

Trajectory planning for overhead crane by trolley acceleration shaping†

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Abstract

This paper proposes a novel off-line trolley trajectory planning method for underactuated overhead cranes. The proposed technique is feasible and efficient for overhead crane operation. Dynamic coupling between trolley motion and payload swing was successfully exploited using a staircase form of trolley acceleration. The payload swings in the constant velocity phase were efficiently suppressed and the trolley reached the desired position using this technique. The reasonable number of stairs can be determined by evaluating the residual oscillation amplitude according to the number of stairs and variation in the natural frequency of the pendulum. The proposed approach was first simulated from the kinematics viewpoint to verify the validity of the trolley trajectory and the swing angle of the payload. The proposed approach was then combined with the dynamics of the overall crane, wherein the robust sliding mode controller was applied to ensure that the trolley tracks the designed trajectory. The numerical simulation results demonstrated superior performance and robustness against parameter uncertainties of the proposed method. The proposed method exhibited potential for application in the control of underactuated systems, such as overhead cranes, single-link flexible-joint manipulators, and flexible Cartesian manipulators.

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Keywords: Underactuated system; Overhead cranes; Anti-swing trajectory planning; Acceleration shaping

1. Introduction

Overhead cranes are widely used in various fields, such as heavy industries, seaports, automotive factories, and construction facilities. The productivity and efficiency of an overhead crane depends on payload weight and velocity, as well as on the capability of the crane to quickly reduce the swing angle of the payload after each operation. Theory and practice have shown that high acceleration and deceleration correspond to large swing angles. This condition leads to hazardous situation and may cause serious accidents when the cargo swing angle becomes extremely large. A large cargo swing angle can break the crane, damage other equipment and infrastructure, or harm nearby people.

The process of operating a crane can be divided into three phases, namely, payload lifting, horizontal transportation, and payload lowering. The second phase presents the most difficult challenge. For this phase, the trolley and the payload need to quickly reach the desired position. Meanwhile, the swing angle of the payload must remain small and must return to zero at the end of each operation. Two main approaches may be applied to satisfy this requirement: the design of an antiswing controller and the design of a reasonable trajectory for the trolley, otherwise known as motion planning. In the second approach, the desired trajectory normally comprises three phases, namely, acceleration, maintaining constant velocity, and deceleration. The time and shape of acceleration in the first and third phases are chosen such that the swing angle of the payload reaches the maximum, and then decreases to zero.

The first approach has drawn interest from various fields of study. A number of control algorithms have been developed for overhead cranes. The simplest approach is designing a linear controller based on the linearized model around its target location. This technique assumes a small payload swing during an operation. The crane dynamics can initially be a linearized model around its target location; subsequently, linear control approaches can be designed according to this model. This technique can be combined with optimal control and extended by gain scheduling [1] and input shaping [2-4].

To improve performance, various nonlinear and advanced control strategies have been tested on the crane control problem. These strategies include partial feedback linearization [5- 9], sliding mode control [10-14], and adaptive control [5, 15- 18]. A combination of the control methods has also been considered, such as adaptive sliding mode control [19-23] and adaptive fuzzy sliding mode control [24].

A design method based on the energy and passivity of the system has recently been investigated. This approach has been successfully applied in the control of underactuated systems

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such as overhead cranes [11, 15, 16, 25]. A controller that includes the passivity of the payload swing by using the payload as an end effector of a manipulator has been developed. H. Aschemann [26] applied energy shaping by interconnection and damping assignment according to a linearized model to increase the damping factor related to the underactuated coordinate.

Most controllers are designed for position regulation. These controllers may be difficult to implement when the trolley must travel over a large distance because of the limitation of the actuating force. Another problem is that these controllers require a sensor to measure trolley motion and another to measure the swing angle for feedback. This requirement increases the cost of the crane and may be difficult to implement.

In the second approach, the swing of the payload is suppressed by designing a reasonable trajectory for the trolley. This technique has also been evaluated. N. Sun et al. [15] applied iterative learning to determine the reasonable trajectory. However, the planned trajectory was obtained numerically. This iterative process needs to be repeatedly performed for each crane operation. H.-H. Lee [27] and N. Sun et al. [16] applied rectangular acceleration profile modifications to eliminate residual oscillations of a payload in the constant or zero velocity phase. Lee [28] used the acceleration in the sine form for the trolley in accelerating and decelerating phases. However, these methods require a precise natural frequency of the payload, which is dependent on the cable length. Thus, this technique is dependent on system parameters. N. Sun et al. [15] used an analytical function to generate an S-shaped trajectory for the trolley and obtain a smooth function for velocity and acceleration. However, with this function, the trolley can only reach its desired location when time approaches infinity.

This paper presents a novel trajectory planning method for overhead cranes, wherein the acceleration in a staircase form is used to eliminate the residual oscillation of the payload in the constant velocity phase and when the trolley reaches its desired location. The constraints on the acceleration amplitudes are given by a set of linear equations that can be easily evaluated. Numerical simulation results have demonstrated the efficiency of the proposed trajectory planning method. In where $x(t)$ denotes the trolley displacement, $\theta(t)$ denotes addition, the planned trajectory is tracked by a simple robust controller that requires only the trolley motion for feedback. These findings suggest that the proposed scheme is efficient and feasible for crane control.

In summary, this paper has the following merits:

(1) The payload swing is theoretically proven to be zero during the constant velocity phase, and no residual swing exists when the trolley reaches the desired location.

against uncertainties in the natural frequency of the payload, $\theta(t)$, which is given by depending on the cable length.

(3) The residual vibration of the payload that depends on the number of stairs is investigated by simulation; this number of stairs should be three or five.

Fig. 1. 2D overhead crane model.

(4) An algorithm for trajectory planning is also presented; this is convenient for practical implementation because of its simple structure.

The remainder of this paper is organized as follows. Section 2 presents the problem formulation. Section 3 presents the method of acceleration shaping, wherein the trolley acceleration in a staircase form is applied to suppress the residual oscillation of the payload. Numerical experiments are demonstrated in Sec. 4. Finally, Section 5 concludes the paper.

2. Problem formulation

Trajectory planning for a crane during the horizontal transportation phase is presented in this paper. The rope has a constant length, and the system has two degrees of freedom. To obtain the dynamic model of the system, the following assumptions are established: (1) the payload is considered a point mass; (2) the mass and stiffness of the hoisting rope are disregarded; and (3) the effects of wind disturbances are disregarded. Based on the Lagrangian formulation [29], the dynamic model of a two-dimensional overhead crane system is expressed as follows: 11 encontrane of this polarity cosons. Second
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Trajectory planning for a crane during the horizontal trans-

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stant le

$$
(mt + mp)\ddot{x} + mpl\cos\theta \ddot{\theta} - mpl\dot{\theta}^2\sin\theta + f_{11}\dot{x} = u
$$
 (1)

$$
m_{p}l\cos\theta \ddot{x} + m_{p}l^{2}\ddot{\theta} + m_{p}gl\sin\theta = 0
$$
 (2)

the payload swing angle (Fig. 1), and u is the force acting on the trolley. In this equation of motion, m_t and m_p represent the trolley mass and the payload mass, respectively; *l* is the length of the rope; *g* is the gravitational acceleration; and f_{11} is the damping coefficient on the trolley. point mass; (2) the mass and surmoss or the noising rope are
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regarded. Based on the Lagrangian formulation [29], the dy-
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ergarded. Based on the Lagrangian formulation [29], the dy-
namic model of a two-dimensional overhead cr *mc* model of a two-dimensional overhead crane system is
 *im*_{*n*} *l <i>l m*_{*n*} *l c*s *g* $\ddot{\theta}$ *- m_n l* $\dot{\theta}$ ² sin θ + f_1 , $\dot{x} = u$ (1)
 m_n l cos θ $\ddot{x} + m_p$ $\frac{2}{r}$ $\ddot{\theta} + m_p$ g

(2) The trajectory planned by acceleration shaping is robust tween trolley acceleration $\ddot{x}(t)$ and the payload swing angle Crane dynamics consists of the actuated part described by Eq. (1) and the unactuated part described by Eq. (2). The second equation captures the kinematic coupling behavior be-

$$
l\ddot{\theta} + g\sin\theta = -\cos\theta \ddot{x}.
$$
 (3)

In practice, the small swing angle of overhead cranes is usu-

$$
\ddot{\theta} + \omega_n^2 \theta = -(\omega_n^2 / g) \ddot{x} \tag{4}
$$

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ally maintained. Thus, the approximations for $\sin \theta \approx \theta$ and
 $\cos \theta \approx 1$ can be applied. Hence, Eq. (3) may be rewritten as
 $\ddot{\theta} + \$ modeled by the pendulum. Eq. (4) clearly shows that trolley acceleration directly influences payload swing. Therefore, a proper trajectory planning method for the trolley significantly affects the reduction or elimination of the payload swing.

Based on the practical operation of overhead cranes, some of the following requirements must be carefully considered through trajectory planning for the design specification of the trolley: the following requirements must be carefully considered x_a

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2. Acceleration and

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stablely reaches the target

P1. The trolley reaches the target position within a finite time t_f ; i.e.,

$$
x(t) = p_d, \forall t \ge t_f. \tag{5}
$$

P2. During movement, the trolley velocity and acceleration must be less than the maximum values

$$
|\dot{x}_d(t)| \le v_{\text{max}}, \qquad |\ddot{x}_d(t)| \le a_{\text{max}} \tag{6}
$$

where v_{max} , a_{max} are the permitted maximum velocity and acceleration, respectively.

P3. The maximum payload swing must be kept within an acceptable domain,

$$
|\theta(t)| \leq \theta_{ub} \tag{7}
$$

P4. No swing should occur during the constant velocity phase, and when the trolley stops at the destination

$$
\theta(t) = 0, \text{ when } \ddot{x}(t) = 0 \text{ and } \theta(t) = 0, \forall t \ge t_f. \tag{8}
$$

signed such that the following conditions are satisfied:

$$
x(0) = 0, \quad \dot{x}(0) = 0, \quad \theta(0) = 0, \quad \dot{\theta}(0) = 0
$$

\n
$$
x(t_f) = p_d, \dot{x}(t_f) = 0, \theta(t_f) = 0, \dot{\theta}(t_f) = 0.
$$
\n(9)

3. Trajectory planning by trolley acceleration shaping

This section presents the design of a trolley trajectory that satisfies Eq. (9). This trajectory planning method is based on $|\theta(t)| \leq \theta_{\text{ns}}$. (1)

Period, then $\theta(\tau_x) = 0$ and $\dot{\theta}(\tau_x) = 0$. Then twist velocity phase is switched at this time

P4. No swing should occur during the constant velocity mined by

phase, and when the trolley stops $|θ(t)| \le \theta_{ab}$.
 PA. No swing should occur during the constant velocity that velocity these is swinched at this time location, the swing
 PA. No swing should occur during the constant velocity and the symbol of the s $|\theta(t)| \leq \theta_{\omega}$. (7)
 **then the swing should occur during the constant velocity masse is switched at this time left

P4.** No swing should occur during the constant velocity minded by
 then angle is maintained at zero *t o t*) is *o*_{*a*}. (7) sature the subset is mathering the constant velocity phase is maintained at zero. The maximum velocity is deter-

phase, and when the trolley stops at the destination
 $\theta(t) = 0$, when $\bar{x}(t$ P4. No swing should occur during the constant velocity mined by

phase, and when the trolley stops at the destination
 $θ(t) = 0$, when $\bar{x}(t) = 0$ and $θ(t) = 0, \bar{y}(t) \ge 0$.

Therefore, the trajectory $x(t)$ of the trolley $t \in [0, \tau_a]$ and decelerating phase phase, and when the trolley stops at the destination
 $\theta(t) = 0$, when $\ddot{x}(t) = 0$ and $\theta(t) = 0, \forall t \ge t_f$.

Therefore, the trajectory $x(t)$ of the trolley must be de-

Eq. (12) suggests that

signed such that the following netefore, the trajectory $x(t)$ of the trolley must be de-

Eq. (12) suggests that the interval of solutions are satisfied:
 $(0) = 0$, $\dot{x}(0) = 0$, $\dot{\theta}(0) = 0$, $\dot{\$ erefore, the trajectory $x(t)$ of the trolley must be de-

Eq. (12) suggests that the in

ed such that the following conditions are satisfied:
 $(0) = 0$, $\dot{x}(0) = 0$, $\dot{\theta}(0) = 0$, $\dot{\theta}(0) = 0$
 $(t_f) = p_a, \dot{x}(t_f) = 0, \dot{\theta}(t_f)$ Therefore, the trajectory $x(t)$ of the trolley must be de-

Eq. (12) suggests that
 $x(0) = 0$, $\dot{x}(0) = 0$, $\dot{\theta}(0) = 0$, $\dot{\theta}(0) = 0$, $\dot{\theta}(0) = 0$
 $x(t_f) = p_a, \dot{x}(t_f) = 0, \dot{\theta}(t_f) = 0$, $\dot{\theta}(0) = 0$

Trajectory planning by Therefore, the trajectory $x(t)$ of the trolley must be de-

Eq. (12) suggests that the

end such that the following conditions are satisfied:
 $x(0) = 0$, $\dot{x}(0) = 0$, $\dot{\theta}(0) = 0$, $\dot{\theta}(0) = 0$
 $x(t_f) = p_a, \dot{x}(t_f) = 0, \dot{\theta}(t$ Sore, the trajectory *x*(*t*) of the trolley must be de-

Eq. (12) suggests that the colluminate the following conditions are satisfied:
 p_{μ} , *x*(*t*) = 0, *θ*(0) = 0, *θ*(*t*_{*r*}) = 0, *θ*(*t_r*) = 0, *θ*(*t_r*

$$
\theta(\tau_a) = 0, \dot{\theta}(\tau_a) = 0
$$

$$
\theta(\tau_f) = 0, \dot{\theta}(\tau_f) = 0.
$$
 (10)

Fig. 2. Acceleration and velocity profile.

For simplicity, the accelerating phase is set to counteract the decelerating phase. Thus, motion planning is presented only in the accelerating phase.

3.1 Rectangular acceleration profile

$$
\theta(t) = \frac{a}{g} (\cos \omega_n t - 1) , \qquad \dot{\theta}(t) = -\omega_n \frac{a}{g} \sin \omega_n t . \tag{11}
$$

 $x(t) = p_x$, $\forall t \ge t_f$.

(2) decelerating phase. Thus, motion planning is presented on

the accelerating phase.

(2) During movement, the trolley velocity and acceleration

(3) decelerating phase.

(3) Eq...,

(3) \mathbb{E}_{\mathbf P2. During movement, the trolley velocity and acceleration **3.1 Rectangular acceleration profile**
 $|\dot{x}_s(t)| \le v_{\text{max}}$, $|\dot{x}_s(t)| \le a_{\text{max}}$
 $|\dot{x}_s(t)| \le a_{\text{max}}$
 $|\dot{x}_s(t)| \le a_{\text{max}}$
 $|\dot{x}_s(t)| \le c_{\text{max}}$
 $|\dot{x}_s(t)| \le c_{\text{max}}$
 $|\$ $\kappa_x(t) \leq v_{\text{max}}$, $|\ddot{x}_d(t) \leq a_{\text{max}}$ (6) Eq. (4) is obtained with $\ddot{x} = a = \text{const}$. Thus,

re v_{max} , a_{max} are the permitted maximum velocity and $\theta(t) = \frac{a}{g}(\cos \omega_a t - 1)$, $\dot{\theta}(t) = -\omega_i \frac{a}{g} \sin \omega_a t$.

eleratio $|\dot{x}_2(t)| \le v_{\text{max}}$, $|\ddot{x}_3(t)| \le a_{\text{max}}$ (6) $\frac{E_0}{t}$ (4) is obtained with $\ddot{x} = a = \text{const}$. Thus,

ere v_{max} , a_{max} are the permitted maximum velocity and $\theta(t) = \frac{a}{g}(\cos \omega_x t - 1)$, $\dot{\theta}(t) = -\omega_x \frac{a}{g} \sin \omega_x t$.
 x_xx(*v*) \Rightarrow x_{max} , a_{max} are the permitted maximum velocity and

elere v_{max} , a_{max} are the permitted maximum velocity and
 $\theta(t) = \frac{a}{g}(\cos \omega_x t - 1)$, $\dot{\theta}(t) = -\omega_x \frac{a}{g} \sin \omega_x$
 \therefore 73. The maximum payload swi $|\ddot{x}_d(t)| \le a_{\text{max}}$ (6)

Eq. (4) is obtained with $\ddot{x} = a = \text{const} \cdot T$

the permitted maximum velocity and
 $\theta(t) = \frac{a}{g}(\cos \omega_n t - 1), \quad \dot{\theta}(t) = -\omega_n \frac{a}{g} \sin \theta$

ply.

payload swing must be kept within an

From Eq. (11), if th ^q ^q = = = = ¹/₂*x*(*x*)₁/2 π_{max} , a_{max} are the permitted maximum velocity and

for, respectively.

for maximum payload swing must be kept within an

From Eq. (11), if the accelerating time is defined

the maximum payl From Eq. (11), if the accelerating time is determined by Fig. 2. Acceleration and velocity profile.

Fig. 2. Acceleration and velocity profile.

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decelerating phase. Thus, motion planning is presented only in

the $\tau_a = kT = 2k\pi/\omega_n$, $k = 1, 2, ...$, multiples of the pendulum period, then $\theta(\tau_a) = 0$ and $\dot{\theta}(\tau_a) = 0$. Thus, when the constant velocity phase is switched at this time location, the swing angle is maintained at zero. The maximum velocity is determined by For simplicity, the accelerating phase is set to counteract the relearting phase. Thus, motion planning is presented only in accelerating phase.
 Rectangular acceleration profile
 $\vec{E}q$. (4) is obtained with \ddot{x} implicity, the accelerating phase is set to counteract the
ating phase. Thus, motion planning is presented only in
elerating phase.

 $\frac{a}{g}$ (cos $\omega_n t - 1$), $\dot{\theta}(t) = -\omega_n \frac{a}{g} \sin \omega_n t$. (11)
 $\frac{a}{g}$ (cos $\omega_n t - 1$), 3.1 Rectangular acceleration profile

Eq. (4) is obtained with $\ddot{x} = a = \text{const}$. Thus,
 $\theta(t) = \frac{a}{g}(\cos \omega_s t - 1)$, $\dot{\theta}(t) = -\omega_a \frac{a}{g} \sin \omega_s t$. (11)

From Eq. (11), if the accelerating time is determined by
 $\tau_s = kT = 2k\pi / \$ $\dot{\theta}(t) = -\omega_n \frac{a}{g} \sin \omega_n t$. (11)

q. (11), if the accelerating time is determined by
 $2k\pi / \omega_n$, $k = 1, 2, ...,$ multiples of the pendulum

m $\theta(\tau_n) = 0$ and $\dot{\theta}(\tau_n) = 0$. Thus, when the con-

nity phase is switched at th $\theta(t) = \frac{a}{g}(\cos \omega_n t - 1)$, $\dot{\theta}(t) = -\omega_n \frac{a}{g}\sin \omega_n t$. (11)

From Eq. (11), if the accelerating time is determined by
 $= kT = 2k\pi/\omega_n$, $k = 1, 2, ...,$ multiples of the pendulum

riod, then $\theta(\tau_n) = 0$ and $\dot{\theta}(\tau_n) = 0$. Thus, orthogonourism behavior and $\sum_{n=0}^{\infty}$ or $\sum_{n=0}^{\infty}$ and $\sum_{n=0}^{\infty}$ and $\sum_{n=0}^{\infty}$ and the integer k must be selected to $\frac{2k\pi}{\omega_n}a$. (12)
ggests that the integer k must be selected to $\frac{2k\pi}{\omega_n}a$. $\epsilon \kappa I = 2k\pi / \omega_s$, $k = 1, 2,...$, muttiples of the pendulum

iod, then $\theta(\tau_s) = 0$ and $\dot{\theta}(\tau_s) = 0$. Thus, when the con-

it velocity phase is switched at this time location, the swing

gle is maintained at zero. The maxi $k = 1, 2, ...,$ multiples of the pendulum
 $t = 0$ and $\dot{\theta}(\tau_a) = 0$. Thus, when the con-

is switched at this time location, the swing

d at zero. The maximum velocity is deter-
 $\frac{1}{4}$.
 $\frac{1}{4}$.
 $\frac{1}{4}$.

(12)

ts $t = 2k\pi/\omega_n$, $k = 1, 2,...$, muttiples of the pendulum

t, then $\theta(\tau_a) = 0$ and $\dot{\theta}(\tau_a) = 0$. Thus, when the con-

relective phase is switched at this time location, the swing

is maintained at zero. The maximum velocity i x, and $v_{\rm v} = v_{\rm max} = v_{\rm iso} = v$. Thus, which reducitly phase is switched at this time location, the swing
is maintained at zero. The maximum velocity is deter-
by
 $= a\tau_a = \frac{2k\pi}{\omega_n} a$. (12)
(12) suggests that the integer

$$
\nu_{\text{max}} = a\tau_a = \frac{2k\pi}{\omega_n}a \,. \tag{12}
$$

Eq. (12) suggests that the integer k must be selected to rived from Eq. (11) as follows:

$$
\theta_{\max} = \theta(\tau/2) = -\frac{2a}{g}.
$$

$$
p_d = x(t_f) = v_{\text{max}} \tau_a + \tau_c v_{\text{max}} = v_{\text{max}} (\tau_a + \tau_c)
$$

= $a_{\text{max}} \tau_a (\tau_a + \tau_c).$ (13)

where τ_c is the time of the constant velocity phase. The acceleration and velocity profiles are shown in Fig. 2.

3.2 Acceleration profile of the stair form

Based on command shaping, which has been successfully applied in control vibration systems, we use a staircase form instead of a rectangular profile to eliminate residual oscillation. Assuming that we can eliminate oscillation in the system after $N+1$ time delay, then the acceleration of the trolley is de-2882 *N. Q. Hoang et al. / Journal of Mechanical Science and Technolog*

instead of a rectangular profile to eliminate residual oscillation. mining the a

Assuming that we can eliminate oscillation in the system after
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1 time delay, then the acceleration of the trolley is d *to a* and the solution of *t* set al. / Journal of Mechanical Science and Technology 28 (7)

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 $t + 1$ time delay, the *N. Q. Hoang et al. / Journal of Mechanical Science and Technology 28 (7) (2l*

of a rectangular profile to eliminate residual oscillation. mining the acceleration

time delay, then the acceleration of the trolley is de-

$$
\ddot{x}(t) = \sum_{i=0}^{N} a_i \sigma(t - \tau_i), \tau_0 = 0, \tau_1 < \tau_2 < \dots < \tau_N \tag{14}
$$

where a_i = const needs to be determined, τ_N is the accel-
constraint equations eration duration, and the step function is defined as follows:

$$
\sigma(t-\tau_i) = \begin{cases} 0, & t < \tau_i \\ 1, & \tau_i \le t \end{cases}
$$
 (15)

Substituting Eq. (14) into Eq. (4), the following equation can be obtained:

Here *a_i* = const needs to be determined, *τ_N* is the accel-
ation duration, and the step function is defined as follows:

$$
\sigma(t - \tau_i) = \begin{cases} 0, & t < \tau_i \\ 1, & \tau_i \le t \end{cases}
$$
These
Substituting Eq. (14) into Eq. (4), the following equation such that
the obtained:

$$
\theta(t) = \frac{1}{g} \sum_{i=0}^{N} a_i [\cos \omega_n (t - \tau_i) - 1] \sigma(t - \tau_i).
$$
 (16)
Thus
From Eq. (16), the residual vibration in the constant veloc-
is a given by

From Eq. (16), the residual vibration in the constant velocity is given by

where *a_i* = const needs to be determined, *r_{si}* is the acceleration direction, and the step function is defined as follows:\n
$$
\begin{aligned}\n\text{In this case, two unknowns need to be determined from four\ncontrol duration, and the step function is defined as follows:\n
$$
\begin{aligned}\na_0 + a_1 = 0, & a_0 + a_1 \cos \omega_1 z = 0 \\
a_1 + a_2 = 0, & a_1 + a_2 \cos \omega_1 z = 0\n\end{aligned}
$$
\n
$$
\begin{aligned}\n\text{Substituting Eq. (14) into Eq. (4), the following equation}\n\end{aligned}
$$
\nSubstituting Eq. (14) into Eq. (4), the following equation\n
$$
\begin{aligned}\n\text{In this case, two unknowns need to be determined from four\nconformal, 15.\n\end{aligned}
$$
\nSubstituting Eq. (14) into Eq. (4), the following equation\n
$$
\begin{aligned}\n\text{Substituting } \text{Eq. } (16) \text{ the residual vibration in the constant velocity}\n\\
\text{Substituting } \text{Eq. } (16) \text{ the residual vibration in the constant velocity}\n\\
\text{This is given by}\n\end{aligned}
$$
\n
$$
\begin{aligned}\n\text{Thus, when } r = 2k\pi / \omega_s, & k = 1, 2, \dots \\
\text{This is given by}\n\end{aligned}
$$
\n
$$
\begin{aligned}\n\text{Thus, the amplitude of the residual oscillation is derived} \\
\begin{aligned}\n\text{In this case, two unknowns need to be determined from four\nusing equations:\n
$$
\begin{aligned}\n\sin \alpha_x \tau = 0.8 \cos \alpha_x \tau = 1 \\
\tau = 2k\pi / \omega_s, & k = \frac{1}{2}, \dots\n\end{aligned}
$$
\n
$$
\begin{aligned}\n\text{Thus, the amplitude of the residual oscillation is derived} \\
\begin{aligned}\n\text{In this case, two unknowns need to be determined from four\nusing equations:\n
$$
\begin{aligned}\n\sin \alpha_x \tau = 0.8 \cos \alpha_x \tau = 1 \\
\tau = 2k\pi / \omega_s, & k = \frac{1}{2}, \dots\n\end{aligned}
$$
\n
$$
\begin{aligned}\n\text{In this case, two unknowns need to be determined from four\nconverfian, 15.\n\end{aligned}
$$
\n
$$
\begin{aligned}\n\text{Substituting } \text{L2} \
$$
$$
$$
$$

Thus, the amplitude of the residual oscillation is derived

$$
= \frac{1}{g} \begin{cases} (a_0 \cos \omega_x \tau_0 + a_1 \cos \omega_x \tau_1 + ... + a_N \cos \omega_x \tau_N) \sin \omega_x t \\ + (a_0 \sin \omega_x \tau_0 + a_1 \sin \omega_x \tau_1 + ... + a_N \sin \omega_x \tau_N) \sin \omega_x t \end{cases}
$$

\nThus, the amplitude of the residual oscillation is derived
\nfrom Eq. (17) by
\n
$$
V(\omega_n) = \sqrt{V_c^2 + V_s^2}
$$

\nwith $V_c = \sum_{i=0}^{N} a_i \cos \omega_x \tau_i$, $V_s = \sum_{i=0}^{N} a_i \sin \omega_x \tau_i$.
\nwith $V_c = \sum_{i=0}^{N} a_i \cos \omega_x \tau_i$, $V_s = \sum_{i=0}^{N} a_i \sin \omega_x \tau_i$.
\n
$$
= \frac{1}{2} \begin{cases} (17) \text{ shows that the following constraints must be satisfied} \\ \text{of the residual oscillation:} \end{cases}
$$

\n
$$
= \frac{1}{2} \begin{cases} (18) \text{ For } \text{Eqs. (19) and (21), the constraints are as} \\ a_0 + a_1 + a_2 = 0 \\ a_0 + a_1 + a_2 = 0 \end{cases}
$$

\n
$$
= \frac{1}{2} \begin{cases} (17) \text{ shows that the following constraints must be satisfied} \\ \text{of } \text{the polynomial oscillation:} \end{cases}
$$

\n
$$
= \frac{1}{2} \begin{cases} 19 \text{ and } (21) \text{ the constraints are as} \\ a_0 + a_1 + a_2 = 0 \\ a_0 \pi_1 + a_2 \sin \omega_x \tau_1 = 0 \end{cases}
$$

\n
$$
= \frac{1}{2} \begin{cases} 19 \text{ and } (19) \text{ and } (11) \text{ the constraints are} \\ a_0 + a_1 + a_2 = 0 \\ a_0 \tau_1 + a_2 \sin \omega_x \tau_2 = 0 \end{cases}
$$

\n
$$
= \frac{1}{2} \begin{cases} 19 \text{ and } (11) \text{ and } (12) \text{ and } (13) \text{ the constraints are} \\ a_0 + a_1 + a_2 = 0 \\ a_0 \tau
$$

Eq. (17) shows that the following constraints must be satisfied to eliminate the residual oscillation:

$$
\sum_{i=0}^{N} a_i = 0
$$
 (19) Two time delays

$$
\sum_{i=0}^{N} a_i = 0
$$
\n
$$
\sum_{i=0}^{N} a_i \cos \omega_n \tau_i = 0, \quad \sum_{i=0}^{N} a_i \sin \omega_n \tau_i = 0.
$$
\n(20)

\nIn addition, the trolley achieves its maximum velocity after
\ncelerating duration τ_N . Thus,

\n
$$
\sum_{i=0}^{N-1} a_i (\tau_N - \tau_i) = v_{\text{max}}.
$$
\n(21)

\nEqs. (19)-(21) form a set of four linear equations for deter-

In addition, the trolley achieves its maximum velocity after accelerating duration τ_N . Thus,

$$
\sum_{i=0}^{N-1} a_i (\tau_N - \tau_i) = v_{\text{max}}.
$$
 (21)

Eqs. (19)-(21) form a set of four linear equations for deter-

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and of a rectangular profile to eliminate residual oscillation. mining the acceleration stairs of the trolley. With

1 time dela *N. Q. Hoang et al. / Journal of Mechanical Science and Technology 28 (7) (2014) 2879-2888*

tead of a rectangular profile to eliminate residual oscillation. mining the acceleration stairs of the trolley. With the suming *N. Q. Hoang et al. / Journal of Mechanical Science and Technology 28 (7) (2014) 2879-2888*

tead of a rectangular profile to eliminate residual oscillation. mining the acceleration stairs of the trolley. With

suming tha (a) the second state of the cost of the troley is de-

in time delay, then the acceleration of the troley is de-

in this case. Is a second of the cost of the co *i* or a constrained points. Commute costain ossumed the residual oscillation is eliminate oscillation in the system after ention profile, the residual oscillation is eliminated by the second that $(2V)$ conditions. Vario Extend on the steading point of the trollowing that we can elliminate oscillation in the system after

suming that we can elliminate oscillation in the system after
 $x(t) = \sum_{r=0}^{N} a_r \sigma(t-r_1)$, $\tau_0 = 0$, $\tau_1 < \tau_2 < ... < \tau_N$ or a recent mannipulation is the schemar of the trollow in the system and so-control section of the transfer conduction is eliminated by

the acceleration of the trolley is de-

stars N are researched below.

Stars (19) mining the acceleration stairs of the trolley. With this acceleration profile, the residual oscillation is eliminated in the nonacceleration phases. Eqs. (19) and (20) are denoted as zerovibration (ZV) conditions. Various cases for the number of stairs *N* are presented below. mining the acceleration stairs of the trolley. With this acceleration profile, the residual oscillation is eliminated in the non-
acceleration phases. Eqs. (19) and (20) are denoted as zero-
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ion profile, the residual oscillation is eliminated in the non-

leration phases. Eqs. (19) and (20) are denoted as zer *Technology 28 (7) (2014) 2879-2888*

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on profile, the residual oscillation is eliminated in the non-

eration phases. Eqs. (19) and (20) are denoted as zero-

t *nd Technology 28 (7) (2014) 2879-2888*
 *ning the acceleration stairs of the trolley. With this accel-

tion profile, the residual oscillation is eliminated in the non-

leleration phases. Eqs. (19) and (20) are denoted ^a a v* ^w ^t ^w ^t ^t + = + = ship the acceleration stairs of the trolley. With this accel-
tion profile, the residual oscillation is eliminated in the non-
electration phases. Eqs. (19) and (20) are denoted as zero-
prainting (ZV) conditions. Various the acceleration stairs of the trolley. With this accel-
profile, the residual oscillation is eliminated in the non-
tion phases. Eqs. (19) and (20) are denoted as zero-
n (ZV) conditions. Various cases for the number of
 g the acceleration stairs of the trolley. With this accel-
 n profile, the residual oscillation is eliminated in the non-
 n acceleration phases. Eqs. (19) and (20) are denoted as zero-
 i on *a k* (2V) conditions increased and scillar in the independent of the trolley. With this accel-

e, the residual oscillation is eliminated in the non-

phases. Eqs. (19) and (20) are denoted as zero-

V) conditions. Various cases for the numbe ming the acceleration stairs of the trolley. With this accel-
tion profile, the residual oscillation is eliminated in the non-
eleration phases. Eqs. (19) and (20) are denoted as zero-
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pration (ZV) conditions. Various cases for the number of
ins N are presented below.
Case 1. $N = 1$
In this case, two unknowns need to be determined from four
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variation (ZV) conditions. Various cases for the number of
\nins *N* are presented below.

\nCase 1. *N* = 1

\nIn this case, two unknowns need to be determined from four
\nnstrain equations

\n
$$
a_0 + a_1 = 0, \qquad a_0 + a_1 \cos \omega_n \tau = 0
$$
\n(22)

\n
$$
a_1 \sin \omega_n \tau = 0, \qquad a_0 \tau = v_{\text{max}}.
$$
\nThese equations can be solved when time τ is chosen
\nsh that

\n
$$
\sin \omega_n \tau = 0 \& \cos \omega_n \tau = 1
$$
\n
$$
\Rightarrow \qquad \tau = 2k\pi / \omega_n, \qquad k = 1, 2, \ldots
$$

\nThus, when $\tau = 2k\pi / \omega_n$, the amplitudes are given by

\n
$$
a_1 = -a_0 = -\frac{v_{\text{max}}}{\tau} = -\frac{\omega_n}{2k\pi} v_{\text{max}}.
$$
\n(23)

\nThe maximum swing angle is determined by

\n
$$
\theta(t < \tau) = \frac{a_0}{g} [\cos \omega_n t - 1]
$$
\n
$$
\Rightarrow \theta_{\text{max}} (t = \tau / 2) = -\frac{2a_0}{g} = \frac{\omega_n}{\pi g} v_{\text{max}}.
$$

\nNotably, the case of *N* = 1 and the case of rectangular acceleration profile are identical.

\nCase 2, *N* − 2

These equations can be solved when time τ is chosen such that hese equations can be solved when time τ is chosen

1 that
 $\sin \omega_n \tau = 0 \& \cos \omega_n \tau = 1$
 $\Rightarrow \qquad \tau = 2k\pi / \omega_n, \qquad k = 1, 2, ...$

hus, when $\tau = 2k\pi / \omega_n$, the amplitudes are given by
 $\omega_1 = -a_0 = -\frac{v_{\text{max}}}{\tau} = -\frac{\omega_n}{2k\pi} v_{\text$

$$
\sin \omega_n \tau = 0 &c \cos \omega_n \tau = 1
$$

\n
$$
\Rightarrow \qquad \tau = 2k\pi / \omega_n, \qquad k = 1, 2, ...
$$

$$
a_1 = -a_0 = -\frac{v_{\text{max}}}{\tau} = -\frac{\omega_n}{2k\pi} v_{\text{max}}.
$$
 (23)

$$
a_0 + a_1 = 0, \t a_0 + a_1 \cos \omega_n \tau = 0
$$
\n
$$
a_1 \sin \omega_n \tau = 0, \t a_0 \tau = v_{\text{max}}.
$$
\nThese equations can be solved when time τ is chosen
\nch that\n
$$
\sin \omega_n \tau = 0 \& \cos \omega_n \tau = 1
$$
\n
$$
\Rightarrow \tau = 2k\pi / \omega_n, \t k = 1, 2, ...
$$
\nThus, when $\tau = 2k\pi / \omega_n$, the amplitudes are given by\n
$$
a_1 = -a_0 = -\frac{v_{\text{max}}}{\tau} = -\frac{\omega_n}{2k\pi} v_{\text{max}}.
$$
\n(23)\n\nThe maximum swing angle is determined by\n
$$
\theta(t < \tau) = \frac{a_0}{g} [\cos \omega_n t - 1]
$$
\n
$$
\Rightarrow \theta_{\text{max}} (t = \tau / 2) = -\frac{2a_0}{g} = \frac{\omega_n}{\pi g} v_{\text{max}}.
$$
\nNotably, the case of $N = 1$ and the case of rectangular ac-
\nleration profile are identical.\n\nCase 2. $N = 2$

Notably, the case of $N=1$ and the case of rectangular acceleration profile are identical.

Case 2. $N = 2$

$$
a_1 = -a_0 = -\frac{v_{\text{max}}}{\tau} = -\frac{\omega_n}{2k\pi} v_{\text{max}}.
$$
\n(23)
\nThe maximum swing angle is determined by
\n
$$
\theta(t < \tau) = \frac{a_0}{g} [\cos \omega_n t - 1]
$$
\n
$$
\Rightarrow \theta_{\text{max}} (t = \tau / 2) = -\frac{2a_0}{g} = \frac{\omega_n}{\pi g} v_{\text{max}}.
$$
\nNotably, the case of $N = 1$ and the case of rectangular acceleration profile are identical.
\nCase 2. $N = 2$
\nFrom Eqs. (19) and (21), the constraints are as follows:
\n
$$
a_0 + a_1 + a_2 = 0
$$
\n
$$
a_0 + a_1 \cos \omega_n \tau_1 + a_2 \cos \omega_n \tau_2 = 0
$$
\n
$$
a_1 \sin \omega_n \tau_1 + a_2 \sin \omega_n \tau_2 = 0
$$
\n
$$
a_0(\tau_2 - 0) + a_1(\tau_2 - \tau_1) = v_{\text{max}}.
$$
\nTwo time delays τ_1, τ_2 (0 $\langle \tau_1 \langle \tau_2 \rangle$) should be chosen to solve Eq. (24). Choosing $\tau_1 = 2k\pi / \omega_n$, $\tau_2 = 2\tau_1$, we obtain
\nthe following:
\n
$$
a_0 + a_1 + a_2 = 0,
$$
\n
$$
a_0 \tau_2 + a_1(\tau_2 - \tau_1) = v_{\text{max}}.
$$
\nThese equations may be expressed in matrix form as follows:
\n
$$
\Phi u = d_1,
$$

$$
a_0 + a_1 + a_2 = 0,
$$

$$
a_0 \tau_2 + a_1 (\tau_2 - \tau_1) = v_{\text{max}}.
$$

These equations may be expressed in matrix form as follows:

$$
\Phi u = d, ,
$$

with

$$
\Phi = \begin{bmatrix} 1 & 1 & 1 \\ \tau_2 & \tau_2 - \tau_1 & 0 \end{bmatrix}, \qquad \mathbf{d} = \begin{bmatrix} 0, v_{\text{max}} \end{bmatrix}^T.
$$

By applying the pseudo-inverse, we obtain the solution
0.2

$$
\mathbf{a} = \mathbf{\Phi}^T (\mathbf{\Phi} \mathbf{\Phi}^T)^{-1} \mathbf{d}
$$

chosen based on the pendulum frequency.

Case 3. $N \geq 3$

In this case, constraint Eqs. (19) and (21) are determined after choosing an arbitrary τ_i , ($i = 1, 2, ..., N$) that satisfies pendulum period of the payload; thus, $T = 2\pi / \omega_n$.

Case 4. $N \geq 5$

For $N \geq 5$, the system consisting of four constraints, that is, least two additional equations. We set the derivative of Eq. (20) with respect to the natural frequency ω_n to zero and thus obtaining the following: y, the time τ , τ_1 , τ_2 in cases $N = 1$ and $N = 2$ is

and on the pendulum frequency.

i. $N \ge 3$

case, constraint Eqs. (19) and (21) are determined af-

the movement distance in the accelerating phase

ease, co Since of residual vibration appearance of variables of residual vibration appearance of variables constraint Eqs. (19) and (21) are determined af-

sin cose, constraint Eqs. (19) and (21) are determined af-

the movement *i n i i i n i ⁿ ⁱ ⁱ ^N ^N i n i i i n i ⁿ ⁱ ⁱ ^a ^a* otably, the time τ , τ _i, τ _i an cases $N = 1$ and $N = 2$ is
 $\frac{0.6}{N} - \frac{0.8}{N} - \frac{1}{N}$ and an absorb the pendulum frequency.
 ase 3. $N \ge 3$

this case, constraint Eqs. (19) and (21) are determined af-

th che absord on the pendulum frequency.

see 3. $N \ge 3$

this case, constraint Eqs. (19) and (21) are determined af-

the movement distance in the accelerating pl

choosing an arbitrary τ_i , $(i = 1, 2, ..., N)$ that satisfies
 ter choosing an arbitrary τ_i , $(i = 1, 2, ..., N)$ that satisfies
 $0 = \tau_0 < \tau_i < \tau_{i\infty}$. The time τ_N should be greater than the

pendulum period of the payload; thus, $T = 2\pi / \omega_s$.
 $\alpha(t) = \sum_{i=0}^{N} a_i \sigma(t - \tau_i)$
 Case 4.

$$
\frac{\partial}{\partial \omega_n} \sum_{i=0}^N a_i \cos \omega_n \tau_i = \sum_{i=0}^N a_i \tau_i \sin \omega_n \tau_i = 0
$$
\n
$$
\frac{\partial}{\partial \omega_n} \sum_{i=0}^N a_i \sin \omega_n \tau_i = \sum_{i=0}^N a_i \tau_i \cos \omega_n \tau_i = 0.
$$
\n(25)\n
\nThe duration is not s

Combining Eqs. (19) , (21) , and (25) leads to a set of six straint equations are denoted as zero-vibration derivative (ZVD) conditions for eliminating the residual oscillation of the payload and rewritten as follows:

(b) With respect to the natural frequency
$$
\omega_n
$$
 to zero and
\nis obtaining the following:
\n
$$
\frac{\partial}{\partial \omega_n} \sum_{i=0}^{N} a_i \cos \omega_n \tau_i = \sum_{i=0}^{N} a_i \tau_i \sin \omega_n \tau_i = 0
$$
\n
$$
\frac{\partial}{\partial \omega_n} \sum_{i=0}^{N} a_i \sin \omega_n \tau_i = \sum_{i=0}^{N} a_i \tau_i \cos \omega_n \tau_i = 0
$$
\n(c) The duration $\tau_c = [(p_d - 2x(\tau_N)] / v_{\text{max}} \ge 0]$.
\nCombining Eqs. (19), (21), and (25) leads to a set of six
\ntion is not satisfied, the maximum
\nduced. Reduction of V_{max} leads
\nfrom the natural equations for N+1 ≥ 6 unknowns. These con-
\ntion is not satisfied, the maximum
\nduced. Reduction of V_{max} leads
\n $\frac{1}{2} \omega_a \sin \omega_a \tau_i = 0$
\n $\frac{1}{2} \omega_a \sin \omega_a \tau_i = 0$
\n $\frac{1}{2} \omega_a \cos \omega_a \tau_i = 0$
\n $\frac{1}{2} \omega_a \sin \$

3.3 Determination of time for the constant velocity phase

The constant velocity duration τ_c must be determined to ensure that the trolley reaches the desired location p_d . First,

Fig. 3. Magnitude of residual vibration dependent on N and ω/ω_n .

the movement distance in the accelerating phase is evaluated by integrating the velocity function from $t = 0$ to $t = \tau_N$

0.6 0.7 0.8 0.9
$$
\frac{1}{1}
$$
 1.1 1.2 1.3 1.4
Fig. 3. Magnitude of residual vibration dependent on *N* and ω/ω_n .
the movement distance in the accelerating phase is evaluated
by integrating the velocity function from $t = 0$ to $t = \tau_N$

$$
a(t) = \sum_{l=0}^{N} a_l \sigma(t - \tau_i)
$$

$$
v(t) = \int_0^t a(\overline{t}) d\overline{t}
$$
(27)

$$
x(\tau_N) = \int_0^{\tau_N} v(\overline{t}) d\overline{t}
$$
(27)
The time for constant velocity phase is then expressed as
follows:

$$
\tau_c = [(\rho_d - 2x(\tau_N)] / \nu_{\text{max}} \ge 0.
$$
(28)
The duration τ_c must be non-negative. When this condi-
tion is not satisfied, the maximum velocity ν_{max} can be re-
duced. Reduction of ν_{max} leads to the reduction of acceleration a_i ; thus, the traveling distance in the accelerating dura-
tion $x(\tau_N)$ is also decreased.
3.4 *Robustness to parameter uncertainties*
The residual oscillation is eliminated with any number of
fairs $N \ge 1$ when the natural frequency ω_n is precisely
determined. However, this requirement may not be satisfied

The time for constant velocity phase is then expressed as follows:

$$
\tau_c = [(p_d - 2x(\tau_N)] / v_{\text{max}} \ge 0. \tag{28}
$$

The duration τ_c must be non-negative. When this condition is not satisfied, the maximum velocity v_{max} can be reduced. Reduction of v_{max} leads to the reduction of acceleration a_i ; thus, the traveling distance in the accelerating duration $x(\tau_N)$ is also decreased.

3.4 Robustness to parameter uncertainties

 $\sum a_i \cos \omega_n \tau_i = 0$ the residual vibration amplitude has been assumed to be de- $\sum a_i \sin \omega_n \tau_i = 0$ ence of the stair number *N*. Fig. 3 shows the magnitude $\sum_{i=0}^{N} a_i \tau_i \sin \omega_n \tau_i = 0$ around the nominal values ω_n for some cases $N = 1,3,5,6$.
For $N = 5$ and 6, cases including both ZV and ZVD condi- $\sum_{i=0} a_i \tau_i \cos \omega_n \tau_i = 0$ the magnitude of the residual vibration with $N = 1$ exhibits the The residual oscillation is eliminated with any number of stairs $N \ge 1$ when the natural frequency ω_n is precisely determined. However, this requirement may not be satisfied. In this section, the number N is selected. The dependence of fined by Eq. (18) on ω_n and N to demonstrate the influ- $\tau_c = [(p_d - 2x(\tau_N)]/v_{max} \ge 0.$ (28)

The duration τ_c must be non-negative. When this condi-

tion is not satisfied, the maximum velocity v_{max} can be re-

duced. Reduction of v_{max} leads to the reduction of accelera-

t $V(\omega_n, N)$ when the natural frequency of the pendulum varies tions are considered. As indicated in Figs. 3(a) and (b), the highest rate of change around the location *n* a_i ; thus, the traveling distance in the accelerating dura-
n $x(x_N)$ is also decreased.
Robustness to parameter uncertainties
The residual oscillation is eliminated with any number of
Iris $N \ge 1$ when the natura

 $\omega/\omega_{n} = 1$, and is depicted as the highest curve. The curve with $N = 3$ is lower than the curves with $N = 5$ and $N = 6$, as shown in Fig. 3(a). With the ZVD constraints considered (Fig. 3(b)), the curves with $N=5$ (ZVD) and 6 (ZVD) are almost the same and lower than the curve with $N = 5$ (ZV) around the nominal values ω_n . Based on these figures, we

3.5 Algorithm for determining trajectory parameters

- Number of stairs $N = 1, 2, 3, ...$
- Time sequence:
- $k = 1, 2, \ldots$.
- If $N = 3, 4, \ldots$: $\tau_0 = 0$, consider freely accelerating du-
-
- Sequence of acceleration step a_i corresponding to time sequence τ_i , $i = 0, 1, ..., N$.
-
-
-
-

4.1 Kinematic simulation

Numerical simulations were conducted using MATLAB to verify the validity and efficiency of the proposed approach. In the simulation, the system parameters were set as $m_i = 2.0$ kg, $m_p = 0.85$ kg, $l = 1.20$ m, and $g = 9.81$ m/s², and the target position of the trolley was set as $x_d = 4$ m. The maximum velocity is selected as $v_{\text{max}} = 0.5 \text{ m/s}$. The pendulum frequency and the corresponding period were determined to be ω _n = 2.8592 rad/s and *T* = 2.1975 s. By applying the algorithm presented in Section 3, the trajectory parameters for some cases of *N* were obtained, as presented in Table 1.

The simulation results for the displacement and acceleration of the trolley, as well as the swing angle of the payload, are shown in Fig. 4. The desired position is achieved after a finite time. The swing angle is kept small in the non-acceleration phase and when the trolley reaches its final position. The maximum velocity can be increased without changing the acceleration duration. However, increasing maximum velocity leads to increases in the acceleration magnitude and the swing angle in the acceleration time.

To compare the robustness between the ZVD conditions and the ZV conditions, we consider the case wherein the cable length is changed to 20% of the nominal one, $l = 1.2l_0$. The swing angles of the payload in some cases are shown in Fig. 5.

Table 1. Calculation results for some cases of *N.*

2884 N. Q. Hoang et al. / Journal of Mechanical Science and Technology 28 (7) (2014) 2879~2888							
conclude that selecting $N = 3$ for the case ZV or $N = 5$ for the case ZVD can improve the robustness of the approach.	Table 1. Calculation results for some cases of N.						
	Case	N and k	\mathbf{a} [m/s ²]	τ [s]	T_c [s]	T_f [s]	$\theta_{\rm max}$ [°]
3.5 Algorithm for determining trajectory parameters	$1-1.$	$N=1$. $k=1$	0.2275 -0.2275	$\tau = \lceil 0 \rceil$ 2.1975]	5.802	10.197	2.654
The proposed trajectory planning method for overhead cranes can be summarized as follows:	$1-2.$	$N=1$, $k = 2$	0.1138 -0.1138	$\tau = \lceil 0 \rceil$ 4.3951]	3.604	12.395	1.325
(1) The cable length l , traveling distance or desired loca- tion p_d , and maximum velocity v_{max} are obtained. (2) The natural frequency of the pendulum $\omega_n = \sqrt{g/l}$ and the period of vibration $T = 2\pi / \omega_n$ are calculated. (3) The number of stairs N and time sequence are selected as follows - Number of stairs $N = 1, 2, 3, $ - Time sequence:	2. $\tau_{N} = 1.5$ T	$N=3$ (ZV)	0.1138 0.1138 -0.1138 -0.1138	$\tau = [0]$ 1.0988 2.1975 3.2963]	4.703	11.296	1.333
	3. $\tau_{N} = 1.5$ T	$N=4$ (ZV)	0.0889 0.1257 -0.0000 -0.1257 -0.0889	$\tau = \lceil 0 \rceil$ 0.8241 1.6482 2.4722 3.2963]	4.703	11.296	1.775
• If $N=1$ or 2: consider $\tau_0 = 0$, $\tau_N = kT = k \cdot 2\pi / \omega_n$, with $k = 1, 2, \dots$. • If $N = 3, 4, $: $\tau_0 = 0$, consider freely accelerating du- ration $\tau_N > T = 2\pi / \omega_n$. (4) The following are also calculated:	4. $\tau_{N} = 1.5 \text{ T}$	$N=5$ (ZV)	0.0822 0.0938 0.0658 -0.0658 -0.0938 -0.0822	$\tau = \lceil 0 \rceil$ 0.6593 1.3185 1.9778 2.6370 3.2963]	4.703	11.296	1.705
- Sequence of acceleration step a_i corresponding to time sequence τ_i , $i = 0, 1, , N$. - Travelling distance in the accelerating and decelerating phases $d = 2 \cdot x(\tau_N)$. - Time duration for constant velocity phase $\tau_c = [p_d - 2 \cdot x(\tau_N)] / v_{\text{max}}$.	5. $\tau_{N} = 1.5$ T	$N=5$ (ZVD)	0.1107 0.0261 0.1268 -0.1268 -0.0261 -0.1107	$\tau = [0]$ 0.6593 1.3185 1.9778 2.6370 3.2963]	4.703	11.296	1.420
- Total time of operation							
$T_f = \tau_c + 2\tau_N = 2\tau_N + [p_d - 2 \cdot x(\tau_N)] / v_{\text{max}}$.	4						
4. Numerical simulations	3 $x \in [m]$ $\overline{2}$					$N = 1$ $=$ 3	
4.1 Kinematic simulation	1					$N = 5$ zv	
Numerical simulations were conducted using MATLAB to				$N = 5$ zvD			

Fig. 4. Kinematic simulation results when $l = 1.2$ m and $x_d = 4$ m.

The residual oscillation of the payload when $N = 5$ (ZVD) is the smallest compared with other cases with $N = 1$ and $N = 5$ (ZV), as shown in Fig. 5. These results verify the robustness

Fig. 5. Simulation results for $N = 5$ ($l = l_0$, $l = 1.2$ l_0 , ZV, and ZVD).

of the approach to the uncertainty in the cable length.

4.2 Dynamic simulation with a sliding mode controller

The trolley must be manipulated to track the designed trajectory to ensure a small swing angle in the constant velocity phase. A sliding mode controller was applied by considering disturbances induced by the payload swing. This controller was developed based on the dynamics of the trolley, and required only the trolley motion for feedback.

The control equation is given by

$$
u = (m_t + m_p)\ddot{x}_r + f_{11}\dot{x} - k s - k_s \text{ sgn}(s)
$$
 (29)

parameter). The control equation denoted as Eq. (29) may cause chattering in the system because of the sgn -function, with $c \gg 1$.
Some simulations were performed with the controller pa-

The trolley must be manipulated to track the designed tra-
johas. A sliding mode controller was applied by considering
disturbances induced by the payload swing. This controller
disturbances induced by the payload swing. displacement of the trolley, swing angle of the load, and control input are shown in Fig. 6. The trolley tracks the desired trajectory and reaches its destination, and the swing angle remains small. The residual oscillations in the cases $N = 3$ (ZV) and $N = 5$ (ZVD) are smaller than those in cases $N = 1$ and $N = 5$ ZV, respectively.

4.3 Experiments

Experiments with the laboratory overhead crane have been conducted to verify the proposed approach. As shown in Fig. 7, the laboratory crane used in the experiment is equipped with three direct current (DC) motors to manipulate the trolley and the bridge motions, as well as hoist the payload. Five incremental encoders with 1024 counts per revolution were used to measure the trolley and the bridge displacements, cargohoisting motion along the cable, and two payload swing angles corresponding to the motion directions of the trolley and the bridge. The crane system was connected to a target personal computer (PC) with two interface cards. An NI PCI-

Fig. 6. Simulation results with a sliding mode controller.

Fig. 7. Laboratory overhead crane system.

6602 card was used to send pulse-width modulation signals to the amplifiers of the DC motors and acquire signals from the encoders. An NI PCI-6025E multifunction card was used to transfer the direction control signals to the motor amplifiers. The target PC was connected to a host PC through RS-232 ports. The overhead crane was controlled by the host PC, which integrated the proposed controller design based on MATLAB/SIMULINK with xPC Target. For purpose of experiment with the proposed controller, the bridge motion and hoist motion are fixed so that the apparatus becomes two-DOF system driven only by one actuator.

The parameters of the laboratory crane are as follows: $m_t = 2.0$ kg, $m_p = 0.85$ kg, $f₁₁ = 10$ Ns/m, $l = 0.6$ m, and $g = 9.81 \text{ m/s}^2$. The target position of the trolley is set as

 $x_d = 0.5$ m. The frequency $\omega_n = 4.044$ s⁻¹ of the pendulum is determined based on the cable length, and the accelera-2886 *N. Q. Hoang et al. / Journal of Mechanical Science and Technology 28 (7) (2014) 2879–288*
 $x_d = 0.5$ m. The frequency $\omega_n = 4.044$ s⁻¹ of the pendulum

is determined based on the cable length, and the accelera-

t number of stairs $N = 3$, the maximum velocity $v_{\text{max}} = 0.2 \text{ m/s}$, $\tau = \begin{bmatrix} 0 & 0.65 & 1.29 & 1.94 \end{bmatrix}^T$, the $\begin{bmatrix} 0.3 \end{bmatrix}$ magnitude of stairs were determined to be $a =$ 2886 *N. Q. Hoang et al. / Journal of Mechanical Science and Technology 28 (7) (2014) 2879-2888*
 $x_y = 0.5$ m. The frequency $\omega_n = 4.044$ s⁻¹ of the pendulum

is determined based on the cable length, and the accelera-
 section, the system parameters were chosen to be large enough to clearly show three phases (accelerating, constant movement, and decelerating); hence, $l = 1.2$ m and $x_d = 4$ m were set. However, in the experiment, these parameters were set to the actual values $l = 0.6$ m and $x_d = 0.5$ m because of the limitation of the system dimension. Therefore, the acceleration and velocity profiles used in the experiment are different from those of the simulation. The obtained velocity profile of the experiment looks like a bell shape denoted as a broken line in Fig. 8. Almost no constant velocity phase exists because of such a small dimension, such that the force trajectory as control input also looks like a bell shape, as shown in Fig. 10. Meanwhile, the force trajectory of the simulation looks like a trapezoid, as shown in Fig. 6. The acceleration profile for the experiment is denoted as a broken line in Fig. 10. With the planned trajectory, the sliding mode controller of Eq. (29) is applied to the real crane system. The sliding mode controller is the same as that of the simulation and the same control parameters: $\lambda = 2$, $k = 10$, and $k_s = 2$ are applied for controlling the crane. However, in the experiment uses parameters were set to the

nactual values $l = 0.6$ m and $x_l = 0.5$ in because of the limitation of the system dimension. Therefore, the acceleration

and velocity profiles used in the expe For the simulation. The obtained we can be carried with the parameters of the simulation, such that the force trajectory as some state when $\frac{1}{2}$ and $\frac{1}{2}$ and $\frac{1}{2}$ are carried to the simulation books like a

Similarly, for $N = 5$ (zvd) and the acceleration time magnitude of stairs are determined as

$$
\tau = \begin{bmatrix} 0 & 0.466 & 0.932 & 1.398 & 1.864 & 2.330 \end{bmatrix}' \qquad \qquad \begin{bmatrix} 8 & 2 \\ 9 & 1 \end{bmatrix} \qquad \begin{bmatrix} 1 & 0.466 & 0.932 & 1.398 & 1.864 & 2.330 \end{bmatrix}'.
$$

Two experiments have been carried out for $N = 3$ and $N = 5$ (zvd), in which the parameters of the controller Eq. (30) are chosen as $\lambda = 2$, $k_s = 2$, and $k = 10$.
The experimental results, including the trolley motion,

swing angle, and control force, are shown in Figs. 8-10. For comparison with the numerical simulation, the responses of the simulation with $l = 0.6$ m and $x_d = 0.5$ m are also shown as dotted lines in Figs. 8-10. The trolley reaches the target position while the swing angle is kept small at about 1° during operation for both the simulation and the experiment. Apparently, the responses of the experiment are quite similar to those of the simulation. The sway angle of the experiment is somewhat different from that of the simulation. However, considering that the maximum value of the sway is small, the difference is minimal. We assumed that the crane has only 2 dimensional motion, but the actual swing motion of the rope would be 3-dimensional in the experiment. Coupled dynamic effect of the 3D motion and un-modeled nonlinear characteristics may cause difference in swing response between the simulation and the experiment. These factors also induce re-

Fig. 8. Displacement and velocity of the trolley.

Fig. 9. Swing angle of the rope.

Fig. 10. Control input on the trolley.

sidual vibration after 5 s in the experiment, as shown in Fig. 9. During the beginning of the acceleration and deceleration periods, 2-dimensional driving effect is dominant in the swing angle; hence, the experimental and simulation responses coincide. These experimental results directly confirm that a small swing angle is maintained given a suitable trajectory for the trolley. It is seen from Fig. 9 that, the swing angle for $N = 5$ is smaller than that for $N = 3$.

5. Conclusions

This paper presents a simple trajectory planning method for underactuated overhead cranes using acceleration in a staircase form. This acceleration profile ensures a small swing angle of the payload during the constant velocity phase and as the trolley reaches its desired location. The efficiency of the proposed approach was verified by numerical simulation in the kinematics viewpoint and in that of crane dynamics. These simulations indicate that the payload swing is suppressed by a reasonable trajectory of the trolley and a simple controller that requires only the trolley motion for feedback. The advantages exhibited potential applications in motion planning for underactuated systems, such as overhead cranes, single-link flexible-joint manipulators, and flexible Cartesian manipulators. The proposed method intends to include three-dimensional overhead cranes in future studies.

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