NONLOCAL DIFFUSION PROBLEMS THAT APPROXIMATE THE HEAT EQUATION WITH DIRICHLET BOUNDARY CONDITIONS

ΒY

CARMEN CORTAZAR AND MANUEL ELGUETA

Departamento de Matemática, Universidad Catolica de Chile Casilla 306, Correo 22, Santiago, Chile e-mail: ccortaza@mat.puc.cl, melgueta@mat.puc.cl

AND

Julio D. Rossi

Departamento de Matemática, FCEyN UBA (1428) Buenos Aires, Argentina e-mail: jrossi@dm.uba.ar

ABSTRACT

We present a model for nonlocal diffusion with Dirichlet boundary conditions in a bounded smooth domain. We prove that solutions of properly rescaled nonlocal problems approximate uniformly the solution of the corresponding Dirichlet problem for the classical heat equation.

1. Introduction

Let $J : \mathbb{R}^N \to \mathbb{R}$ be a nonnegative, radial, continuous function with $\int_{\mathbb{R}^N} J(z)$, dz = 1. Assume also that J is strictly positive in B(0, d) and vanishes in $\mathbb{R}^N \setminus B(0, d)$. Nonlocal evolution equations of the form

(1.1)
$$u_t(x,t) = (J * u - u)(x,t) = \int_{\mathbb{R}^N} J(x-y)u(y,t) \, dy - u(x,t),$$

and variations of it, have been widely used to model diffusion processes. As stated in [12] equation (1.1) models a "random walk" continuous in time where

Received June 1, 2007

the probability distribution of jumping from location y to location x is given by J(x-y). For recent references on nonlocal diffusion see, [2], [1], [3], [4], [5], [6], [8], [12], [14], [15] and references therein.

In this article we propose the following nonlocal "Dirichlet" boundary value problem: Given g(x,t) defined for $x \in \mathbb{R}^N \setminus \Omega$ and $u_0(x)$ defined for $x \in \Omega$, find u(x,t) such that

(1.2)
$$\begin{cases} u_t(x,t) = \int_{\mathbb{R}^N} J(x-y)(u(y,t) - u(x,t))dy, & x \in \Omega, \ t > 0, \\ u(x,t) = g(x,t), & x \notin \Omega, \ t > 0, \\ u(x,0) = u_0(x), & x \in \Omega. \end{cases}$$

In this model we prescribe the values of u outside Ω which is the analogue of prescribing the so-called Dirichlet boundary conditions for the classical heat equation. However, the boundary data is not understood in the usual sense as we will see in Remark 2.1, below. As explained before the integral $\int J(x-y)(u(y,t)-u(x,t)) dy$ takes into account the individuals arriving or leaving position $x \in \Omega$ from or to other places while we are prescribing the values of u outside the domain Ω by imposing u = g for $x \notin \Omega$. When g = 0 we get that any individuals that leave Ω , die, this is the case when Ω is surrounded by a hostile environment. See [11] for a similar model.

Existence and uniqueness of solutions of (1.2) is proved by a fixed point argument in Section 2, where a comparison principle is also obtained.

Let us consider the classical Dirichlet problem for the heat equation,

(1.3)
$$\begin{cases} v_t(x,t) - \Delta v(x,t) = 0, & x \in \Omega, \ t > 0, \\ v(x,t) = g(x,t), & x \in \partial\Omega, \ t > 0, \\ v(x,0) = u_0(x), & x \in \Omega. \end{cases}$$

The nonlocal Dirichlet model (1.2) and the classical Dirichlet problem (1.3) share many properties, among them the asymptotic behavior of their solutions as $t \to \infty$ is similar as was proved in [7].

The main goal of this article is to show that the Dirichlet problem for the heat equation (1.3) can be approximated by suitable nonlocal problems of the form of (1.2).

More precisely, for a given J and a given $\varepsilon > 0$ we consider the rescaled kernel

(1.4)
$$J_{\varepsilon}(\xi) = C_1 \frac{1}{\varepsilon^N} J\left(\xi/\varepsilon\right), \quad \text{with} \ C_1^{-1} = \frac{1}{2} \int_{B(0,d)} J(z) z_N^2 \, dz.$$

55

Here C_1 is a normalizing constant in order to obtain the Laplacian in the limit instead of a multiple of it. Let $u^{\varepsilon}(x,t)$ be the solution of

(1.5)
$$\begin{cases} u_t^{\varepsilon}(x,t) = \int_{\Omega} \frac{J_{\varepsilon}(x-y)}{\varepsilon^2} (u^{\varepsilon}(y,t) - u^{\varepsilon}(x,t)) dy, & x \in \Omega, t > 0, \\ u(x,t) = g(x,t), & x \notin \Omega, t > 0, \\ u(x,0) = u_0(x), & x \in \Omega. \end{cases}$$

Our main result now reads as follows.

THEOREM 1.1: Let Ω be a bounded $C^{2+\alpha}$ domain for some $0 < \alpha < 1$. Let $v \in C^{2+\alpha,1+\alpha/2}(\overline{\Omega} \times [0,T])$ be the solution to (1.3) and let u^{ε} be the solution to (1.5) with J_{ε} as above. Then, there exists C = C(T) such that

(1.6)
$$\sup_{t \in [0,T]} \|v - u^{\varepsilon}\|_{L^{\infty}(\Omega)} \le C\varepsilon^{\alpha} \to 0, \quad \text{as } \varepsilon \to 0.$$

Related results for the Neumann problem where recently obtained in [10].

Note that the assumed regularity of v is a consequence of regularity assumptions on the boundary data g, the domain Ω and the initial condition u_0 , see [13].

The rest of the paper is organized as follows: in Section 2, we prove existence, uniqueness and a comparison principle for our nonlocal equation and in Section 3 we prove the convergence result.

2. Existence, uniqueness and a comparison principle

Existence and uniqueness of solutions is a consequence of Banach's fixed point theorem. We look for $u \in C([0,\infty); L^1(\Omega))$ satisfying (1.2). Fix $t_0 > 0$ and consider the Banach space $X_{t_0} = \{w \in C([0,t_0]; L^1(\Omega))\}$ with the norm

$$|||w||| = \max_{0 \le t \le t_0} \|w(\cdot, t)\|_{L^1(\Omega)}.$$

We will obtain the solution as a fixed point of the operator $\mathcal{T}: X_{t_0} \to X_{t_0}$ defined by

$$\mathcal{T}_{w_0}(w)(x,t) = w_0(x) + \int_0^t \int_{\mathbb{R}^N} J(x-y) \left(w(y,s) - w(x,s) \right) dy \, ds,$$

where we impose

$$w(x,t) = g(x,t), \quad \text{for } x \notin \Omega.$$

LEMMA 2.1: Let $w_0, z_0 \in L^1(\Omega)$, then there exists a constant C depending on J and Ω such that

$$|||\mathcal{T}_{w_0}(w) - \mathcal{T}_{z_0}(z)||| \le Ct_0|||w - z||| + ||w_0 - z_0||_{L^1(\Omega)}$$

for all $w, z \in X_{t_0}$.

Proof. We have

$$\begin{split} &\int_{\Omega} |\mathcal{T}_{w_0}(w)(x,t) - \mathcal{T}_{z_0}(z)(x,t)| \, dx \le \int_{\Omega} |w_0 - z_0|(x) \, dx \\ &+ \int_{\Omega} \left| \int_0^t \int_{\mathbb{R}^N} J\left(x - y\right) \left[(w(y,s) - z(y,s)) - (w(x,s) - z(x,s)) \right] dy \, ds \right| \, dx. \end{split}$$

Hence, taking into account that w and z vanish outside Ω ,

$$|||\mathcal{T}_{w_0}(w) - \mathcal{T}_{z_0}(z)||| \le ||w_0 - z_0||_{L^1(\Omega)} + Ct_0|||w - z|||,$$

as we wanted to prove.

THEOREM 2.1: For every $u_0 \in L^1(\Omega)$ there exists a unique solution u, such that $u \in C([0,\infty); L^1(\Omega))$.

Proof. We check first that \mathcal{T}_{u_0} maps X_{t_0} into X_{t_0} . Taking $z_0 \equiv 0$ and $z \equiv 0$ in Lemma 2.1 we get that $\mathcal{T}_{u_0}(w) \in C([0, t_0]; L^1(\Omega))$ for any $w \in X_{t_0}$.

Choose t_0 such that $Ct_0 < 1$. Now taking $z_0 \equiv w_0 \equiv u_0$ in Lemma 2.1 we get that \mathcal{T}_{u_0} is a strict contraction in X_{t_0} and the existence and uniqueness part of the theorem follows from Banach's fixed point theorem in the interval $[0, t_0]$. To extend the solution to $[0, \infty)$ we may take as initial data $u(x, t_0) \in L^1(\Omega)$ and obtain a solution up to $[0, 2t_0]$. Iterating this procedure we get a solution defined in $[0, \infty)$.

Remark 2.1: Note that in general a solution u with $u_0 > 0$ and g = 0 is strictly positive in $\overline{\Omega}$ (with a positive continuous extension to $\overline{\Omega}$) and vanishes outside $\overline{\Omega}$. Therefore, a discontinuity occurs on $\partial\Omega$ and the boundary value is not taken in the usual "classical" sense, see [7].

We now define what we understand by sub and supersolutions.

Definition 2.1: A function $u \in C([0,T); L^1((\Omega))$ is a supersolution of (1.2) if

(2.1)
$$\begin{cases} u_t(x,t) \ge \int_{\mathbb{R}^N} J(x-y)(u(y,t)-u(x,t))dy, & x \in \Omega, \ t > 0, \\ u(x,t) \ge g(x,t), & x \notin \Omega, \ t > 0, \\ u(x,0) \ge u_0(x), & x \in \Omega. \end{cases}$$

As usual, subsolutions are defined analogously by reversing the inequalities.

LEMMA 2.2: Let $u_0 \in C(\overline{\Omega})$, $u_0 \ge 0$, and $u \in C(\overline{\Omega} \times [0,T])$ a supersolution to (1.2) with $g \ge 0$. Then, $u \ge 0$.

Proof. Assume to the contrary that u(x,t) is negative in some point. Let $v(x,t) = u(x,t) + \varepsilon t$ with ε so small such that v is still negative somewhere. Then, if (x_0, t_0) is a point where v attains its negative minimum, there holds that $t_0 > 0$ and

$$v_t(x_0, t_0) = u_t(x_0, t_0) + \varepsilon > \int_{\mathbb{R}^N} J(x - y)(u(y, t_0) - u(x_0, t_0)) \, dy$$
$$= \int_{\mathbb{R}^N} J(x - y)(v(y, t_0) - v(x_0, t_0)) \, dy \ge 0$$

which is a contradiction. Thus, $u \ge 0$.

COROLLARY 2.1: Let $J \in L^{\infty}(\mathbb{R}^N)$. Let u_0 and v_0 in $L^1(\Omega)$ with $u_0 \geq v_0$ and $g, h \in L^{\infty}((0,T); L^1(\mathbb{R}^N \setminus \Omega))$ with $g \geq h$. Let u be a solution of (1.2) with $u(x,0) = u_0$ and Dirichlet datum g and v be a solution of (1.2) with $v(x,0) = v_0$ and datum h. Then, $u \geq v$ a.e.

Proof. Let w = u - v. Then, w is a supersolution with initial datum $u_0 - v_0 \ge 0$ and datum $g - h \ge 0$. Using the continuity of solutions with respect to the data and the fact that $J \in L^{\infty}(\mathbb{R}^N)$, we may assume that $u, v \in C(\overline{\Omega} \times [0, T])$. By Lemma 2.2 we obtain that $w = u - v \ge 0$. So the corollary is proved.

COROLLARY 2.2: Let $u \in C(\overline{\Omega} \times [0,T])$ (resp., v) be a supersolution (resp., subsolution) of (1.2). Then, $u \geq v$.

Proof. It follows from the proof of the previous corollary.

3. Convergence to the heat equation

In order to prove Theorem 1.1, let \tilde{v} be a $C^{2+\alpha,1+\alpha/2}$ extension of v to $\mathbb{R}^N \times [0,T]$.

Let us define the operator

$$\tilde{L}_{\varepsilon}(z) = \frac{1}{\varepsilon^2} \int_{\mathbb{R}^N} J_{\varepsilon}(x-y) \big(z(y,t) - z(x,t) \big) dy.$$

Then \tilde{v} verifies

(3.1)
$$\begin{cases} \tilde{v}_t(x,t) = \tilde{L}_{\varepsilon}(\tilde{v})(x,t) + F_{\varepsilon}(x,t) & x \in \Omega, \ (0,T], \\ \tilde{v}(x,t) = g(x,t) + G(x,t), & x \notin \Omega, \ (0,T], \\ \tilde{v}(x,0) = u_0(x), & x \in \Omega. \end{cases}$$

where, since $\Delta v = \Delta \tilde{v}$ in Ω ,

$$F_{\varepsilon}(x,t) = -\tilde{L}_{\varepsilon}(\tilde{v})(x,t) + \Delta \tilde{v}(x,t).$$

Moreover as G is smooth and G(x,t) = 0 if $x \in \partial \Omega$ we have

 $G(x,t) = O(\varepsilon),$ for x such that $\operatorname{dist}(x,\partial\Omega) \leq \varepsilon d.$

We set $w^{\varepsilon} = \tilde{v} - u^{\varepsilon}$ and we note that

(3.2)
$$\begin{cases} w_t^{\varepsilon}(x,t) = \tilde{L}_{\varepsilon}(w^{\varepsilon})(x,t) + F_{\varepsilon}(x,t) & x \in \Omega, \ (0,T], \\ w^{\varepsilon}(x,t) = G(x,t), & x \notin \Omega, \ (0,T], \\ w^{\varepsilon}(x,0) = 0, & x \in \Omega. \end{cases}$$

First, we claim that, by the choice of C_1 , the fact that J is radially symmetric and $\tilde{u} \in C^{2+\alpha,1+\alpha/2}(\mathbb{R}^N \times [0,T])$, we have that

(3.3)
$$\sup_{t\in[0,T]} \|F_{\varepsilon}\|_{L^{\infty}(\Omega)} = \sup_{t\in[0,T]} \|\Delta \tilde{v} - \tilde{L}_{\varepsilon}(\tilde{v})\|_{L^{\infty}(\Omega)} = O(\varepsilon^{\alpha}).$$

In fact,

$$\Delta \tilde{v}(x,t) - \frac{C_1}{\varepsilon^{N+2}} \int_{\mathbb{R}^N} J\left((x-y)/\varepsilon \right) \left(\tilde{v}(y,t) - \tilde{v}(x,t) \right) \, dy$$

becomes, under the change variables $z = (x - y)/\varepsilon$,

$$\Delta \tilde{v}(x,t) - \frac{C_1}{\varepsilon^2} \int_{\mathbb{R}^N} J(z) \left(\tilde{v}(x-\varepsilon z,t) - \tilde{v}(x,t) \right) dz$$

and hence (3.3) follows by a simple Taylor expansion. This proves the claim.

We proceed now to prove Theorem 1.1.

Proof of Theorem 1.1. In order to prove the theorem by a comparison we first look for a supersolution. Let \overline{w} be given by

(3.4)
$$\overline{w}(x,t) = K_1 \varepsilon^{\alpha} t + K_2 \varepsilon$$

For $x \in \Omega$ we have, if K_1 is large,

(3.5)
$$\overline{w}_t(x,t) - \tilde{L}(\overline{w})(x,t) = K_1 \varepsilon^{\alpha} \ge F_{\varepsilon}(x,t) = w_t^{\varepsilon}(x,t) - \tilde{L}_{\varepsilon}(w^{\varepsilon})(x,t).$$

Since

$$G_{\varepsilon}(x,t) = O(\varepsilon)$$
 for x such that $\operatorname{dist}(x,\partial\Omega) \leq \varepsilon$

choosing K_2 large, we obtain

(3.6)
$$\overline{w}(x,t) \ge w^{\varepsilon}(x,t)$$

for $x \notin \Omega$ such that $\operatorname{dist}(x, \partial \Omega) \leq \varepsilon d$ and $t \in [0, T]$. Moreover it is clear that

(3.7)
$$\overline{w}(x,0) = K_2 \varepsilon > w^{\varepsilon}(x,0) = 0.$$

By (3.5), (3.6) and (3.7) we can apply the comparison result and conclude that

(3.8)
$$w^{\varepsilon}(x,t) \le \overline{w}(x,t) = K_1 \varepsilon^{\alpha} t + K_2 \varepsilon.$$

In a similar fashion we prove that $\underline{w}(x,t) = -K_1 \varepsilon^{\alpha} t - K_2 \varepsilon$ is a subsolution and hence,

(3.9)
$$w^{\varepsilon}(x,t) \ge \underline{w}(x,t) = -K_1 \varepsilon^{\alpha} t - K_2 \varepsilon.$$

Therefore

(3.10)
$$\sup_{t \in [0,T]} \|u - u^{\varepsilon}\|_{L^{\infty}(\Omega)} \le C(T)\varepsilon^{\alpha}.$$

This proves the theorem.

ACKNOWLEDGMENTS. The authors were supported by Universidad de Buenos Aires under grant X066, by ANPCyT PICT No. 03-13719, by CONICET (Argentina) and by FONDECYT Project 1070944 and Coop. Int. 7050118 (Chile).

References

- P. Bates and A. Chmaj, A discrete convolution model for phase transitions, Archives for Rational Mechanics and Analysis 150 (1999), 281–305.
- [2] P. Bates and A. Chmaj, An integrodifferential model for phase transitions: stationary solutions in higher dimensions, Journal of Statistical Physics 95 (1999), 1119–1139.
- [3] P. Bates, P. Fife, X. Ren and X. Wang, *Travelling waves in a convolution model for phase transitions*, Archives for Rational Mechanics and Analysis **138** (1997), 105–136.
- [4] P. Bates and J. Han, The Dirichlet boundary problem for a nonlocal Cahn-Hilliard equation, Journal of Mathematical Analysis and Applications 311 (2005), 289–312.
- [5] P. Bates and J. Han, The Neumann boundary problem for a nonlocal Cahn-Hilliard equation, Journal of Differential Equations 212 (2005), 235–277.
- [6] C. Carrillo and P. Fife, Spatial effects in discrete generation population models, Journal of Mathematical Biology 50 (2005), 161–188.
- [7] E. Chasseigne, M. Chaves and J. D. Rossi, Asymptotic behavior for nonlocal diffusion equations, Journal de Mathématiques Pures et Applquées 86 (2006), 271–291.
- [8] X Chen, Existence, uniqueness and asymptotic stability of traveling waves in nonlocal evolution equations, Advances in Differential Equations 2 (1997), 125–160.
- [9] C. Cortazar, M. Elgueta and J. D. Rossi, A non-local diffusion equation whose solutions develop a free boundary, Annales Henri Poincaré 6 (2005), 269–281.
- [10] C. Cortazar, M. Elgueta, J. D. Rossi and N. Wolanski, How to approximate the heat equation with Neumann boundary conditions by nonlocal diffusion problems, Archives of Rational and Mechanical Analysis, 187 (2008), 137–156.
- [11] P. Fife, Clines and material interfaces with nonlocal interaction, in Partial Differential Equations and Applications (P. Marcellini, G. Talenti and E. Visentinti, eds.), Lecture Notes Pure Appl. Math., Vol. 177, Marcel Dekker, 1996.
- [12] P. Fife, Some nonclassical trends in parabolic and parabolic-like evolutions, in Trends in Nonlinear Analysis, Springer, Berlin, 2003, pp. 153–191.
- [13] A. Friedman, Partial Differential Equations of Parabolic Type, Prentice-Hall, Englewood Cliffs, NJ, 1964.
- [14] X. Wang, Metaestability and stability of patterns in a convolution model for phase transitions, Journal of Differential Equations 183 (2002), 434–461.
- [15] L. Zhang, Existence, uniqueness and exponential stability of traveling wave solutions of some integral differential equations arising from neuronal networks, Journal of Differential Equations 197 (2004), 162–196.