

Gated trapped ion mobility spectrometry coupled to fourier transform ion cyclotron resonance mass spectrometry

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Abstract Analysis of molecules by ion mobility spectrometry coupled with mass spectrometry (IMS-MS) provides chemical information on the three dimensional structure and mass of the molecules. The coupling of ion mobility to trapping mass spectrometers has historically been challenging due to the large differences in analysis time between the two devices. In this paper we present a modification of the trapped ion mobility (TIMS) analysis scheme termed “Gated TIMS” that allows efficient coupling to a Fourier Transform Ion Cyclotron Resonance (FT-ICR) analyzer. Analyses of standard compounds and the influence of source conditions on the TIMS distributions produced by ion mobility spectra of labile ubiquitin protein ions are presented. Ion mobility resolving powers up to 100 are observed. Measured collisional cross sections of ubiquitin ions are in excellent qualitative and quantitative agreement to previous measurements. Gated TIMS FT-ICR produces results comparable to those acquired using TIMS/time-of-flight MS instrument platforms as well as numerous drift tube IMS-MS studies published in the literature.

Keywords Fourier transform ion cyclotron resonance mass spectrometry · Trapped ion mobility spectrometry · High resolution · Ubiquitin

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Introduction

When ion mobility spectrometry (IMS) was developed in the 1960s [1], it was first used to study gas phase ion chemistry. Whereas standalone IMS systems have been used primarily as a chemical sensors in homeland security applications [2], the coupling of IMS to mass spectrometry (MS) creates a versatile instrument that provides information on both the size and the molecular weight of the analyzed species [3]. Accordingly, IMS-MS has been more recently used in a wide variety of chemical [4–9], physical [10–14], and biological applications [15–33].

Ion mobility spectrometry was initially performed in a drift tube. As a dispersive ion mobility technique, ion separation in drift tube IMS involves guiding ions with a weak electric field through a tube filled with neutral buffer gas where separation is based on both size and charge. Though drift tubes have the advantage that they can be used to determine mobility based on first principles, drift tubes suffer from radial ion losses due to the thermal diffusion and Coulomb repulsion. This principally limits IMS sensitivity when coupled to MS.

The ion funnel was developed in the 1990s [34] to improve the transfer of ions generated from soft ionization sources, such as electrospray ionization (ESI), as they passed from high- to low-pressure regions, and ultimately into the vacuum of the mass analyzer. Ion funnels operate by focusing ions through a series of stacked ring electrodes of decreasing inner diameter. A radiofrequency potential is applied 180° out of phase to adjacent electrodes, thus creating a radially confining pseudopotential [34]. The focusing ability of the ion funnel has been successfully implemented in drift tube instruments, to recapture radially diffuse ion clouds after mobility separation, for efficient transfer into the mass spectrometer. Placement of an ion funnel between the drift tube and the mass spectrometer significantly improves the sensitivity of the

hybrid technique by focusing the radially diffuse ion packets through conductance limiting apertures [35, 36].

Recently, the technique of Trapped Ion Mobility Spectrometry (TIMS) was developed by Park and coworkers [37], wherein ion mobility separation is performed entirely inside an ion funnel. Inverse to drift tube IMS, TIMS operates by passing gas through swarms of stationary ions that are repelled by a repulsive electric field gradient. Advantages of TIMS include: (1) high transmission efficiencies in MS and IMS-MS modes, (2) compact design allowing integration into many instrument platforms, (3) determination of reduced mobility or collisional cross section from a simple calibration [38, 39], (4) low operating potentials, (5) flexibility to alter ion mobility resolving power and duty cycle, and (6) superior mobility resolving power that can exceed 250 [37, 39, 40]. Recently, TIMS has been shown to produce mobility distributions for proteins and peptides (whose conformations are known to depend on a number of physicochemical parameters) that are virtually identical to those measured by drift tube IMS operating below the low field limit where ion temperatures are known to be near thermal [39, 40].

One of the challenges in coupling dispersive IMS to MS is the need for a mass analyzer with a spectral acquisition rate sufficiently fast to sample the mobility distribution. Typically, time-of-flight (TOF) mass analyzers are employed, owing to their ability to efficiently sample mobility-separated ions on the millisecond timescale. However, a few IMS-MS instruments have dispersive IMS techniques coupled to a trapping mass analyzer [41–43]. Alternatively, approaches have been designed to decrease the dispersive IMS timescale, but none of these experimental techniques have been applied to hybridization with trapping MS analyzers [44, 45]. IMS coupling to Fourier Transform Ion Cyclotron Resonance mass spectrometry (FT-ICR MS) is particularly challenging due to its relatively slow spectral acquisition rate (~1 s per spectrum).

To overcome this principal limitation, TIMS was first coupled to an FT-ICR analyzer using a modified electric field gradient that allowed for increased ion accumulation times. This embodiment, termed Selective Accumulation-TIMS (SA-TIMS), has been successfully used to analyze complex mixtures. The combination of ion mobility and accurate mass measurement has been used to identify compounds from complex mixtures, in applications ranging from hydrocarbon analysis in petroleum to the identification of endocrine disrupters from complex environmental matrices [46, 47].

Since the inception of TIMS, a number of significant improvements have been recently implemented. The application of “parallel accumulation” allows a near 100 % duty cycle [48], modification of the field curvature of the TIMS trapping field allows for improved ion storage capacity and resolving power [49], and a fundamental theory for TIMS operation has been developed and validated by simulations and experiments [50, 51]. These improvements cannot be directly applied to

SA-TIMS due to differences in the experimental method, and results produced by SA-TIMS cannot be easily compared with results produced by TIMS-TOF MS. With these recent advancements in TIMS and with the desire to align TIMS analyzers across all MS platforms, we evaluate here a new operational scheme for TIMS, termed Gated TIMS, as an alternative method for coupling to FT-ICR MS.

Materials and methods

Ubiquitin (Sigma-Aldrich, St. Louis, MO) was diluted to a concentration of 1 μ M in 70:30:0.1 water:acetonitrile:formic acid. Low concentration Tune Mix (Agilent Technologies, Santa Clara, CA), a series of perfluoroalkylphosphazenes, was used without dilution. Samples were ionized via static nanospray using a glass capillary pulled in-house. The acquisition parameters for the TIMS and FTMS instruments described below are listed in the [Appendix](#).

Experiments were performed on a modified 12T FT-ICR MS (solariX, Bruker Daltonics, Bremen, Germany). The production level ion funnel that conventionally occupies the first pumping region of the instrument vacuum system was replaced with a TIMS funnel. This design incorporates an entrance funnel, a tunnel, and an exit funnel. A schematic of the relevant parts of the experimental setup and the various steps in its operation are shown in [Fig. 1](#).

[Figure 1a](#) is a cross-sectional view of the TIMS funnel. Ions are produced in the ESI source via static nanospray. Ions entrained in nitrogen gas are introduced into the first pumping region via the capillary exit. As shown in [Fig. 1a](#), the gas flow is directed across the inlet of the entrance funnel. Application of a repulsive potential to the deflection plate deflects ions out of the gas stream, into the entrance funnel, and ultimately into the TIMS tunnel. In MS-only operation, ions are allowed to pass through the TIMS tunnel and are simply transmitted downstream to the mass analyzer.

However, for TIMS-MS experiments, ions are accumulated in the TIMS tunnel. Application of an RF potential (~250 Vpp, 900 kHz) to the TIMS funnel electrodes confines the ions radially throughout the TIMS funnel. DC potentials applied to the funnel and tunnel electrodes generate the field that pushes ions along the funnel-tunnel axis. The strength of this DC electric field as a function of position in the TIMS funnel at various points in the Gated TIMS FT-ICR MS experiment is depicted in [Fig. 1b–e](#). In these plots, a negative value indicates the field is pushing ions downstream, towards the mass analyzer, whereas a positive value indicates it is pushing the ions upstream.

To perform a TIMS experiment, some of the nitrogen from the capillary exit is diverted to flow through the TIMS tunnel and is then pumped away laterally in the exit funnel region. The gas flow pushes ions in the tunnel downstream. To

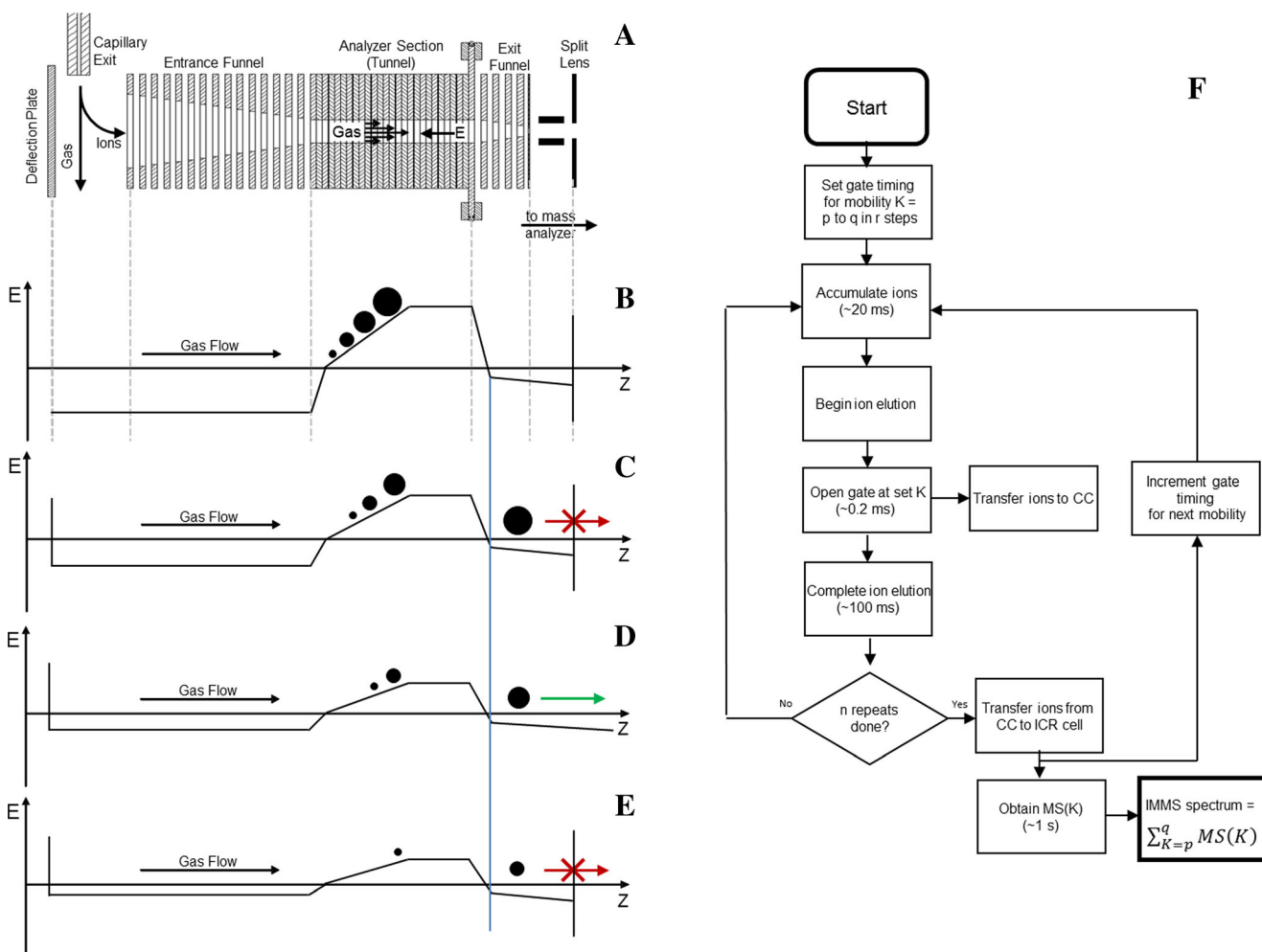


Fig. 1 Schematic of the TIMS device and experiment timing, with ion motion moving left to right along the z-axis. **a** Cross section of the TIMS analyzer with relevant ion optics labeled. **b** Beginning of a TIMS experiment where ions are accumulated in the TIMS analyzer and undergo mobility separation. **c** Beginning of ion elution out of the

accumulate ions in the tunnel, the DC electric field is set as depicted in Fig. 1b. In such a case, ions are deflected into the entrance funnel. The ions are directed further through the entrance funnel and into the TIMS tunnel by the DC field. Finally, the ions are trapped in the tunnel radially by the confining RF field, and axially by the balancing forces exerted by the gas flow and the DC field.

As shown in Fig. 1b, the axial electric field in the tunnel increases in strength with position, with the slope of the electric field set by a resistor-divider chain on the ion funnel. This allows for accumulation of ions having a range of mobilities. Ions of low mobility, depicted as a large circle, are pushed downstream more strongly by the gas and therefore accumulate at a position where the retarding DC field is greater. On the other hand, ions of high mobility, depicted as a small circle, come to rest at a position where the retarding DC field is weaker.

Once a group of ions has been accumulated, the potential on the deflection plate is set to an attractive potential, as

TIMS analyzer, ions cannot pass to the collision cell due to blocking by the gate lens. **d** Ions continue to elute, ions with the desired mobility are allowed to pass the gate lens. **e** Ions continue to elute and are blocked from entering the collision cell by the gate lens. **f** Timing flow chart of the Gated TIMS experiment

indicated in Fig. 1c–e, such that no additional ions from the ESI source can enter the TIMS analyzer. The group of ions that are trapped in the TIMS tunnel can then be eluted. As depicted in Fig. 1b–e, this is done by decreasing the strength of the electric field gradient profile over time. Because the low mobility ions are pushed most strongly by the gas flow, these ions will be the first to overcome the electric field barrier represented by the plateau (Fig. 1c). As the amplitude of the electric field gradient profile continues to be lowered, moderate (Fig. 1d) and then high mobility ions (Fig. 1e) will overcome the barrier and elute.

Unlike previous TIMS analyzers and methods, the experimental setup in the present work includes an ion gate downstream of the TIMS tunnel (see Fig. 1a). This allows for the elimination/selection of ions, based on mobility, as they elute from the TIMS analyzer. All ions outside a selected range or ranges of mobility are eliminated by deflection, followed by collision with the gate electrodes. The gate is comprised of a

metal plate that has been split into two halves—i.e. a “split lens”. To transmit ions, the two halves are biased to the same DC potential such that ions are focused into the downstream optics (Fig. 1d). To block the transmission of ions, the two halves are biased at opposite polarities such that the ions are deflected orthogonal to the ion optical beam path resulting in their collision with the split lens elements (Fig. 1c and e). Tests of the split lens gate showed that it can select/deselect ions with a time resolution of better than 0.2 ms.

The incorporation of an ion gate as described above allows for the hybridization of the TIMS analyzer with slower mass analyzers—i.e. mass analyzers that otherwise would not be able to keep up with the rapid elution of ions from the TIMS analyzer. The method associated with the hybridization of gated TIMS with an FT-ICR mass analyzer is further detailed in Fig. 1f. In the experiment exemplified here, the sample is analyzed over the mobility range from mobility p to mobility q in r steps. The result is a set of r FT-ICR mass spectra, each of which corresponds to a segment of the mobility range p to q . These mass spectra, when taken together, constitute a multidimensional IMMS spectrum of the original sample.

As discussed above, the individual gated TIMS experiment consists of the steps of ion “accumulation” (Fig. 1b), and “elution” (Fig. 1c–e) with gating of the ions. The ions selected in the elution/gating step are transmitted to, and collected in, a collision cell an RF driven hexapole, operated at a pressure of about 10^{-3} mbar of N_2 . As indicated in Fig. 1f, an individual gated TIMS cycle lasts on the order of 100 ms, however, the FT-ICR mass analysis may require as long as 1 s. This allows for the performance of several gated TIMS cycles per mass analysis. Thus, as indicated in Fig. 1f, the gated TIMS cycle is repeated n times, with the resultant ions of selected mobility/mobilities accumulated all together in the collision cell. In the experiments performed here, all of the ions accumulated in the collision cell via n experiments are of the same mobility, K .

At the end of n gated TIMS cycles, the contents of the collision cell are transferred to the ICR cell for m/z analysis. While the FT-ICR mass analyzer is producing a mass spectrum of these ions, the timing of the gated TIMS is incremented to the next mobility in the range and the next group of ions is accumulated in the collision cell. In this manner the TIMS gate and associated timing are incremented through the mobility range—from p to q —in r steps producing a FT-ICR mass spectrum at each step. All the resulting mass spectra are combined to form the multidimensional IMMS spectrum.

Results and discussion

The Gated TIMS device was initially characterized by performing ion mobility separations on Agilent tune mix, as

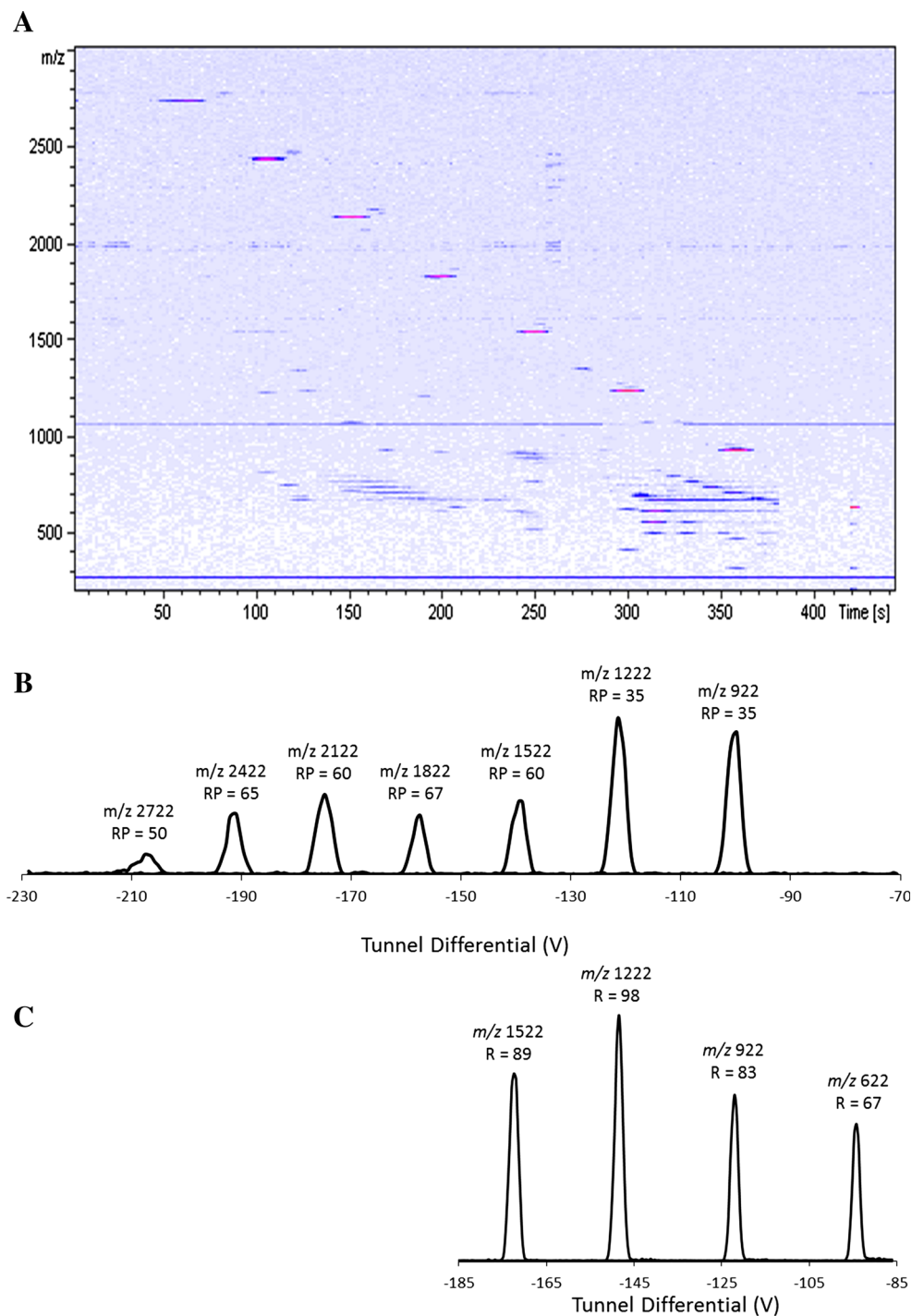
shown in Fig. 2. The heatmap (Fig. 2a) shows baseline separation in both m/z space (y-axis) and mobility space (x-axis) for ions m/z 922– m/z 2722. Extracted ion “mobiligrams” for the tune mix ions are shown in Fig. 2b. Resolving powers around 60 are observed for ions over much of this mobility range, as shown above each peak in the figure.

Unique to the TIMS method is the flexibility to adjust mobility resolving power, mobility range, and the speed of the experiment, according to the application. This is done, in part, by controlling the potential range over which the tunnel DC is scanned and the rate at which it is scanned. For the data presented in Fig. 2a and b, the potential on the tunnel was scanned quickly (3 V/ms) over a broad range of mobilities, producing a spectrum having a relatively modest resolving power. In contrast, the spectrum shown in Fig. 2c was produced by scanning the tunnel voltage more slowly (0.96 V/ms) and over a narrower voltage range. As a result, the observed resolving powers were significantly higher: 80 to near 100 for ions with m/z 622–1522. Achieving higher mobility resolution ion mobility spectra requires not only slow elution from the TIMS analyzer but also narrower mobility steps—i.e., more “ r ”—in a given range (see Fig. 1f)—and therefore a longer time spent acquiring the IMMS spectrum. This is exemplified in the experiment of Fig. 2, where the low resolving power data in Fig. 2b were acquired in approximately 450 s whereas the higher resolving power data in Fig. 2c were acquired in approximately 1650 s.

As described previously with respect to TIMS-TOF MS results [40], data like those in Fig. 2 can be used to calibrate the TIMS analyzer via the simple relation— $K \propto 1/V$. Using the known mobility for each component of tune mix and the data in Fig. 2b, one can calibrate the analyzer for reduced mobility. As with previous TIMS-TOF MS results, a plot of the reduced mobilities of the tune mix species vs. the inverse of their elution voltages revealed a highly linear relationship ($R^2=0.9999$, data not shown). The resultant calibration was subsequently used as an external calibration for the ubiquitin analysis (below). The mobilities measured for ubiquitin ions were then readily converted to collisional cross sections, resulting in the values reported in Figs. 3 and 4.

Ions of the 7+ charge state of ubiquitin display (m/z 1224) a number of different conformers (compact, partially folded, and elongated) that have been extensively studied by ion mobility [52, 53]. In the present work, mobiligrams of the 7+ charge state were acquired under varying experimental conditions using Gated TIMS. These results were then compared to data previously acquired by TIMS-TOF MS (Fig. 3). The mobility spectra shown in Fig. 3a and e were obtained from a prototype TIMS-TOF MS instrument whereas the data in Fig. 3b–d were obtained using the Gated TIMS FT-ICR instrument described above. The plots in Fig. 3 are shown with the x-axis in Å^2 and the cross-section values were calculated from the tune mix calibration.

Fig. 2 **a** Heatmap of m/z vs. elution time from a low ion mobility resolution analysis of Agilent Tune Mix. **b** Extracted ion mobiligrams of each component of tune mix with their m/z and resolving power. **c** High resolution extracted ion mobiligrams from the high ion mobility resolution analysis of the Agilent Tune Mix



Mobiligrams acquired using similar TIMS conditions with m/z analysis by TOF [40] and FT-ICR MS analysis of the 7+ charge state are shown in Fig. 3a and b, respectively. Not surprisingly, these two data sets exhibit similar features. For example, little or no compact conformers, a small amount of partially folded species, and abundant elongated forms are observed in all cases. Also, three abundant peaks corresponding to elongated forms are observed between 1880 and 1950 Å², and an abundant peak is observed at 1800 Å². The

elongated form appears more abundant in the Gated TIMS data than in the TIMS-TOF MS data, but this may be due to an increase in the Funnel 1 DC gradient (discussed below). Overall, very similar mobiligrams were generated from two different mass spectrometers equipped with the same TIMS analyzer, demonstrating the cross-platform capabilities of the TIMS analyzer.

The 7+ charge state of ubiquitin is also ideal for studying the influence of source conditions on the observed

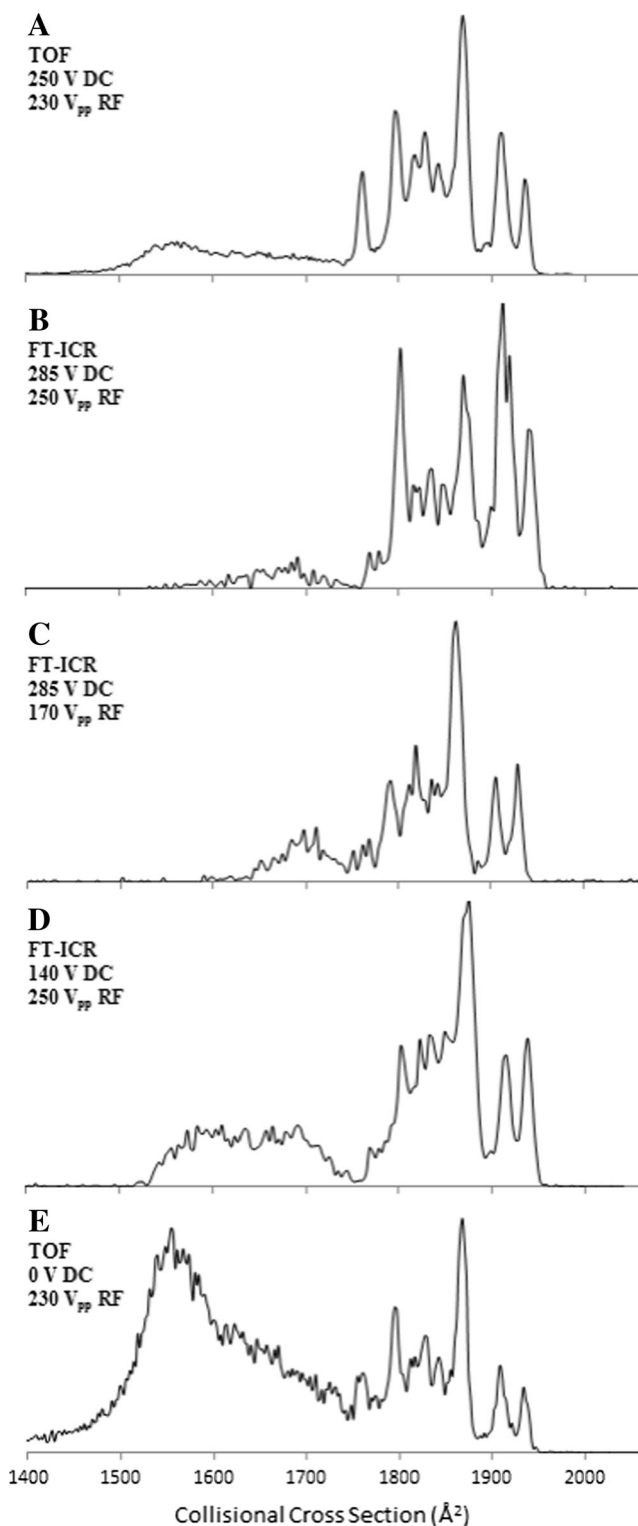


Fig. 3 Conformational distribution of the 7+ charge state of ubiquitin showing, partially folded, and elongated structures as a function of voltage applied across the entrance funnel and the funnel RF amplitude. Plots A and E were generated on a TIMS-TOF MS and reproduced with permission by the Royal Society of Chemistry from [40], plots B–D were generated with Gated TIMS-FT-ICR MS

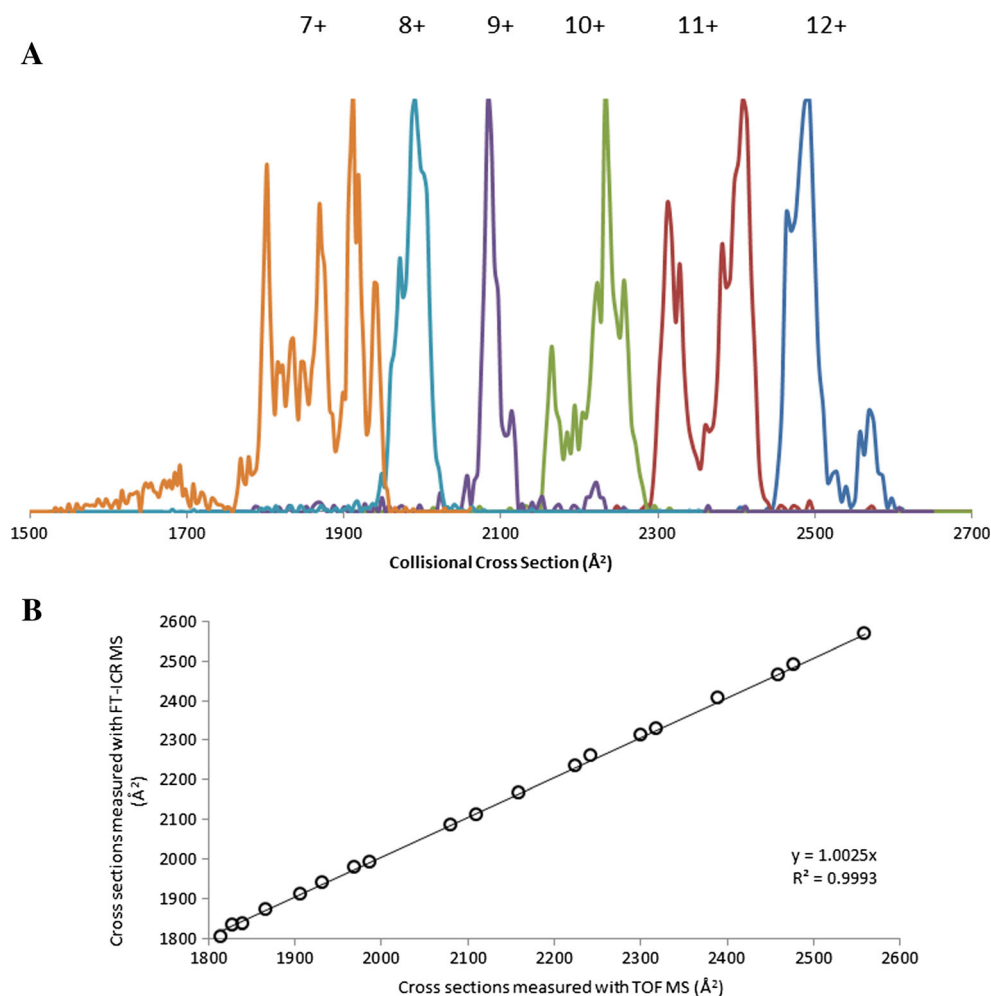
conformations, as it is sensitive to heating and will undergo conformational changes from compact to elongated forms if the source or analyzer is too energetic [54, 55]. For example, ion mobility devices employing RF have the potential to heat ions due to energetic collisions with the background gas, causing annealing of protein ions to gas phase structures [40, 56]. To explore the influence of the RF and DC potentials of Gated TIMS on the conformers of the 7+ charge state, mobiligrams were acquired as the RF amplitudes on the tunnel and funnel and the DC potential difference between the ion funnel 1 entrance and exit lenses were varied, as shown in Fig. 3b–d. The potentials of these two components of the TIMS analyzer have been shown to greatly influence the measured mobiligram [40].

Figure 3 is plotted with conditions being most energetic at Fig. 3a and least energetic at Fig. 3e. Figure 3b and c show the conformational distribution of the 7+ ions of ubiquitin measured with Gated TIMS with the peak-to-peak RF on the ion funnel set at 250 and 170 Vpp, respectively. As shown, decreasing the RF has some small influence on the conformations. For example, the signal height in the partially folded region (1650–1750 Å²) is greater and the signal height in the elongated region (1800–1950 Å²) is reduced with lower RF (Fig. 3c), suggesting slight heating effects due to the applied RF.

Comparing Fig. 3b and d, the DC potential difference across ion funnel 1 has a more significant influence on the observed gas phase structures. For example, when low DC potentials are placed across ion funnel 1 (Fig. 3d), an abundant partially folded state of the 7+ ions is observed around 1500–1700 Å². As the voltage gradient across ion funnel 1 is increased (Fig. 3b), evidence for a significant conversion of partially folded to elongated conformations is observed as a decrease in abundance in the 1500–1700 Å² region, suggesting heating of the ions through collisions with the background gas as they pass through ion funnel 1 and enter the tunnel. Similar behavior is observed in TIMS-TOF measurements, as exemplified in Fig. 3a and e. Decreasing the voltage gradient across ion funnel 1, from 250 to 0 V causes an increase of the partially folded form (1500–1700 Å²) relative to the elongated form (1750–1950 Å²). Overall, similar behaviors with respect to both the RF and DC potentials have been observed in the TIMS-TOF MS and Gated TIMS FT-ICR MS analyses of the 7+ ions of ubiquitin [40], demonstrating the tuning capabilities and performance similarities of these TIMS analyzers.

A compilation of the *m/z* selected Gated TIMS mobiligrams of the observed charge states of ubiquitin is shown in Fig. 4a. Baseline separation of nearly all the charge states is observed, and the data show a direct correlation between the ion charge state and its collisional cross section. Specifically, the collisional cross section increases as the charge state increases, suggesting the preference for the

Fig. 4 **a** TIMS distributions for selected m/z values in the ESI mass spectra of ubiquitin. Patterns are the same for FT-ICR MS and TOF MS. **b** Comparison of measured collisional cross sections of the most abundant peaks for all charge states in the ubiquitin mass spectra generated by Gated TIMS-FT-ICR MS and TIMS-TOF MS measurements



formation of elongated structures [57]. These extracted ion mobiligrams are comparable to what has been previously measured by TIMS-TOF MS. Largely the same number of conformations, as well as the same cross sections, are observed for any given ion charge state by Gated TIMS FT-ICR as by TIMS-TOF MS.

A comparison of the abundant peaks observed in TIMS-TOF [40] to the Gated TIMS FT-ICR data is made in the form of a plot in Fig. 4b. Here the measured collisional cross sections of the most abundant peaks for each charge state generated from TIMS-TOF MS are plotted on the x-axis, and the corresponding results from cross section measurements generated by Gated TIMS FT-ICR MS are plotted on the y-axis. A perfect match between these data sets would generate a line with a slope of one. The data of Fig. 4 form a line with a slope of 1.0025 and an R^2 value of 0.9993, showing excellent correlation between the two sets of data. The average error between Gated TIMS-FT ICR and TIMS-TOF collision cross sections is 0.3 % (about 7\AA^2) which is similar in magnitude to the error associated with day-to-day stability of TIMS-TOF instrument platforms (0.6 %, data not shown).

Conclusions

In the present work we have detailed and demonstrated the experimental setup and operation of Gated TIMS in conjunction with an FT-ICR mass analyzer. Overall, the results are consistent with the expectation that Gated TIMS operates on the same principles as TIMS without a gate. In particular with Gated TIMS FT-ICR, we observed a linear relationship between mobility and $1/V$ ($R^2=0.9999$). This relationship was used to calibrate the Gated TIMS analyzer and then to determine collisional cross sections of ubiquitin ions. Importantly, the observed mobility resolutions (up to about 100) and their dependence on various instrument parameters, for example the scan speed, are also consistent with basic TIMS theory. The measurement of ubiquitin conformers and their cross sections show excellent agreement with previous results ($R^2=0.9993$). Gated TIMS with FT-ICR MS produced spectra showing ubiquitin conformations with cross sections within about 0.3 % on average to those measured by TIMS-TOF MS. The combination of high mobility resolution Gated TIMS and high mass resolution FT-ICR MS offers exceptional prospects for the analysis of complex mixtures such as petroleum,

carbohydrates, lipids, and a variety of other biological samples.

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Appendix

Table 1 Acquisition parameters

	Low resolution tune mix	High resolution tune mix	High resolution ubiquitin
Pressure at funnel entrance	2.6 mbar	3.1 mbar	3.1 mbar
Pressure in P2	1.4 mbar	1.4 mbar	1.4 mbar
TIMS start voltage (V)	−200	−195	−160
TIMS end voltage (V)	−50	−75	−80
TIMS ramp time (ms)	50	125	125
TIMS ramp speed (V/ms)	3	0.96	0.64
TIMS ion accumulation time	10	25	50
Number of TIMS accumulations in collision cell	20	40	40
TIMS gate width (ms)	0.5	1	1
TIMS gate shift (ms)	0.2	0.5	0.5
Number of averaged mass spectra	250	250	250
Total experiment time (s)	450	1650	1900

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