

# Correlation of Impact Conditions, Interface Reactions, Microstructural Evolution, and Mechanical Properties in Kinetic Spraying of Metals: A Review

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In the past, most studies into kinetic spraying technology focused on basic research, but a large portion of current research is devoted to industrial applications of the technology. To advance, however, studies about industrial applications of kinetic spraying require profound understanding of the scientific foundations of the kinetic spray process. Nevertheless, no one has yet provided a well-organized summary of the correlations among impact conditions, interface reactions, microstructural evolution, and mechanical properties across the whole field of kinetic spraying technology. This paper provides such an overview of these correlations for kinetic spraying of metals. For each correlation, the interactions between the given conditions and the material properties of the metal feedstock powder are the most influential. These interactions are so complicated that it is difficult to systematically classify all cases into certain types. Nonetheless, we try to explain and summarize the critical factors and their roles in each relationship.



# 1. Introduction

Kinetic spraying, or the cold gas dynamic spray process, is a method for deposition of micron-sized powder onto a target by accelerating particles to supersonic speeds (300– 1200 m/s) using high-pressure gas with a convergent–divergent de Laval nozzle (Ref [1-8](#page-24-0)). Since the feedstock powder is deposited primarily through the kinetic energy of the supersonic acceleration without melting, coating layers with specific properties that provide various advantages over the thermal spray process can be produced (Ref [3](#page-24-0), [9-18](#page-25-0)). Also, intimate bonding between the impact pairs can be induced through the creation of adiabatic shear instability at interface regions (Ref [19](#page-25-0)-[23](#page-25-0)). Therefore, kinetic spraying technology has received much

attention for use in various industrial fields (i.e., automotive, aviation, defense, and energy) (Ref [24-27](#page-25-0)).

Research on kinetic spraying has been conducted since the 1990s. In the early days, most studies focused on determining the deposition mechanism of metals (Ref [19-](#page-25-0) [23](#page-25-0), [28-42\)](#page-25-0). It was revealed that the creation of adiabatic shear instability at the periphery of the interface region played a dominant role in bond formation via severe plastic deformation (Ref [19-23](#page-25-0), [28](#page-25-0), [29](#page-25-0)). On this basis, the minimum particle velocity required for successful deposition was defined as the critical velocity of kinetic spraying, resulting from competition between the bonding (plastic deformation) and debonding (elastic recovery) energies (Ref [19](#page-25-0)). Early research defined the *rebound velocity* as the velocity at which successful deposition becomes impossible due to excessive elastic recovery (Ref [20,](#page-25-0) [36](#page-25-0)). Subsequently, the effects of process and powder conditions (impact conditions) on the deposition and interface behavior at the moment of impact were investigated (Ref [43](#page-25-0)-[70](#page-26-0)). The factors related to particle velocity were found to be the most influential, although other factors, such as the relative hardness of the impact pairs, the particle/substrate temperature, the degree of oxidation, and the impact angle were also found to be important. Additionally, the characteristic microstructural evolution in kineticsprayed coating layers was elucidated, particularly for the interface region of the deposit (Ref [71-](#page-26-0)[102](#page-27-0)). In many cases, refined structure was observed along the bonding interface, which seemed to be formed by grain refinement and static/dynamic restoration (Ref [74,](#page-26-0) [77-84](#page-26-0)). Nanoscale phase transformation or formation of intermetallic compounds appeared in local parts of the interface region (Ref [85-87\)](#page-26-0). Various researchers also investigated the

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properties of kinetic-sprayed metal deposits (Ref [14](#page-25-0), [16,](#page-25-0) [103](#page-27-0)[-141](#page-28-0)). Specifically, some studies reported that the nanoscale mechanical properties of the interface region were superior to those of bulk metals (Ref [111](#page-27-0), [112,](#page-27-0) [114](#page-27-0)). Of course, at the microscale, enhanced bond strength/density/ hardness and high-density coating layers of metals were achieved without phase transformation or oxidation. However, the overall physical and chemical properties of a kinetic-sprayed coating layer cannot exceed the properties of bulk metals because of microdefects such as pores, cracks, and unbonded interface (Ref [12\)](#page-25-0).

Thus, until recently, most research focused on fundamental investigation of kinetic spraying technology, but a large portion of current research is devoted to industrial applications of the technology. Diverse industrial fields are interested in kinetic spraying, and there is even applicability of the technology as a method for three-dimensional (3D) forming (additive manufacturing) (Ref [142-146](#page-28-0)). However, industrial applications of kinetic spraying require profound understanding of the physical foundations of the process. When such understanding is achieved, the requirements of many industrial fields could be satisfied and a database of various materials and conditions systematically prepared. Nevertheless, no one has yet provided a well-organized summary of the correlations among impact conditions, interface reactions, microstructural evolution, and mechanical properties across the whole field of kinetic spraying technology.

This paper provides such an overview of these correlations for kinetic spraying of metals, discussing each relationship sequentially in detail. The overall discussion focuses on understanding the critical factors and their roles in each relationship.

# 2. Effect of Impact Conditions on Interface Reactions

# 2.1 Outline

Many factors can be considered as impact conditions, but it is unnecessary to examine them one by one. For the kinetic spray process, the most influential impact conditions for interface reactions are particle velocity, the particle and substrate temperatures, the relative hardness of the particles and substrate, the oxidation state, and the impact angle.

#### 2.2 Particle Velocity

Generally, in a kinetic spray process, feedstock material can be deposited on a target when the particle velocity is greater than the critical velocity. However, if particles are excessively accelerated, above the rebound velocity, the impacting particles rebound and successful deposition cannot be achieved. Thus, the interface reactions vary greatly according to the velocity of the particles in flight. The critical and rebound velocities result from the competition between the plastic deformation energy and the recoverable elastic energy. As is well known, the plastic deformation of the impact pair is the dominant factor contributing to bonding in the kinetic spray process, whereas the recoverable elastic force is the major hindrance to bonding. The variation of the plastic deformation and recoverable elastic energies with particle velocity is shown in Fig. 1 (Ref  $36$ ). In the first stage of this plot, the plastic deformation energy  $(E_P)$  is smaller than the recoverable elastic energy  $(E_R)$ , but the rate of increase of  $E_P$  is greater than that of  $E_R$ . However, at a certain point (the critical velocity), the situation reverses. This situation of greater  $E_P$  than  $E_R$  is not maintained continuously, because the rate of increase of  $E_P$  falls below that of  $E_R$ above the critical velocity. Thus, the rebound velocity is the point at which the predominance between  $E_P$  and  $E_R$ is re-reversed ( $E_P \leq E_R$ ). In other words, since the critical and rebound velocity are the primary turning points in terms of deposition behavior, the interface reactions at the moment of impact can be classified into three cases: (1) below critical velocity, (2) between critical velocity and rebound velocity, and (3) above rebound velocity.

In case I, an insufficient amount of impact energy is provided by the particles. The degree of plastic deformation does not induce sufficient thermal energy for satisfactory interfacial bonding. As a result, the particles bounce off due to the recoverable elastic force stored in the particles and substrate. In terms of energy, this phenomenon occurs because the plastic deformation energy (bonding  $E$ ) is smaller than the recoverable elastic energy (debonding  $E$ ). In this case, the effects of strain and strainrate hardening are more dominant than those of thermal softening, and adiabatic shear instability does not occur. Thus, there is no jetting (Ref [19](#page-25-0)) of the interface edge region and no sudden increase of strain or temperature [no thermal boost-up zone  $(Ref 21)$  $(Ref 21)$  $(Ref 21)$ ].

In case II, the particles are traveling fast enough to produce sufficient impact energy. A significant amount of plastic deformation, with consequent thermal energy, is



Fig. 1 Calculated plastic deformation (bonding) and recoverable elastic (debonding) energies for various sized Al–Si particles impacting onto a mild-steel substrate (Ref [36](#page-25-0))

<span id="page-2-0"></span>induced at the impact interface. Even though the shaperestoring phenomenon via the stored recoverable elastic energy at the interface region continues to occur in common with case I, the particles do not bounce off due to their intimate bonding with the substrate or other particles. In other words, in case II, the plastic deformation energy (bonding  $E$ ) is greater than the recoverable elastic energy (debonding  $E$ ). The thermal softening effect outweighs the strain and strain-rate hardening, and adiabatic shear instability is generated at the periphery of the interface region. Here, jetting occurs at the interface edge, with rapid increases in strain and temperature (occurrence of thermal boost-up zone).

In case III, plastic deformation of the impact pair is adequately induced by the impact of the rapidly accelerated particles. However, although the interfacial bond is well generated instantly, the bonded particles cannot withstand the subsequent rebound phenomenon caused by the excessive accumulated recoverable elastic energy in the impact pair. In case III, the amount of recoverable

elastic energy converted from the impact energy is greater than the plastic deformation energy. This implies that the plastic deformation  $E$  (bonding  $E$ ) again becomes smaller than the recoverable elastic energy (debonding  $E$ ). The jetting and strain/temperature increase induced by the adiabatic shear instability can be observed at the interface region, but successful deposition cannot be achieved.

### 2.3 Particle and Substrate Temperatures

To obtain a proper interface reaction for successful deposition, use of a particle velocity faster than the critical velocity is the most important requirement in kinetic spraying. To induce an interface reaction similar to that of case II but at a relatively slow particle velocity, the critical velocity of the chosen feedstock powder can be lowered by enhancing the plastic deformability of the feedstock powder and substrate. Doing so generates an appropriate interface reaction for successful deposition. Figure 2 shows the effect of powder preheating on the deposition



Fig. 2 Temperature distribution resulting from a multiparticle impact simulation of pure Ti at 600 °C and 2.5 MPa using N<sub>2</sub> gas (a) without preheating and (b) with preheating, and (c) the temperature and strain alteration vs. time for (a) and (b) (Ref [22](#page-25-0))

behavior of pure Ti in kinetic spray deposition (Ref [22](#page-25-0)). As shown in Fig.  $2(a)$  $2(a)$ , without preheating of the feedstock powder, the degree of particle plastic deformation was generally insignificant and no jetting occurred at the edge of the interface regions. In contrast, when the feedstock powder was preheated (Fig. [2](#page-2-0)b), the particles were severely plastic-deformed and some particles showed the jetting phenomenon. The temperature increase at the interface region became remarkable following preheating of the feedstock powder (Fig. [2c](#page-2-0)). In addition, the strain in the vicinity of the particle interface region was greatly enhanced by preheating the feedstock powder. This indicates that preheating of the feedstock powder can be helpful to minimize the critical velocity of a given material.

### 2.4 Relative Hardness of Particles and Substrate

Even for the same impact energy and identical feedstock powder, the interface reaction will vary depending on the target material. Figure 3 shows four kinds of impact behavior classified according to the relative hardness of the particles and substrate: soft–soft, hard–hard, soft–hard, and hard–soft (Ref [21](#page-25-0)). Naturally, the degree of plastic deformation is the same for similar pairs (soft–soft and hard– hard) (Fig. 3a, b), but it is biased toward the soft side for dissimilar pairs (soft–hard and hard–soft) (Fig. 3c, d). Thus, the division of impact energy consumption varies between similar and dissimilar pairs. For similar pairs, the impact energy is almost evenly consumed by the impacting and impacted parts, so it is somewhat difficult to heat either side of the interface to the melting point (comparatively high critical velocity). On the other hand, for dissimilar pairs, the impact energy is primarily consumed by the soft part, which is therefore more easily heated to the melting point (relatively low critical velocity). Practically, according to Bae et al. (Ref [21\)](#page-25-0), the critical velocity of Al on Al  $(-750 \text{ m/s})$ , soft–soft case) and Ti on Ti (~850 m/s, hard–hard case) is higher than that of Al on mild steel (~350 m/s, soft–hard case) or Ti on Al (~650 m/s, hard–soft case).

These four cases also differ with respect to high-temperature areas and aspects of energy development (Ref [21](#page-25-0)). As shown in Fig.  $4(a)$  $4(a)$  and (d), the soft–soft and hard– soft cases show the widest high-temperature areas, which means the widest chemical/metallurgical bonding area assisted by thermal energy. In both cases, the bonding (plastic deformation) energy is more dominant than the debonding (recoverable elastic) energy. It is obvious that the coating quality of these cases is usually excellent, as long as proper impact conditions are used. On the other hand, for the soft–hard case, the high-temperature area is only wide on the soft side (Fig. [4c](#page-4-0)). The bonding energy is not much higher than the debonding energy, unlike the soft–soft and hard–soft cases. Thus, the bonding quality is inferior to that of the prior cases in most situations. The hard–hard pair has the narrowest high-temperature area among the four pairs. Also, the bonding energy is lower than the debonding energy, regardless of particle velocity (Fig. [4](#page-4-0)d). Hence, successful deposition with high bonding quality is difficult to achieve for the hard–hard case.

#### 2.5 Oxidation State

To obtain successful bonding in the kinetic spray process, the amount of thermal energy converted from plastic deformation at the moment of impact should be higher than a certain threshold. If the impact energy is consumed unnecessarily, particle–particle and particle–substrate bonding cannot be achieved. Unnecessary use of impact energy in kinetic spraying occurs primarily in the oxide layer on the surface of the feedstock powder, which is naturally formed during storage. The oxide layer obstructs sound plastic deformation by consuming impact energy as it is broken. Thus, it is ideal to remove the oxide layer before deposition, although complete removal is practically impossible. Surface oxidation of metal powders is unavoidable because metals react easily with atmospheric oxygen.

Figure [5](#page-4-0) shows the effect of surface oxidation on the critical velocity of copper, 316L stainless steel, and nickel alloy (Ref [54\)](#page-26-0). With increasing oxygen content of the feedstock powder, the critical velocity increases continuously. For copper, the critical velocity rapidly develops from



Fig. 3 Four kinds of impact behavior classified according to the relative hardness of the impacting particles and impacted substrate: (a) soft–soft, (b) hard–hard, (c) soft–hard, and (d) hard–soft (Ref [21](#page-25-0))

<span id="page-4-0"></span>

Fig. 4 Temperature distribution induced by impact for the four cases: (a) soft–soft, (b) hard–hard, (c) soft–hard, and (d) hard–soft, with corresponding adhesive/rebound energy variations vs. particle velocity (Ref [21\)](#page-25-0)



Fig. 5 Effect of surface oxidation on critical velocity of copper, 316L stainless steel, and nickel alloy (Ref [54\)](#page-26-0)

310 m/s (0.02 wt.% oxygen) to 610 m/s (0.38 wt.% oxygen). Kang et al. also presented similar results in a study of the oxidation dependence of the critical velocity of aluminum (Ref [55\)](#page-26-0). The flattening ratio of the impacting Al particles decreased as the oxygen content increased (from 1.95 at 0.001 wt.% to 1.74 at 0.045 wt.%) (Fig. [6](#page-5-0)). This indicates that hindrance of plastic deformation at the interface region was intensified with increasing amount of oxides. Also, according to Ichikawa et al., when CoNiCrAlY was deposited on Inconel 625 using kinetic spraying, the deposition efficiency was degraded from 54.8 to 0.4% as the substrate surface became covered with an oxide layer (Ref [57\)](#page-26-0).

In addition, the oxide layer disturbs interfacial bonding by preventing direct contact between the impact pair. It is well known that debris from broken oxides is emitted from the contact surface through the jetting phenomenon at the periphery of the interface. However, complete emission of oxide debris is impossible, and the remaining oxides interrupt bonding between the impact pair. A schematic sequence of oxide breakage at the moment of impact is shown in Fig.  $7$  (Ref  $56$ ). Complete extrusion of the broken oxides cannot be achieved, and some debris remains in the form of inclusions at the bonding interface. Figure [8](#page-6-0) presents the bonding state between an Al particle and its substrate achieved in practice (Ref [55](#page-26-0)). A hollow gap was formed between the Al particle and the substrate in the

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Fig. 6 Cross-sectional scanning electron microscopy (SEM) images of single deposited Al particles with oxygen content of (a) 0.001 wt.%, (b) 0.012 wt.%, (c) 0.023 wt.%, and (d) 0.045 wt.% (Ref [55\)](#page-26-0)

form of pores with remaining aluminum oxides. Thus, aluminum oxides might discourage interfacial bonding by blocking contact with the bare surface.

Dynamic oxidation of the surface can also be induced during the deposition stage. Kim et al. (Ref [58](#page-26-0)) reported that the fresh surface of kinetic-sprayed Ti particles was dynamically reoxidized after breakage and ejection of the preexisting oxide layer through jet formation with adiabatic shear instability. Figure [9](#page-7-0) shows the interfacial state of the subsequently deposited Ti particles. As shown in Fig.  $9(f)$  $9(f)$  and (g), an amorphous TiO<sub>2</sub> layer tens of nanometers thick was present in the interface region. This  $TiO<sub>2</sub>$  layer was newly formed rather than preexisting, because it was continuous without any trace of impact damage (Fig. [9](#page-7-0)e). It is reasonable to believe that kineticsprayed Ti particles could be dynamically oxidized during the deposition stage, considering that Ti and Al possess high oxygen affinity. Nonetheless, despite the oxide layer at the interface, the bonding of two subsequently deposited particles was very intimate (Fig. [9](#page-7-0)b, g). Kim et al. insisted that thin amorphous oxides do not hinder bonding of deposited particles due to their high cohesive energy to

crystalline material with a mismatched lattice (Ref [147,](#page-28-0) [148](#page-28-0)). However, a clear bonding mechanism was not discussed in their research.

# 2.6 Impact Angle

Generally, in the kinetic spray process, particles impact the target at an angle of  $\sim 90^\circ$ , since the deposition efficiency decreases as the impact angle deviates from 90°; That is, at other angles, some portion of the impact energy is used in processes other than plastic deformation, preventing the development of adiabatic shear instability in the interface region. Of course, the tangential behavior of the impacting particles induces generation of thermal energy via friction, but its effects on bonding are inferior to those of plastic deformation with adiabatic shear insta-bility. Figure [10](#page-7-0) shows the degradation of the deposition efficiency and porosity for pure Ti as the impact angle deviates from  $90^{\circ}$  (Ref [60](#page-26-0)). The deposition efficiency and porosity are not degraded much until  $70^{\circ}$ , but worsen rapidly thereafter.

<span id="page-6-0"></span>

Fig. 7 Schematic diagram of break-up and extrusion of oxide layer at moment of impact in kinetic spraying (Ref [56](#page-26-0))



Fig. 8 Bright-field transmission electron microscopy (TEM) image of interface between kinetic-sprayed Al particles with 0.045 wt.% oxygen content (Ref [55\)](#page-26-0)

Deviation of the impact angle from  $90^\circ$  also results in inhomogeneous particle deformation (Ref  $60$ ). As shown in Fig. [11](#page-8-0), jetting at the particle interface edge became biased to one side as the impact angle deviated from 90. In this case, the bonding quality was inevitably worsened because the interfacial bonding on the side with a relatively small degree of plastic deformation was poor. Thus, particles were bonded on one side, but not the other.

# 2.7 Interaction Between Impact Condition and Feedstock Properties

In previous sections, the effect of several impact conditions on the interface reactions in kinetic spraying was discussed. However, practically, the interface reactions depend on the interaction between the properties of the feedstock powder and the given impact conditions; For instance, criteria based on conditions such as the critical/

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Fig. 9 Kinetic-sprayed Ti particles: (a) scanning transmission electron microscopy (STEM) bright-field image, (b) magnified image of white boxed area in (a), results of element mapping for (c) tungsten, (d) titanium, and (e) oxygen, and (f, g) high-resolution (HR) images of areas marked in (b) (Ref [58](#page-26-0))



Fig. 10 Degradation of deposition efficiency and porosity for pure Ti as impact angle deviates from  $90^{\circ}$  (Ref [58\)](#page-26-0)

rebound velocity, relative hardness, and degree of oxidation are determined by the material properties of the impact pair. Thus, it is helpful to understand the interactions between the impact conditions and material properties. Although all such interactions are important, only interactions between the particle velocity and material properties are discussed in this section, because the others (i.e., for relative hardness and oxidation state) are simple and obvious.

As is well known, the critical and rebound velocities result from the competition between the bonding (plastic deformation) and debonding (recoverable elastic) energies. Therefore, the interaction with respect to the particle velocity can be distinguished based on the elastic/plastic behaviors of the impact pair, which are determined by their material properties. Applying the Johnson–Cook plasticity model in a particle impact simulation using the finite-element method (FEM) allows comparison of the elastic/plastic behaviors of impact pairs, based on their elastic moduli and material-specific parameters in the equation. The Johnson–Cook plasticity model is as follows (Ref [149](#page-28-0)):

$$
\sigma = \left[A + B\varepsilon_{\mathrm{p}}^{n}\right]\left[1 + C\mathrm{ln}\left(\frac{\dot{\varepsilon}_{\mathrm{p}}}{\dot{\varepsilon}_{0}}\right)\right][1 - (T^{*})^{m}],
$$

where  $\sigma$  is the equivalent flow stress,  $\varepsilon_p$  and  $\varepsilon_p$  are the equivalent plastic strain and strain rate,  $\dot{\epsilon}_0$  is the normalizing reference strain rate, and  $T^*$  is the normalized temperature.  $A$ ,  $B$ ,  $C$ ,  $n$ , and  $m$  are characteristic material parameters. A is the yield stress in a quasistatic simple tension or compression test,  $B$  is the strainhardening coefficient, C is the dimensionless strain-rate hardening coefficient, and  $n$  and  $m$  are the power exponents of the strain-hardening and thermal-softening terms.

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Fig. 11 Morphology of Ti particles after impact on low-carbon steel substrate (1000  $\degree$ C, 4 MPa process gas) at various spray angles: (a) 90°, (b) 85°, (c) 70°, and (d) 45° (Ref [58\)](#page-26-0)

Table 1 Numerical values for Ni, Ti, and Ta for particle impact simulation using FEM

	Ni	Ti	Тя
Density (kg m <sup><math>-3</math></sup> )	8890	4510	16,650
Young's modulus (GPa)	207	116	186
Poisson's ratio	0.31	0.34	0.35
Heat capacity (J kg <sup>-1</sup> K <sup>-1</sup> )	456	528	153
Melting point $(K)$	1726	1923	3269
$A$ (MPa)	163	806.57	340
B(MPa)	648	481.61	260
n	0.33	0.319	0.32
C	0.006	0.0194	0.042
m	1.44	0.655	0.88
Ref. temp. (K)	298	298	300
Ref. strain rate $(s^{-1})$	1		4

Here, Ni, Ti, and Ta are considered as representative metals. Table 1 presents their numerical values applied in a particle impact simulation using FEM (Ref [21](#page-25-0)). If the ability to be heated to near the melting point at the interface region is considered as a criterion for successful bonding, each material property can be distinguished according to whether or not it is advantageous. The influence of each material property is summarized in Table 2. Whether corresponding material property is advantageous or not is marked with symbol of 'O' in the table.

Compared with Ti and Ta, Ni has the benefit of possessing low yield strength and melting point, and a high thermal-softening exponent. However, its elastic modulus





and strain-hardening coefficient/exponent are higher than those of the other metals, which is disadvantageous for plastic deformation. Figure [12](#page-9-0) shows kinetic-sprayed pure Ni for various process conditions. The Ni was hard to deposit successfully at 600 °C and 2.5 MPa using  $N_2$  gas (condition C1) (Fig.  $12a$  $12a$ , b) (Ref  $52$ ). The point to note here is that deposition of pure Ni was not that successful even at  $600 \degree C$  and 1.5 MPa with He gas (condition C2) (Fig. [12c](#page-9-0)). Also, the surface of the Ni coating layer was severely eroded by subsequent impacts from accelerated Ni particles (Fig. [12d](#page-9-0)). This indicates that the particle velocity of the later particles approached the rebound velocity of Ni, despite the relatively small difference in

<span id="page-9-0"></span>

**Fig. 12** Cross-sectional SEM images of kinetic-sprayed pure Ni at (a) 600 °C and 2.5 MPa with N<sub>2</sub> gas (condition C1) and (c) 600 °C and 1.5 MPa with He gas (condition C2), and magnified images of the top regions (b, d) shown in (a) and (c), respectively (Ref [52\)](#page-26-0)

particle velocity between conditions C1 and C2 (~600 m/s and ~770 m/s, respectively). Such a narrow gap between the critical and rebound velocities implies that the effect of debonding (recoverable elastic) energy was rapidly strengthened relative to that of the bonding (plastic deformation) energy with an increase in particle velocity, indicating a powerful hardening effect. The primary cause of such a severe hardening effect is the high strain-hardening coefficient/exponent of Ni. Thus, when using Ni, it is more efficient to investigate other ways to widen the window of sprayability (WS) (Ref [20\)](#page-25-0). Considering that pure Ni has a relatively high thermal-softening exponent but a low melting point, preheating the feedstock powder could be more effective than increasing the particle velocity. Figure [13](#page-10-0) presents the stress and temperature distributions for an impact pair when using 600 °C and 2.5 MPa with  $N_2$ gas, with and without preheating, and when using  $600 \degree C$ and 1.5 MPa with He gas (Ref [52](#page-26-0)). The stress distribution is much more severe without preheating (Fig. [13](#page-10-0)a, c) than with preheating (Fig. [13b](#page-10-0)). Also, the temperature distribution in the interface region is highest among the three cases when powder preheating is applied. In practice, the coating thickness (deposition efficiency) and bond strength were dramatically improved as the preheating temperature was increased (Fig. [14](#page-10-0)). In Fig. 14(e), the coating thickness and bond strength of each pre-heating condition are presented in the form of curve and bars, respectively.

Compared with Ni and Ta, pure Ti presents almost no advantageous features, but suffers from various disadvantages (i.e., high heat capacity and yield strength, low density and thermal-softening exponent). Nevertheless, deposition of Ti is not as difficult as it may seem, since a pure Ti layer can be fabricated even at particle velocity lower than the critical velocity reported in literature. Figure  $15(a)$  $15(a)$  and (d) show a Ti layer above a certain degree of thickness, despite use of a particle velocity lower than the critical velocity reported in Ref [22](#page-25-0). Surprisingly, the deposition efficiency was about 80%, which is the value generally obtained under optimal conditions in most cases of kinetic spraying. According to Bae et al., this result was induced by the higher adiabacity of pure Ti relative to other metals (i.e. Cu, Al, Ni, and Ta). Ti retains its high adiabacity, since its thermal conductivity is lower than that of other metals (Ref [150\)](#page-28-0). However, as presented in Fig.  $15(a)$  $15(a)$  and (d), the quality of the Ti coating layer was poor, even though deposition was achieved. To achieve a Ti layer with low density and adequate bonding, the particles should be accelerated above the critical velocity, or the critical velocity of Ti must be lowered by powder preheating. The results with powder preheating (at  $600 °C$ ) and enhanced process conditions (changing from 600 °C and 2.5 MPa with  $N_2$ gas to  $600 °C$  and  $1.5 MPa$  with He) are shown in Fig.  $15(b, e)$  $15(b, e)$  and  $(c, f)$ . The fabricated Ti layers became denser and the bonding state was improved in both cases. Here, the improvement in coating quality by enhancing the particle acceleration was more obvious than that achieved by powder preheating, in contrast to the case of Ni. For pure Ni, increasing the particle velocity was not optimal because of the high strain-hardening coefficient/exponent and elastic modulus of Ni. However, Ti does not possess outstanding strain and strain-rate

<span id="page-10-0"></span>

Fig. 13 von Mises stress and temperature distributions of the Ni–mild steel impact pair at 600 °C and 2.5 MPa with N<sub>2</sub> gas (a) with and (b) without preheating and at (c)  $600 °C$  and 1.5 MPa with He gas (Ref [52](#page-26-0))



Fig. 14 Cross-sectional SEM images of kinetic-sprayed pure Ni with powder preheating at temperature of (a) 300 °C, (b) 400 °C, (c) 500 °C, and (d) 600 °C, and (e) variation in coating thickness and bond strength with powder preheating temperature (Ref [52\)](#page-26-0)

coefficient/exponent. Also, the thermal-softening effect is unremarkable, because of the quite low thermal-softening exponent of Ti (0.655) relative to that of other metals (generally  $\sim$ 1.0). Figure [16](#page-12-0) clearly suggests that increasing the particle acceleration is better than powder preheating for improvement of the bonding of kineticsprayed pure Ti (Ref [22\)](#page-25-0). The degree of strain and the temperature increase at the interface region were more noticeable with particle acceleration than with powder preheating.

<span id="page-11-0"></span>

**Fig. 15** Cross-sectional SEM images of kinetic-sprayed pure Ti at 600  $^{\circ}$ C and 2.5 MPa with N<sub>2</sub> gas (a) with and (b) without preheating at 600 °C and at (c) 600 °C and 1.5 MPa with He gas and (d–f) their magnified images (Ref [22\)](#page-25-0)

Pure Ta has some advantages, such as high density and low heat capacity and strain-hardening coefficient, relative to Ni and Ti. However, its melting point and strain-rate hardening coefficient are higher than those of Ni and Ti. In particular, the extremely high melting point of Ta  $(3269 K)$  compared with Ni  $(1726 K)$  or Ti  $(1923)$  is the most important obstacle to increasing the interfacial temperature to near its melting point. Nevertheless, the critical velocity of pure Ta is  $\sim$ 550 m/s, including it in the group of low-critical-velocity metals compiled by Bae et al. (Ref [21](#page-25-0)). Practically speaking, pure Ta was well deposited at 550 °C and 3.0 MPa with  $N_2$  gas when the particle velocity [numerically calculated by equation (Ref [4,](#page-24-0) [47](#page-25-0))] was  $\sim$ 555 m/s (Fig. [17a](#page-12-0)). This indicates that the advantageous features of Ta (i.e., high density, low heat capacity and strain-hardening coefficient) dominate its disadvantageous features (i.e., high melting point and strain-rate hardening coefficient), resulting in a relatively low critical velocity. However, although a certain degree of deposition was achieved, the deposition efficiency of pure Ta developed from ~50 % to more than 80 % on enhancement of the process conditions (from 550  $\degree$ C and 3.0 MPa with  $N_2$  gas to 550 °C and 2.0 MPa with He gas) and powder preheating (Fig. [17f](#page-12-0)). Here, in Fig. [17](#page-12-0)(f), bars are corresponding to the wt.% of O and N elements, and upper and lower curves are relevant to the deposition efficiency and hardness, respectively. The point to note here is that powder preheating was more effective than enhancing the process conditions, probably because of the high strain-rate hardening coefficient of Ta.

In summary, whether an interaction that is beneficial for successful deposition will be induced is determined by the impact conditions and material properties of the feedstock powder. Even though these interactions cannot be obviously classified into certain types for all metals, some predictions can be deduced using the materialspecific parameters of the Johnson–Cook plasticity model and the general properties (i.e., density, heat capacity/conductivity, melting point, and elastic modulus) of the metal feedstock powder.

### 2.8 Summary

As discussed above, interface reactions in kinetic spray technology are strongly affected by the impact conditions (i.e., particle velocity, particle and substrate temperatures, relative hardness of particles and substrate, oxidation state, and impact angle) and feedstock properties.

In terms of particle velocity, the interface reaction at the moment of impact can be classified into three types: (1) below critical velocity, (2) between critical velocity and rebound velocity, and (3) above rebound velocity. In case I, particles bounce off because of insufficient plastic deformation (bonding  $E <$  debonding  $E$ ). In case II, interfacial bonding is generated and is assisted by adiabatic shear instability (bonding  $E \ge$  debonding  $E$ ). In case III, bonded particles cannot withstand the rebound phenomenon due to the dominance of recoverable elastic energy (bonding  $E <$  debonding  $E$ ).

Preheating the particles and substrate is helpful for generating a good bonding state with relatively low particle velocity. The critical velocity for certain materials can be reduced by improving the plastic deformability by preheating the particles and substrate.

The relative hardness of the impact pair (particles and substrate) is also an influential factor and can be divided

<span id="page-12-0"></span>

**Fig. 16** Temperature distribution resulting from multiparticle impact simulation of pure Ti at (a) 600 °C and 2.5 MPa with N<sub>2</sub> gas with preheating at 600 °C and at (b) 600 °C and 1.5 MPa with He gas without preheating, and (c) the temperature and strain variation vs. time for (a) and (b)  $(Ref 22)$  $(Ref 22)$ 



Fig. 17 Cross-sectional SEM images of kinetic-sprayed pure Ta under various process conditions: at 550 °C and 3.0 MPa with N<sub>2</sub> gas with (a) no preheating, (b) 400 °C preheating, (c) 500 °C preheating, and (d) 600 °C preheating, and (e) at 550 °C and 1.5 MPa with He gas, and (f) deposition efficiency according to nitrogen/oxygen content for each process condition

<span id="page-13-0"></span>into four cases: soft–soft, hard–hard, soft–hard, and hard– soft. The division of impact energy consumption differs between similar pairs (soft–soft and hard–hard) and dissimilar pairs (soft–hard and hard–soft). It is easier for dissimilar pairs to be heated to the melting point at the interface region than for similar pairs. Also, the areas of high temperature and energy development are distinct in these four cases. The area of high temperature is wider with dominance of bonding (plastic deformation) energy in the following sequence: (1) soft–soft and hard–soft cases, (2) soft–hard case, and (3) hard–hard case.

Furthermore, the effect of an oxide layer should be considered because it naturally forms on the surface of the feedstock powder during storage. The oxide layer functions as an obstacle to plastic deformation by consuming impact energy as it breaks. Also, ejection of broken oxides by jetting cannot be completely achieved, so some debris of broken oxides remaining at the bonding interface can disturb interfacial bonding by preventing direct contact between the particles and substrate.

As the impact angle deviates from  $90^\circ$ , the deposition efficiency degrades because of the decrease in particle impact velocity. Deviation of the impact angle from 90° also leads to inhomogeneous particle deformation, which reduces the bonding quality.

Practically, interface reactions depend on the interaction between the properties of the feedstock powder and the given impact conditions. These interactions are fairly complicated, but it is possible to compare the elastic/plastic behaviors of impact pairs based on their material-specific parameters and general properties (i.e., density, heat capacity/conductivity, melting point, and elastic modulus).

# 3. Dependence of Microstructural Evolution on Interface Reactions and Material Properties

### 3.1 Outline

The previous section presented the relationships among various impact conditions and interface reactions for different material properties. In the relationships between interface reactions and microstructural evolution, it is most important whether the induced interface reaction properly satisfies the requirements for successful bonding. With inappropriate interface reactions, bonding between impact pairs does not occur properly, and no coating layer is formed. In this section, microstructural evolution refers not to the coating layer but to the individual particles. However, since the microstructure considered in this section corresponds to a successfully fabricated coating layer, the discussion is limited to cases of interface reactions suitable to induce adequate bonding. Also, several review papers have already focused on the microstructure of the kinetic spray process (Ref  $151$ ,  $152$ ), and they can provide good understanding of the basic mechanics involved.

#### 3.2 Inhomogeneity of Kinetic-Sprayed Deposit

In kinetic spraying, reactions directly related to bonding intensively appear at the interface region. Thus, there is always a microstructural difference between internal particles and the interface region, as visible at a glance from Fig. [18,](#page-14-0) which presents cross-sectional images of pure Cu feedstock and a kinetic-sprayed Cu deposit (Ref [74](#page-26-0)). The regions marked A-C correspond to internal particle areas, while D and E are at the interface region. The grains in the internal particle regions (A-C) are relatively well preserved compared with the feedstock particle (Fig. [18](#page-14-0)a), but those in the interface regions (D and E) are refined (Fig. [18](#page-14-0)b). Such grain refinement at the interface region must be induced by large amounts of plastic deformation and thermal energy. In most cases, grains are refined by static recovery/recrystallization (SRV/SRX) or dynamic recovery/recrystallization (DRV/ DRX) (Ref [74,](#page-26-0) [77-84\)](#page-26-0). Detailed discussion on grain refinement is presented in section 3.3.

Even within the interface region, the microstructure is highly distinctive according to the location (i.e., south pole and intermediate and edge parts of the particle interface region) because all the critical reactions related to bonding are induced at the periphery of the interface region. In particular, the thermal energy and stored strain energy generated at the moment of impact are specifically concentrated in the edge parts of the bonding interface due to the occurrence of adiabatic shear instability. On the other hand, at the south pole region, thermal energy is scarcely induced, and the effect of debonding (recoverable elastic) energy is dominant over the bonding (plastic deformation) energy. Figure [19](#page-14-0) presents high-magnification images of the parts corresponding to the periphery (marked D in Fig.  $18$ ) and south pole (marked E in Fig.  $18$ ) of the interface region (Ref [74](#page-26-0)). In the periphery of the interface region (Fig. [19a](#page-14-0)), intimate bonding is generated, as seen from the indistinct particle–substrate boundary. A thin layer of nanoscale grains seems to be induced along the bonding interface by the interaction between the large amounts of thermal and stored strain energy. On the other hand, the south pole region shows a bimodal structure with an upper nanograin region and lower submicrongrain region with deformation twins (Fig. [19b](#page-14-0)). Perhaps the thermal energy was insufficient to generate recrystallization at the lower part of the south pole region, where the degree of deformation must be larger than in the upper part.

## 3.3 Microstructural Alteration Primarily Induced by Thermal Energy

This section mainly focuses on the microstructural evolution at the periphery of the interface region. The main point to consider here is how the feedstock material responds to the energies induced by supersonic impact. In the case of kinetic spraying, the major energies that influence the microstructural evolution are the stored strain energy and thermal energy. At the moment of deposition, the strain energy is first stored in the interface

<span id="page-14-0"></span>

Fig. 18 Cross-sectional SEM images of (a) feedstock Cu particles and (b) kinetic-sprayed Cu particles (Ref [72\)](#page-26-0)



Fig. 19 Magnified cross-sectional SEM images of (a) region D and (b) region E in Fig. [16](#page-12-0) (Ref [72\)](#page-26-0)

region of the deposited particles, then thermal energy is induced in accordance with the progress of plastic deformation. Simultaneously, the feedstock powder responds to the given thermal and stored strain energy in one of two ways: (i) reacting with thermal energy or (ii) reacting with thermal and stored strain energy.

Generally, the edge of the interface region heats instantly to the melting point at the moment of successful deposition. Although the absolute quantity of induced thermal energy is insufficient to produce a reaction through a large area at the interface region, it is sufficient to generate nanoscale phase transformation and the formation of intermetallic compounds at highly localized parts of the interface. Figure [20](#page-15-0) shows nanoscale amorphization of Al at the interface (Ref  $85$ ). A 3-4-nm amorphous Al layer (white box numbered 1) was created with an indistinguishable transient boundary from the amorphous to crystalline region (white box numbered 2), which seemed to be stabilized by oxygen after its formation. This amorphization primarily resulted from the rapid temperature increase to the melting point, followed by rapid quenching. Apparently, nucleation sites were crystallographic planes weakened by high-strain and strainrate plastic deformation.

Similar to the amorphization phenomenon, localized crystallization can occur at the periphery of the interface region when amorphous material is used as the feedstock powder in kinetic spraying. Yoon et al. observed <span id="page-15-0"></span>nanocrystallization of the amorphous phase in a kinetic-sprayed deposit of CuNiTiZr (Fig. [21\)](#page-16-0) (Ref [86\)](#page-26-0). Crystallized CuNiTiZr appeared at the center of the amorphous region in nanoscale, as shown in Fig.  $21(a)$  $21(a)$  and (b). The main cause of the crystallization was induced thermal energy, assisted by a lowered activation energy resulting from high-strain and strain-rate plastic deformation. The degree of crystallization was enhanced as the process gas pressure was increased (Fig. [21c](#page-16-0), d), providing more energy to the interface region of the CuNiTiZr deposit.

In addition, formation of nanoscale intermetallic compounds can be induced at the interface region when the kinetic spray deposition is conducted using a dissimilar impact pair (e.g., Al particles on Cu substrate). Figure [22](#page-16-0) shows Al-Cu intermetallic compounds created at the interface in nanoscale when Al particles were successfully deposited on Cu substrate (Ref  $87$ ). Al<sub>4</sub>Cu<sub>9</sub> and Al<sub>2</sub>Cu phases with size of ~400 nm were identified by fast Fourier transform (FFT) pattern analysis in the region marked in Fig. [22](#page-16-0). The sequence of intermetallic compound formation was as follows: (1) chemical/metallurgical bonding was created at the interface of the impact pair, (2) diffusion of elements was induced through the bonding interface, and (3) reaction occurred between the elements in the diffusion layer, assisted by the thermal energy provided by severe plastic deformation.

# 3.4 Microstructural Alteration Caused by Thermal and Stored Strain Energies

3.4.1 Grain Refinement and Static/Dynamic Recovery/ Recrystallization. Alterations in microstructure induced primarily by thermal energy were discussed above. There is assistance from stored strain energy, but the thermal energy effect is much more dominant in the aforementioned microstructural reactions (i.e., phase transformation and intermetallic compound formation). In this section, various microstructural evolutions dependent on both the thermal and stored mechanical energies are discussed.

As mentioned repeatedly, severe plastic deformation is induced at the interface region of kinetic-sprayed deposits, assisted by adiabatic shear instability. This plastic deformation converts a large amount of the impact energy into strain energy that is stored at the interface region in the form of defects, such as dislocations and shear bands, during the progress of deposition. This means that the interface region becomes unstable because of the high level of stored strain energy. Thus, some microstructural phenomena, such as grain refinement and static/dynamic recovery/recrystallization, stabilize the interface region, when sufficient thermal energy is available.

Grain refinement is induced dynamically during the deposition stage. In this process, countless dislocations are generated at the interface region by high-strain and strainrate plastic deformation. Since a specific arrangement is more stable than a random dispersion, dislocation cells are created via rearrangement of dislocations. In this way, numerous subgrains are formed in the grains near the bonding interface. The sequence of grain refinement is summarized in Fig. [23](#page-17-0), as suggested by Zou et al. (Ref [79](#page-26-0)). As shown in Fig. [23](#page-17-0), subgrains are mainly formed along grain boundaries, where dislocations generally accumulate. The subgrains have elongated shape because of the grain refinement and plastic deformation of particles,



Fig. 20 TEM images at Al particle–substrate bonding interface: (a) sampled by focused ion beam (FIB), (b) high-resolution electron microscopy (HREM) image at Al particle–substrate bonding interface, and (c) highly magnified image of boxed region in (b) with fast Fourier transform (FFT) patterns of selected area (Ref [83](#page-26-0))

<span id="page-16-0"></span>

Fig. 21 HREM images of crystalline area in kinetic-sprayed CuNiTiZr deposit: (a) HREM image of nanosize crystals in Cu-based bulk metallic glass (BMG) coating (under gas pressure of 1.5 MPa), (b) highly magnified HREM image and FFT pattern in selected area of (a), (c) HREM image of nanosize crystals in Cu-based BMG coating (under gas pressure of 3.0 MPa), and (d) highly magnified HREM image and FFT pattern in selected area of (c) (Ref [84\)](#page-26-0)



Fig. 22 TEM images and FFT pattern analysis of Al-Cu intermetallic compounds at the interface region: (a) TEM image of interface region and (b) highly magnified HREM images with FFT pattern of each intermetallic compound for the boxed area in (a) (Ref [85](#page-26-0))

<span id="page-17-0"></span>which can induce subsequent elongation of subgrains after their creation.

Furthermore, relief of the stored strain energy results from static/dynamic recovery/recrystallization at the interface region. Simply, DRV/DRX can be induced concurrently with deformation during the deposition stage, and SRV/SRX can be generated at the interface region with additional thermal energy provided by subsequent particle impacts.

DRV is the phenomenon of dislocation dipole annihilation. Here, a dislocation dipole corresponds to two dislocations of opposite sign located on different slip planes. Annihilation of dislocation dipoles is primarily induced by interactions between stored network dislocations and mobile dislocations. Specifically, dislocation dipoles are annihilated by dislocation climb that results from nonzero elastic forces perpendicular to their slip planes (Ref [153-](#page-28-0) [156](#page-28-0)). In the case of kinetic spraying, DRV can occur actively at the interface region because the dislocation climb phenomenon is sufficiently activated by the thermal energy induced by the severe plastic deformation.

In DRX, relief of stored strain energy is achieved by formation of dislocation-free grains. DRX induction is divided into two types according to its mechanism: migrational and rotational. Migrational DRX is primarily accomplished by migration of high-angle grain boundaries, whereas rotational DRX occurs through rotation of subgrains. For the kinetic spray process, the latter mechanism is more likely to apply than the former, because the mobility of high-angle grain boundaries is very low under high-strain-rate conditions (above  $4.0 \times 10^4$  s<sup>-1</sup>) (Ref [157](#page-28-0)). When particles impact a target in kinetic spraying, the strain rate of plastic deformation at the interface region definitely exceeds the strain-rate criterion for mobility of high-angle grain boundaries. Thus, it is reasonable to explain DRX at the interface region in a kinetic spray process using the theory of rotational DRX. Zou et al. explained the detailed sequence of rotational DRX in kinetic spraying (Fig. [24](#page-18-0)) (Ref [79](#page-26-0)). Numerous subgrains are created by the grain refinement phenomenon near grain boundaries (Fig. [24](#page-18-0)a). Afterwards, the created subgrains are broken into much smaller pieces, as the plastic deformation of the impacted particles progresses, even after the formation of subgrains (Fig. [24](#page-18-0)b). Consequently, the broken pieces of subgrains are rotated by continuous exertion of shear forces, and strain-free grains are formed with high-angle grain boundaries (Fig. [24c](#page-18-0)).

SRV and SRX in the kinetic spray process are generated by residual heat or additional supply of thermal energy from subsequent particle impacts. In the case of residual heat, SRV/SRX progresses during the cooling stage after adiabatic heating of the interface at the moment of deposition. In this case, the residual heat needs to be maintained for a sufficiently long time. Therefore, occurrence of SRV/SRX is more likely in high-adiabacity (low thermal conductivity) materials such as Ti. Here, the adiabacity indicates how long a material retains the generated thermal energy in a specific region. Thus, for SRV/ SRX via additional supply of thermal energy, heat is provided by additional plastic deformation of a predeposited layer through successive particle impacts or through transfer from the upper deposited layer of subsequent particles. In this case, SRV and SRX can occur even in materials with high thermal conductivity (e.g., Al and Cu). The additional plastic deformation from subsequent particle impacts has to be sufficient in terms of quantity and persistence. According to Kang et al., SRV and SRX occur in sequence (Fig. [25\)](#page-18-0) (Ref [77](#page-26-0)). As shown in Fig.  $25(a)$  $25(a)$ , SRV and SRX are not induced in the moment preceding deposition, and only some parts of the lower bonding interface experience DRV or DRX. However, when deposition of subsequent particles occurs, heat is generated by additional plastic deformation of the predeposited layer and is transferred from the upper deposited layer (Fig. [25b](#page-18-0)). As a result, static restoration (SRV/SRX) occurs in regions with strain energy induced by the impact of subsequent particles above a certain threshold (Fig. [25c](#page-18-0)).

3.4.2 Different Microstructural Responses According to Material Properties. As explained above, grain refinement and static/dynamic restoration phenomena can arise at the interface region if it retains a high degree of strain energy induced by high-strain and strain-rate plastic deformation. However, the definitive microstructural evolution is determined by the interaction of the given thermal and stored strain energies with the material properties of the particular feedstock powder. All such interactions cannot be classified completely into specific cases, but one point of obvious importance is whether the energy requirements are satisfied to induce an alteration in the microstructure of a given feedstock material. Hence, it is helpful to examine practical cases to understand the microstructural relationship and predict the microstructural changes.

Figure [26](#page-19-0) shows the microstructure near the interface region of kinetic-sprayed Ni, Ti, and Ta layers (Ref [78,](#page-26-0) [79](#page-26-0)). In all cases, refined microstructures appeared along the bonding interface. Although refinement of the



Fig. 23 Sequence of subgrain formation: (a) nondeformed microstructure with low dislocation density, (b) accumulation of dislocations with shearing, and (c) resultant grain refinement with formation of elongated subgrains (Ref [77\)](#page-26-0)

<span id="page-18-0"></span>

Fig. 24 Sequence of rotational DRX: (a) formation of elongated subgrains, (b) break-up of elongated subgrains, and (c) rotation of broken pieces from subgrains (Ref [77\)](#page-26-0)



Fig. 25 Sequence of SRV and SRX by impact of subsequent particles: (a) microstructural state after deposition of first layer, (b) additional heating by impact of subsequent particles, and (c) occurrence of SRV and SRX using provided thermal energy (Ref [75\)](#page-26-0)

microstructure can be induced by both static and dynamic restoration, it is obvious that such grain refinement was caused by DRV/DRX rather than SRV/SRX. Almost no strain or defects appear in newly created refined grains produced by SRV and SRX. However, it is obvious that a considerable number of defects remained in the refined grains nearby, which were apparently generated after grain refinement (left side of Fig. [26a](#page-19-0)). In the case of the grains created by DRV and DRX, generation of additional defects in the grains was inevitable because plastic deformation proceeds simultaneously with DRV/DRX, as shown clearly in the pattern quality map of kineticsprayed Ni obtained by electron backscatter diffraction (EBSD) (right side of Fig. [26a](#page-19-0)) (Ref [79\)](#page-26-0). The pattern quality of the newly formed nanograins was extremely poor, which means there were a significant number of defects in the refined grains near the interface region. Strain-free grains were also rare in the kinetic-sprayed Ti and Ta deposits (left side of Fig. [26b](#page-19-0) and c) (Ref [78](#page-26-0)). However, particularly for Ti, some dislocation-free nanograins (~10–100 nm) were found very near the bonding interface (right side of Fig. [26](#page-19-0)b) (Ref [78\)](#page-26-0). Such nearly nondefective grains might have been created by static restoration. As explained in section [3.4.1](#page-15-0), SRV and SRX can be induced in materials that possess high adiabacity, which can retain the thermal energy resulting from plastic deformation for a sufficiently long time during the cooling stage of the deposited particles.

The microstructures in the interface region for facecentered cubic (FCC) metals (Al, Ni, and Cu) are pre-sented in Fig. [27](#page-20-0) (Ref [83](#page-26-0)). Although Al, Ni, and Cu share the same crystal structure, the freedom of dislocation movement varies among them and critically influences strain restoration. Thus, the stacking fault energy (SFE) can significantly influence the microstructural evolution in the kinetic spray process, since the dislocation movement and interaction among dislocations vary according to the SFE of the feedstock powder. Therefore, the microstructures of Al, Ni, and Cu are distinctive from one another, as shown in Fig. [27.](#page-20-0) In the Al deposit, an equiaxed nanocrystalline structure was created (marked C in Fig. [27](#page-20-0)a), which must have been the result of static/dynamic restoration at the moment of deposition. However, recovery and recrystallization have a competitive relationship, because both are driven by consumption of stored strain energy. It is therefore unavoidable that one will be inferior when the other is superior. In the case of Al, the feedstock metal with the highest SFE, creation of a stacking fault plane can increase the energy state considerably (Ref [83](#page-26-0)). Compared with Ni and Cu, dislocations generated in Al are difficult to dissociate in particle dislocations due to the high SFE, guaranteeing freedom of dislocation movement. Thus, in Al, recovery is preferred over recrystallization, because of its dependence on dislocation movement. It is more reasonable to conclude that the grain refinement in the kinetic-sprayed Al deposit was induced by recovery rather than recrystallization. A tangled structure corresponding to regions A and B in Fig. [27](#page-20-0)(a) was also observed.

On the other hand, Ni possesses lower SFE than Al, so recrystallization becomes more competitive with recovery than for Al. The refined structure on the left side of Fig. [27](#page-20-0)(b) might have been created by both recovery and recrystallization, although Borchers et al. insisted that the restoration was mainly caused by recrystallization (Ref [83](#page-26-0)). Interestingly, coffee-bean-like structures appear on the right side of Fig. [27\(](#page-20-0)b). According to Borchers et al., this characteristic microstructure is formed when intrinsic/ extrinsic dislocation loops are generated in the grains (Ref [158](#page-28-0)). Intrinsic/extrinsic dislocation loops are created by agglomeration of point defects and are assisted by dislocation climbing.

In the case of kinetic-sprayed Cu deposits, recovery is restrained by its low SFE (relative to Ni and Al). In this case, phenomena that depend on dislocation movement (i.e., cross-slip and recovery) rarely arise because of the

<span id="page-19-0"></span>

Fig. 26 Microstructural images of kinetic-sprayed Ni, Ti, and Ta coating layers: (a) EBSD characterization of Ni deposit; Euler angle map (left) and pattern quality map (right) (Ref [77](#page-26-0)), (b) TEM images of kinetic-sprayed Ti deposit and the particle–particle interface region (left), and magnified bright-field images with dislocation-free small grains marked with black dotted arrow (right) (Ref [76\)](#page-26-0), and (c) TEM images of kinetic-sprayed Ta deposition and the particle–substrate interface region (left) and tangled dislocations (right)

frequent dissociation of dislocations that involve creation of stacking fault planes (Ref [83](#page-26-0)). Thus, recrystallization must have been the primary phenomenon in the Cu deposit. However, since recrystallization is sensitive to local conditions such as dislocation density/orientation/arrangement and temperature, the microstructural evolution is quite distinctive with respect to the local conditions even within the same interface region. Figure  $27(c)$  $27(c)$  shows the complicated organization of microstructures in a kinetic-sprayed Cu deposit: large grains whose walls consist of countless dislocations (A), elongated refined grains (B), equiaxed nanocrystalline grains (C), and micrograins with recrystallization twins.

Thus, the microstructural evolution differs case by case. In accordance with local conditions and material properties, a dominant mechanism of restoration that forms characteristic microstructures, such as the coffee-bean-like structure in Ni, can be determined. Other distinct microstructures can also be induced in the kinetic-sprayed coating layer (Ref [158-160](#page-28-0)), but those details are not dealt with here. In the next section, the effects of various microstructures in kinetic spraying on the mechanical properties are discussed in order of importance.

### 3.5 Summary

As discussed above, the microstructural features in kinetic-sprayed layers result primarily from the generation of high-strain and strain-rate deformation, which instantly induces large amounts of thermal energy at the interface region. Basically, microstructural differences always exist between internal particle and interface regions in kineticsprayed deposits because plastic deformation and the rise of thermal energy are limited to the interface region. However, even within the interface region, the microstructure varies considerably by location (i.e., south

<span id="page-20-0"></span>

Fig. 27 TEM images of kinetic-sprayed Al, Cu, and Ni deposits: (a) particle–particle boundary region of Al deposit, (b) particle–particle boundary region of Ni deposit with interparticle boundaries marked with black arrows, and (c) triple joint of deposited Cu particles with interparticle boundaries marked by black arrows (Ref [81\)](#page-26-0)

pole, intermediate and edge parts of the particle interface region), since most of the critical reactions related to bonding are induced only at the periphery of the interface; For instance, at the periphery of the interface region, a thin layer of many nanoscale grains can appear along the bonding interface, while the south pole region shows a bimodal structure or relatively coarse, microscale structures.

Specifically, microstructural alterations can occur in two ways in accordance with the type of energy required: (i) thermal energy and (ii) thermal and stored strain energy. In case I, the microstructure is altered when the edge part of the interface region is instantly heated to the melting point. Although the absolute quantity of induced thermal energy is insufficient to produce a reaction across a large area of the interface region, it is sufficient to generate several reactions (i.e., phase transformation and formation of intermetallic compounds) in highly localized areas in nanoscale. In case II, some microstructural phenomena, such as grain refinement and static/dynamic recovery/recrystallization, stabilize the interface region when sufficient thermal energy is provided to relieve the high level of stored strain energy that results from the severe plastic deformation. Specific forms of microstructural evolution are determined by the interactions between the given thermal and stored strain energies and the material properties of the feedstock powder. Not all interactions can be classified into specific cases, but whether the specific energy requirements are satisfied for certain materials is obvious.

# 4. Distinctive Features of Mechanical Properties Induced by Characteristic Microstructural Evolution at the Interface Area

As explained in section [3](#page-13-0), a remarkable characteristic of the microstructures in kinetic-sprayed metal deposits is that nanoscale grains are created by grain refinement and static/dynamic restoration. When a nanocrystalline structure is created, the mechanical behavior of the material changes from that of a microscale structure. In kinetic spraying, the hardness is greatly improved by grainboundary strengthening (Hall–Petch relationship) due to the generation of ultrafine (100-500 nm) grains with a certain amount of internal dislocations that result from DRV and DRX. Such enhancement of the hardness is assisted by the high density of dislocations far from the bonding interface that is not relieved by static/dynamic restoration because of insufficient thermal energy supply. Zou et al. (Ref [84\)](#page-26-0) and Bae et al. (Ref [52\)](#page-26-0) confirmed this hardness increase at the interface region. According to Zou et al. (Ref [84\)](#page-26-0), the hardness of the interface region was much higher than that of the internal region. For kinetic-sprayed Ni deposit, the average hardness at the interface region was 4.3 GPa, whereas that of the internal region was 3.3 GPa (Fig. 28). They calculated the numerical value of hardness using  $H = H_i + H_{\Delta GB} +$ municrical value of naturess using  $H = H_1 + H_{\Delta GB} + H_{\Delta \rho} = H_1 + H_{\Delta GB} + 3\sqrt{3} \alpha \mu b \sqrt{\Delta \rho}$ , to confirm that the primary reason for the increase of hardness was the grain boundary and strain hardening phenomena. Although  $H_{\Lambda GB}$  and  $H_{\Lambda \rho}$  were not precisely estimated, the hardness at the interface region calculated assuming  $H_{\text{AGB}} = -1$  GPa (Ref [159,](#page-28-0) [160](#page-28-0)) and  $\Delta \rho = 2 \times 10^{14} \text{ m}^{-2}$  (Ref [161](#page-28-0)) (i.e.,  $4.6 \pm 0.2$  GPa) was similar to the experimental result (4.3 GPa). Additionally, Bae et al. (Ref [52\)](#page-26-0) verified that the maximum nanohardness of a kinetic-sprayed Ni deposit (~5.5 GPa) was adequately high, similar to that of electrodeposited ~20-nm nanocrystalline Ni (5.1-6.0 GPa) (Ref [162](#page-28-0)), though they both had lower hardness than  $\sim$ 6-nm nanocrystalline Ni (7.6-9.7 GPa) (Ref [163](#page-28-0)).

Nanocrystalline (~30-100 nm) grains smaller than ultrafine (100-500 nm) grains can be created at the interface region by SRV and SRX. At this scale, plastic deformation is encouraged by behavior involving grain boundaries rather than by dislocation movement (i.e., grain boundary sliding, GBS) (Ref [164](#page-28-0)). Dislocations cannot be generated in nanograins, since the equilibrium distance between dislocations is larger than the size of the grains, which is called dislocation pile-up breakdown. Bae et al. (Ref [78\)](#page-26-0) reported evidence of GBS in a kineticsprayed specimen. Figure [29](#page-22-0) shows the morphology of a transparticle fracture surface in a kinetic-sprayed Ni layer.

Obviously, many nanosized dimples resulted from GBS among high-angle grain boundaries (HABs) during plastic deformation. In this case, the material can acquire superplasticity, where extraordinary ductility appears during plastic deformation within a certain range of strain rate and elevated temperature. However, no research has yet focused on the superplasticity behavior of kineticsprayed metal deposits.

The most interesting characteristic of kinetic-sprayed metal coating layers is the coexistence of ultrafine (100- 500 nm) grains with nanocrystalline (~30-100 nm) structures at the interface region. As presented in Fig. [30](#page-22-0) (marked with black dotted circle) (Ref [78](#page-26-0)), a bimodal structure of ultrafine grains and nanocrystalline structures shows both excellent hardness and ductility (Ref [165-167](#page-28-0)).

However, these advanced mechanical properties are only reasonable in the interface region. Within the whole coating layer, the bonding interface is usually imperfect because of defects such as porosity and cracks induced by the supersonic impacts. These flaws significantly degrade the mechanical properties of a kinetic-sprayed metal deposit by serving as starting points for cracks. Table [3](#page-22-0) presents the mean hardness of kinetic-sprayed Ti deposits under various process conditions as measured by the indentation method (Ref [111](#page-27-0)). All the specimens possessed higher nanohardness than bulk Ti, but only a few of them presented advanced microhardness relative to bulk Ti when the particle velocity was above a certain threshold. Here, the nanohardness values were calibrated using Eq 1 to minimize the effect of indentation load and size.

$$
H = H_0 \sqrt{1 + \frac{h^*}{h}},\tag{Eq 1}
$$

where  $H$  is the nanohardness,  $h$  is the indentation depth,  $H_0$  is the material hardness at infinite indentation depth, and  $h^*$  is the characteristic length scale. The primary reason for the different tendencies exhibited by the nanoand microhardness is hindrance due to microscale defects,



Fig. 28 EBSD characterization of Ni coating after nanoindentation: (a) Euler angle map, (b) image quality (IQ) map, and (c) IQ map with local hardness (Ref [82](#page-26-0))

<span id="page-22-0"></span>such as pores and debonding of particles. For the nanohardness, the evaluation is almost unaffected by such microscale defects because the measuring range is confined to the nanoscale and is assisted by a low load (1- 20 mN). However, evaluation of microhardness cannot be



Fig. 29 SEM image of fracture surface in kinetic-sprayed Ti deposit with ductile dimple fracture mode (Ref [76\)](#page-26-0)



Fig. 30 Highly magnified bright-field TEM image near interparticle boundary of kinetic-sprayed Ni deposit (Ref [52\)](#page-26-0)

free from the influence of microscale defects (i.e., pores, interparticle cracks, and interparticle debonding). Figure [31](#page-23-0) shows indents caused by microhardness measurements for bulk Ti and Ti deposits (Ref [111\)](#page-27-0). In the specimen with poor microhardness (Table 3), there is obvious distortion of the indents, which seems to be caused by debonding between particles (Fig. [31c](#page-23-0), d). Thus, pores and debonding of particles contributed to the reduction of microhardness. However, as shown in Fig.  $31(e, f)$  $31(e, f)$  and  $(g, h)$ , the distortion in the Ti deposits was minimized when using enhanced process conditions (increased particle velocity), corresponding to the higher microhardness of the well-formed Ti deposits compared with bulk Ti (Table  $3$ ); That is, to achieve a kineticsprayed metal deposit of good quality, the bonding quality must be maximized by using optimized process conditions. However, some limited parts of the deposit (i.e., the interface region) usually possess superior properties, regardless of the process conditions.

Furthermore, even for the case of Cu (Ref [12\)](#page-25-0), which is one of the materials easily deposited by kinetic spraying, the tensile strength of the as-sprayed layer  $(<100$  MPa) is much lower than that of cold-rolled plate (~390 MPa). The elongation to failure was also poor for as-sprayed Cu deposits  $( $0.1\%$ ) relative to cold-rolled Cu plate$ (~2.6%). The reason for such inferior tensile strength and ductility in the kinetic-sprayed specimens can be determined by observing the fracture surface of the cold-rolled and as-sprayed Cu (Fig.  $32$ ). As shown in Fig.  $32(b)$ , brittle fracture was generated along the interparticle boundary. No dimples appeared on the fracture surface, which is evidence of intimate bonding in the kinetic spray process. Thus, the degradation in the evaluated mechanical properties of the kinetic-sprayed Cu deposit occurred because of weakening caused by imperfect bonding between particles.

However, if optimized process conditions are applied in fabrication of metal deposits via a kinetic spray process, metal coatings with mechanical properties almost equivalent to those of the bulk metal can be formed, although it is impossible to achieve perfect interface bonding throughout the whole layer. Kim et al. (Ref [122](#page-27-0)) reported that a kinetic-sprayed thick Al deposit showed superior mechanical properties relative to bulk Al fabricated by equal-channel angular pressing, which is a fabrication method involving severe plastic deformation (Ref [168,](#page-28-0) [169](#page-28-0)).

Table 3 Mean hardness and splat adhesion strength of kinetic-sprayed Ti deposits under various process conditions, measured by indentation method

	<b>Nanohardness</b> (GPa)	Microhardness (GPa)	<b>Splat adhesion (MPa)</b>
608 m/s ( $N_2$ at 300 °C and 3 MPa)	$2.92 \pm 0.12$	$1.14 \pm 0.03$	$32 + 52$
648 m/s ( $N_2$ at 500 °C and 3 MPa)	$2.86 \pm 0.13$	$1.32 \pm 0.09$	$92 + 51$
688 m/s (N <sub>2</sub> at 600 $^{\circ}$ C and 3 MPa)	$3.06 \pm 0.07$	$1.85 \pm 0.04$	$128 \pm 86$
785 m/s ( $N_2$ at 750 °C and 4 MPa)	$3.15 \pm 0.20$	$2.21 \pm 0.08$	$152 \pm 16$
805 m/s ( $N_2$ at 800 °C and 4 MPa)	$2.98 \pm 0.10$	$2.15 \pm 0.35$	$238 + 33$
1173 m/s (He at $350$ °C and 4 MPa)	$2.60 \pm 0.08$	$2.57 \pm 0.04$	$254 + 21$
Bulk Ti	$2.19 \pm 0.04$	$1.97 \pm 0.06$	$\cdots$

<span id="page-23-0"></span>

Fig. 31 Optical microscopy images of 0.1-N and 5-N indents in: (a, b) bulk Ti plate, and as-sprayed Ti layers deposited with particle velocity of (c, d) 608 m/s, (e, f) 805 m/s, and (g, h) 1173 m/s (Ref [109\)](#page-27-0)

<span id="page-24-0"></span>

Fig. 32 Microstructure and morphology of fracture surface in (a) cold-rolled Cu plate and (b) as-sprayed Cu deposit fabricated by kinetic spray process (Ref [12](#page-25-0))

# 5. Summary

Throughout this paper, the correlations among impact conditions, interface reactions, microstructural evolution, and mechanical properties have been presented. Although these relationships cannot be systematically classified into specific cases due to their complexity, there are critical factors for each relationship. For interface reactions, the most important factor is the interactions between given impact conditions (i.e., particle velocity, particle and substrate temperatures, relative hardness of particles and substrate, oxidation state, and impact angle) and the material properties. It is possible to predict the interface reactions for given impact conditions using material-specific parameters from the Johnson–Cook plasticity model and the general properties of the metal feedstock powder. The microstructural response of the feedstock material is primarily determined by the interactions between its material properties and the given thermal/stored strain energy. In most cases, nanocrystalline structure is formed along the bonding interface through grain refinement and static/dynamic restoration. Accordingly, the local mechanical properties (i.e., ductility, hardness, and toughness) of the interface region of the coating layer are superior to those of the bulk metal, whereas the microscale properties are usually inferior because of microdefects, such as pores and interparticle debonding, unless optimized process conditions are applied.

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