

# Active vibration control based on modal controller considering structure-actuator interaction†

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### **Abstract**

Active vibration control to suppress structural vibration of the flexible structure is investigated based on a new control strategy considering structure-actuator interaction. The experimental system consists of a clamped-free rectangular plate, a controller based on modal control switching, and a magnetostrictive actuator utilized for suppressing the vibrations induced by external excitation. For the flexible structure, its deformation caused by the external actuator will affect the active control effect. Thus interaction between structure and actuator is considered, and the interaction model based on magnetomechanical coupling is incorporated into the control system. Vibration reduction strategy has been performed resorting to the actuator in optimal position to suppress the specified modes using LQR (linear quadratic regulator) based on modal control switching. The experimental results demonstrate the effectiveness of the proposed methodology. Considering structure-actuator interaction (SAI) is a key procedure in controller design especially for flexible structures.

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*Keywords*: Flexible structures; Vibration suppressing; Structure-actuator interaction; Modal control switching

## **1. Introduction**

As is well known, an effective approach in making some structures more efficient is to reduce their weight especially in the area of aerospace. As a result, these structures are typically made thin. However, some components of these structures are prone to experiencing excessive vibrations such as flutter and fatigue failure. In order to solve these vibration problems, both passive and active vibration control approaches have been developed and implemented. For instance, the inertial actuator as a usual vibration control method [1, 2], is a mass supported on a spring and damper. Due to its structural features, it sometimes is not as lightweight as is required. Subsequently, active actuators made of smart material are applied onto the flexible structures. Many investigators have implemented experimental research for active vibration control employing smart material actuators [3-8].

Compared to piezoelectric ceramic material (PZT) and shape memory alloy (SMA), giant magnetostrictive material (GMM) has some advantages such as generation of large forces, high load bearing capability, high magnetomechanical coupling, and rapid reaction. In view of the advantages of GMM, giant magnetostrictive actuators are a promising technology in active vibration control [9, 10]. Moon and Lim have designed a linear magnetostrictive actuator (MSA) using Terfenol-D, and analyzed the characteristics of the actuator by implementing a series of experimental and numerical tests which confirmed that the linear MSA had a good control performance [5]. Braghin and Cinquemani introduced a linear model of magnetostrictive actuators that was valid in a range of frequencies below 2 kHz, and the model was verified through experiments [9]. Zhou and Zheng presented a nonlinear constitutive model-based vibration control system for giant magnetostrictive actuators, and the effectiveness of a real control system for suppressing a vibration was demonstrated by a case study with negative velocity feedback [11]. However, a more widespread application for giant material actuators (GMA) is restrained by an inherent property of GMM. Therefore, researchers have started to focus on the characteristics of GMA, such as nonlinear hysteresis phenomena [12-15] which can cause the instability of the system. The nonlinear behaviors can be modeled by a range of approaches [16-20]. Liu and Zhang constructed a dynamic model in a given rate range for a rate-dependent hysteresis system, and a PID control combined with feedforward compensation was employed on a giant magnetostrictive actuator. The results gotten by simulations and experiments indicated the effectiveness of the proposed methods [21]. Aljanaideh and Rakheja proposed a phenomenological hysteresis model for the hysteresis nonlinearities of

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Fig. 1. GMA model.

a magnetostrictive actuator, and the model effectively describe the nonlinear hysteresis properties of the actuator [22]. From previous investigations of vibration control systems, mostly the nonlinear behaviors of actuators are considered by using compensation algorithms, with little consideration of the relationship between actuator and structure that is also important for the whole system. In Ref. [23], computational models considering the interaction between the GMA and structure were developed, and the results demonstrated that consideration of the CSI and the dynamics of the GMA can improve the performance of a controller significantly. In this paper, the main purpose is to demonstrate the control strategy with considering the interaction between the actuator and the structure for a flexible structure, thus the nonlinear behavior for GMA like rate-dependency hysteresis will not be described.

For flexible structures, application of an external force will generate a reactive force to the object producing the external force. For smart material such as piezoelectric/piezomagnetic actuators, this interaction may affect the performance of these actuators. In this work, in order to validate the proposed method, an interaction model between structure and actuator is derived and constructed. Based on the model, the experiment is implemented utilizing the proposed control strategy. The experiment results demonstrate the effectiveness of the control model.

## **2. Model establishment of giant magnetostrictive actuator**

In order to improve the control efficiency, the actuating



Fig. 2. Equivalent circuit of GMA.

principle of the actuator should be studied. A giant magnetostrictive actuator (GMA) is used in this paper as shown in Fig. 1 and the parameters of GMM (giant magnetostrictive material) are shown in Table 1. *m M M m* **<b>***m m MARR m* 

# *2.1 Positive effect of GMA*

When the voltage is applied to the ends of a coil, GMM will produce a deformation along the axial direction due to the Joule effect. This process is also named the "positive effect" and is shown in Fig. 2.

As is well known, magnetomotive forces (briefly named MMF) can be defined by their forms, and the two of them are given as follows

$$
\begin{cases}\nMMF = NI \\
MMF = R_m \Phi\n\end{cases} (1)
$$

2. Equivalent circuit of GMA.<br>
ciple of the actuator should be studied. A giant magne-<br>
cictive actuator (GMA) is used in this paper as shown in<br>
1 and the parameters of GMM (giant magnetostrictive<br>
reial) are shown in Ta . 2. Equivalent circuit of GMA.<br>
nciple of the actuator should be studied. A giant magne-<br>
trictive actuator (GMA) is used in this paper as shown in<br>
1. 1 and the parameters of GMM (giant magnetostrictive<br>
terial) are sho where *N* is the number of coil turns; *I* is the coil current;  $R_m$  is magnetic reluctance; and  $\Phi$  is the magnetic flux. Here eddy current effect is neglected. According to Eq. (1), we can get that When the voltage is applied to the ends of a coil, GMM will<br>duce a deformation along the axial direction due to the<br>le effect. This process is also named the "positive effect"<br>is shown in Fig. 2.<br>As is well known, magneto As is well known, magnetomotive forces (briefly named<br>
MF) can be defined by their forms, and the two of them are<br>
en as follows<br>  $\int MMF = NI$  (1)<br>  $MMF = R_m \Phi$  (1)<br>
ere *N* is the number of coil turns; *I* is the coil current;<br> is well known, magnetomotive forces (briefly named<br>  $\gamma$ ) can be defined by their forms, and the two of them are<br>
as follows<br>  $MMF = NI$  (1)<br>  $MMF = N_{m}\Phi$  (1)<br>  $MF = R_{m}\Phi$  (1)<br>  $MF = R_{m}\Phi$  (1)<br>
is magnetic reluctance; and  $\Phi$  is

$$
\Phi = \frac{NI}{R_n} \,. \tag{2}
$$

Furthermore, the output force of GMM can be expressed by

$$
F = \Phi K_F = \Phi \frac{1}{d_{33}} \tag{3}
$$

where  $K_F$  is the interaction coefficient between force and flux and  $d_{33}$  is piezomagnetic strain constant. Substituting Eq. (2) into Eq. (3), a new form with respect to the current is written as ere *N* is the number of coil turns; *I* is the coil current;<br>
is magnetic reluctance; and Φ is the magnetic flux.<br>
re eddy current effect is neglected. According to Eq. (1),<br>
can get that<br>  $\Phi = \frac{NI}{R_m}$ . (2)<br>
Furthermor *Reflered According to Eq.* (1), the magnetic reluctance; and  $\Phi$  is the magnetic flux.<br> *R* eddy current effect is neglected. According to Eq. (1), the magnet that<br>  $\frac{N}{R_m}$ . (2)<br>  $\frac{N}{R_m}$ . (2)<br>  $\Phi K_F = \Phi \frac{1}{d_{33}}$ 

$$
F = \frac{NIK_F}{R_m} = \frac{NI}{R_m d_{33}} = \frac{NU}{RR_m d_{33}}.
$$
\n(4)

## *2.2 Inverse effect of GMA*

Aside from the positive effect of GMM, there also exists an inverse effect named the "Villari effect" shown in Fig. 3. When a force acts on the end of GMM, the voltage of the coil



will be changed.

For GMM, constitutive equation is given by

$$
\begin{cases}\n\varepsilon = \frac{\sigma}{E} + d_{33}H \\
B = d_{33}\sigma + \mu^{\sigma}H\n\end{cases}
$$
\n(5)  $\mu$ 

*B d H*<sup>s</sup> <sup>s</sup> <sup>ï</sup> = + <sup>î</sup> where  $\varepsilon$ ,  $\sigma$  and *E* are strain, stress and elastic modulus, respectively. *B*, *H* and  $\mu^{\sigma}$  are magnetic induction, magnetic field strength and magnetic conductivity, respectively. Moreover, according to the Faraday theorem, the induction voltage can be described in the following form 3. Equivalent circuit of inverse effect.<br>
3. Equivalent circuit of inverse effect.<br>
as<br>
be changed.<br>  $\epsilon = \frac{\sigma}{E} + d_{ss}H$ <br>  $\epsilon = \frac{\sigma}{E} + d_{ss}H$ <br>
(5)  $E$ ,  $d_{\text{in}}$  a 3. Equivalent circuit of inverse effect.<br> **as**<br>  $\Delta U = \left(\frac{Nd_x SE}{l_a}\right) \frac{dx}{dt} + \left(L \frac{R_a}{NK_F}\right) \frac{1}{c}$ <br>
For GMM, constitutive equation is given by<br>  $\begin{cases}\n\varepsilon = \frac{\sigma}{E} + d_x H & (5) \\
B = d_y \sigma + \mu^c H & (5) \\
B = d_y \sigma + \mu^c H & (6) \\
\end{cases}$ <br>
enere  $\vare$ 

$$
\Delta U = N d_{33} S E \frac{d \varepsilon}{dt} + L \frac{dI}{dt}
$$
\n<sup>(6)</sup>

where *S* is the cross-sectional area of the GMM rod, and *L* is the induction coefficient. From Eq. (6), it can be seen that the induction voltage has a relation with the strain change ratio and the current change ratio.

## **3. Structure-actuator interaction model**

For an active control system with an intelligent material actuator, theoretically only the actuator has a control effect on the structure, and the structure (i.e. the controlled object) has no influence on the actuator. However, especially for intelligent material, the structure will produce an opposing force to the actuator. The control stability is then destroyed due to the inverse effect. It is important that structure-actuator interaction (SAI) should be considered when a control system is designed including intelligent materials such as piezoelectricity and piezomagnetism. SAI will be described as follows. *s* in induction order the relation with the strain change ratio<br> *AU* =  $Z_1 \frac{dx_1}{dt} + Z_2$ <br>
induction voltage has a relation with the strain change ratio<br>
of an active control eystem and intelligent material ac-<br>
for an tor, theoretically only the actuator has a control effect on<br>
structure c i.e. the controlled object) has<br>
influence on the structure (i.e. the controlled object) has<br>
influence on the actuator. However, especially for in structure, and the structure (i.e. the controlled object) has<br>influence on the actuator. However, especially for inelli-<br>in order to indicate<br>at that actuator. The control stability is then destroyed due to the<br>actuator.

When GMM is driven by voltage, it will generate axial displacement which satisfies the material mechanics theory

$$
x = \varepsilon l_m \tag{7}
$$

where *x* is axial displacement;  $l_m$  is the length of the GMM ( rod. According to Eq. (7), the change rate of axial displacement can then be expressed by

$$
\frac{dx}{dt} = l_m \frac{d\varepsilon}{dt} \,. \tag{8}
$$

In addition, Eq. (4) can be shown in the form of current



Fig. 4. System model considering SAI.

changing with respect to time, thus Eq. (6) also can be written as

$$
\Delta U = \left(\frac{Nd_{33}SE}{l_m}\right)\frac{dx}{dt} + \left(L\frac{R_m}{NK_F}\right)\frac{dF}{dt}.
$$
\n(9)

Due to the significant changes for some parameters such as *E*,  $d_{33}$  and  $\mu^{\sigma}$  under the action of pressure and temperature, correction coefficients will be introduced into Eq. (9), thus changing it to

$$
\Delta U = \gamma_1 \left( \frac{Nd_{33}SE}{l_m} \right) \frac{dx}{dt} + \gamma_2 \left( L \frac{R_m}{NK_F} \right) \frac{dF}{dt}
$$
(10)

**1**<br> **l**  $\frac{d}{dx}$  **l**  $\$  $\frac{1}{\sqrt{2}}$   $\frac{1}{\sqrt{2}}$  where  $\gamma_1$  and  $\gamma_2$  are correction factors. From Eq. (10), it can be seen that  $\gamma_1 \left( \frac{Nd_{33}SE}{l_m} \right)$  and  $\gamma_2 \left( L \frac{R_m}{NK_F} \right)$  only have idering SAI.<br>
to time, thus Eq. (6) also can be written<br>  $+\left(L\frac{R_m}{NK_F}\right)\frac{dF}{dt}$ . (9)<br>
t changes for some parameters such as<br>
er the action of pressure and tempera-<br>
cients will be introduced into Eq. (9),<br>  $\frac{k_x}{dt} + \gamma_2 \$ msidering SAI.<br>
t to time, thus Eq. (6) also can be written<br>  $\frac{x}{t} + \left(L \frac{R_m}{NK_F}\right) \frac{dF}{dt}$ . (9)<br>
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ader the action of pressure and tempera-<br>
ficients will be introduced into Eq. (9),<br>  $\frac{dx}{dt} + \$ 6) also can be written<br>
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F<br>
F<br>
(10)<br>
rs. From Eq. (10), it<br>  $\left(L\frac{R_m}{NK_F}\right)$  only have<br>
ers of GMM, so Eq.<br>
(11) relations with the characteristic parameters of GMM, so Eq. (10) is simplified as  $U = \left(\frac{Nd_{33}SE}{l_{\pi}}\right) \frac{dx}{dt} + \left(L \frac{R_n}{NK_r}\right) \frac{dF}{dt}$ . (9)<br>
we to the significant changes for some parameters such as<br>  $d_{33}$  and  $\mu''$  under the action of pressure and tempera-<br>
correction coefficients will be introduc  $\frac{d}{dx} + \left(\frac{V}{c} \frac{V}{NK_F}\right) \frac{d}{dt}$ . (9)<br>
e significant changes for some parameters such as<br>  $d \mu^{\sigma}$  under the action of pressure and tempera-<br>
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ge it to<br>  $\frac{Nd_{33}SE}{l_m}\frac{dx}{dt} + \gamma_2 \left(L \frac{R_m}{NK_F}\right) \frac{dF}{dt}$  (10)  $\Delta U = \left(\frac{Nd_{33}S_E}{l_m}\right) \frac{dx}{dt} + \left(L\frac{R_m}{NK_F}\right) \frac{dt}{dt}$ . (9)<br>
Due to the significant changes for some parameters such as<br>  $d_{33}$  and  $\mu^{\sigma}$  under the action of pressure and tempera-<br>
e, correction coefficients will be in o the significant changes for some parameters such as<br>
and  $\mu^{\sigma}$  under the action of pressure and tempera-<br>
rection coefficients will be introduced into Eq. (9),<br>
nging it to<br>  $\gamma_1 \left( \frac{Nd_{33}SE}{l_m} \right) \frac{dx}{dt} + \gamma_2 \left( L \frac$ the significant changes for some parameters such as<br>
and  $\mu^{\alpha}$  under the action of pressure and tempera-<br>
tion coefficients will be introduced into Eq. (9),<br>
ing it to<br>  $\left(\frac{Nd_{33}SE}{l_m}\right) \frac{dx}{dt} + \gamma_2 \left(L \frac{R_n}{NK_F}\right) \frac{dF$  $u^{\sigma}$  under the action of pressure and tempera-<br>
a coefficients will be introduced into Eq. (9),<br>
to<br>  $t_{\frac{1}{2}s}E\left(\frac{dx}{dt} + r_2\left(L\frac{R_m}{NK_F}\right)\frac{dF}{dt}$  (10)<br>  $\gamma_2$  are correction factors. From Eq. (10), it<br>
hat  $\gamma_1\left(\frac$ <sup>33</sup>C<sub>m</sub> and  $\gamma_2 \left( L \frac{R_m}{NK_F} \right)$  only have<br>
eristic parameters of GMM, so Eq.<br>
(11)<br>
and  $Z_2 = \gamma_2 \left( L \frac{R_m}{NK_F} \right)$ . Note  $Z_1$ <br>
stants due to the correction factors,<br>
nly applied to the current case and for some parameters such as<br>
on of pressure and tempera-<br> *R*<sub>*m*</sub>  $\frac{R_m}{NK_F}$   $\frac{dF}{dt}$  (10), it<br>
and  $\gamma_2 \left( L \frac{R_m}{NK_F} \right)$  only have<br>
parameters of GMM, so Eq.<br>
(11)<br>  $Z_2 = \gamma_2 \left( L \frac{R_m}{NK_F} \right)$ . Note  $Z_1$ <br>
the to r some parameters such as<br>
of pressure and tempera-<br>
e introduced into Eq. (9),<br>  $\frac{dF}{dr}$  (10)<br>  $\frac{dF}{dt}$  (10)<br>  $\frac{dF}{d\tau}$  (10), it<br>  $\frac{d\gamma_2 \left(L \frac{R_m}{NK_F}\right)}{NK_F}$  only have<br>
ameters of GMM, so Eq.<br>
(11)<br>  $\gamma_2 \left(L \frac{R$ 

$$
\Delta U = Z_1 \frac{dx}{dt} + Z_2 \frac{dF}{dt}
$$
\n(11)

where 
$$
Z_1 = \gamma_1 \left( \frac{Nd_{33}SE}{l_m} \right)
$$
 and  $Z_2 = \gamma_2 \left( L \frac{R_m}{NK_F} \right)$ . Note  $Z_1$ 

that  $Z_2$  and are not constants due to the correction factors, and the two values are only applied to the current case and gotten by identifying trial results.

In order to indicate that some parameters of GMM influence the parameter  $K_F$ , another correction coefficient is introduced. Substituting Eq. (11) into Eq. (4), the output force is given by considering structure-actuator interaction that  $\gamma_1 \left( \frac{M_{\text{d,s}} - M_{\text{d,s}}}{I_m} \right)$  and  $\gamma_2 \left( L \frac{r_{\text{d,s}}}{N K_F} \right)$  only have<br>
th the characteristic parameters of GMM, so Eq.<br>
fified as<br>  $\frac{k}{\pi} + Z_2 \frac{dF}{dt}$  (11)<br>  $\gamma_1 \left( \frac{Nd_{\text{d,s}} S E}{I_m} \right)$  and  $Z_2 = \gamma_2 \left($ ations with the characteristic parameters of GMM, so Eq.<br> *NU* =  $Z_1 \frac{dx}{dt} + Z_2 \frac{dF}{dt}$  (11)<br>
ere  $Z_1 = \gamma_1 \left( \frac{Nd_{33}SE}{l_m} \right)$  and  $Z_2 = \gamma_2 \left( L \frac{R_m}{NK_F} \right)$ . Note  $Z_1$ <br>  $Z_2$  and are not constants due to the correcti  $\left(\frac{l_m}{l_m}\right)^{2m\omega}$   $\left(\frac{2}{N}K_F\right)^{2m\omega}$  and  $\frac{2}{l_m}$  and  $\frac{Z_2}{l_m} = \gamma_2 \left(L\frac{R_m}{N K_F}\right)$ . Note  $Z_1$ <br>
t constants due to the correction factors,<br>
are only applied to the current case and<br>
t trail results.<br>
te that ons with the characteristic parameters of GMM, so Eq.<br>
s simplified as<br>  $J = Z_1 \frac{dx}{dt} + Z_2 \frac{dF}{dt}$  (11)<br>  $\frac{R_a}{L} = \gamma_1 \left( \frac{Nd_{33}SE}{l_m} \right)$  and  $Z_2 = \gamma_2 \left( L \frac{R_m}{N K_F} \right)$ . Note  $Z_1$ <br>  $Z_2$  and are not constants due to t and the two values are only applied to the current case and<br>gotten by identifying trial results.<br>In order to indicate that some parameters of GMM influ-<br>roce the parameter  $K_F$ , another correction coefficient is in-<br>trodu

$$
F = \gamma_3 \frac{NIK_F}{R_m} = \frac{\gamma_3 \left( U - Z_1 \frac{dx}{dt} - Z_2 \frac{dF}{dt} \right) NK_F}{RR_m} \,. \tag{12}
$$

From Eq. (12), we can see that the coefficient  $\frac{\gamma_3 N K_F}{R}$  $\frac{NK_F}{RR_m}$ *RR*

only has relation with the material characteristic parameters of GMA. Therefore the coefficient is replaced here by  $Z_3$ , and

$$
F = Z_3 U - Z_1 Z_3 \dot{x} - Z_2 Z_3 \dot{F}.
$$
\n(13)

According to the equations shown above, the system model is shown in Fig. 4.

## **4. Control algorithm**

In this paper, a hybrid modal space control method is proposed, and the process of this method is derived by following equations. Here, the governing equation of motion for a structure with piezoelectric actuators is written as

$$
M\ddot{x} + Kx = u \tag{14}
$$

*Mx Mx Mx* is the set of model space of this method is propertional apple to the process of this method is propertion in this paper, a hybrid modal space control method is pro-<br>
sed, and the process of this method is where *u* represents the load due to actuation, *M* is the mass matrix,  $K$  is the stiffness matrix, and  $x$  contains node displacements. In order to realize decoupling, the displacement of nodes in physical space will be transformed to modal space according to expansion theorem *x x* **And the plotted in the matter of model serieve and the discuss of units and the conservant of the extention of metal of the conservant of the conservant of the extent of node to actuation,** *M* **is the given by even by eve** Let a represents the load due to actuation, *M* is the given by<br>
ss matrix, *K* is the stiffness matrix, and *x* contains<br>
declinged decoupling, the dis-<br>
declined is  $\ddot{q}_i + \sum_{i=1}^{N} h_u \dot{q}_i + \sum_{i=1}^{N} (g_u + \dot{q}_i)$ <br>

$$
x = \phi q \tag{15}
$$

where  $\phi$  is the vector of mode shape, and *q* is the corresponding modal coordinate. Substituting Eq. (15) into Eq. (14), it can be gotten by *i*  $M\phi\ddot{q} + k\dot{\phi}q = u$ .<br> *i*  $\ddot{q} + \alpha y^2 q = \phi_i^T u = F_i$ <br> *i i*  $\ddot{q} + k\dot{\phi}q = u$ .<br> **ii**  $\ddot{q} + k\dot{\phi}q = u$ .<br> **ii**  $i\theta = \begin{cases} 15 & \text{if } i \neq s \\ 1, & \text{if } i = s \end{cases}$ <br> **ii**  $i\theta = 0$ .<br> **ii**  $i\theta = 0$  iii  $i\theta = 0$  iii  $i\theta = 0$  a

$$
M\phi\ddot{q} + k\phi q = u.\tag{16}
$$

In this paper, normalization of mass is employed and Eq. (16) can be written in the following form as

$$
\ddot{q} + \Omega q = \phi^T u. \tag{17}
$$

Then, for any order natural frequency, the governing equation is given as follows

$$
\ddot{q}_i + \omega_i^2 q_i = \phi_i^T u = F_i \tag{18}
$$

where  $\omega_i$  is the natural frequency and  $F_i$  is the actuating force.

#### *4.1 Independent modal space control*

According to the active control principle, the control force can be written as

$$
F_i = -g_i q_i - h_i \dot{q}_i \qquad i = 1, 2, \cdots, N \tag{19}
$$

where  $g_i$  is displacement gain and  $h_i$  is speed gain. Then the closed loop formed in modal space is expressed by

$$
\ddot{q}_i + h_i \dot{q}_i + \left(g_i + \omega_i^2\right) q_i = 0. \tag{20}
$$

Then, for any order natural frequency, the governing equa-<br>  $\ddot{q}_i + \omega_i q_i = \phi_i^{\dagger} u = F_i$  (18)<br>
where  $A = \begin{bmatrix} 0 & I \\ -\Omega & 0 \end{bmatrix}, B = \begin{bmatrix} 0 \\ \phi^{\dagger} \end{bmatrix}$ .<br>
ere  $\omega_i$  is the natural frequency and  $F_i$  is the actuating<br>
In m From Eq. (20), it can be seen that the governing equation of each order is in a separate state. It is easy to design the control law, and the real control signal can be written as *i* formed in modal space is expressed by<br>  $g_i + \omega_i^2 q_i = 0.$  (20)<br>
(20)<br>
(20)<br>
(30), it can be seen that the governing equation of<br>
in a separate state. It is easy to design the control<br>
(al control signal can be written a *i* **h** *i alh, + B,u,*  $(i = 1, 2, \dots, n)$ <br>According to the active control principle, the control force<br>A<sub>i</sub> =  $\begin{bmatrix} 0 & I \\ -\omega_i^2 & 0 \end{bmatrix}$ ,  $B_i = \begin{bmatrix} 0 \\ \theta_i^T \end{bmatrix}$ .<br>be written as<br> $F_i = -g_i q_i - h_i \dot{q}_i$   $i = 1, 2, \dots, N$ <br>(19) *h*,  $= A_i h_i + B_i u_i$  ( $i = 1$ <br>
coording to the active control principle, the control force<br>
where  $A_i = \begin{bmatrix} 0 & I \\ -\omega_i^2 & 0 \end{bmatrix}$ ,  $I$ <br>
where  $A_i = \begin{bmatrix} 0 & I \\ -\omega_i^2 & 0 \end{bmatrix}$ ,  $I$ <br>
where  $A_i = \begin{bmatrix} 0 & I \\ -\omega_i^2 & 0 \end{bmatrix}$ ,  $I$ <br>
w

$$
u = -\sum_{i=1}^{N} M \phi_i (g_i q_i + h_i \dot{q}_i)
$$
 (21)

## *4.2 Dependent modal space control*

For dependent modal space control, let  $F_i$  be

1 Technology 32 (8) (2018) 3515–3521

\n2 Department modal space control

\nFor dependent modal space control, let 
$$
F_i
$$
 be

\n
$$
F_i = -\sum_{s=1}^{N} (g_{is}q_s + h_{is}\dot{q}_s) \quad i = 1, 2, \cdots, N
$$
\n22)

\n18 The closed loop form considering the external control is

\n29

\n20

and the closed loop form considering the external control is given by

For dependent modal space control, let 
$$
F_i
$$
 be  
\n
$$
F_i = -\sum_{s=1}^{N} (g_{is}q_s + h_{is}\dot{q}_s) \quad i = 1, 2, \dots, N
$$
\n(22)  
\nd the closed loop form considering the external control is  
\nven by  
\n
$$
\ddot{q}_i + \sum_{s=1}^{N} h_{is}\dot{q}_s + \sum_{s=1}^{N} (g_{is} + \omega_i^2 \delta_{is})q_s = 0
$$
\n(23)  
\nhere  $h_s$  and  $g_{is}$  are the speed gain and the displacement

*Technology 32 (8) (2018) 3515-3521*<br> **Dependent modal space control**<br>
For dependent modal space control, let  $F_i$  be<br>  $F_i = -\sum_{s=1}^{N} (g_n q_s + h_n \dot{q}_s)$   $i = 1, 2, \dots, N$  (22)<br>
4 the closed loop form considering the external c Technology 32 (8) (2018) 3515-3521<br> **Dependent modal space control**<br>
For dependent modal space control, let  $F_i$  be<br>  $F_i = -\sum_{i=1}^{\infty} (g_{i0}q_i + h_{i0}q_i)$   $i = 1, 2, \dots, N$  (22)<br>
d the closed loop form considering the external where  $h_{is}$  and  $g_{is}$  are the speed gain and the displacement gain in modal space, respectively. In addition, here  $\delta_{is}$  can be given by  $\delta_{is} = \begin{cases} 0, & i \neq s \\ 1, & i = s \end{cases}$ . The real control signal is 8) (2018) 3515-3521<br> **nodal space control**<br>
modal space control, let  $F_i$  be<br>  $x + h_n \dot{q}_x$   $i = 1, 2, \dots, N$  (22)<br>
oop form considering the external control is<br>  $\sum_{s=1}^{N} (g_{is} + \omega_i^3 \delta_{is}) q_s = 0$  (23)<br>  $g_{is}$  are the speed ga *i* = 1, 2, ..., *N* (22)<br> *i* = 1, 2, ..., *N* (22)<br> **i** considering the external control is<br>  $\omega_i^2 \delta_{is} g_g = 0$  (23)<br>
the speed gain and the displacement<br>
spectively. In addition, here  $\delta_{is}$  can<br>  $\vec{i} = s$  . The real co shown in the form as  $g_{is}$  are the speed gain and the displacement<br>pace, respectively. In addition, here  $\delta_{is}$  can<br> $\delta_{is} =\begin{cases} 0, & i \neq s \\ 1, & i=s \end{cases}$ . The real control signal is<br>m as<br> $\int_{1}^{t} M \phi_{i} F_{i} = -\sum_{i=1}^{N} \sum_{s=1}^{N} M \phi_{i} F_{i} (g_{is}$  $\sum_{s=1}^{\infty} (g_{is} + \omega_i^2 \delta_{is}) q_s = 0$  (23)<br>  $g_{is}$  are the speed gain and the displacement<br>
space, respectively. In addition, here  $\delta_{is}$  can<br>  $\delta_{is} = \begin{cases} 0, & i \neq s \\ 1, & i = s \end{cases}$ . The real control signal is<br>
rm as<br>  $\sum_{i=1$ are the speed gain and the displacement<br> *i*, respectively. In addition, here  $\delta_s$  can<br>
<sup>0</sup>,  $i \neq s$ <br>
<sup>1</sup>,  $i = s$  The real control signal is<br>  $i_f F_i = -\sum_{i=1}^N \sum_{s=1}^N M \phi_i F_i (g_s q_s + h_s \dot{q}_s)$ . (24)<br> **ace control**  $F_i = -\sum_{s=1}^{N} (g_s q_s + h_s \dot{q}_s)$   $i = 1, 2, \dots, N$  (22)<br> **d** the closed loop form considering the external control is<br>
cen by<br>  $\ddot{q}_i + \sum_{s=1}^{N} h_s \dot{q}_s + \sum_{s=1}^{N} (g_{si} + \omega_i^2 \delta_{si}) q_s = 0$  (23)<br>
ere  $h_s$  and  $g_a$  are the speed behaviology 32 (8) (2018) 3515-3521<br> **Dependent modal space control**<br>
dependent modal space control, let F<sub>c</sub> be<br>  $= -\sum_{r=1}^{\infty} (g_{\alpha}q_{x} + h_{\alpha}q_{x}) \quad i = 1, 2, \dots, N$  (22)<br>
the closed loop form considering the external cont given by<br>  $\ddot{q}_i + \sum_{k=1}^{N} h_{ik} \dot{q}_k + \sum_{k=1}^{N} (g_{ik} + \omega_i^2 \delta_{ik}) q_i = 0$  (23)<br>
where  $h_{ik}$  and  $g_{ik}$  are the speed gain and the displacement<br>
gain in modal space, respectively. In addition, here  $\delta_{ik}$  can<br>
be given by For  $h_n$  and  $g_n$  are the speed gain and the displacement<br> *n* in modal space, respectively. In addition, here  $\delta_n$  can<br>
given by  $\delta_n = \begin{cases} 0, & i \neq s \\ 1, & i = s \end{cases}$ . The real control signal is<br>
win in the form as<br>  $u = M\phi F$ foodal space, respectively. In addition, here  $\delta_n$  can<br>by  $\delta_n =\begin{cases} 0, & i \neq s \\ 1, & i=s \end{cases}$ . The real control signal is<br>the form as<br> $\delta F = \sum_{i=1}^N M \phi_i F_i = -\sum_{i=1}^N \sum_{i=1}^N M \phi_i F_i (g_n q_i + h_n \dot{q}_i)$ . (24)<br>d modal space contro expectively. In addition, here  $\delta_a$  can<br>  $i \neq s$ . The real control signal is<br>  $i = s$ . The real control signal is<br>  $i = -\sum_{i=1}^{N} \sum_{s=1}^{N} M \phi_i F_i (g_a q_s + h_a \dot{q}_s)$ . (24)<br> **control**<br>
2 hybrid modal space control, let<br>
2 hybri

$$
u = M\phi F = \sum_{i=1}^{N} M\phi_{i} F_{i} = -\sum_{i=1}^{N} \sum_{s=1}^{N} M\phi_{i} F_{i} \left(g_{s} q_{s} + h_{s} \dot{q}_{s}\right).
$$
 (24)

## *4.3 Hybrid modal space control*

(17)  $h = [q \t q]$ , for any order natural frequency, the governing equal<br>  $\ddot{q} + \Omega q = \phi^r u$ .<br>
(17)  $h = [q \t q]$ ,  $\Gamma$ ,  $\Gamma$ , (17) can be written in the space<br>
Then, for any order natural frequency, the governing equa-<br>  $\dot{n}$ In order to describe hybrid modal space control, let space *h h*  $H = M\phi F = \sum_{i=1}^{N} M\phi_i F_i = -\sum_{i=1}^{N} \sum_{i=1}^{N} M\phi_i F_i (g_s g_s + h_s \dot{q}_s)$ . (24)<br> *Hybrid modal space control*<br>
in order to describe hybrid modal space control, let<br>  $\begin{bmatrix} 0 & \dot{q} \end{bmatrix}$ , Eq. (17) can be written in th  $F = \sum_{i=1}^{n} M\phi_{i}F_{i} = -\sum_{i=1}^{n} \sum_{i=1}^{n} M\phi_{i}F_{i}(g_{n}g_{i} + h_{n}g_{i}).$ (24)<br> **I modal space control**<br>
f to describe hybrid modal space control, let<br>  $\int_{0}^{T}$ , Eq. (17) can be written in the form of the state<br>  $+ Bu$ (2  $\sum_{r=1}^{n} \sum_{s=1}^{n} M\phi_{r} F_{i}(g_{s}g_{s} + h_{s}g_{s}).$  (24)<br> **ontrol**<br>
hybrid modal space control, let<br>
be written in the form of the state<br>
(25)<br>  $\begin{bmatrix} 0 \\ \phi^T \end{bmatrix}.$ <br>
ing to Eq. (25), we can get that<br>  $\begin{bmatrix} \cdot \cdot \cdot \cdot \cdot n \\ \phi^{$ 

$$
h = Ah + Bu \tag{25}
$$

where  $A = \begin{bmatrix} 0 & I \\ -\Omega & 0 \end{bmatrix}$ ,  $B = \begin{bmatrix} 0 \\ \phi^T \end{bmatrix}$ .  $A = \begin{bmatrix} 0 & I \\ -\Omega & 0 \end{bmatrix}, B = \begin{bmatrix} 0 \\ \phi^T \end{bmatrix}.$ 

In modal space, according to Eq. (25), we can get that

$$
\dot{h}_i = A_i h_i + B_i u_i \quad (i = 1, 2, \cdots, n)
$$
\n(26)

where  $A_i = \begin{bmatrix} 0 & I \\ -\omega_i^2 & 0 \end{bmatrix}, B_i = \begin{bmatrix} 0 \\ \phi_i^T \end{bmatrix}.$ 

For dependent modal space control, the linear quadratic in dex can be written as

$$
i = Ah + Bu
$$
\n(25)\n  
\nhere  $A = \begin{bmatrix} 0 & I \\ -\Omega & 0 \end{bmatrix}, B = \begin{bmatrix} 0 \\ \phi^T \end{bmatrix}.$ \n  
\nIn modal space, according to Eq. (25), we can get that\n
$$
\vec{h}_i = A_i \vec{h}_i + B_i \vec{u}_i \quad (i = 1, 2, \dots, n)
$$
\n(26)\n  
\nhere  $A_i = \begin{bmatrix} 0 & I \\ -\omega_i^2 & 0 \end{bmatrix}, B_i = \begin{bmatrix} 0 \\ \phi_i^T \end{bmatrix}.$ \nFor dependent modal space control, the linear quadratic in-  
\nx can be written as\n
$$
J = \frac{1}{2} \int_0^\infty (h^T Q h + u^T R u) dt
$$
\n(27)\n  
\nhere *Q* is a diagonal matrix and *R* is a symmetric posi-  
\ne definite matrix.\nSimilarly to independent modal space control, the linear  
\nandratio index can be written as

where  $Q$  is a diagonal matrix and  $R$  is a symmetric positive definite matrix.

Similarly to independent modal space control, the linear quadratic index can be written as

$$
\dot{n}_i = A_i h_i + B_i u_i \quad (i = 1, 2, \dots, n)
$$
\n(26)  
\nhere  $A_i = \begin{bmatrix} 0 & I \\ -\omega_i^2 & 0 \end{bmatrix}, B_i = \begin{bmatrix} 0 \\ \phi_i^T \end{bmatrix}.$   
\nFor dependent modal space control, the linear quadratic in-  
\nex can be written as  
\n
$$
J = \frac{1}{2} \int_0^\infty (h^T Q h + u^T R u) dt
$$
\n(27)  
\nhere *Q* is a diagonal matrix and *R* is a symmetric posi-  
\nve definite matrix.  
\nSimilarly to independent modal space control, the linear  
\nandratio index can be written as  
\n
$$
J_i = \frac{1}{2} \int_0^\infty (h_i^T Q_i h_i + u_i^T r_i u_i) dt \quad (i = 1, 2, \dots, n)
$$
\n(28)

where  $Q_i$  is the second order diagonal matrix and  $r_i$  is a constant.

In accordance with optimal control theory, the optimal control force is expressed as

$$
\begin{cases}\nF_c = -R^{-1}B^T G(t)h \\
F_{ci} = -r_i^{-1}B_i^T G_i(t)h_i\n\end{cases}
$$
\n(29)

as shown in Eq. (30).

$$
\begin{cases}\n- A^T G(t) - G(t) A + G(t) B R^{-1} B^T G(t) - Q = 0 \\
- A_i^T G_i(t) - G_i(t) A + G_i(t) B_i r_i^{-1} B_i^T G_i(t) - Q_i = 0\n\end{cases}
$$
\n(30)

in addition, Eq. (30) is easy to be solved by MatLab.

In the process of the closed loop control, a modal tracking strategy is taken to determine control-switching, which is described as follows. According to the principle of mode super position, a definition with respect to modal coordinate is given by  $\left\{\begin{aligned}\nF_{ci} &= -r_i^{-1}B_i^T G_i(t)h_i \\
\text{there } G_i(t) &= \text{and } G_i(t) \text{ are obtained from Riccati equations} \\
\text{shown in Eq. (30).} \n\end{aligned}\right.$ Freedback<br>
shown in Eq. (30).<br>  $\left\{\begin{aligned}\n- A^T G_i(t) - G_i(t) A + G_i(t) B_i r_i^{-1} B_i^T G_i(t) - Q = 0 \\
-A_i^T G_i(t) - G_i(t) A + G_i(t) B_i r_i^{-1} B_i^T G_i(t) - Q = 0\n\end{aligned}\right.$ <br>
add  $\int -A^T G(t) - G(t)A + G(t)BR^{-1}B^T G(t) - Q = 0$ <br>  $\left[-A_i^T G_i(t) - G_i(t)A + G_i(t)Br_i^{-1}B_i^T G_i(t) - Q_i = 0\right]$ <br>
in addition, Eq. (30) is easy to be solved by MatLab.<br>
In the process of the closed loop control, a model tracking<br>
strategy is taken to

$$
W_h = \sum_{i=1}^n \|h_i\| \tag{31}
$$

where  $W_h$  represents the aggregate weighting of the controlled modes. Moreover, one may define

$$
\Gamma_i = \frac{\left\|h_i\right\|}{W_h} \tag{32} \qquad \qquad 5.1 \text{ Ex}
$$

where  $\Gamma_i$  denotes the contribution of *i*-th mode to the aggregate weighting. Obviously, if the controlled mode needs to be determined, one condition should be satisfied by

$$
\Gamma_j = \max(\Gamma_1, \Gamma_2, \cdots, \Gamma_n)
$$
\n(33)

where  $\Gamma_j$  indicates the chosen mode that is used to determine feedback gain. Fig. 5 shows the program flow chart of hybrid modal space control.

Here, the control program shown in Fig. 5 is called modal control switching. In control process, the actuator output will produce a appropriate signal according to the modal tracking given by  $Z_1 = 1200, Z_2 = 0.045, Z_3 = 0.204$ . Besides, it is also strategy.

#### **5. Experiments and results**

With the actuator model completed and control algorithm considering SAI designed, it is of importance to experimentally evaluate the proposed method for active vibration control. In this section, an experimental setup is described, and ex perimental results are presented and discussed.



Fig. 5. The program flow chart of hybrid modal space control.



Fig. 6. Diagrammatic sketch of the experimental system.

## *5.1 Experimental setup*

In this experiment, the actuator is positioned near the end of the rectangular plate and is mounted on a support. The sensor is placed on the opposite side of the actuator. Fig. 6 shows a diagrammatic sketch of the experimental system.

The system shown in Fig. 6 is realized by adopting three means. Firstly, the relative parameters of actuator are acquired by some trials, the real actuator is seen in Fig. 7. Secondly, the rectangular plate with one side clamped is constructed by using virtual prototyping technology (VPT). Thirdly, the hybrid modal space control algorithm is implemented using Matlab/Simulink.

From Eq. (13), note that SAI needs three parameters  $Z_1$ ,  $Z_2$  and  $Z_3$ . Through addressing the trial results, these are Fig. 6. Diagrammatic sketch of the experimental system.<br>
5.*I Experimental setup*<br>
In this experiment, the actuator is positioned near the end of<br>
the rectangular plate and is mounted on a support. The sensor<br>
is placed noted that Eq. (13) can be divided into two models according to the computation requirements for a control system. The first model is the full expression of SAI given in Eq. (13). The second model is an ideal linear actuator, which neglects SAI. in this experiment, the actuator is positioneal near the end of<br>the rectangular plate and is mounted on a support. The sensor<br>is placed on the opposite side of the actuator. Fig. 6 shows a<br>diagrammatic sketch of the exper This model,  $F = Z<sub>3</sub>U$ , is usually used. In subsequent simulations, the vibration suppression results will be given considering these two models respectively.

In addition, some relative parameters of the plate structure are listed in Table 2.

Parameter	Value
Length (mm)	100
Width (mm)	50
Thickness (mm)	
Density ( $kg/m3$ )	2800
Young's ratio (GPa)	70
Damping ratio	0.0002

Table 2. Parameters of the plate structure.

Table 3. Natural frequencies of the plate structure.

Number of order		-		
Natural frequency (Hz)	24.2	38.2	61.8 81.3	96.1



Fig. 7. Giant material actuator.

## *5.2 Experimental results and discussion*

For the aim of active vibration control, only the first five modes are considered, and their natural frequencies are given in Table 3.

In this section, several experimental results are presented for illustrating the influence of interaction models on the control system. For exciting the mode shapes of the plate structure, a concentrated load is applied at a central location on the free end of the plate, which is a sine sweep excitation signal with an amplitude of 100 N and a frequency of 0~200 Hz. In the first experimental results, the first two modes are controlled and the remaining three modes are regarded as residual modes, as shown in Fig. 8.

In the second set of experimental results, the first three modes are controlled and the remaining two modes are con sidered as residual modes, which is shown in Fig. 9.

As can be seen in these two figures, the vibrations are suppressed by using an additional actuator. In Figs. 8 and 9, it can be seen that the plate with SAI considered has more "stiffer" than the one without SAI considered. When the first two modes are controlled, the vibrations are suppressed by using the control method without SAI. But when the first three modes are controlled, for the response of the third order frequency, the control method without SAI does not act. The



Fig. 8. Frequency response of the plate.



Fig. 9. Frequency response of the plate.

proposed approach, meanwhile, effectively suppresses vibration suppressing effectiveness whether in either the first ex perimental results or in the second experimental results.

## **6. Conclusions**

This work presented in this paper indicates that the proposed active vibration control with SAI (structure-actuator interaction) considered can be utilized successfully to suppress vibrations of a flexible plate structure. In this investigation, the SAI model is constructed according to magnetomechanical coupling, based on which the controller is designed. In order to improve the control effect, a modal control switching strategy is employed. According to the proposed methodology, an experimental setup is established to verify the feasibility of the method. The experimental results show that the stiffness effect of the plate structure could be increased and the amplitudes of the plate vibration are reduced more effectively. Furthermore, the present methodology can be extended to other flexible structures.

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