 + 10m_1^2\theta_1^2)}{4.2}$ <br>  $\therefore$   $\frac{x^* H^2(m_1m_1\theta_1^2/65 + 5x^$  $\frac{d}{dt} = \int_{\pi}^{T} \int_{\pi$ +  $X^*(H^Tm,m_1+2H^Tm_1^2\theta)/T^2 + X'H^Tm_1^2/252$ <br>
+  $X^*(1m,m_2\theta)^2 + m_2\theta^2 + 0 \tan^2\theta^2/252$ <br>
+  $X^*(1m,m_2\theta)^2 + N^2(H^2m_1\theta^2/252)$ <br>
+  $X^*(Hm_1\theta^2/\theta + 2 \cot^2\theta) + Y(Hm_1^2m_1\theta^2/\theta^2/2$ <br>
+  $X^*(Hm_1\theta^2/\theta^2 + 0 \cot^2\theta) + X^*(Hm_1^2m_1\theta^2/\theta$ 1. = :<br>
1. **Particle swarm optimization (PSO)**<br>
1. **h** final temperature distribution in the fin is given as<br>
by Kennedy and Ebenhat [21], which works quite<br>
22) finals and ficted by the social behavior of birds, PSO is Inspired by the social behavior of birds, PSO is developed by Kennedy and Eberhart [21], which works quite well in such scenarios which involve nonlinearity as well as noisy fitness landscapes and fetches the near optimum values in less computational effort. It is a non-calculus based optimization technique to solve the global optimization problems. The algorithm starts by creating random particles and assigning ran dom positions to these particles. At every iteration personal previous best or p-best and global best or g-best play a crucial role in improving the particle's positions. The velocity and hence position of each particle is thus updated based on the following equations: (28)<br>  $0.0001 \leq Da \leq 0.01$ . (28)<br> **Particle swarm optimization (PSO)**<br>
Inspired by the social behavior of birds, PSO is developed<br>
Kennedy and Eberhart [21], which works quite well in<br>
the senarios which involve nonlinear **Particle swarm optimization (PSO)**<br>**particle swarm optimization (PSO)**<br>**is Kennedy and Eberhart [21]**, which works quite well in<br>**Kennedy and Eberhart [21]**, which works quite well in<br>the senarios which involve nonlinear

$$
v_{i+1} = w v_i + c_1 r_1 (pbest_i - x_i) + c_2 r_2 (gbest_i - x_i)
$$
 (29)

$$
x_{i+1} = x_i + v_{i+1} \tag{30}
$$

 $(T_b - T_a) \left[ q_{ideal} \right]$   $\left( {2Nu(1-\phi)} \right)/\psi + RaDa/\psi + Nu$  [0,1]. There are three main parts of Eq. (29) which are the **and fin efficiency**<br>
computational effort. It is a non-calculus based c<br>
technique to solve the global optimization problem<br>
intim stars by ereating random particles and assessment<br>
and positions to these particles. At e *F* and  $\overline{f} = \frac{1}{\left[\frac{2hL}{dt} + W\right)^{1/2} + V_1 + W_2 + W_1}$  and  $\overline{f} = \frac{1}{\left[\frac{2hL}{dt} + W\right)(1 - \phi)(T_1 - T_1) + hW(T_1 - T_2)}$  (22) and ideal heat transfer rate is calculated by applying Four-<br> *Actual heat transfer rate and fin eff* **Example the the transfer rate and fine fiftiency<br>
<b>Computational** effort. It is a non-calculus based optimization<br>
the actual heat transfer rate is calculated by applying Fou-<br>
the metal of the spoke the global optimizat (22) finess landscapes and fetches the near optimum value<br>
(22) finess landscapes and fetches the mear optimum value<br>
computational effort. It is a non-calculus based optir<br>
technique to solve the global optimization prob efficiency<br>
computational effort. It is a non-calculus based optimization<br>
technique to solve the global optimization problems. The algo-<br>
ulated by applying Fou-<br>
timm starts by creating random particles and saigining ra  $\sigma = \lim_{\mu \to \infty} y + y_1 + y_2 + y_3 + ...$  (22) thus bandscapes and felches the near optimum values in less<br>
computational effort it is a non-calculus based optimization<br>  $A$ -tutal heat transfer rate is calculated by applying Fou-<br> e  $(q_a)$  and ideal heat transfer rate  $(q_i)$  in dimensional form proving the interpretion of each proving the reduce position of each proving  $\begin{bmatrix} q_a \\ q_i \end{bmatrix} = \begin{bmatrix} k_{af}W\left[dT/dx\right]_{at} \\ 2hL(t+W)(1-\phi)(T_s - T_s) + \dot{m}C_s(T_s - T_s) + hW(T_s -$ The velocity of the particles gets updated as per Eq. (29) while the position gets changed by Eq. (30). Here the acceleration constants  $c_1$  and  $c_2$  are cognitive and social factors respectively which collectively help in changing the velocity during iteration. The random numbers  $r_i$  and  $r_j$  are in the range inertia, the cognitive and the social components respectively. A good choice of ' *w* ' which is associated with the momentum of the particles helps to maintain a balance between ex ploitation and exploration. A modification in original version of PSO considering variable inertia weight instead of a con-



Fig. 2. Significance analysis plot for  $Q_a$  vs  $\psi$  with variation of *Nu*.



Fig. 3. Significance analysis plot for  $Q_a$  vs  $\phi$  with variation of *Da*.

stant one can yield better results [22]. In this work with the help of reflecting boundary condition [23] the particles are prevented from flying out of the feasible space.

## **5. Significant variable analysis**

A significance analysis has been carried out to gauge the most important or highly sensitive variables out of the seven variables involved in the present problem. This analysis helps to select the most sensitive variables in a problem that influ ences the output to a greater extent and thus reducing the computational effort and time by omitting the variables which has negligible or insignificant role in the overall output. In the current study the change in heat transfer rate with respect to the change in the values of *Ra*, *Da*, *Nu*,  $\phi$ , *T<sub>a</sub>*, *T<sub>b</sub>* and  $\psi$  are The temperature noted respectively and based on the gradient the four significant variables ( $Nu, \psi, \phi$  and  $Da$ ) are selected which has been finally optimized by PSO to achieve greater heat transfer rate by judicial combination of these variables.

Figs. 2 and 3 are generated to show the individual effect of these parameters on heat transfer rate. It is seen from Fig. 2

Table 1. Comparison of temperature distribution between the present analytical and numerical results  $(Ra = 10^4, Da = 0.0001, Nu = 20, k_f =$ 0.02 Wm<sup>-1</sup>K<sup>-1</sup>,  $k_s = 200$  Wm<sup>-1</sup>K<sup>-1</sup>,  $T_b = 373K$ ,  $T_a = 300K$ ,  $\phi = 0.4$ ,  $w = 0.05$ .

X	Ĥ	Difference	
	Numerical	<b>HPM</b>	
0.0	0.52307	0.52293	$-0.00014$
0.2	0.53810	0.53976	0.00166
0.4	0.58854	0.59101	0.00247
0.6	0.67992	0.68090	0.00098
0.8	0.81388	0.81428	0.00040
1.0			

that the heat transfer rate enhances with the increase in *Nu* due to more convection current through the fin surfaces. For each *Nu*, there is an optimum value of  $\psi$  at which the heat transfer rate from the fin attains a maximum value and the optimum value of  $\psi$  increases by increasing *Nu*. Again heat transfer rate is a function of porosity parameter, thermal conductivities,  $\psi$  and gradient at the base. With the increase in  $\psi$ , gradient is decreased and thus at a particular value of  $\psi$  heat transfer becomes maximum. On the other hand, heat transfer rate is decreased by increasing porosity parameter after a certain value. In porous fins, the effective surface area increases, which tends to increase the heat transfer rate, but simultaneously the effective thermal conductivity decreases, which results in decreasing the heat transfer rate. Thus, depending upon the other thermo-physical parameters, the heat transfer rate from the porous fin may becomes less than the solid fin. However, actual heat transfer rate is increased with the in crease in Darcy number as seen in Fig. 3. With increasing *Da*, the fin becomes more permeable which facilitates better flow of fluid through the pores and this enhances the actual heat transfer rate.

## **6. Results and discussions**

The analytical model is validated by comparing the obtained temperature distribution with that obtained from a nu merical model, as seen in Table 1. For the numerical result, Eq. (6) is first discretised by central difference scheme of fourth order accuracy followed by solving the difference equations by Gauss Seidel iterative procedure, satisfying the boundary conditions and a convergence criterion of  $10^{-6}$ . As seen in Table 1, the temperature distribution from the present analytical model show a high accuracy with that obtained from the numerical model.

and  $\psi$  are The temperature along the fin length is slightly lower when the value of Darcy number (*Da*) is higher as seen in Fig. 4. This is because with the increase in *Da*, the permeability in creases which allows more heat to flow out of the fin through the pores thereby reducing the temperature. Again, the tem perature along the fin length drops as the porosity of the fin increases. The reason for this trend is two folds. Firstly, at



Fig. 4. Temperature distribution with the variation of *Da* and  $\phi$ .



Fig. 5. Efficiency plots as a function of  $Nu$ ,  $\psi$  and  $\phi$ .

high porosity the effective conductivity suffers due to the removal of excess material and secondly due to higher pore density ambient cool air mixes well with the hot air in fin surface. As seen in Figs. 5 and 6 as  $\psi$  increases, fin efficiency increases. It is because at a particular thickness, as the length decreases the conduction resistance reduces. Again, in Fig. 5, it is seen that efficiency decreases with increasing  $\phi$  value. This is because as the pores increases, the heat transfer drops due to the removal of metal which in turn reduces the conductive heat transfer.

Darcy number which is related with permeability improves the efficiency at lower values as seen in Fig. 6. It is because at higher *Da* the increase in permeability comes at the cost of increase in ideal heat transfer which nullifies the effect of increased value of actual heat transfer to the extent of lowering the efficiency. The same justification can be given for the fall in efficiency in the same figure with increase in *Ra* values. As the buoyant forces increases, the ideal heat transfer increases at a rate higher than the actual heat transfer. The vari-



Fig. 6. Efficiency plots as a function of *Ra*, y and *Da.*



Fig. 7. Convergence plot of significant variable  $\phi$  with three different tip temperatures.

ables are varied to understand their importance in temperature distribution and efficiency. This being done, the four important parameters are chosen for optimization to determine a sufficiently good heat transfer rate.

Figs. 7-10 show the convergence plots of four parameters with iteration number. Each parameter is studied for three different fin tip temperatures under consideration. PSO being a metaheuristic technique converges faster to near optimum or good values. It however cannot guarantee the best or optimum results [24] due to randomness present in the algorithm which is an integral identity of metaheuristics. However the aim here is to reach a seemingly good near optimum value. Fig. 7 shows the plot of porosity whose feasible region lies from 0.4 to 0.6. The graph shows fluctuating tendency of the variable in the initial iterations which slowly converges to a good point as shown in figure. Similar trend is seen in the other parameters (<sup>y</sup> and *Nu* and *Da*) in Figs. 8-10.

The initial fluctuations seen in all the four variables in three



Fig. 8. Convergence plot of  $\psi$  with different tip temperatures.



Fig. 9. Convergence plot of *Nu* with different tip temperatures.



Fig. 10. Convergence plot of *Da* with different tip temperatures.

different tip temperatures can be understood from PSO's nature of scanning the domain for the best possible combination of values of the variables that yield the maximum value of

0.10  $\begin{bmatrix} 1 & \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$  ----  $\theta_1 = 0.2500$   $\begin{bmatrix} -\frac{1}{2} & \frac{1}{2} & \frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & \frac{1}{2} \\ -\frac{1}{2} & \frac{1}{2} & \frac{1}{2} \end{bmatrix}$   $-\theta = 0.5000$  |  $T_b = 373 \text{ K}, T_a = 300 \text{ K}.$ Table 2. Optimized values of the variables obtained from PSO for three

Tip tem- perature $\theta$ .	Φ	$\psi$	Da	Nu	$O_a$ (max)
0.25	0.4699	0.05345	0.001	31.5119	1098.591
0.50	0.48861	0.06172	0.001	35.9544	1243.739
0.75	0.53133	0.04949	0.001	31.2791	2697.160



Fig. 11. Convergence pattern of fitness function (heat transfer rate) by varying Rayleigh number with three different tip temperatures.

objective function without disobeying the constraint. Once a good combination is obtained, the variables converge to their individual best or near best values so as to produce an en hanced heat transfer rate. Table 2 shows the optimized values of the variables obtained from PSO for three different tip tem peratures.

Rayleigh number plays a crucial role in heat transfer at higher tip temperatures in convective medium. This dimen sionless parameter which is associated with buoyant forces, results in better heat transfer through the fin at higher values. This happens because at higher values of Rayleigh number the buoyancy increases which in turn encourages convective heat transfer to take place at a greater scale. The trend obtained from the convergence analysis of the fitness function i.e., actual heat transfer rate is in accordance with the said logic. In Fig. 11, the convergence pattern is studied with respect to the number of iterations for two different Rayleigh number for three tip temperatures. The heat transfer rate increases with increasing tip temperature due to the increasing temperature 0 20 40 60 80 100 120 140 difference. Also for each tip temperature, the heat transfer rate is more for higher Rayleigh number because of more driving force due to buoyancy effect.

#### **7. Conclusions**

A rectangular porous fin is studied analytically (using HPM) to obtain the temperature distribution and efficiency

under convective fin tip conditions. For three different tip temperatures four important variables are optimized by PSO to maximize the actual heat transfer rate. The key findings of this study are summarized as follows:

- Fin tip temperature is decreased with the increase in *Da* and  $\phi$ .
- Actual heat transfer rate is increased by increasing *Nu* and *Da*. However, with the increase in  $\phi$  heat transfer rate decreases due to mainly conduction loss.
- The variables after being optimized by PSO shows in creased heat transfer rate for all the three different tip temperatures and the heat transfer rate enhances with the increase in tip temperature.
- The actual heat transfer increases marginally with the in crease in *Ra* for low tip temperatures whereas the change is somewhat appreciable for high tip temperature.

#### Nomenclature--

- $C_P$  : Specific heat of air at constant pressure (Jkg<sup>-1</sup>K<sup>-1</sup>) g  $\therefore$  Acceleration due to gravity (ms<sup>-2</sup>) *h* : Convective heat transfer coefficient  $(W \text{ m}^{-2} K^{-1})$ *kR* : Thermal conductivity ratio *L* : Fin length (m) *K* : Permeability of porous fin  $(m^2)$  $q_{act}$  : Actual heat transfer rate per unit width  $(\text{Wm}^{-1})$ *Q<sup>a</sup>* : Dimensionless actual heat transfer rate per unit width  $q_{ideal}$  : Ideal heat transfer rate per unit width  $(Wm^{-1})$ *Ra* : Rayleigh number *T* : Local fin surface temperature (°C)
- *T<sup>a</sup>* : Ambient temperature (°C)
- 
- $T_b$  : Fin base temperature ( ${}^{\circ}$ C)
- $T_t$  : Fin tip temperature (°C)<br>  $\overline{U}$  : Fin volume per unit wide
- $\overline{U}$  : Fin volume per unit width  $(m^2)$
- *U* : Dimensionless fin volume per unit width
- *x* : Axial length measured from fin tip (m)
- *X* : Dimensionless distance, *x/L*

## *Greek symbols*

- $\phi$  : Porosity
- $\beta$  : Coefficient of thermal expansion  $(K^{-1})$
- $\gamma$  : Kinematic viscosity (m<sup>2</sup>s<sup>-1</sup>)
- $\psi$  : Thickness to length ratio
- $\eta$ : Fin efficiency
- $\theta$  : Dimensionless temperature,  $(T-T_a)/(T_b-T_a)$

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