

# Characteristics analysis of aero-engine whole vibration response with rolling bearing radial clearance<sup>†</sup>

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(Manuscript Received July 7, 2016; Revised November 28, 2016; Accepted January 7, 2017)

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#### **Abstract**

For the influence of rolling bearing radial clearance on the whole vibration in the aero-engine whole system, a real engine rotor bearing-casing whole model is established. The rotor and casing systems are modeled by means of FEM; the support systems are modeled by lumped-mass model; rolling bearing radial clearance and strong-nonlinearity of Hertz contact force at four different supports are considered. The coupled system response is obtained by the numerical integral method. The characteristics of the whole vibration re sponse are analyzed. For rolling bearing at a typical support, the rotor, outer ring of rolling bearing and casing response characteristics at different rotating speeds are analyzed. The changing law of contact forces for each ball and the global contact forces at different speeds are analyzed. The influence of the radical clearance on the contact forces on the whole vibration is analyzed. The results show that the contact forces will be larger and the acceleration amplitude jumps obviously when the radial clearance is increased, and due to the variable stiffness of the rolling bearing, the natural frequency will appear when the stiffness changes fiercely, that is *frequency-locked* phe nomenon. When the radical clearance is larger and the rotating speed is between two critical speeds, the rotor squeezes the outer ring now and then. Reducing the radical clearance can reduce the whole vibration and increas

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*Keywords*: Radial clearance; The whole vibration; The finite model; Hertz contact force; Unbalanced force

# **1. Introduction**

Rolling bearing is an important part in aero-engine. Due to the strong nonlinear characteristics in rolling bearing, it has great influence on the whole vibration response. Scholars at home and abroad have made many studies on the clearance and contact force between the ball and the race. Ehrich [1, 2] studied the dynamics of rotors with bearing clearance. Sinou [3] studied the non-linear dynamic response of a flexible rotor supported by ball bearings by using the Harmonic balance method. Harsha et al. [4, 5] analyzed non-linearity causing the transition from contact to no contact state due to the bearing clearance. Tiwari et al. [6, 7] analyzed the non-linear rotor dynamics numerically and experimentally. Jang et al. [8] studied the vibration due to ball bearing waviness in a rotating system. Chen et al. [9] established a new rotor-support-stator coupling system dynamic model for practical aero-engine considering the nonlinear Hertzian contact force between balls and races, and the varying compliance vibration with rubbing coupling faults. The numerical method was applied to gain responses, and bifurcation and chaotic motions were analyzed. Bai et al. [10] established a general dynamic model with surface waviness. Wu et al. [11] studied the non-synchronous vibration of a Jeffcott rotor due to internal radial clearance in roller bearing.

As can be seen from the current literatures, the existing rolling bearing models are complete, however, considering char acteristics analysis of the aero-engine whole vibration response is less. The existing models are based on the simple rotor model; the rolling bearing models are not applied in the real engine model. The ball bearing clearance is considered in the existing results, which leads to nonlinear response of the rotor system, however, the contact force for each ball, the global contact forces on the rolling bearing, and the rotor operation laws are not considered.

In this paper, a real aero-engine whole vibration model is established. The rolling bearing radial clearance is considered in this model. Based on the casing acceleration signal and the rotor displacement signal, the real engine vibration response characteristics are analyzed. The forces of the rolling element and the rotor operation laws under different speeds are analyzed. The influence of the radial clearance on the contact forces on whole vibration is studied.

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<sup>†</sup>Recommended by Associate Editor Junhong Park

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Fig. 1. Rotor-support-casing coupling model of a type of aero-engine (unit: *mm*).



Fig. 2. The finite element rotor model.

# **2. A dynamic model for certain type aero-engine**

# *2.1 The structure sketch map for certain real engine*

A certain real engine model can be referred in Refs. [15, 16]. Its structure is described in detail. Its sketch map is shown in Fig. 1. The symbols in Fig. 1 are described as follows:  $P_1-P_5$ denote different disks; the  $C_1$  denotes the casing;  $G_1-G_3$  denote different gear couplings; *S*1–*S*4 denote different support point bearings;  $I_1-I_2$  denote different front installation nodes;  $k_g$  is the mesh stiffness of a gear pump;  $k_{\text{fl}}-k_{\text{fl}}$  denote different support stiffness;  $k_c$  is the connection stiffness.

The rotor-support-casing coupling dynamic model for a real engine is established. Its concrete method for modeling is as follows.

#### *2.2 Dynamics model*

# *2.2.1 Rotor model and casing model*

Beam element for rotor model and casing model is established by FEM. The mass, inertias, and gyroscopic moments of the disks are considered in rotor model; however, these parameters are not considered in casing model. The casing model is considered beam element model without considering rotating effect. The finite element rotor dynamic model is shown in Fig. 2.

The symbols in Fig. 2 are described in Ref. [15]. A coordinate system is shown in Fig. 2, and the O-XYZ is a fixed coordinate system, the four vibration directions of the ro-[15].



Fig. 3. Rotor-casing support.



Fig. 4. Rolling bearing model.

$$
(\mathbf{M}_s)\ddot{\mathbf{q}}_s + (\mathbf{C}_s - \omega \mathbf{G}_s)\dot{\mathbf{q}}_s + \mathbf{K}_s \mathbf{q}_s = \mathbf{Q}_s \tag{1}
$$

is the external force vector;  $M<sub>s</sub>$  is the mass matrix;  $G_s$  is gyroscopic matrix;  $K_s$  is the stiffness matrix; **C***<sup>s</sup>* is the damping matrix.

In this paper,  $C_s = \alpha_0 \mathbf{M}_s + \alpha_1 \mathbf{K}_s$ , of which,  $\alpha_0$  and  $\alpha_1$ are supposed to be constants.

is [15]. The coupling connection between rotors, and Elastic support between casing and base modeling can be referred in Ref.

#### *2.2.2 Discrete support model*

supports are considered in exist pare considered in Ref. [15]. A particle single considered in the consistent of the considered in the roto-support-staing coupling dynamic model for a real mass  $m_8$  is the bearing botwen The connection modeling between rotor and casing can be referred in Ref. [15]. Suppose that,  $m_{wi}$  is the outer bearing mass;  $m_{bi}$  is the bearing housing mass;  $k_{ti}$  is the support stiffness;  $c_t$  is the damping coefficient;  $k_f$ ,  $c_f$  are respectively the support stiffness and damping coefficient. As shown in Fig. 3,  $F_{yRi}$  and  $F_{xRi}$  are the force of rotor acting on the support RC<sub>*i*</sub>, and the  $F_{yCi}$  and  $F_{xCi}$  are the force of casing acting on support RC*i*. In this section, the modeling process of rolling bearings is described in detail. and **y** and **y c** is the conding power from an easing can be referred in Ref. [15]. Suppose that,  $m_{wi}$  is the surport stiff-<br>ness;  $c_{u}$  is the damping coefficient;  $k_{fb}$ ,  $c_{jl}$  are respectively the support stiffne

The rolling bearing model is shown in Fig. 4, when the ball bearing works, with the contact position periodically varying, the total stiffness and compliance of the bearing will periodically vary, and the Varying compliance (VC) of the bearing is a parametric excitation of a rotor-balling bearing coupling system, finally, a so-called VC vibration is generated. VC vibration is an inherent vibration. it always exists even if the

bearing is new and does not have any faults. Assume the displacements of the *<sup>m</sup>*th node of rotor are *xRm* et al. [13], the ball force can be expressed as:



Fig. 5. Solving flow for rotor-support-casing coupling dynamics.

Newtonant- $\beta$	method	Newtonant- $\beta$		
Response of casings	Response of cross and casings	Response of rators		
g. 5. Solving flow for rotor-support-casing coupling dynamics.	Parameters	Disk $P_1$	Disk $P_1$	Disk $P_2$
$\int_{\text{Pak}} = \sum_{j=1}^{N} C_b(x \cos \theta_j + y \sin \theta_j - r_0)^{3/2} H(x \cos \theta_j + y \sin \theta_j - r_0) \sin \theta_j$	Cross-polar inertia inertia $J_{dd}(kg m^2)$	0.015	0.	
$\int_{\text{Psi}} = \sum_{j=1}^{N} C_b(x \cos \theta_j + y \sin \theta_j - r_0)^{3/2} H(x \cos \theta_j + y \sin \theta_j - r_0) \sin \theta_j$	Table 3. Parameters of fan rotor.			

In the formula,  $C<sub>b</sub>$  is the Hertzian contact stiffness and  $F_{\text{sub}} = \sum_{j=1}^{N} C_{\text{b}}(x \cos \theta_{j} + y \sin \theta_{j} - r_{0})^{3/2} H(x \cos \theta_{j} + y \sin \theta_{j} - r_{0}) \cos \theta_{j}$ <br>  $F_{\text{sub}} = \sum_{j=1}^{N} C_{\text{b}}(x \cos \theta_{j} + y \sin \theta_{j} - r_{0})^{3/2} H(x \cos \theta_{j} + y \sin \theta_{j} - r_{0}) \sin \theta_{j}$ .<br>
Table 3. Parameters of fan rotor<br>
In the formu  $\frac{2\pi}{N_{\rm b}}(j-1)$ ,  $j = 1, 2, \cdots N_{\rm b}$ , where,  $\pi$  , and the contract of  $\overline{\phantom{a}}$  $N<sub>b</sub>$  is the number of balls.  $\omega_{\text{Case}}$  is the cage's speed. Suppose the outer race radius is  $R$ , the inner ring radius is Fig. 5. Solving flow for rotor-support-easing coupling dynamics.<br>
Fig. 5. Solving flow for rotor-support-easing coupling dynamics.<br>  $\begin{bmatrix}\nF_{\text{vis}} = \sum_{j=1}^{N} C_s(x \cos \theta_j + y \sin \theta_j - r_0)^{3/2} H(x \cos \theta_j + y \sin \theta_j - r_0) \cos \theta_j \\
F_{\text{ys}} = \sum_{j=$  $\omega_{\text{Case}} = \frac{\omega \times r}{R+r}$ , and the  $\omega$  is shaft rotating angular velocity. (•) is the rieaviside tunction.  $\theta_j$  is the *j* fin the article tunction,  $\theta_j$  is the *j* c  $\theta_{\text{cusp}} \times t + \frac{2\pi}{N_b} (j-1)$ ,  $j = 1, 2, \dots, N_b$ , where,  $\frac{4}{5}$  is the number of balls.  $\omega_{\text{cusp}}$  is the cage's speed. Supp •) is the Fieaviside function.  $\theta_j$  is the *f* th balls angle po-<br>
on, that is  $\theta_j = \omega_{\text{Cusp}} \times t + \frac{2\pi}{N_b} (j-1), j = 1, 2, \cdots N_b$ , where,<br>
is the number of balls.  $\omega_{\text{Cusp}}$  is the cage's speed. Suppose<br>  $\theta_{\text{Cusp}} = \frac{\omega N}{f$ 

$$
\omega_{\rm vc} = \omega_{\rm Cage} \times N_{\rm b} = \omega_{\rm Rotor} \times B_{\rm N}
$$
\n(3)

*dxi* í = && <sup>L</sup> && (4)

$$
F_{dxi} = c_{ti} (\dot{x}_{wi} - \dot{x}_{bi}), \quad F_{dyi} = c_{ti} (\dot{y}_{wi} - \dot{y}_{bi}). \tag{5}
$$

# *2.3 Solution of rotor-support-casing model*

The solution can be referred in Ref. [15], which is shown in Fig. 5. The support looseness is considered between bearing housing and casing in Ref. [15], however, the rolling bearing radial clearance is considered at different supports in this pa per. The merits of combination method have: The enormous matrix need not to be formed, therefore, the computation efficiency is improved greatly.

Table 1. Unit number of rotors and casings.

Fan rotor	Compressor rotor	Turbine rotor	Casing



	ran folof	Compressor rotor		I urbine rotor		Casing
Implicit integration Explicit integration (Zhai	11	10		11		24
(Newmark- $\beta$ ) method) Responses of Responses of rotors and casings	Table 2. Parameters of rotors.					
rotors	Parameters	Disk $P_1$	Disk $P_2$	Disk $P_3$	Disk $P_4$	Disk $P_5$
	Mass $m_n$ (kg)	3.88	1.41	5.17	10.28	10.28
ow for rotor-support-casing coupling dynamics.	Cross-polar inertia $J_{\rm pd}$ (kg·m <sup>2</sup> )	0.03	0.003	0.03	0.05	0.05
$x\cos\theta_i + y\sin\theta_i - r_0^{3/2}H(x\cos\theta_i + y\sin\theta_i - r_0)\cos\theta_i$	Cross-equator inertia $J_{dd}$ (kg·m <sup>2</sup> )	0.015	0.0015	0.015	0.025	0.025
$\cos\theta_i + y\sin\theta_i - r_0^{3/2}H(x\cos\theta_i + y\sin\theta_i - r_0)\sin\theta_i$ .						
	Table 3. Parameters of fan rotor.					
(2)	Number	Coordinate/mm		Outer diameter/mm Inner diameter/mm		
alla, $Cb$ is the Hertzian contact stiffness and		$\mathbf{0}$		30		$\mathbf{0}$
eaviside function. $\theta_i$ is the <i>j</i> th ball's angle po-	$\overline{2}$	30		48		$\mathbf{0}$
	3	71		30		$\mathbf{0}$
$\theta_j = \omega_{\text{Cage}} \times t + \frac{2\pi}{N} (j-1), \ j = 1, 2, \cdots N_{\text{b}}$ , where,	4	90		30		$\theta$
er of balls. $\omega_{\text{Case}}$ is the cage's speed. Suppose	5	114		36		$\mathbf{0}$
	6	162.6		24		$\mathbf{0}$
e radius is $R$ , the inner ring radius is						

Table 4. Parameters of compressor rotor.

$H(\bullet)$ is the Heaviside function. $\theta_i$ is the <i>fth</i> ball s angle po-	3	71	30	$\mathbf{0}$
sition, that is $\theta_j = \omega_{\text{Cage}} \times t + \frac{2\pi}{N_{\text{b}}} (j-1)$ , $j = 1, 2, \dots N_{\text{b}}$ , where,	$\overline{4}$	90	30	$\mathbf{0}$
	5	114	36	$\boldsymbol{0}$
$N_{\rm b}$ is the number of balls. $\omega_{\rm Caee}$ is the cage's speed. Suppose	6	162.6	24	$\mathbf{0}$
the outer race radius is $R$ , the inner ring radius is	$\overline{7}$	230	24	$\mathbf{0}$
$r, \omega_{\text{cage}} = \frac{\omega \times r}{R + r}$ , and the $\omega$ is shaft rotating angular velocity. Because the inner race is fixed to the shaft, $\omega_{\text{inner}} = \omega_{\text{rotor}}$ . If		Table 4. Parameters of compressor rotor.		
$Nb$ is the number of balls, then VC frequency can be given by	Number	Coordinate/mm	Outer diameter/mm Inner diameter/mm	
	1	230	37.6	24
$\omega_{\rm vc} = \omega_{\rm Case} \times N_{\rm b} = \omega_{\rm Rotor} \times B_{\rm N}$ (3)	$\overline{c}$	248.5	37.6	24
	3	286.5	37.6	24
where $B_N = N_b \times r/(R+r)$ . Obviously, the number $B_N$ depends on	$\overline{4}$	307.5	37.6	24
the bearing size. In this paper, $B_{N1} = 5.5$ , $B_{N2} = 4.85^{\circ}$ .	5	329.5	37.6	24
The bearing outer race's differential equation is	6	374.5	37.6	24
	$\tau$	408.5	37.6	24
$\begin{cases} m_{\rm wi} \ddot{x}_{\rm wi} + k_{\rm ti} (x_{\rm wi} - x_{\rm bi}) + F_{\rm div} = F_{\rm xRi} \\ m_{\rm wi} \ddot{y}_{\rm wi} + k_{\rm ti} (y_{\rm wi} - y_{\rm bi}) + F_{\rm div} = F_{\rm vRi} - m_{\rm wi} g \end{cases} \quad i = 1, 2, \cdots N$ (4)	8	427.5	37.6	24
		Table 5. Parameters of turbine rotor.		
where $F_{dx}$ and $F_{dy}$ are damping forces, suppose the	Number	Coordinate/mm		Outer diameter/mm Inner diameter/mm
damping is viscous, then,	1	427.5	30	18
	2	463.5	49	36.9
$F_{dxi} = c_{ii}(\dot{x}_{wi} - \dot{x}_{bi}), \quad F_{dyi} = c_{ii}(\dot{y}_{wi} - \dot{y}_{bi}).$ (5)	3	593.5	49	36.9
	4	625.5	49	36.9
	5	670.5	49	36.9
2.3 Solution of rotor-support-casing model				

Table 5. Parameters of turbine rotor.



# **3. The whole vibration simulation with the radial clearance**

#### *3.1 Dynamic parameters*

The parameters of the rotor, the casing and the connection parameters of system are shown in Tables 1-13. Mass of out  $\min_{w} m_{w}$  is 0.2 kg, and Bearing carrier mass  $m_{b}$  is 10 kg.

Table 6. Parameters of rotors.



Table 7. Parameters of casing.



Table 8. Parameters of casings.

Elastic modulus $E$ (Pa)	Density $\rho$ (kg/m <sup>3</sup> )	Poisson's ratio $\mu$	Proportion $\alpha_0$	Proportion damping ratio damping ratio $\alpha_1$
$2.07\times10^{11}$	$7.8 \times 10^3$	0.3		$1.35 \times 10^{-5}$

Table 9. Parameters of ball bearing.

Rolling bearing	Outer race- Inner race- way radius R/mm	way radius r/mm	Ball number $N_{\rm h}$	Contact s tiffness $C_{\rm b} / (N / m^{3/2})$	Bearing clearance $r_0$ /um
$S_1$	39.5	29	13	$12.4\times10^{9}$	
$S_2$	39.5	29	13	$12.4\times10^{9}$	
S <sub>3</sub>	32	17	14	$11.9\times10^{9}$	
$\mathrm{S}_4$	32		14	$11.9\times10^{9}$	

Table 10. Support parameters of rotor-casing.

Sup- ports	Node of Casing rotor	(node)	$k_t$ (N/m)	$c_t$ (N.s/m)	$k_f(N/m)$	$c_f$ (N.s/m)	
RC <sub>1</sub>		2	$1\times10^8$	2000	$1\times10^8$	1000	
RC <sub>2</sub>		9	$1\times10^8$	2000	$1\times10^8$	1000	
RC <sub>3</sub>	11	16	$1\times10^8$	2000	$1\times10^8$	1000	
RC <sub>4</sub>	8	22	$1\times10^8$	2000	$1\times10^8$	1000	

Table 11. Collection parameters of rotor-casing.

Collection Node of Casing $k_{gx}$ (N/m) $c_{gx}$ (N/m) $k_{gy}$ (N/m) $c_{gy}$ (N/s/m)				
RK1		$1\times10^8$	$1\times10^8$	

Table 12. Collection parameters of casing-base.







# *3.2 Calculation condition*

(1) The radial clearances at the supports  $S_1$ ,  $S_2$ ,  $S_3$  and  $S_4$  are considered.

(2) The output is the casing acceleration response and the vibration displacement response at the supports  $S_1$ ,  $S_2$ ,  $S_3$  and *S*<sup>4</sup>

.(3) The rotating speed range is 5000-40000 rpm.

# *3.3 The critical speed analysis under different radical clear ances*

Figs. 6(a)-(d) show the amplitude-speed curves of the casing lateral acceleration at the ninth node at different supports  $S_1$ ,  $S_2$ ,  $S_3$  and  $S_4$  under different radical clearances that is 0  $\mu$ m, 10  $\mu$ m, 30  $\mu$ m and 50  $\mu$ m. As can be seen from the figures, the first three order critical speeds are 16100 rpm, 20600 rpm and 33800 rpm, respectively. Figs. 6(b)-(d) show the clear ances has little influence on the casing response. In Fig. 6(a), At the support  $S_1$ , when the clearance is large, the amplitude  $r_0$ /um jumping phenomenon is obvious; when the rotating speed is close to the second order critical speed, the unbalance force plays a major role and the nonlinearity is less, so the amplitude jumping phenomenon is less obvious; when the rotating speed is near the third order critical speed, the nonlinearity plays a major role and the unbalance force is less, so the am plitude jumping phenomenon is obvious. Fig. 7 is the enlargement of Fig. 6(a). Fig. 8 is the result of Ref. [17]. Figs. 7 and 8 show the obvious amplitude jumping phenomenon, which is caused by the strength of nonlinearity of rolling bearing.

### *3.4 Calculating modal analysis*

The first node of fan rotor is applied the lateral transient force 100 N, and the frequency response of the casing lateral acceleration is obtained, which is used to simulate hammering method, as shown in Fig. 9.

# *3.5 Casing-rotor response characteristics at different rotating speeds*

Figs. 10(a)-(e) show the cascade plot of the support  $S_1$ ,  $S_2$ , *S*<sup>3</sup> and *S*4 lateral acceleration response at different rotating speeds with the clearances 30  $\mu$ m, 30  $\mu$ m, 30  $\mu$ m and 30  $\mu$ m, respectively. In these figures, the frequency changing in the cascade plot of the support  $S_1$  is not obvious; the second order natural frequency  $f_{n2}$ , the sixth order natural frequency  $f_{n6}$  and the combined frequency of rotational frequency and the sec-



(a) The casing lateral acceleration at the support *S*<sup>1</sup>



(b) The casing lateral acceleration at the support  $S_2$ 





(d) The casing lateral acceleration at the support *S*<sup>4</sup>

Fig. 6. Amplitude-speed curve of casing acceleration response under different clearances at different supports.

ond order natural frequency  $f_{n2}$  appear in the cascade plot of the support  $S_2$ ; the same phenomenon is in the cascade plot of the support  $S_3$ ; the sixth order natural frequency  $f_{\text{no}}$  appears in the cascade plot of the support  $S_4$ . The phenomenon is called v *whirl phenomenon* or *frequency-locked phenomenon*. The whip has a constant frequency not depending on the rotating speed [17]. Fig. 11 is from Ref. [17].

Figs. 12(a)-(d) show the bifurcation diagram of the rotor



Fig. 7. The enlargement of Fig. 6(a).



Fig. 8. Result of Ref. [17].



Fig. 9. Frequency response at the ninth node of the casing.

displacements at different supports. In these figures, when the rotating speed is near the critical speed, the bifurcation dia gram at different supports shows periodic motion; when the rotating speed is far away from the critical speed, different levels of quasi-periodic and chaotic motion appear; when the rotating speed is higher than the second order critical speed at the supports  $S_1$  and  $S_2$ , the phenomenon of quasi-periodic and chaotic motion is obvious.

# *3.6 Orbits of rotor-out ring and characteristics of casing response at typical rotating speeds*

Rotor and casing response characteristics at the support  $S_2$ with the radial clearance  $30 \mu m$  are analyzed at different speeds, which reveal the rotor operation law and casing response characteristics.

Figs. 13(a)-(d) show the orbits of the rotor and out ring, the bifurcation diagram of the rotor, the wave and frequency





(a) Cascade plot of the casing acceleration response at different speeds without regard to the radical clearance



Frequency (Hz)

(b) Cascade plot of the casing acceleration response at different speeds considering the radical clearance of the  $S_1$ 



(c) Cascade plot of the casing acceleration response at different speeds considering the radical clearance of the  $S_2$ 



(d) Cascade plot of the casing acceleration response at different speeds considering the radical clearance of the  $S_3$ 



(e) Cascade plot of the casing acceleration response at different speeds considering the radical clearance of the S<sup>4</sup>

Fig. 10. Cascade plot of casing acceleration response with radial clearance at different supports.

Fig. 11. The result of Ref. [11].



(a) Bifurcation diagram at the support  $S<sub>1</sub>$ 



(b) Bifurcation diagram at the support  $S_2$ 



(c) Bifurcation diagram at the support S<sup>3</sup>



(d) Bifurcation diagram at the support S<sup>4</sup>

Fig. 12. Bifurcation diagram of the rotor displacements at different supports.





Fig. 13. Orbits of rotor-outer ring, bifurcation diagram of the rotor and wave and spectrum of casing acceleration when the rotating speed is 2000 rpm.

spectrum of casing acceleration at the compressor rotor support *S*2 when the rotating speed is 2000 rpm. In Fig. 13(a), due to the effect of many supports and the larger radial clearance, pears in high frequency. the compressor rotor is whirling at the bottom of the rolling bearing, so the out ring is not squeezed. In Fig. 13(b), the rotor shows weak chaotic motion; In Figs. 13(c) and (d), the rotational frequency, the combined frequency of frequency multi-

Fig. 14. Orbit of rotor-outer ring, bifurcation diagram of the rotor and wave and spectrum of casing acceleration when the rotating speed is 5000 rpm.

plication and the second order natural frequency  $f_{n2}$  appear in frequency spectrum; the sixth order natural frequency  $f_{\text{no}}$  ap-

Figs.  $14(a)-(d)$  show the orbits of the rotor and out ring, the bifurcation diagram of the rotor, the wave and frequency spectrum of casing acceleration at the compressor rotor support  $S_2$  when the rotating speed is 5000 rpm. In Fig. 14(a), the



Fig. 15. Orbits of rotor-outer ring, and wave and frequency spectrum of casing acceleration when the rotating speed is 8000 rpm.

compressor rotor is squeezing the out ring. In Fig. 14(b), the rotor shows period-doubling motion; in Figs. 14(c) and (d), the rotational frequency, frequency multiplication and the sixth order natural frequency  $f_{\text{no}}$  appear in frequency spectrum.

Figs. 15(a)-(d) show the orbits of the rotor and out ring, the wave and frequency spectrum of casing acceleration at the



Fig. 16. Orbits of rotor-outer ring, and wave and frequency spectrum of casing acceleration when the rotating speed is 12000 rpm.

compressor rotor support  $S_2$  when the rotating speed is 8000 rpm. In Figs. 15(a) and (b), the compressor rotor is squeezing the out ring and whirling along the race, and because the contact force is large, the vibration of the out ring is big, which appear like wavy shape. In Figs. 15(c) and (d), the rotational frequency, the combined frequency of frequency multiplication and the sixth order natural frequency  $f_{\text{no}}$  appear in frequency spectrum; the fifth order natural frequency  $f_{\text{no}}$  appears in high frequency.<br>Figs.  $16(a)-(c)$  show the orbits of the rotor and out ring, the

wave and frequency spectrum of casing acceleration at the compressor rotor support  $S_2$  when the rotating speed is 12000 rpm. In Fig. 16(a), the compressor rotor is squeezing the out ring and whirling along the race, and because the contact force is larger, the vibration of the out ring is bigger. In Figs. 16(b) and (c), the rotational frequency and the VC1 frequency which is  $B_{N1}$  times of the rotating frequency, appear in frequency spectrum.



Fig. 17. Orbits of rotor-outer ring, and wave and frequency spectrum of casing acceleration when the rotating speed is 16100 rpm.



Fig. 18. Orbits of rotor and outer ring when the rotating speed is 20600 rpm.

Figs. 17(a)-(c) show the orbits of the rotor and out ring, the wave and frequency spectrum of casing acceleration at the compressor rotor support  $S_2$  when the rotating speed is 16100 rpm, that is, the first order critical speed. In Fig. 17(a), the compressor rotor is squeezing the out ring and whirling along



Fig. 19. Orbits of rotor-outer ring, bifurcation diagram of the rotor and wave and spectrum of casing acceleration when the rotating speed is 22500 rpm.

the race, and because the contact force is larger, the vibration of the out ring is bigger, which appears like ellipsoid. In Figs. 17(b) and (c), when the rotating speed is about the first order critical speed, because the unbalanced force is bigger, the rotational frequency appears in frequency spectrum.

 $0.0<sub>0</sub>$ 



Fig. 20. Orbits of rotor-outer ring, bifurcation diagram of the rotor and wave and spectrum of casing acceleration when speed is 30000 rpm.

Fig. 18 shows the orbits of the rotor and out ring at the compressor rotor support  $S_2$  when the rotating speed is 20600 rpm, that is, the second order critical speed. In Fig. 18, because the unbalanced force continues to increase, the vibration of out ring is further increase.

Figs. 19(a)-(d) show the orbits of the rotor and out ring, the bifurcation diagram of the rotor, the wave and frequency spectrum of casing acceleration at the compressor rotor support  $S_2$ when the rotating speed is 22500 rpm. In Fig. 19(a), the orbits



Fig. 21. Orbits of rotor-outer ring, and wave and spectrum of casing acceleration when the rotating speed is 35000 rpm.

of the rotor and the out ring appear overlap, that is, the rotor squeezes or doesn't squeezes the out ring now and then. In Fig. 19(b), the bifurcation diagram of the rotor appears like a cloud, that is, chaotic motion. In Figs. 19(c) and (d), because the rotational frequency is near the corresponding frequency of the second order critical speed, the continuous spectrum appears, and the wave of the casing acceleration appears beat vibration phenomenon. Due to the instability of the rotor, the stiffness variation fiercely, the second order natural frequency *f*n2 is excited.

Figs. 20(a)-(d) show the orbits of the rotor and out ring, the bifurcation diagram of the rotor, the wave and frequency spectrum of casing acceleration at the compressor rotor support *S*2 when the rotating speed is 30000 rpm. In Fig. 20(a), after the rotating speed passed the second order critical speed, due to the larger speed, the contact force increases, the vibration of the out ring is bigger. In Fig. 20(b), the bifurcation diagram of the rotor appears closed loop, that is, chaotic mo-



Fig. 22. The changing law of the contact forces for each ball at differ ent rotating speeds.

tion. In Figs. 20(c) and (d), the discrete spectrum appears in the casing acceleration frequency spectrum, that is, quasi periodic phenomenon.

Figs. 21(a)-(c) show the orbits of the rotor and out ring, wave and frequency spectrum of the casing acceleration at the compressor rotor support  $S_2$  when the rotating speed is 35000 rpm. In Fig. 21(a), due to the larger rotating speed, the contact force is larger, and the vibration of the out ring is bigger. In Figs. 21(b) and (c), the rotational frequency appears in frequency spectrum.

# *3.7 The nonlinear contact forces for each ball and the evolution of contact force for the rolling bearing*

Figs. 22(a)-(f) show the evolution of contact forces for each ball when the radial clearance is  $30\mu$ m at the support  $S_2$ , the contact time is one period of revolution of the rolling bearing cage and the ball number is 13. The evolution of contact forces for each ball changing with time and ball number can be seen in these figures. Figs.  $22(a)-(d)$  show the rotor squeeze two balls, three balls, four balls and five balls every moment, respectively. Figs. 22(e) and (f) show the rotor squeezes the out ring now and then. It is found that the time of the rotor squeezing each ball is same in one period of revolution of the rolling bearing cage before the first order critical speed, the number is increasing with the increase of the rotating speed and the movement of the rolling bearing is periodic in Figs.  $22(a)$ -(d), which belongs to the condition that the rotor always squeezes the out ring. Between the two critical speeds, the rotor squeezes the out ring now and then, the time of the rotor squeezing the out ring is different in one period of revolution of the rolling bearing cage, and the movement of the rolling bearing is quasi periodic and chaos in Figs. 22(e) and (f), that is instability phenomenon.

Figs. 23(a)-(f) show the evolution of the contact forces for the rolling bearing  $S_2$ . The evolution of the contact forces for  $\alpha$ the rolling bearing  $S_2$  can be seen in these figures. Figs. 23(a)-(d) show the contact forces are increasing gradually and the



Fig. 23. The changing law of the contact forces for the rolling-bearing at different speeds.



radial clearances.



Fig. 25. The changing law of the global contact forces for the rolling bearing in different radial clearances.

contact forces for the rolling bearing  $S_2$  are periodicity. Figs. 23(e) and (f) show the condition the rotor squeezes the out ring now and then and the contact forces for the rolling bearing  $S<sub>2</sub>$  are quasi periodicity and chaos.

# *3.8 The influence of the radial clearance on the contact forces*

In order to study the influence of the radial clearance on the contact forces, the evolution of contact forces for each ball and the global restoring forces for the rolling bearing  $S_2$  are shown in Figs. 24 and 25 when the clearances are  $0 \mu m$ , 30  $\mu m$  and 50  $\mu$ m, respectively. Figs. 24(a)-(c) show the rotor squeeze five balls, three balls, and two balls every moment, respectively. The number is decreasing with the increase of the ra dial clearances. Figs. 25(a)-(c) show the fluctuation of the global contact forces is high by the increase of the radial clearances.

#### **4. Conclusion**

In this paper, certain type real engine whole vibration model with the radial clearance of the rolling bearing is established, the strong nonlinearity of the radial clearance and the Hertz contact forces are considered in this model. The response of coupled system is obtained by numerical integration method. Some results are obtained as follows:

(1) The changing law of the casing acceleration amplitude with the rotating speeds is analyzed at different radial clear ances. It is found that the jumping phenomenon is caused by the strength of nonlinearity of rolling bearing with the radial clearance. It shows that it can reduce the whole vibration and improve the rotor's stability by reducing the radial clearance.

(2) The characteristics of the casing acceleration response are analyzed with the radial clearance. In high rotating speed, duo to the variable stiffness of the rolling bearing, when the stiffness changes obviously, the rotor squeezes the out ring now and then, so the natural frequencies of the system are excited, that is *frequency-locked phenomenon*.

(3) The evolution of contact forces for each ball and the global contact forces at different rotating speeds are analyzed, and the rotor and the out ring load characteristics are analyzed. When the out ring is period loaded, the whole vibration is stabile. When the out ring is loaded now and then, the whole vibration is unstable.

(4) The influence of the radial clearance on the contact forces is analyzed. The contact forces will larger when the radial clearance is increased.

## **Acknowledgment**

This work is supported by Funding of Jiangsu Innovation Program for Graduate Education KYLX\_0295.

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