ORIGINAL RESEARCH



Uniform convergence analysis of a new adaptive upwind finite difference method for singularly perturbed convection-reaction-diffusion boundary value problems

Quan Zheng¹ · Zhongli Liu²

Received: 14 November 2023 / Revised: 9 December 2023 / Accepted: 19 December 2023 / Published online: 14 January 2024 © The Author(s) under exclusive licence to Korean Society for Informatics and Computational Applied Mathematics 2024

Abstract

In this paper, a new adaptive upwind finite difference method based on the arc-length equidistribution principle is studied for solving the general linear singularly perturbed convection-reaction-diffusion two-point boundary value problem. Under the discrete comparison principle and its derived properties, the a posteriori error estimate of the upwind finite difference scheme on an arbitrary mesh is obtained by using a linear interpolation, the existence of the solution of the adaptive upwind finite difference method is proved without the presupposition, and the boundedness of the arc-length of the numerical solution and then the uniform first-order convergence of the adaptive upwind finite difference method are achieved by the discrete Green's function. Finally, the proposed method is verified to be uniformly first-order convergent and is compared with the existing method in numerical examples

Keywords Singularly perturbed problem \cdot Upwind finite difference scheme \cdot Arc-length monitor function \cdot A posteriori error estimate \cdot Adaptive/moving grid \cdot Uniform first-order convergence

Mathematics Subject Classification $\,65L10\cdot 65L11\cdot 65L12\cdot 65L20$

1 Introduction

The numerical methods for solving singularly perturbed problems are often concerned and have been developed for decades (see [1-5]). Besides fitted operator numeri-

Quan Zheng zhengq@ncut.edu.cn
 Zhongli Liu shtzhongli@buu.edu.cn

¹ College of Sciences, North China University of Technology, Beijing 100144, China

² Institute of Mathematics and Physics, Beijing Union University, Beijing 100101, China

cal methods on the uniform mesh (see [1, 2, 5, 6]), two strategies are commonly used to solve these problems on the non-uniform mesh. One strategy is to use layeradapted grid methods, i.e., numerical methods on a priori layer-adapted meshes, such as Bakhvalov-type meshes and Shishkin-type meshes (see [1, 2, 4, 5, 7-10]). Another strategy is to use adaptive grid methods, i.e., numerical methods on adaptive meshes. The application of the adaptive grid gets rid of the limitation of a priori knowledge of the exact solution, only needs to find a monitor function to help optimize the grid continually (see [1, 3, 5, 11-28]). The monitor function is defined in terms of the local error (or the local residual) or the arc-length of the solution and so on. Because the construction of the adaptive grid (or the moving mesh) does not depend on the properties of the exact solution, the progresses on the convergence analysis of the adaptive grid method based on the arc-length monitor function have been made gradually almost only for the special cases (see [14-21]), but seldom for the general linear non-conservative case (see [22-26]).

We consider the general linear singularly perturbed convection-reaction-diffusion problem in the non-conservative form:

$$\begin{cases} Tu(x) := -\varepsilon u''(x) - p(x)u'(x) + q(x)u(x) = f(x), \ x \in (0, 1), \\ u(0) = 0, \ u(1) = 0, \end{cases}$$
(1)

where $0 < \varepsilon \ll 1$ is a sufficiently small positive perturbation parameter, p, q and f are sufficiently smooth functions, $0 < \beta \le p(x) \le \beta^*$, $|p'(x)| \le C_p$, $0 \le q(x) \le \gamma^*$ and $|q'(x)| \le C_q$ for $x \in [0, 1]$, $\beta, \beta^*, \gamma^*, C_p$ and C_q are constants, and C is a generic positive constant, independent of the perturbation parameter ε . The problem has a unique solution u that has an exponential boundary layer at x = 0.

This problem (1) was solved by the upwind finite difference scheme on the arclength equidistribution mesh as follows (see [22–25]):

$$\begin{cases} T^{N}u_{i}^{N} := -\varepsilon DD^{-}u_{i}^{N} - p_{i}D^{+}u_{i}^{N} + q_{i}u_{i}^{N} = f_{i}, i = 1, ..., N - 1, \\ u_{0}^{N} = 0, u_{N}^{N} = 0, \end{cases}$$
(2)

where $h_i = x_i - x_{i-1}$, $D^- u_i = \frac{u_i - u_{i-1}}{h_i}$, $D^+ u_i = \frac{u_{i+1} - u_i}{h_{i+1}}$ and $DD^- u_i = \frac{2(D^- u_{i+1} - D^- u_i)}{h_{i+1} + h_{i+1}}$.

The development of related adaptive methods and their convergence analysis are reviewed in the following.

(i) In the early days, for the non-conservative homogeneous linear case $-\varepsilon u''(x) - p(x)u'(x) = 0$, the convergence of the upwind finite difference scheme on the grid equidistributing the monitor function $M(x, u(x)) = |u'(x)|^{1/m}$ and $\sqrt{1 + (u'(x))^2}$ of the exact solution *u* was obtained by using the a priori truncation error in [11–14]. For the non-homogeneous case $-\varepsilon u''(x) - p(x)u'(x) = f(x)$, the convergence of the upwind finite difference scheme by using $M(x, u(x)) = \alpha + |u''(x)|^{1/m}$ of the exact solution *u* was obtained by the a priori truncation error in [15]. For the reaction-diffusion case $-\varepsilon u''(x) - b(x)u(x) = f(x)$, the convergence of the central finite difference scheme using $M(x, u(x)) = \alpha + |w''(x)|^{1/m}$ (*w* is the layer component of *u*) was obtained by the a priori truncation error in [16].

For the conservative quasilinear case $-\varepsilon u''(x) - b(x, u)' = f(x)$, the first-order convergence of the simple upwind finite difference scheme using monitor functions based on a priori information of the exact solution was also obtained in [17].

- (ii) In fact, to devise a rigorous analysis of any adaptive algorithm, the first ingredient one needs is an a posteriori bound on the error of the computed solution (see §I.2.5 in [1]). For the conservative quasilinear case $-\varepsilon u''(x) - b(x, u)' = f(x)$, the convergence of the simple upwind finite difference scheme on the adaptive mesh equidistributing the numerical arc-length monitor function $M(x, u^N(x)) = \sqrt{1 + ((u^N(x))')^2}$ was firstly analyzed by the a posteriori truncation error in [18]. Its uniform first-order error bound was obtained by using the Green's function, since the conservative equation could easily be reduced to the first-order equation (see [1, 18, 19]). The numerical solution of the above problem was also calculated by equidistributing a monitor function of discrete second-order derivatives in [20].
- (iii) Afterwards, the special case $-\varepsilon u''(x) p(x)u'(x) = 0$ was solved by the simple upwind scheme on the adaptive mesh equidistributing the numerical arc-length monitor function $M(x, u^N(x))$. By using some techniques developed in [18], the existence of the adaptive solution was proved with a strong grid assumption $h_i/h_{i+1} \le 1$, and the uniform first-order convergence was obtained by the a posteriori error estimation and using a quadratic interpolation of the numerical solution (see [21]).
- (iv) However, the uniform first-order convergence of (2) for the general case $-\varepsilon u''(x) p(x)u'(x) + q(x)u(x) = f(x)$ was proved in [22] by the a priori truncation error and the grid properties obtained in [14] for the special case $-\varepsilon u''(x) p(x)u'(x) = 0$, but not the a posteriori error estimate and not free from the properties of the exact solution. The uniform first-order convergence of (2) for singularly perturbed mixed boundary value problems was derived in [23] also by a priori truncation error, the properties of the exact solution and the grid properties obtained in [14].
- (v) In the later, by developing the techniques in [21], the existence of the numerical solution of (2) for the general problem (1) on the adaptive mesh equidistributing the numerical arc-length monitor function $M(x, u^N(x))$ was proved by a slightly weakened grid restriction $h_i/h_{i+1} \leq C$, and the uniform first-order convergence was obtained by the a posteriori error estimation and also the quadratic interpolation of the numerical solution for (1) in [24].
- (vi) Further, the quadratic interpolation was replaced by a linear interpolation in the error analysis of (2) for (1) in [25] and also of a similar scheme for a singularly perturbed differential difference equation with small delay in [26]. However, the uniform first-order convergence in [25, 26] was obtained by citing [24] and [18, 21] respectively, where the quadratic interpolation and the impractical condition $h_i/h_{i+1} \leq C$ or the conservative equation were used, but without any proof of the existence of numerical solution and the boundedness of the arc-length of the numerical solution.
- (vii) Recently, for the reaction-diffusion singularly perturbed boundary value problem $-\varepsilon u''(x) + a(x)u(x) = f(x)$ of robin-type, the spline difference method on the a posteriori layer adapted mesh by using the monitor function $M(x, u(x)) = \alpha + |w''(x)|^{1/m}$ (see also [16] was discussed in [27]. A posteriori error analyses

for the system of reaction-diffusion problems, the nonlinear singularly perturbed system of delay differential equations and the singularly perturbed nonlinear parameterized problem were also discussed in [28–30]. The algorithm and its extension for the system of mixed type reaction diffusion problems and several singularly perturbed parabolic problems were developed in [31–35].

The stability of T in (1) was derived in Corollary 4.2 in [24] as follows:

Lemma 1.1 (see [24]) Let $F(x) = \int_x^1 f(s)ds + c_0$ be a bounded piecewise continuous function. Then, there exists a unique weak solution $u \in H_0^1(0, 1)$ for the problem (1), such that

$$\|u\|_{\infty} \le C_1 \|Tu\|_{*},\tag{3}$$

where $C_1 = C(2\exp(\gamma^*/\beta) + \beta^*/\beta)/\beta^2 + C$, and

$$\|u\|_{\infty} = ess \sup_{x \in [0,1]} |u(x)| \quad and \quad \|u\|_{*} = \min_{U:U'=u} \|U(x)\|_{\infty}.$$
 (4)

Actually, the stability (3) can also be proved under the condition $q \ge 0$ and $q + p' \ge 0$ by the derivation for Theorem 1.7 in §I.1.1.2 in [1].

In this paper, we investigate a new upwind finite difference scheme different from (2) on the adaptive mesh equidistributing the numerical arc-length monitor function $M(x, u^N(x))$ to solve the general non-conservative non-homogeneous convection-reaction-diffusion problem (1). In Section 2, the a posteriori error estimate for the upwind finite difference scheme on an arbitrary mesh is obtained by the stability of T and a linear interpolation of u^N . In Section 3, the adaptive upwind finite difference method is proposed and the existence of the solution is proved by the fixed point theorem without the grid restriction $h_i/h_{i+1} \leq 1$ or C. In Section 4, the boundedness of the arc-length of the numerical solution and then the uniform first-order convergence are derived by using the property of the discrete Green's function. Finally, in Section 5, the adaptive method is compared with the existing method and shown to be effective and uniformly first-order convergent in the numerical examples.

2 The upwind finite difference scheme and its a posteriori error estimate

For any nodes on interval [0, 1], $0 = x_0 < x_1 < ... < x_n = 1$, we investigate a new simple upwind finite difference scheme as follows:

$$\begin{cases} T^{N}u_{i}^{N} := -\varepsilon D^{+}(D^{-}u_{i}^{N}) - p_{i}D^{+}u_{i}^{N} + q_{i}u_{i}^{N} = f_{i}, i = 1, ..., N - 1, \\ u_{0}^{N} = 0, \ u_{N}^{N} = 0, \end{cases}$$
(5)

where $h_i = x_i - x_{i-1}$, $D^-u_i = \frac{u_i - u_{i-1}}{h_i}$, $D^+u_i = \frac{u_{i+1} - u_i}{h_{i+1}}$ and then $D^+(D^-u_i)$ = $\frac{D^-u_{i+1} - D^-u_i}{h_{i+1}}$. The simple upwind scheme (5) when $h_i = h$ gives the simple upwind scheme (2.12) in §I.2.1.2 in [1] and is much similar to the scheme $-D^+(\varepsilon D^-u_i^N + \varepsilon D^-u_i^N)$ $b(x_i, u_i^N) = f_i$ in [17–20] and the scheme in [8–10], but different from the usually adopted classical simple upwind scheme (2) in [7, 11–15, 21–26]. And it is also different from the central difference scheme to avoid oscillations of it, see e.g. the central difference scheme (2.3) in §I.2.1.1 and the discussion in §I.2.1.2 in [1].

Since $T^N u_i^N = -\frac{\varepsilon}{h_i h_{i+1}} u_{i-1}^N + (\frac{\varepsilon}{h_i h_{i+1}} + \frac{\varepsilon}{h_{i+1} h_{i+1}} + \frac{p_i}{h_{i+1}} + q_i) u_i^N - (\frac{\varepsilon}{h_{i+1} h_{i+1}} + \frac{p_i}{h_{i+1}}) u_{i+1}^N$, the matrix of T^N of the scheme (5) is diagonally dominant and has non-positive off-diagonal entries. Thus, the matrix is an M-matrix. Let $\Psi_i^{\pm} = \frac{1}{\beta} (1 - x_i) \|T^N u_i^N\|_{\infty} \pm u_i^N$. Then

$$T^{N}\Psi_{i}^{\pm} = \frac{p_{i}}{\beta} \|T^{N}u_{i}^{N}\|_{\infty} + \frac{q_{i}}{\beta}(1-x_{i})\|T^{N}u_{i}^{N}\|_{\infty} \pm T^{N}u_{i}^{N} \ge 0,$$

 $\Psi_0^{\pm} \ge 0$ and $\Psi_N^{\pm} \ge 0$. So, by the discrete comparison principle, $\Psi_i^{\pm} \ge 0$, i.e., T^N is uniformly stable (cf. Lemma 2.11 in §I.2.1 in [1]):

$$|u_i^N| \le \frac{1}{\beta} \|T^N u_i^N\|_{\infty} \le \frac{1}{\beta} \|f_i\|_{\infty}, \quad i = 0, 1, ..., N.$$
(6)

For any $x \in I_i = (x_i, x_{i+1})$, i = 0, ..., N - 1, we have a piecewise constant function $f^N(x) = f_i$ and a piecewise linear interpolation

$$u^{N}(x) = \frac{1}{h_{i+1}} (u^{N}_{i+1}(x - x_{i}) + u^{N}_{i}(x_{i+1} - x)),$$
(7)

which satisfies $u^N(x_i) = u_i^N$, $u^N(x_{i+1}) = u_{i+1}^N$ and $(u^N(x))' = D^+ u_i^N$. Alternatively, a quadratic interpolation was used in [21–24].

For any $x \in I_i = (x_i, x_{i+1}), i = 0, ..., N - 1$, we have

$$\begin{split} Tu^{N}(x) &- f^{N}(x) \\ &= -\varepsilon(u^{N}(x))'' - p(x)(u^{N}(x))' + q(x)u^{N}(x) \\ &+ \varepsilon D^{+}(D^{-}u_{i}^{N}) + p_{i}D^{+}u_{i}^{N} - q_{i}u_{i}^{N} \\ &= (-\varepsilon(u^{N}(x))' + \int_{x}^{1} p(t)(u^{N}(t))' - q(t)u^{N}(t)dt \\ &- \sum_{j=i+1}^{N-1} \int_{I_{j}} \varepsilon D^{+}(D^{-}u_{j}^{N}) + p_{j}D^{+}u_{j}^{N} - q_{j}u_{j}^{N}dt \\ &- \int_{x}^{x_{i+1}} \varepsilon D^{+}(D^{-}u_{i}^{N}) + p_{i}D^{+}u_{i}^{N} - q_{i}u_{i}^{N}dt + \varepsilon D^{-}u_{N}^{N})' \\ &= (-\varepsilon D^{+}u_{i}^{N} + \sum_{j=i+1}^{N-1} \int_{I_{j}} p(t)D^{+}u_{j}^{N} - q(t)u^{N}(t)dt \\ &+ \int_{x}^{x_{i+1}} p(t)D^{+}u_{j}^{N} - q(t)u^{N}(t)dt \\ &- \varepsilon D^{-}u_{N}^{N} + \varepsilon D^{-}u_{i+1}^{N} - \sum_{j=i+1}^{N-1} \int_{I_{j}} p_{j}D^{+}u_{j}^{N} - q_{j}u_{j}^{N}dt \\ &+ \int_{x}^{x_{i+1}} f_{i}dt + \varepsilon D^{-}u_{N}^{N})' \\ &= (\sum_{j=i+1}^{N-1} \int_{I_{j}} (p(t) - p_{j})D^{+}u_{j}^{N} + (q_{j}u_{j}^{N} - q(t)u^{N}(t))dt \\ &+ \int_{x}^{x_{i+1}} (p(t) - p_{i})D^{+}u_{i}^{N} + (q_{i}u_{i}^{N} - q(t)u^{N}(t)) \\ &+ p_{i}D^{+}u_{i}^{N} - q_{i}u_{i}^{N} + f_{i}dt)'. \end{split}$$

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And owing to Lemma 1.1, (1) and (4), we have

$$\begin{split} \|u^{N}(x) - u(x)\|_{\infty} &\leq C \|Tu^{N}(x) - Tu(x)\|_{*} \\ &= C \|Tu^{N}(x) - f^{N}(x) + f^{N}(x) - f(x)\|_{*} \\ &\leq C \|\sum_{j=i+1}^{N-1} \int_{I_{j}} (p(t) - p_{j})D^{+}u_{j}^{N} + (q_{j}u_{j}^{N} - q(t)u^{N}(t))dt \\ &+ \int_{x}^{x_{i+1}} (p(t) - p_{i})D^{+}u_{i}^{N} + (q_{i}u_{i}^{N} - q(t)u^{N}(t)) + p_{i}D^{+}u_{i}^{N} - q_{i}u_{i}^{N} + f_{i}dt \\ &- \int_{x}^{1} f^{N} - fdt\|_{\infty} \\ &\leq C \max_{0 \leq i \leq N-1} \{\sum_{j=i}^{N-1} \int_{I_{j}} |(p(t) - p_{j})D^{+}u_{j}^{N}| + |(q(t) - q_{j})u_{j}^{N}| \\ &+ |q(t)(u^{N}(t) - u_{j}^{N})| + |f_{j} - f|dt + \int_{I_{i}} |p_{i}D^{+}u_{i}^{N} - q_{i}u_{i}^{N} + f_{i}|dt \}. \end{split}$$

Since the items of the right-hand side satisfy:

$$\begin{split} &\int_{I_j} |(p(t) - p_j) D^+ u_j^N | dt \leq C_p \int_{I_j} |(t - x_j) D^+ u_j^N | dt \leq C h_{j+1}^2 | D^+ u_j^N |, \\ &\int_{I_j} |(q(t) - q_j) u_j^N | dt \leq C \int_{I_j} |q(t) - q_j| dt \leq C C_q h_{j+1}^2, \\ &\int_{I_j} |q(t) (u^N(t) - u_j^N) | dt \leq \gamma^* \int_{I_j} |\int_{x_j}^t (u^N(s))' ds | dt \\ &\leq \gamma^* \int_{I_j} \int_{x_j}^t |D^+ u_j^N | ds dt \leq C h_{j+1}^2 | D^+ u_j^N |, \\ &\int_{I_j} |f_j - f(t)| dt \leq C h_{j+1}^2, \\ &\int_{I_i} |p_i D^+ u_i^N - q_i u_i^N + f_i | dt \leq \beta^* h_{i+1} | D^+ u_i^N | + C h_{i+1}, \end{split}$$

we have

$$\|u^{N}(x) - u(x)\|_{\infty} \le C \max_{0 \le i \le N-1} \{ \sum_{j=i}^{N-1} (h_{j+1}^{2} | D^{+} u_{j}^{N} | + h_{j+1}^{2}) + h_{i+1} | D^{+} u_{i}^{N} | + h_{i+1} \},$$

and then

$$\|u^{N}(x) - u(x)\|_{\infty} \le C \max_{0 \le i \le N-1} (h_{i+1}|D^{+}u_{i}^{N}| + h_{i+1}).$$
(8)

Theorem 2.1 Let u(x) be the solution of (1) and u_i^N be the solution of (5) on an arbitrary mesh. Then there exists a constant C such that:

$$||u_i^N - u(x_i)||_{\infty} \le C \max_{1 \le i \le N} h_i \sqrt{1 + (D^- u_i^N)^2}.$$

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Proof Let $u^N(x)$ be the piecewise linear interpolation of u_i^N . From (8), we have

$$\begin{aligned} \|u_i^N - u(x_i)\|_{\infty} &\leq \|u^N(x) - u(x)\|_{\infty} \leq C_1 \max_{1 \leq i \leq N} (h_i |D^- u_i^N| + h_i) \\ &\leq C \max_{1 \leq i \leq N} h_i \sqrt{1 + (D^- u_i^N)^2}. \end{aligned}$$

3 The adaptive upwind FDM and the existence of its solution

The problem (1) is solved iteratively by the adaptive upwind finite difference method:

- Step 1. Set an initial mesh $\{x_i^{(0)}\}$, e.g., a uniform mesh $\{0, \frac{1}{N}, \frac{2}{N}, ..., 1\}$ or a Shishkintype mesh.
- Step 2. For k = 0, 1, ..., compute the numerical solution $\{u_i^{(k)}\}$ by (5) on the current mesh $\{x_i^{(k)}\}$, and set $l_i^{(k)} = \sqrt{(x_i^{(k)} - x_{i-1}^{(k)})^2 + (u_i^{(k)} - u_{i-1}^{(k)})^2}$ and $L^{(k)} =$ $\sum_{j=1}^{N} l_i^{(k)}.$ Step 3. Take a user-chosen constant $C_0 \ge 1$. If the condition

$$\frac{\max_{1 \le i \le N} l_i^{(k)}}{L^{(k)}} \le \frac{C_0}{N}$$
(9)

is satisfied, then set $\{x_i^*\} = \{x_i^{(k)}\}, u^* = u^{(k)}$ and stop. Otherwise, continue to the next step.

Step 4. Set the next mesh $0 = x_0^{(k+1)} < x_1^{(k+1)} < ... < x_N^{(k+1)} = 1$ by equidistributing the arc-length of the polygonal curve of the numerical solution $u^{(k)}(x)$, such that

$$l_i^{(k+1)} = \sqrt{(x_i^{(k+1)} - x_{i-1}^{(k+1)})^2 + (u^{(k)}(x_i^{(k+1)}) - u^{(k)}(x_{i-1}^{(k+1)}))^2} = \frac{L^{(k)}}{N}$$

Return to step 2. Similar to the contents in [18, 21, 24-26], but the equation (1) is different from those in [18, 21, 26] and the scheme (5) is different from that in [18, 21, 24–26], the above adaptive method for (1) constructs a mapping, and its existence theorem also follows the fixed point theorem, which only demands the following newly proved preliminary lemma without using the presupposition $h_i/h_{i+1} \leq 1$ or C used in [21, 24–26].

Lemma 3.1 The solution $\{u_i^N\}$ by (5) for (1) on an arbitrary mesh $\{x_i\}$ satisfies

$$|D^{-}u_{i}^{N}| \leq C(N + \varepsilon^{-1})$$
 for $i = 1, 2, ..., N$,

where C a constant independent of ε and N.

Proof From (6), $|u_i^N| \le C$ for each *i*. So, $h_k \ge N^{-1}$ and $|D^-u_k^N| \le CN$ for some *k*. For i = k, k + 1, ..., N - 1, from (5) we have

$$\begin{split} &-\varepsilon(D^{+}u_{j}^{N}-D^{-}u_{j}^{N})-p_{j}h_{j+1}D^{+}u_{j}^{N}+q_{j}h_{j+1}u_{j}^{N}=h_{j+1}f_{j},\\ &(\varepsilon+p_{j}h_{j+1})D^{+}u_{j}^{N}=\varepsilon D^{-}u_{j}^{N}+q_{j}h_{j+1}u_{j}^{N}-h_{j+1}f_{j},\\ &\varepsilon|D^{+}u_{j}^{N}|\leq\varepsilon|D^{-}u_{j}^{N}|+Ch_{j+1},\\ &\varepsilon(|D^{+}u_{i}^{N}|-|D^{-}u_{k}^{N}|)\leq C\sum_{j=k}^{i}h_{j+1}\leq C,\\ &|D^{-}u_{i+1}^{N}|\leq|D^{-}u_{k}^{N}|+C\varepsilon^{-1}\leq C(N+\varepsilon^{-1}). \end{split}$$

For i = 1, 2, ..., k - 1, from (5) we have

$$\begin{split} \varepsilon(D^{-}u_{i}^{N}-D^{-}u_{k}^{N}) &= \sum_{j=i}^{k-1} \{p_{j}h_{j+1}D^{+}u_{j}^{N}-q_{j}h_{j+1}u_{j}^{N}+h_{j+1}f_{j}\} \\ &= p_{i}(u_{i+1}^{N}-u_{i}^{N})+\ldots+p_{k-1}(u_{k}^{N}-u_{k-1}^{N})-\sum_{j=i}^{k-1} \{q_{j}h_{j+1}u_{j}^{N}-h_{j+1}f_{j}\} \\ &= -p_{i}u_{i}^{N}+(p_{i}-p_{i+1})u_{i+1}^{N}+\ldots+(p_{k-2}-p_{k-1})u_{k-1}^{N}+p_{k-1}u_{k}^{N} \\ &-\sum_{j=i}^{k-1} \{q_{j}h_{j+1}u_{j}^{N}-h_{j+1}f_{j}\}, \\ &\varepsilon|D^{-}u_{i}^{N}| \leq \varepsilon|D^{-}u_{k}^{N}|+C+CC_{p}\sum_{j=i}^{k-2}h_{j+1}+C\sum_{j=i}^{k-1}h_{j+1}. \end{split}$$

We also have

$$\varepsilon |D^- u_i^N| \le \varepsilon CN + C$$
, i.e., $|D^- u_i^N| \le C(N + \varepsilon^{-1})$.

We need to prove that there exists a solution by (5) satisfying the equations:

$$l_i = \sqrt{(x_i - x_{i-1})^2 + (u_i - u_{i-1})^2} = \frac{\sum_{j=1}^N l_j}{N}, \ 1 \le i \le N.$$
(10)

Theorem 3.2 (Existence of a solution of the adaptive upwind FDM) For any $\varepsilon \in (0, 1)$ and every positive integer N, there exists a mesh equidistributing the arc-length (10) along the piecewise linear interpolation of the solution of (5).

Proof Let us regard steps 2 and 4 of the algorithm as a mapping: $\Psi : (h_1, h_2, ..., h_N) \rightarrow (\tilde{h}_1, \tilde{h}_2, ..., \tilde{h}_N)$, where h_i and \tilde{h}_i are of the mesh before and after moving it. Define a set

$$S_Q = \{(h_1, h_2, ..., h_N) \in \mathbf{R}^N : h_i \ge Q, 1 \le i \le N, \sum_i^N h_i = 1\}$$

where $Q = Q(\varepsilon, N)$ satisfies $0 < Q < \frac{1}{N}$. Clearly, S_Q is not empty. Let u_i^N be the numerical solution from (5) on the mesh with the step size h_i and set

 $l_i = h_i \sqrt{1 + |D^- u_i^N|^2}$ (i = 1, 2, ..., N). From Lemma 3.1 the slope of each segment of $u^N(x)$ is at most $C(N + \varepsilon^{-1})$, so the arc-length of $u^N(x)$ on each interval of the new mesh with length \tilde{h}_i for every *i* is at most $\tilde{h}_i \sqrt{1 + C(N + \varepsilon^{-1})^2}$. Since the new mesh is taken to split the piecewise linear function $u^N(x)$ into *N* pieces, we have

$$\tilde{h}_i \sqrt{1 + C(N + \varepsilon^{-1})^2} \ge \frac{\sum_{j=1}^N \tilde{h}_j}{N} \ge \frac{\sum_{j=1}^N \tilde{h}_j}{N} = \frac{1}{N},$$

which leads to

$$\tilde{h}_i \ge Q := \frac{1}{N\sqrt{1 + C(N + \varepsilon^{-1})^2}}, 1 \le i \le N.$$

Hence $0 < Q < \frac{1}{N}$ and then Ψ maps S_Q into itself.

The nonempty set S_Q is compact and convex, and Ψ is obviously continuous. So, Ψ has a fixed point in S_Q by the Brouwer fixed point theorem, i.e., there is a mesh, as well as a computed solution on it, such that $l_i = l_j$ for all i and j.

4 The uniform first-order convergence of the adaptive upwind FDM

Define the discrete Green's function $G(x_i, x_j)$ of T^N in (5) as follows:

$$T^N G_{i,j} = \delta^N_{i,j}, \quad i = 1, 2, \dots, N-1,$$
 (11)

and $G_{0,j} = G_{N,j} = 0$, for j = 1, 2, ..., N - 1, where

$$\delta_{i,j}^{N} = \begin{cases} \frac{1}{h_{i+1}}, & i = j, \\ 0, & \text{otherwise.} \end{cases}$$

Thus, the numerical solution u_i^N has the expression

$$u_i^N = \sum_{j=1}^{N-1} h_{j+1} G_{i,j} f_j, \qquad (12)$$

which satisfies (5).

Similar to the property of the discrete Green's functions in [1, 16, 18, 21, 24], we can prove the following property:

Lemma 4.1 The discrete Green's function $G(x_i, x_j)$ in (11) satisfies:

$$0 \le G_{i,j} \le \beta^{-1}$$
, for $i, j = 0, ..., N$,

and

$$\sum_{i=1}^{N} |G_{i,j} - G_{i-1,j}| \le \frac{2}{\beta}, \text{ for } j = 1, \dots, N-1.$$
(13)

Proof Let the barrier function for $G_{i,j}$ be (see, e.g., [1, 16, 18]):

$$B_{j}^{i} = \begin{cases} \beta^{-1}, & \text{for } i = 0, ..., j, \\ \beta^{-1} \prod_{k=j+1}^{i} (1 + \frac{\beta h_{k}}{\varepsilon})^{-1}, & \text{for } i = j+1, ..., N \end{cases}$$

For i = 1, ..., j - 1, we have

$$T^N B^i_j = q_i \beta^{-1} \ge 0.$$

For i = j, we have

$$T^{N}B_{j}^{j} = -\frac{\varepsilon}{h_{j+1}} \left(\frac{B_{j}^{j+1} - B_{j}^{j}}{h_{j+1}} - \frac{0}{h_{j}}\right) - p_{j}\frac{B_{j}^{j+1} - B_{j}^{j}}{h_{j+1}} + q_{j}B_{j}^{j}$$
$$= -\left(\frac{\varepsilon}{h_{j+1}} + p_{j}\right)\frac{B_{j}^{j}}{h_{j+1}}\left[\left(1 + \frac{\beta h_{j+1}}{\varepsilon}\right)^{-1} - 1\right] + q_{j}B_{j}^{j}$$
$$= \frac{\varepsilon + p_{j}h_{j+1}}{(\varepsilon + \beta h_{j+1})h_{j+1}} + q_{j}\beta^{-1} \ge \frac{1}{h_{j+1}}.$$

For i = j + 1, ..., N - 1, we have

$$T^{N}B_{j}^{i} = -\frac{\varepsilon}{h_{i+1}}\left(\frac{B_{j}^{i+1}-B_{j}^{i}}{h_{i+1}} - \frac{B_{j}^{i}-B_{j}^{i-1}}{h_{i}}\right) - p_{i}\frac{B_{j}^{i+1}-B_{j}^{i}}{h_{i+1}} + q_{i}B_{j}^{i}$$
$$= \frac{\varepsilon}{h_{i+1}}\left(\frac{\varepsilon}{\varepsilon+\beta h_{i+1}} - \frac{\beta}{\varepsilon}\right)B_{j}^{i} + p_{i}\frac{\beta}{\varepsilon+\beta h_{i+1}}B_{j}^{i} + q_{i}B_{j}^{i}$$
$$= \frac{(p_{i}-\beta)\beta}{\varepsilon+\beta h_{i+1}}B_{j}^{i} + q_{i}B_{j}^{i} \ge 0.$$

Hence

$$T^N B^i_j \ge \delta_{i,j} = T^N G_{i,j}.$$

Since $B_j^0 \ge G_{0,j}$ and $B_j^N \ge G_{N,j}$, the bounds in the following are established by the discrete comparison principle

$$0 \le G_{i,j} \le B_j^i \le \beta^{-1}$$
, for $i, j = 0, ..., N$.

As we know that $G_{i,j}$ increases for i = 0, 1, ..., j and $G_{i,j}$ decreases for i = j, j + 1, ..., N, we have

$$\sum_{i=1}^{N} |G_{i,j} - G_{i-1,j}| = 2G_{j,j} \le \frac{2}{\beta}, \text{ for } j = 1, \dots, N-1.$$

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Furthermore,

$$\mathcal{L}^{N} = \sum_{i=1}^{N} h_{i} \sqrt{1 + (D^{-} u_{i}^{N})^{2}} \le \sum_{i=1}^{N} h_{i} [1 + |D^{-} u_{i}^{N}|] = 1 + \sum_{i=1}^{N} |u_{i}^{N} - u_{i-1}^{N}|, \quad (14)$$

where \mathcal{L}^N is the arc length of discrete solution $\{u_i^N\}$. From (5), (12)-(14), we have

$$\mathcal{L}^{N} \leq 1 + \sum_{i=1}^{N} \sum_{\substack{j=1\\j=1}}^{N-1} h_{j+1} |f_{j}| \cdot |G_{i,j} - G_{i-1,j}|$$

$$\leq 1 + \|f\|_{\infty} \sum_{j=1}^{N-1} h_{j+1} \sum_{i=1}^{N} |G_{i,j} - G_{i-1,j}|$$

$$\leq 1 + \frac{2\|f\|_{\infty}}{\beta}.$$
 (15)

Theorem 4.2 The solution u_i^N of (5) on the grid $\{x_i\}$ of arc-length equidistribution satisfies:

$$||u_i^N - u(x_i)||_{\infty} \le CN^{-1},$$

where *C* is a constant independent of ε and *N*.

Proof According to Theorem 2.1, Theorem 3.2, (9) and (15), the solution u_i^N of (5) for (1) on the grid $\{x_i\}$ of arc-length equidistribution satisfies:

$$\|u_i^N - u(x_i)\|_{\infty} \le C_1 \max_{1 \le i \le N} h_i \sqrt{1 + (D^- u_i^N)^2} \le C_1 C_0 \mathcal{L}^N / N \le C N^{-1}.$$

5 Numerical examples

The numerical examples for the problem (1) are solved by the old upwind scheme (2) and new (5) respectively on the uniform mesh and the adaptive mesh. The maximum node errors are denoted as e_{old}^N and e_{new}^N , respectively. The numerical convergence orders are denoted as r_{old}^N and r_{new}^N , respectively. The number of iterations are denoted as *Iter*_{old} and *Iter*_{new}, respectively. The numerical convergence order is computed by

$$r^{N} = \log_2 \frac{\max |u_i - u_i^{N}|}{\max |u_i - u_i^{2N}|}.$$

Example 1 Consider the singularly perturbed two-point boundary value problem:

$$\begin{cases} -\varepsilon u''(x) - (1+x)u'(x) + u(x) = f(x), & x \in (0,1), \\ u(0) = u(1) = 0, \end{cases}$$

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N	16	32	64	128	256	512
$\varepsilon = 10^{-2}$						
Iter _{old}	4	6	6	4	3	3
e_{old}^N	1.1613e-01	7.5945e-02	4.6723e-02	2.6277e-02	1.4325e-02	7.6502e-03
r_{old}^N		0.6127	0.7008	0.8303	0.8752	0.9050
Iternew	7	3	3	2	2	2
e_{new}^N	5.4921e-02	3.2226e-02	1.7497e-02	1.1045e-02	5.0560e-03	2.2858e-03
r_{new}^N		0.7691	0.8811	0.6638	1.1273	1.1453
$\varepsilon = 10^{-5}$						
Iter _{old}	17	19	38	10	6	10
e_{old}^N	1.3675e-01	8.0317e-02	5.1264e-02	3.1368e-02	1.9866e-02	1.1102e-02
r_{old}^N		0.7677	0.6478	0.7087	0.6590	0.8395
Iternew	12	14	5	4	31	12
e_{new}^N	6.0147e-02	3.2131e-02	1.7137e-02	9.4281e-03	4.1500e-03	2.2783e-03
r_{new}^N		0.9045	0.9068	0.8621	1.1839	0.8651
$\varepsilon = 10^{-8}$						
Iter _{old}	15	18	16	18	20	35
e_{old}^N	1.3655e-01	8.0134e-02	5.1230e-02	3.1182e-02	1.8343e-02	1.0922e-02
r_{old}^N		0.7690	0.6454	0.7163	0.7655	0.7480
Iternew	15	9	10	6	5	21
e_{new}^N	6.1471e-02	3.3082e-02	1.6479e-02	8.6102e-03	4.7830e-03	2.2665e-03
r_{new}^N		0.8939	1.0054	0.9365	0.8481	1.0775

Table 1 The numerical results of Ex.1

where f(x) is chosen so that $u(x) = e^{-\frac{x}{\varepsilon}} - e^x + (e - e^{-\frac{1}{\varepsilon}})x$. The numerical results are shown in Table 1, Fig. 1 and Fig. 2.

Example 2 Consider the singularly perturbed two-point boundary value problem:

$$\begin{cases} -\varepsilon u''(x) - u'(x) + u(x) = 0, & x \in (0, 1), \\ u(0) = 0, & u(1) = 1. \end{cases}$$

The exact solution is $u(x) = \frac{\exp(m_1 x) - \exp(m_2 x)}{\exp(m_1) - \exp(m_2)}$, where $m_1 = \frac{-1 + \sqrt{1 + 4\varepsilon}}{2\varepsilon}$, $m_2 = \frac{-1 - \sqrt{1 + 4\varepsilon}}{2\varepsilon}$.

The numerical results are shown in Table 2 and Fig. 3.

Example 3 Consider the singularly perturbed two-point boundary value problem:

$$\begin{cases} -\varepsilon u''(x) - (1 + x(1 - x))u'(x) = f(x), & x \in (0, 1), \\ u(0) = u(1) = 0, \end{cases}$$

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Fig. 1 The solutions of Ex. 1 when N = 32 for $\varepsilon = 10^{-2}$



Fig. 2 The errors of Ex. 1 when N = 32 for $\varepsilon = 10^{-2}$

N	32	64	128	256	512	1024		
$\varepsilon = 10^{-2}$								
Iter _{old}	5	4	3	3	2	2		
e_{old}^N	3.4408e-02	2.0047e-02	1.0891e-02	5.7997e-03	3.1689e-03	1.5717e-03		
r_{old}^N		0.7794	0.8802	0.9091	0.8720	1.0117		
Iternew	4	3	3	2	2	2		
e_{new}^N	1.0738e-02	5.7549e-03	3.0807e-03	1.6893e-03	8.3495e-04	4.1298e-04		
r_{new}^N		0.8998	0.9015	0.8668	1.0167	1.0156		
$\varepsilon = 10^{-4}$								
Iter _{old}	17	11	17	10	7	6		
e_{old}^N	5.1225e-02	3.1884e-02	1.9278e-02	1.1334e-02	6.4437e-03	3.5708e-03		
r_{old}^N		0.6840	0.7259	0.7663	0.8147	0.8516		
Iternew	6	23	19	13	7	9		
e_{new}^N	1.3833e-02	6.7360e-03	3.6323e-03	1.8955e-03	9.7837e-04	4.9868e-04		
r_{new}^N		1.0382	0.8910	0.9383	0.9541	0.9723		
$\varepsilon = 10^{-6}$								
Iter _{old}	16	47	35	27	166	31		
e_{old}^N	5.1613e-02	3.2332e-02	1.9512e-02	1.1622e-02	6.8388e-03	3.8243e-03		
r_{old}^N		0.6747	0.7286	0.7475	0.7651	0.8385		
Iternew	9	18	16	28	13	23		
e_{new}^N	1.1767e-02	6.7545e-03	3.6409e-03	1.9009e-03	9.8197e-04	5.0078e-04		
r_{new}^N		0.8008	0.8915	0.9376	0.9530	0.9715		

Table 2 The numerical results of Ex. 2

where f(x) is chosen such that $u(x) = \frac{1-e^{-\frac{x}{\varepsilon}}}{1-e^{-\frac{1}{\varepsilon}}} - \cos(\frac{\pi}{2}(1-x))$. The numerical results are shown in Table 3 and Fig. 4.

From the rows of iteration numbers, the errors and the numerical convergence orders for each ε in the tables, and from the ordinates of errors in the figures, we can see that the two adaptive methods both achieve better accuracy when the mesh is finer and are close to the first-order convergence uniformly with respect to ε . Moreover, the proposed adaptive upwind finite difference method with new scheme (5) usually uses less iterations to satisfy the iteration tolerance and obtains better precision to resolve the boundary layers than the existing adaptive upwind finite difference method does.

6 Conclusion

In this work, the rigorous error analysis of the upwind finite difference scheme (5) on the grid that equidistributes the arc-length monitor function $M(x, u^N(x)) = \sqrt{1 + ((u^N(x))')^2}$ of the numerical solution u^N to solve the general non-conservative convection-reaction-diffusion problem (1) is finally accomplished through new efforts.



Fig. 3 The errors of Ex. 2 when N = 32 for $\varepsilon = 10^{-2}$



Fig. 4 The errors of Ex. 3 when N = 32 for $\varepsilon = 10^{-2}$

N	32	64	128	256	512	1024	
$\varepsilon = 10^{-2}$							
Iter _{old}	7	7	6	4	3	3	
e_{old}^N	6.3718e-02	4.0954e-02	2.3724e-02	1.3043e-02	6.9076e-03	3.5914e-03	
r_{old}^N		0.6377	0.7877	0.8630	0.9171	0.9436	
Iter _{new}	6	4	3	2	2	2	
e_{new}^N	4.5076e-02	2.3682e-02	1.2122e-02	7.0324e-03	3.2826e-03	1.5496e-03	
r_{new}^N		0.9286	0.9662	0.7855	1.0992	1.0830	
$\varepsilon = 10^{-5}$							
Iter _{old}	22	23	23	21	69	15	
e_{old}^N	6.5623e-02	4.2844e-02	2.6595e-02	1.5934e-02	9.4057e-03	5.9246e-03	
r_{old}^N		0.6151	0.6880	0.7390	0.7605	0.6668	
Iter _{new}	19	17	16	4	9	14	
e_{new}^N	5.5522e-02	2.8054e-02	1.4075e-02	7.3192e-03	3.5138e-03	1.8164e-03	
r_{new}^N		0.9849	0.9951	0.9434	1.0586	0.9520	
$\varepsilon = 10^{-8}$							
Iter _{old}	26	27	26	26	71	53	
e_{old}^N	6.5639e-02	4.2843e-02	2.6599e-02	1.5937e-02	9.3793e-03	5.9478e-03	
r_{old}^N		0.6155	0.6877	0.7390	0.7649	0.6571	
Iternew	20	19	17	18	59	23	
e_{new}^N	5.5577e-02	2.8057e-02	1.4082e-02	7.0495e-03	3.5168e-03	1.8209e-03	
r_{new}^N		0.9861	0.9945	0.9982	1.0033	0.9496	

Table 3 The numerical results of Ex.3

By contrast, considering adaptive methods, M(x, u(x)) was of the exact solution u in [11–16], the equation was of conservative form in [17–20], of reaction-diffusion form with no first derivative term in [16, 27, 28, 31–33] and of convection-diffusion form with no zero derivative term in [11–15, 21], the a priori truncation error was from a priori knowledge of the exact solution in [22, 23], and the condition $h_i/h_{i+1} \leq 1$ or C was used in [21, 24–26]. In fact, the a posteriori error estimate of the upwind finite difference scheme, the existence of the solution of the adaptive upwind finite difference method, and its ε -independent first-order convergence are developed step by step with new progresses in this paper. The numerical examples confirm its uniform first-order convergence and show its advantage in precision, too. The works for solving nonlinear singularly perturbed problems, singularly perturbed mixed BVPs, problems having interior layers, systems of singularly perturbed BVPs, singularly perturbed two-dimensional elliptic problems, fractional order integro-parabolic equations and so on are expected to be extended for future study (see e.g. [27–39]).

Acknowledgements The authors thank the anonymous reviewers for their valuable comments which help to improve the presentation. This paper is supported by National Natural Science Foundation (Grant No. 11471019) of China.

Declarations

Conflict of interest statement The authors declare that they have no conflict of interests.

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