



STAR FORMATION

Statistical analysis of dynamical evolution of open clusters

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Abstract. We present the dynamical evolution of 10 open clusters, which were part of our previous studies. These clusters include both young and intermediate-age open clusters with ages ranging from 25 ± 19 Myr to 1.78 ± 0.20 Gyr. The total mass of these clusters ranges from 356.18 ± 142.90 to $1811.75 \pm 901.03 M_{\odot}$. The Galactocentric distances to the clusters are in the range of 8.91 ± 0.02 – 11.74 ± 0.18 kpc. The study is based on the ground-based UBVR data supplemented by the astrometric data from the Gaia archive. We studied the minimum spanning tree of the member stars for these clusters. The mass segregation in these clusters was quantified by mass segregation ratios calculated from the mean edge length obtained through the minimum spanning tree. The clusters NGC 2360, NGC 1960, IC 1442, King 21 and SAI 35 have Γ_{MSR} to be 1.65 ± 0.18 , 1.94 ± 0.22 , 2.21 ± 0.20 , 1.84 ± 0.23 and 1.96 ± 0.25 , respectively, which indicate moderate mass segregation in these clusters. The remaining five clusters are found to exhibit weak or no mass segregation. We used the ratio of half mass radius to the tidal radius i.e., R_h/R_t to investigate the effect of the tidal interactions on the cluster structure and dynamics. The ratios of half mass radii to tidal radii are found to be positively correlated with the Galactocentric distances with a linear slope of 0.06 ± 0.01 having linear regression coefficient r -square = 0.93 for the clusters.

Keywords. Open clusters—mass segregation—tidal interactions.

1. Introduction

The star clusters evolve dynamically as the age passes, after their birth from molecular clouds. Open clusters go through two-body relaxation and their core approaches spherical structure due to the dynamical relaxation. This dynamical relaxation causes low-mass stars to preferentially settle in the outer part of the clusters. These outskirt stars are subjected to dynamic ejection from the cluster, which lowers the gravitational potential and loosens the cluster (Gieles & Baumgardt 2008; Tang *et al.* 2019). The external perturbations like tidal effects, disc crossing and differential rotation also strengthen the disintegration process of open clusters. These internal and external factors affect the dynamical evolution of the clusters, which manifests in the morphology,

shape and spatial distribution of the stars of the clusters. Thus, open clusters with accurately determined physical parameters are useful in understanding the impact of the internal and external effects on the dynamical evolution of the cluster.

The spatial distribution of stars in the clusters has been found to be dependent on the stellar masses. The massive stars are found to be preferentially concentrated in the inner region of the clusters compared to low-mass stars (Dib *et al.* 2018; Joshi *et al.* 2020a; Maurya *et al.* 2021). This mass segregation phenomenon has been explained by two theories that attribute mass segregation either to the star formation process itself or dynamical evolution. The theory considering the star formation process responsible for the mass segregation proposes that massive stars are preferentially formed in the inner region of cluster (Dib *et al.* 2008). However, the dynamical evolution theory of mass segregation suggests that the segregation happens due to an internal

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Table 1. Derived values of physical parameters of the clusters. The name, right ascension, declination, completeness limit in the V band, mass of the most massive star, mass completeness limit of the data, number of member stars up to the completeness limit of the data, reddening, logarithmic age and distance calculated using the methods described by [Bailer-Jones et al. \(2018\)](#) are given in columns 1–10, respectively.

Cluster	RA	Dec	V_{lim} (mag)	M_{upper} (M_{\odot})	M_{lower} (M_{\odot})	N	$E(B-V)$ (mag)	$\log(\text{age})$ (Myr)	D_{BJ} (pc)
NGC 381	01:08:19.6	+61:35:18.2	20	2.80	0.61	134	0.36 ± 0.04	8.65 ± 0.05	1147 ± 38
NGC 2360	07:17:43.5	−15:38:39.8	18	2.21	0.68	276	0.07 ± 0.03	8.95 ± 0.05	1072 ± 44
Berkeley 68	04:44:30.0	+42:05:55.8	19	1.75	0.97	229	0.52 ± 0.04	9.25 ± 0.05	3206 ± 199
NGC 1960	05:36:20.2	+34:08:06	19	7.18	0.72	253	0.24 ± 0.02	7.44 ± 0.02	1169 ± 54
IC 1442	22:16:03.7	+53:59:29.4	18	9.02	1.43	205	0.54 ± 0.04	7.40 ± 0.30	3492 ± 230
King 21	23:49:55.0	+62:42:18.0	19	6.89	1.21	238	0.76 ± 0.06	7.70 ± 0.20	2953 ± 174
Trumpler 7	07:27:23.8	−23:56:56.4	19	5.12	0.76	146	0.38 ± 0.04	7.85 ± 0.25	1700 ± 98
SAI 35	04:10:46.8	+46:52:33.2	19	3.34	1.02	156	0.61 ± 0.04	8.50 ± 0.10	2826 ± 266
SAI 44	05:11:10.5	+45:42:10.2	19	2.47	1.00	171	0.34 ± 0.04	8.82 ± 0.10	3670 ± 184
SAI 45	05:16:29.4	+45:35:35.9	19	2.14	0.79	79	0.34 ± 0.02	9.07 ± 0.10	1668 ± 47

two-body dynamical relaxation process ([Allison et al. 2009a](#)). The debate on the origin of mass segregation has not concluded and requires further analysis ([Dib et al. 2018](#)).

The comprehensive study of the cluster parameters like core radii, tidal radii, half-mass radii and ages helps in understanding the complete scenario of the interplay among structural evolution, dynamical evolution and tidal interactions. Open clusters generally become more centrally concentrated due to dynamical evolution, which makes mass loss less likely, which is caused by tidal interactions in these clusters. In this two-body dynamical relaxation, massive stars sink towards the central region, while low-mass stars preferentially shift towards the outer region, which evaporate from the cluster with passing time.

Impact of the tidal field on the cluster size can be characterized by using the ratio of half-mass radius to the tidal radius of the cluster i.e., R_h/R_t . Previously, it has been found that the R_h/R_t ratio parameter is also related to the survival of the cluster as it has been found that the more compact clusters have smaller R_h/R_t values, which favor the survival of the clusters located at smaller Galactocentric distances by reducing the mass loss due to tidal fields ([Angelo et al. 2021](#)).

To probe the role of dynamical evolution in defining the shape and morphology of the clusters, we present the analysis of the dynamical evolution of the 10 open clusters. These clusters are parts of our previous studies ([Joshi et al. 2020a](#); [Maurya & Joshi 2020](#); [Maurya et al. 2020, 2021](#)) and their physical parameters, such as reddening, age and distance are listed in Table 1. In the current study, we applied a statistical approach to study

the dynamical evolution of this homogeneous sample of open clusters.

2. Data

The observations for the clusters were taken using the 1.3-m Devasthal Fast Optical Telescope (DFOT) at Devasthal and 1.04-m Sampurnanand Telescope (ST) at Nainital. The telescopes were equipped with $2k \times 2k$ CCDs having a field-of-view of $\sim 18' \times 18'$ and $\sim 13' \times 13'$ for the DFOT and ST telescopes, respectively. The plate scales were 0.54 and 0.75 arcsec pixel $^{-1}$ for the DFOT and ST telescopes, respectively. The observations were taken in the nights of 30 November 2010, 24 and 25 March 2017, 21 October 2017 and 13 January 2018. We used the Image Reduction and Analysis Facility (IRAF) for cleaning the photometric data. We used the PSF technique to obtain the instrumental magnitudes using DAOPHOT II software packages. The instrumental magnitudes were calibrated to standard magnitudes through the process described by [Stetson \(1992\)](#). The conversion formula used for the standardization of instrumental magnitudes is given in our previous paper ([Maurya & Joshi 2020](#)). The membership of these clusters was calculated using proper motions and parallaxes from Gaia DR3 archives. We first plotted the vector-point diagram (VPD) of these clusters using proper motions. We noticed over-density regions in the VPDs of the clusters, which were found to be formed by probable cluster members on the basis of location of the stars belonging to these regions on the Hertzsprung–Russell (H–R) diagrams. To quantify the membership probabilities of stars of the clusters, we

calculated the membership probabilities from the proper motions of the stars using a statistical approach originally suggested by [Sanders \(1971\)](#). The membership probability for the i th star was calculated as follows:

$$P(i) = \frac{n_c \cdot \phi_c^v(i)}{n_c \cdot \phi_c^v(i) + n_f \cdot \phi_f^v(i)},$$

where n_c and n_f denote the number of stars belonging to field and cluster regions in the normalized form such that $n_c + n_f = 1$. The $\phi_c^v(i)$ and $\phi_f^v(i)$ in the above equation are the frequency distribution functions for the stars belonging to cluster and field populations, respectively. We calculated the $\phi_c^v(i)$ and $\phi_f^v(i)$ from the i th stars' proper motions ($\mu_{\alpha^*i}; \mu_{\delta i}$), errors in the proper motions ($\epsilon_{\alpha^*i}; \epsilon_{\delta i}$) and the mean proper motion of probable cluster members or field stars ($\mu_{\alpha^*c}; \mu_{\delta c}$) with their dispersions ($\sigma_{\alpha^*c}; \sigma_{\delta c}$) as described in [Maurya & Joshi \(2020\)](#). We also used Gaia parallaxes to remove the possible contamination of member stars from field stars with similar proper motions. The number of stars, N , identified as member stars, up to the completeness limit of the data for each cluster is given in [Table 1](#). The ADDSTAR routine of DAOPHOT packages was utilized to estimate the completeness of the data by adding artificial stars in the original images of the clusters. The completeness of the data was measured based on the ratio of the number of recovered stars to the number of artificially added stars in each magnitude bin. The completeness determination process is described in our previous study ([Maurya & Joshi 2020](#)). The completeness limit of the data in the V band is denoted by V_{lim} in [Table 1](#).

The physical parameters of the studied clusters are derived through the samples of member stars identified by us. The photometric data using our observations are complemented by Gaia DR3 and Pan-STARSS data for this study. The physical parameters derived in our previous studies, such as reddening, extinction law, distance, ages and completeness limits were used for the present study to be uniform in our sample as well as our approach ([Joshi et al. 2020a; Maurya & Joshi 2020; Maurya et al. 2020, 2021](#)). Thus, we prefer to select these clusters for the present work.

The reddening of all the clusters except SAI 35, SAI 44 and SAI 45, were estimated through fitting zero-age main sequence isochrones given by [Schmidt-Kaler \(1982\)](#) on the (U–B)/(B–V) color–color diagrams. In the absence of U band data, we calculated reddening of the clusters SAI 35, SAI 44 and SAI 45 using a 3D reddening map given by [Green et al. \(2019\)](#). We calculated E(B–V) values for the clusters from the E(g–r)

values provided by the reddening map using the extinction ratio relations of [Wang & Chen \(2019\)](#). The clusters included in this study are associated with relatively low reddening for the detection of lower mass regime member stars. The total-to-selective extinction values for the clusters were determined using (V– λ)/(B–V) two-color diagrams, where $\lambda = R, I, J, H, K$ bands' magnitudes. The near-IR bands' magnitudes were obtained from 2MASS archives. The distances to the clusters were estimated using Gaia DR3 parallaxes through the parallax inversion method as described in our previous study ([Maurya et al. 2020](#)). The distances to the clusters are calculated and are in the range of 1072 ± 44 – 3670 ± 184 parsec. The ages of the clusters were estimated by fitting [Marigo et al. \(2017\)](#) isochrones on the color–magnitude diagrams for derived reddening and distance modulus. The ages of the clusters were estimated in the range of 25 ± 19 Myr to 1.78 ± 0.20 Gyr. The methods for the derivation of the physical parameters are briefly described in [Maurya et al. \(2020, 2021\)](#) and their values are given in [Table 1](#).

3. Results

3.1 Mass segregation

The mass segregation can be attributed to the escape of the low-mass stars from the cluster besides the concentration of massive stars in the central region of the cluster. The mass segregation effect can be triggered due to the dynamical evolution through the equipartition of energy or may be imprint of the star formation process itself ([Sagar et al. 1988](#)). The cumulative distribution of stars with a radius for various mass ranges is often used to study the mass segregation in star clusters. To study the effect of dynamical evolution and mass segregation, we determined cumulative radial distributions of member stars for different mass ranges in our previous studies ([Joshi et al. 2020a; Maurya & Joshi 2020; Maurya et al. 2020](#)). However, the method based on the cumulative distribution of stars with radial distance depends on the size of mass bins and cumulative radii, which may give misleading results. Therefore, we used a method given by [Allison et al. \(2009b\)](#), which is based on mass segregation ratio (MSR) for the study of mass segregation scenario in the clusters.

The MSR is calculated from the mean edge length, γ , using a minimum spanning tree (MST) for the member stars. The MST for the member stars is the shortest path that connects all the member stars barring closed loops ([Prim 1957](#)). We have also shown the MST for these clusters in [Figure 1](#). The MST was generated using

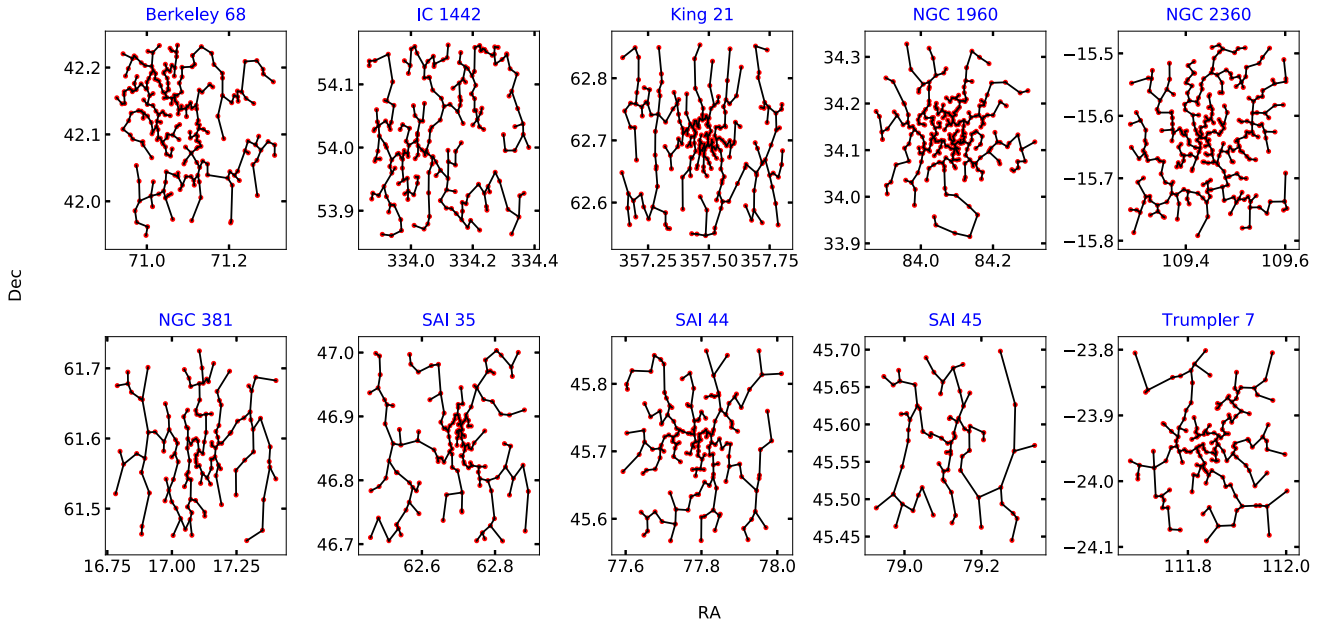


Figure 1. Plot of minimum spanning tree for the clusters. The vertices (dots) in the plots represent member stars. The lines connecting these vertices are generally referred as edges, which are parts of a spanning tree. The x and y axes in the plot represent the right ascension and declination for the stars associated with the clusters.

the Python package provided by Naidoo (2019). The densely located stars would have smaller edge lengths in the MST. These MST plots are helpful in the visual inspection of spatial distribution of the member stars in clusters. The clusters King 21 and NGC 1960 show the concentration of vertices in their central region, however, this is related to the distribution of stars in general and cannot be directly linked to mass segregation as the MSTs shown in Figure 1 do not have stellar mass information.

The γ is calculated for n most massive stars γ_{mm} and n random stars chosen from all members of the cluster γ_{rand} , separately. Then, we repeated the same process 500 times to calculate the mean of the γ_{rand} i.e., $\langle \gamma_{rand} \rangle$. The MSR value, Γ_{MSR} , was calculated using the formula provided by Olczak *et al.* (2011), which is given below:

$$\Gamma_{MSR} = \frac{\langle \gamma_{rand} \rangle}{\gamma_{mm}}.$$

We calculated the standard deviation $\Delta \Gamma_{MSR}$ value in the mass segregation ratio Γ_{MSR} as follows:

$$\Delta \Gamma_{MSR} = \Delta \gamma_{rand}.$$

The MSR method works on the principle that mass segregation will cause massive stars to have closer spatial distribution than low-mass stars. The Γ_{MSR} value for a cluster having similar spatial distribution for both the massive and low-mass stars will be ~ 1 , which means the absence of mass segregation. The value of

$\Gamma_{MSR} > 1$ would be interpreted as the relatively closer spatial distribution of the most massive stars than the rest of the stars, which means the presence of mass segregation.

We used only those member stars of these clusters to study the mass segregation, whose magnitudes are within the completeness limit in the V band as mentioned in Table 1. The clusters included in this study have ages > 20 Myr and are not embedded anymore in the star-forming regions. Hence, we took star magnitudes as the proxy for their masses. In our calculation of the Γ_{MSR} , we started with a minimum of 10 most massive stars in the clusters to have statistically significant calculations, then we repeated the calculation of the Γ_{MSR} by increasing the number of massive stars in steps of one in each iteration up to 30% of the total stars. The mass range, denoted by $M_{30\%}$, corresponding to the most massive stars, whose numbers are up to 30% of total stars, is given in Table 2. We have shown the plots for the obtained values of Γ_{MSR} vs. the number of massive stars $N_{massive}$ in Figure 2.

We have reported the maximum value of Γ_{MSR} obtained for the clusters in Table 2. We found that the clusters NGC 2360, NGC 1960, IC 1442, King 21 and SAI 35 exhibit moderate signatures of mass segregation having Γ_{MSR} to be 1.65 ± 0.18 , 1.94 ± 0.22 , 2.21 ± 0.20 , 1.84 ± 0.23 and 1.96 ± 0.25 , respectively. The $\Gamma_{MSR} \sim 2.0$ values for IC 1442 are only for $N_{massive}$ values of 11, 12 and 13. After these values, the MSR drops sharply to around 1.0, which indicates that IC 1442 exhibits mass

Table 2. Name of the clusters, total mass of the clusters, mass range of 30% most massive stars, R_h , R_t , R_{gc} , R_h/R_t and Γ_{MSR} are given in columns 1 to 8, respectively.

Cluster	$M_C (M_\odot)$	$M_{30\%} (M_\odot)$	R_h (pc)	R_t (pc)	R_{gc} (pc)	R_h/R_t	Γ_{MSR}
NGC 381	467.13 ± 187.41	2.80–1.86	1.50 ± 0.05	10.30 ± 1.36	8.91 ± 0.02	0.15 ± 0.02	1.26 ± 0.18
NGC 2360	887.54 ± 356.08	2.21–1.17	1.65 ± 0.07	12.75 ± 1.69	8.93 ± 0.03	0.13 ± 0.02	1.65 ± 0.18
Berkeley 68	1396.11 ± 552.29	1.75–1.27	5.41 ± 0.34	17.06 ± 2.23	11.29 ± 0.19	0.32 ± 0.05	1.16 ± 0.13
NGC 1960	618.94 ± 248.32	7.18–1.51	2.21 ± 0.10	11.65 ± 1.54	9.36 ± 0.05	0.19 ± 0.03	1.94 ± 0.22
IC 1442	1811.75 ± 901.03	9.02–2.38	6.30 ± 0.41	16.77 ± 2.76	9.52 ± 0.12	0.38 ± 0.07	2.21 ± 0.20
King 21	1454.19 ± 654.43	6.89–4.22	3.18 ± 0.19	15.93 ± 2.37	9.86 ± 0.12	0.20 ± 0.03	1.84 ± 0.23
Trumpler 7	508.00 ± 203.81	5.12–1.52	1.78 ± 0.10	10.80 ± 1.43	9.21 ± 0.06	0.16 ± 0.02	1.04 ± 0.19
SAI 35	817.47 ± 327.96	3.34–1.38	3.78 ± 0.36	13.93 ± 1.85	10.81 ± 0.25	0.27 ± 0.04	1.96 ± 0.25
SAI 44	881.70 ± 353.73	2.47–1.39	4.59 ± 0.23	15.01 ± 1.99	11.74 ± 0.18	0.31 ± 0.04	1.19 ± 0.18
SAI 45	356.18 ± 142.90	2.14–1.84	2.09 ± 0.06	9.98 ± 1.32	9.80 ± 0.05	0.21 ± 0.03	1.14 ± 0.24

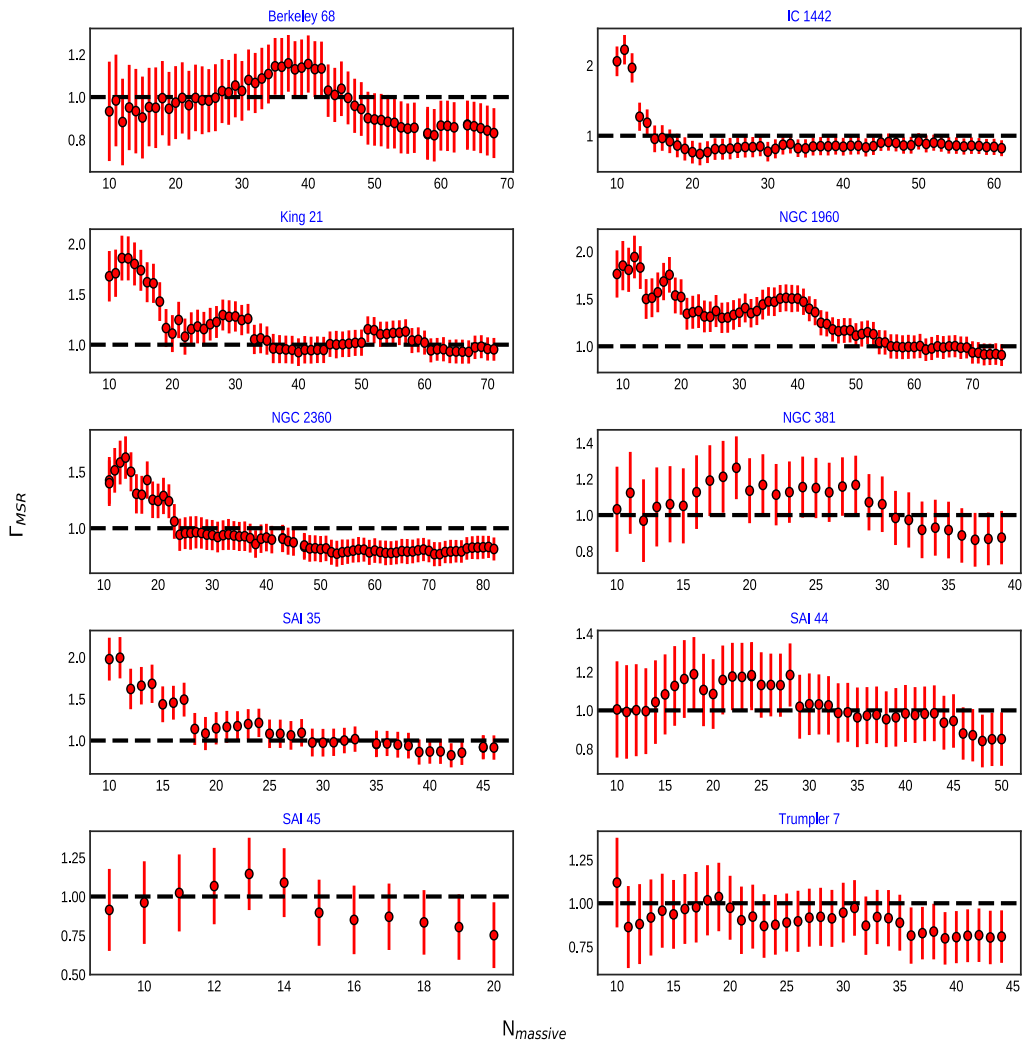


Figure 2. Plot of Γ_{MSR} vs. $N_{massive}$ for the clusters. The error bar in the plot denotes the standard deviation $\Delta\Gamma_{MSR}$ values for each value of $N_{massive}$.

segregation for only a small number of the most massive stars. The value of $\Gamma_{\text{MSR}} = 1.26 \pm 0.18$ for cluster NGC 381, which indicates the presence of weak mass segregation in the cluster. The remaining four clusters namely, Berkeley 68, Trumpler 7, SAI 44 and SAI 45 exhibit Γ_{MSR} values within uncertainties to be 1, thus, we do not find any evidence of mass segregation in these clusters. There is a bump around N_{massive} values from 30 to 42 for the cluster Berkeley 68, which could be due to the presence of subgroups of massive stars (Yu *et al.* 2017). However, the MSR for cluster Berkeley 68 is very low with $\Gamma_{\text{MSR}} = 1.16 \pm 0.13$.

3.2 Half mass radius, tidal radius and cluster structure

It has been found that tidal interactions can influence the structure and dynamical evolution of the open clusters (Chumak *et al.* 2010). The ratio of half mass radius to the tidal radius, R_h , can be used as an indicator of the disruption of a cluster caused by tidal forces (Angelo *et al.* 2021). The radial distance from the cluster center containing half of the total mass of the clusters is defined as half mass radius. The tidal radius, R_t , is defined as the radial distance from the cluster center, where the tidal field of the Galaxy is balanced by the gravitational field of the cluster. We used the relation given by Kim *et al.* (2000) to calculate the tidal radius as follows:

$$R_t = \left(\frac{M_C}{2M_G} \right)^{1/3} \times R_{gc}.$$

The symbols M_C and R_{gc} in the above equation denote the total mass of the cluster and the distance to the cluster from the Galactic center, respectively. The mass of the Galaxy contained within the Galactocentric distance of the cluster is denoted by M_G and calculated by using Genzel & Townes (1987) relation given below:

$$M_G = 2 \times 10^8 M_\odot \left(\frac{R_{gc}}{30 \text{ pc}} \right)^{1.2}.$$

We used the above two equations to calculate the tidal radius of the clusters. The R_{gc} values for the clusters were calculated using the relation $R_{gc}^2 = R_\odot^2 + (d \cos b)^2 - 2R_\odot d \cos b \cos l$. The d and R_\odot represent distances to clusters and the distance of the Sun from the Galactic center, respectively, whereas l and b denote the Galactic longitude and latitude, respectively. We used R_\odot as 8.2 ± 0.10 kpc for the calculations as given by Bland-Hawthorn *et al.* (2019). The values of R_t and R_{gc} are given in Table 2.

The stellar initial mass function peaks at $0.5 M_\odot$, whereas mass completeness limits are above $0.5 M_\odot$ for all the clusters studied in this analysis. Therefore,

we estimated the total stellar mass, M_C , of the clusters utilizing the stellar mass function. The mass function slopes for these clusters obtained in our previous studies for the stars above $1 M_\odot$ were mostly in agreement with the Kroupa (2001) mass-function slopes within the uncertainty values (Joshi *et al.* 2020b; Maurya & Joshi 2020; Maurya *et al.* 2020, 2021). So, we used Kroupa (2001) mass function of multiple-part power law form to estimate the total mass of the clusters including stellar populations up to $0.08 M_\odot$. We applied a similar method as described by Snider *et al.* (2009). We are giving a brief description here.

The number of stars in the mass range m_1 – m_2 will be:

$$N = A \times \int_{m_1}^{m_2} M^{-\alpha} dM.$$

Using the above equation for the known number of stars from our samples, we calculated the values of the normalization constant A for the α corresponding to the mass range of $M/M_\odot \geq 1$ given by Kroupa (2001). The total mass can be obtained using the following relation:

$$M_{\text{tot}} = \int_0^N M dN = A \times \int_{m_1}^{m_2} M^{1-\alpha} dM.$$

Using the above equation, we calculated the total mass of the clusters in the mass ranges of $M/M_\odot \geq 1.0$, $0.5 \leq M/M_\odot < 1.0$, $0.08 \leq M/M_\odot < 0.5$. The normalization constant values for the other mass ranges namely, $0.5 \leq M/M_\odot < 1.0$ and $0.08 \leq M/M_\odot < 0.5$ were calculated using the relations of the normalization constant values for different mass ranges given by Maschberger (2013). The estimated total mass values, M_C , through this method up to the lower stellar mass limit of $0.08 M_\odot$ for the clusters are given in Table 2.

The half-mass radii, R_h , of the clusters were calculated from the individual masses of member stars of the clusters. The individual masses of the member stars were determined by fitting the solar metallicity isochrones of Marigo *et al.* (2017) on the V/(B–V) color–magnitude diagrams of the clusters. These isochrones were corresponding to the derived age, reddening and distance of the clusters. The half-mass radii were estimated corresponding to the observed total masses of the clusters. We derived the half-mass radii of the clusters by including stars with magnitudes within the completeness limits of our data mentioned in Table 1. Generally, the half-mass radii estimated from deep photometric data tend to become larger compared to the half-mass radii estimated from shallow photometric data (Bonatto & Bica 2008). However, half-mass

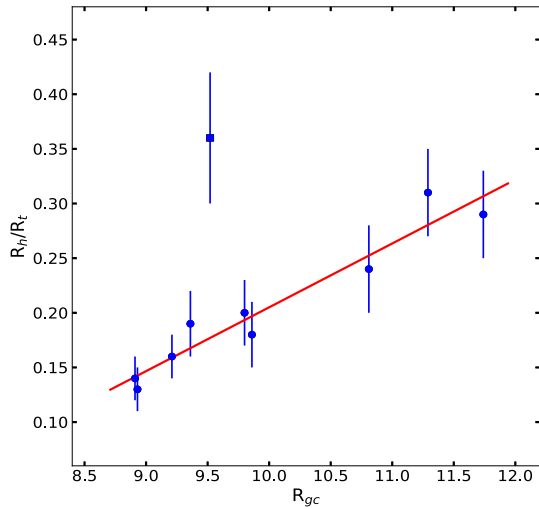


Figure 3. Plot of R_h/R_t vs. R_{gc} for the clusters. The point shown by the square is an outlier. The outlier point belongs to cluster IC 1442.

radii and other structural parameters become insensitive to photometric depth in case of uniform mass function slopes i.e., absence of mass segregation (Bonatto & Bica 2008). As the studied clusters do not show strong mass segregation and the Γ_{MSR} values reach around 1.0 even for most massive stars, which are 30% of the total stars, the half-mass radii derived by us are reasonably good estimates. However, a larger sample having clusters with a high number of member stars would be better to constrain the correlation shown in Figure 3. The values of R_h and R_h/R_t are given in Table 2.

We plotted R_h/R_t vs. R_{gc} for the clusters as shown in Figure 3. We found that the R_h/R_t ratios are positively correlated with R_{gc} values as seen in the figure. The R_h/R_t value of 0.38 ± 0.07 with $R_{gc} = 9.52 \pm 0.12$ for the cluster IC 1442 is an outlier in the above plot. The binary clusters have been found to show relatively larger R_h/R_t values (Angelo *et al.* 2021). The slope of the linear fit was found to be 0.06 ± 0.01 with linear regression coefficient r -square = 0.93 in the range of R_{gc} from 8.9 to 11.7 kpc. The clusters located at the larger Galactocentric distances might be subjected to weaker tidal fields, which could be a possible reason for these clusters to have larger R_h/R_t ratios (Angelo *et al.* 2021).

4. Discussion

We studied a sample of 10 open clusters with photometric data mostly complete up to 19 mag in the V band to understand the dynamical evolution of the clusters. These clusters were chosen for being associated

with comparatively lower reddening to assist the detection of fainter stars. We utilized precise astrometric data from Gaia DR3 to ascertain the membership of the stars and accurate measurements of the distances, which help in the precise estimation of distance modulus essential for mass determination through isochrone fitting on the CMDs of the clusters.

We used a method based on the minimum spanning tree of the member stars as suggested by Allison *et al.* (2009b) to study mass segregation scenarios in the clusters. We found that the clusters NGC 2360, NGC 1960, IC 1442, King 21 and SAI 35 exhibit evidence of the presence of moderate mass segregation having MSR to be 1.65 ± 0.18 , 1.94 ± 0.22 , 2.21 ± 0.20 , 1.84 ± 0.23 and 1.96 ± 0.25 . The remaining clusters show weak or no mass segregation with Γ_{MSR} around 1 within the uncertainty level. The mass segregation in the open clusters has been thought to increase with age (Dib *et al.* 2018). In the present study, we could not find such a trend between mass segregation and the age of the clusters. Similarly, Tarricq *et al.* (2022) did not find any clear trend between the MSR calculated from the 10 most massive stars and the clusters' ages using a sample of 389 open clusters.

The ratio of the half-mass radius to the tidal radius, R_h/R_t , is a good parameter to study the influence of the tidal interactions on the dynamical evolution and shapes of the clusters (Baumgardt *et al.* 2010; Angelo *et al.* 2020). This ratio measures the fraction of the tidal volume filled by the half-mass contents. We estimated the total mass of the clusters for the calculation of R_t by using the multiple-part power law form of mass function given by Kroupa (2001) to include the masses of fainter stars up to $0.08 M_{\odot}$. The ratios R_h/R_t are found to be positively correlated to R_{gc} with a slope of 0.06 ± 0.01 and linear regression r -square coefficient of 0.93. This indicates a tendency that the clusters located at larger R_{gc} had larger R_h/R_t ratios. In a similar study with a larger sample of open clusters, Angelo *et al.* (2020) could not find any such trend, however, they noticed that the clusters located at $R_{gc} > 9$ kpc had at least 50% larger dispersion in R_h/R_t values. It has been suggested that the clusters located at larger R_{gc} face lower external gravitational forces, which allow the internal stellar content to fill a larger fraction of tidal volume without tidal disruption, so these clusters exhibit larger R_h/R_t values (Angelo *et al.* 2020, 2021). The correlation between R_h/R_t and R_{gc} , found here, is based on a small sample of open clusters of a wide age range, which is an important caveat of the present study. Therefore, a study based on larger samples of open clusters of similar age across different galactic locations would

give more insight into the dynamical evolution of the clusters.

5. Conclusions

We investigated the dynamical evolution of 10 open clusters, whose physical parameters were previously determined by us using a homogeneous approach and data sets. In the present study, we identified member stars of the clusters using Gaia DR3 astrometric data. Our findings can be summarized as follows:

- We found moderate mass segregation with Γ_{MSR} values to be 1.65 ± 0.18 , 1.94 ± 0.22 , 2.21 ± 0.20 , 1.84 ± 0.23 and 1.96 ± 0.25 for the clusters NGC 2360, NGC 1960, IC 1442, King 21 and SAI 35, respectively.
- The cluster NGC 381 exhibits a weak signature of mass segregation with Γ_{MSR} is 1.26 ± 0.18 .
- We found no evidence of mass segregation in the clusters Berkeley 68, Trumpler 7, SAI 44 and SAI 45.
- The ratios R_h/R_t are found to be positively correlated to R_{gc} with a slope of 0.06 ± 0.01 and linear regression r -square coefficient of 0.93.

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