



Research on a High-Threshold-Voltage AlGaIn/GaN HEMT with P-GaN Cap and Recessed Gate in Combination with Graded AlGaIn Barrier Layer

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Abstract

Due to the enhanced-mode (E-mode) operation, AlGaIn/GaN high-electron-mobility transistors (HEMTs) are considered to be safer for circuit operation. In order to improve the threshold voltage (V_{th}) of the device, this work provides a hybrid gate structure HEMT by embedding a P-GaN cap on the etched graded AlGaIn barrier layer. Through simulation calculations, the P-GaN cap (thickness of P-GaN = 50 nm, concentration of P-type = $2 \times 10^{18} \text{ cm}^{-3}$) and the aluminum (Al) composition (Al:0.3 \rightarrow 0.24), in the graded AlGaIn barrier layer were optimized. Although simulation calculations show that the optimized P-GaN layer can significantly increase the device's V_{th} to 8.6 V and transconductance (g_m) to 94.7 mS/mm, the device exhibits a lower saturation current (I_{sat}). Therefore, to improve the output characteristics of the devices, the addition of an N-well in the GaN channel layer of such structures was proposed. It can increase the device's source–drain current while maintaining a steady V_{th} . Compared with the HEMT structure/combined P-GaN cap with recessed gate and a graded AlGaIn barrier layer, the device with the added N-well exhibits a significant improvement of 11.2% in the saturation current ($I_{sat} = 718 \text{ mA/mm}$). The results demonstrate that HEMT structures combining recessed gates and P-GaN with N-well have promising applications in next-generation high-power devices.

Keywords E-mode · AlGaIn/GaN HEMT · P-GaN cap · graded AlGaIn barrier layer · N-well

Introduction

Gallium nitride (GaN) has emerged as an excellent candidate for high-power and high-temperature applications due to its wide bandgap. It also possesses high breakdown field strength and high electron saturation velocity.^{1–3} These properties make GaN-based high-electron-mobility transistors (HEMTs) potential candidates for next-generation power electronic devices in applications such as power supplies, photovoltaic inverters, and motor drives.⁴ HEMT devices formed by AlGaIn/GaN heterostructures at the two-dimensional electron gas (2DEG) are naturally in the conducting state, which increases power losses and poses challenges to circuit safety.⁵ However, enhancement-mode HEMTs⁶ have

been developed to address these issues. They not only reduce the complexity and cost of circuit design in applications such as radio frequency (RF) and low-noise power amplifiers, but also significantly improve circuit safety and system reliability in high-power switching applications. Moreover, enhancement-mode HEMTs can alleviate the P-channel deficiency in nitride semiconductors. As a result, a growing number of researchers are exploring enhancement-mode HEMTs. Currently, there are several methods to achieve device enhancement mode: recessed gate structures,⁷ technology of P-GaN cap,⁸ F⁻ implantation technology,⁹ and polarization engineering techniques.¹⁰ However, the temperature reliability of devices using fluorine implantation is not suitable for practical applications.^{11,12} Polarization engineering techniques have extreme process complexity and cost, making them impractical. The recessed gate technique, while effective, results in a thinner barrier layer due to etching of the AlGaIn barrier, causing significant gate leakage current and affecting gate swing and device reliability. On the other hand, P-GaN gate technology employs a P-GaN layer with an activated hole concentration that effectively depletes the

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2DEG at the AlGa_N/Ga_N interface, rendering the gate non-conductive under zero bias voltage.¹³ Su et al.¹⁴ successfully achieved enhancement-mode HEMT devices by employing heavily doped P-GaN cap layers, resulting in a threshold voltage as high as 4.3 V and a maximum drain–source current of 208 mA/mm. However, the elevated doping concentration of Mg in the P-GaN layer introduces defects,^{15–21} leading to a decrease in device performance. Moreover, many related studies have been reported focusing on individual device parameters, and achieving high threshold voltages and drain–source currents simultaneously still poses significant challenges. Therefore, in this paper, we propose a hybrid gate structure combining P-GaN with a recessed gate, and determine the optimal P-GaN cap layer thickness and concentration, leading to a significant improvement in the V_{th} of the HEMT device. In addition, the Al component in the graded AlGa_N barrier layer is optimized to reduce the remote scattering effect²² and improve the electron mobility in the channel. Finally, an N-well is added to the Ga_N

channel layer, effectively increasing the source–drain current of the device. The structure with a P-GaN gate grown on an etched AlGa_N barrier layer and an N-well inserted in the Ga_N channel was realized using Silvaco TCAD (technology computer-aided design) simulation software. This not only significantly increases the V_{th} and g_m of the device but also enhances the maximum drain–source (saturation) current.

Device Structure

Figure 1a, b, and c illustrate the architecture of structure A, consisting of a conventional MIS-recessed gate HEMT; structure B, the recessed gate HEMT covered with a P-GaN layer ($Al_{0.3}Ga_{0.7}N/GaN$); and structure C, the P-GaN cap and recessed gate HEMT with a graded AlGa_N layer ($Al_{0.30 \rightarrow 0.24}Ga_{0.70 \rightarrow 0.76}N/GaN$). The mole fraction of Al in the AlGa_N barrier layer of structure A is 0.24. The mole fraction of Al in the AlGa_N barrier layer of structures B and C is 0.3, and the P-GaN cap layer

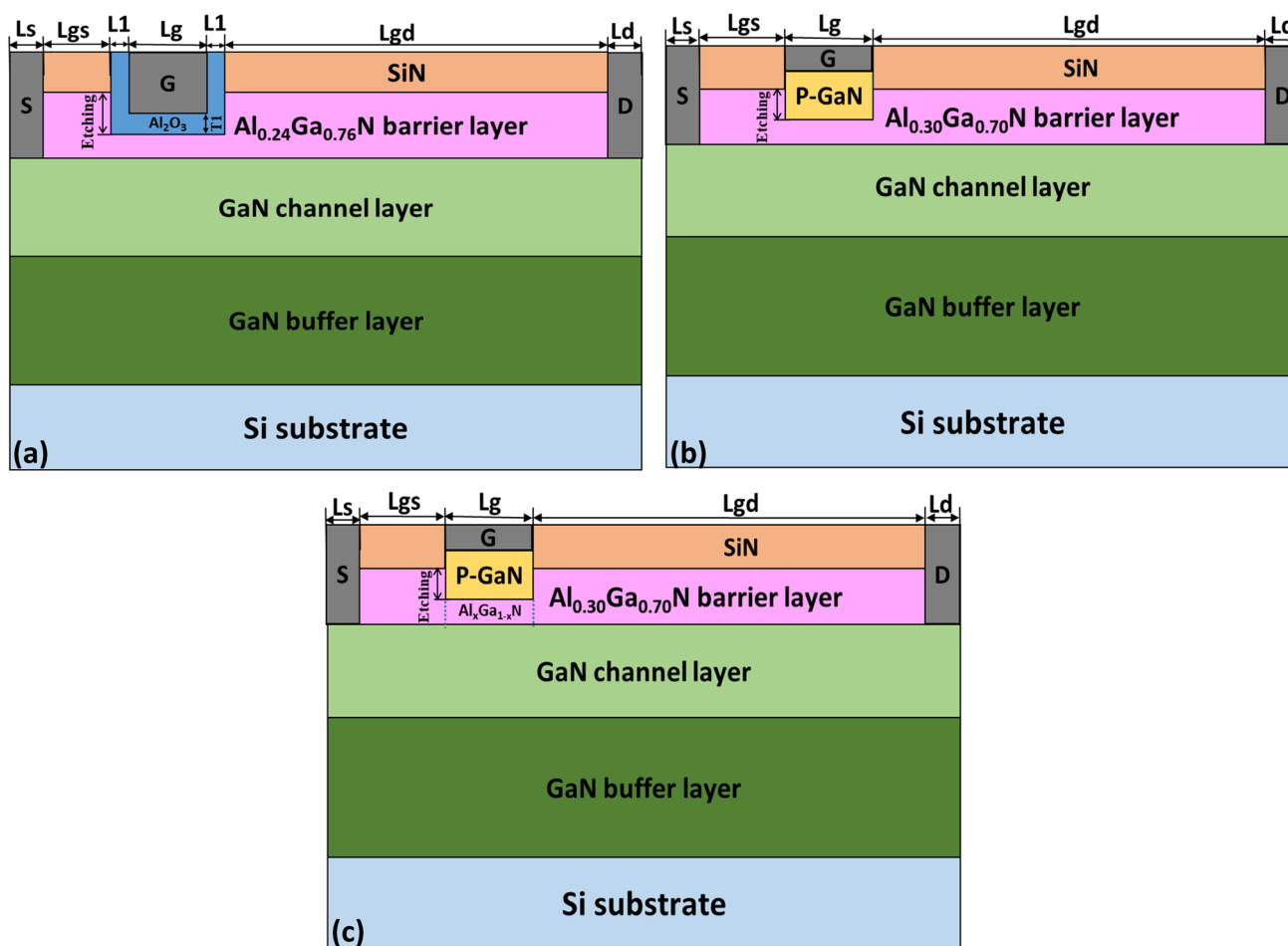


Fig. 1 The structure cross sections of the three devices: (a) structure A: a conventional MIS-recessed gate HEMT, (b) structure B: the recessed gate HEMT covered with a P-GaN layer HEMT

($Al_{0.3}Ga_{0.7}N/GaN$), and (c) structure C: the P-GaN cap and recessed gate HEMT with the graded AlGa_N layer HEMT structure ($Al_{0.30 \rightarrow 0.24}Ga_{0.70 \rightarrow 0.76}N/GaN$).

is doped with a *P*-type concentration of $1 \times 10^{18} \text{ cm}^{-3}$. Among them, the GaN buffer layer is unintentionally doped. Table I summarizes the detailed information of the three main structural parameters (a), (b), and (c).

Simulation Model

In this study, we employed the Silvaco ATLAS TCAD tool for 2D simulations.²³ The numerical solution was performed for the Poisson equation, electron continuity equation based on the drift–diffusion (DD) model.²⁴ The transport equations were solved using the drift–diffusion model,²⁵ and all calculations in the simulations were based on Fermi–Dirac statistics. Additionally, the Shockley–Read–Hall (SRH) model was utilized to simulate the impact of traps in the bandgap on carrier recombination,²⁵ and a polarization model was employed to simulate the polarization charges along the heterojunction interface. Considering that the electron mobility in the GaN channel varies with the electric field strength, we considered the effect of the electric field on the electron mobility of the nitride in our simulations and chose a specific field-dependent mobility model. We used the Silvaco ATLAS simulation software to simulate the transfer characteristics of the conventional AlGaIn/GaN HEMT and compared the results with previously reported experimental data.²⁶ As shown in Fig. 2, the results demonstrate excellent agreement between the simulation and experimental data, indicating that the physical parameters chosen in our simulations are consistent with the actual conditions.

Results and Discussion

The polarization model is crucial for the simulation of GaN-based HEMT devices. In Fig. 3, the polarization charge distribution at the AlGaIn/GaN heterojunction interface is illustrated for structure B and structure C. The spontaneous polarization arises from the asymmetry along the interface of the wurtzite crystal, while the piezoelectric polarization results from the stress and strain caused by the lattice mismatch between the GaN and AlGaIn layers. In the absence of an external electric field, the total polarization intensity

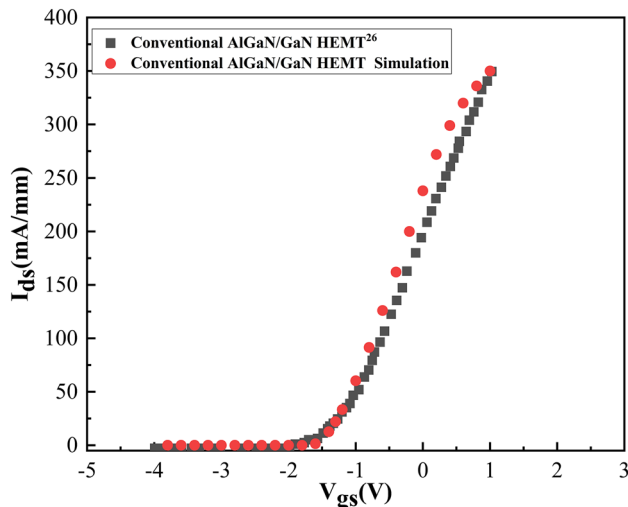


Fig. 2 I_d – V_g characteristic comparison of conventional AlGaIn/GaN HEMT²⁶ with TCAD simulation for validation at $V_{ds}=2 \text{ V}$, $V_g=-4 \sim 1 \text{ V}$.

Table I The detailed information on the three main structural parameters (a), (b), and (c)

Parameters	Unit	Value of structure A	Value of structure B	Value of structure C
Source length (Ls)	μm	1	1	1
Distance from gate to source (Lgs)	μm	3	3	3
Gate length (Lg)	μm	1.99	2	2
Thickness of the Al ₂ O ₃ dielectric layer (T1)	nm	1	–	–
Width of the Al ₂ O ₃ dielectric layer (L1)	nm	5	–	–
Distance from gate to drain (Lgd)	μm	9	9	9
Drain length (Ld)	μm	1	1	1
Length of the P-GaN layer (Lp)	μm	–	2	2
Thickness of the P-GaN layer	nm	–	50	50
Al component of the AlGaIn barrier layer under the gate	mol	0.24	0.30	0.3 → 0.24
Etching depth of the AlGaIn barrier layer	nm	17	10	10
Thickness of AlGaIn barrier layer	nm	20	20	20
Thickness of the GaN channel layer	μm	0.4	0.4	0.4
Thickness of the high-resistance GaN buffer layer	μm	2.6	2.6	2.6

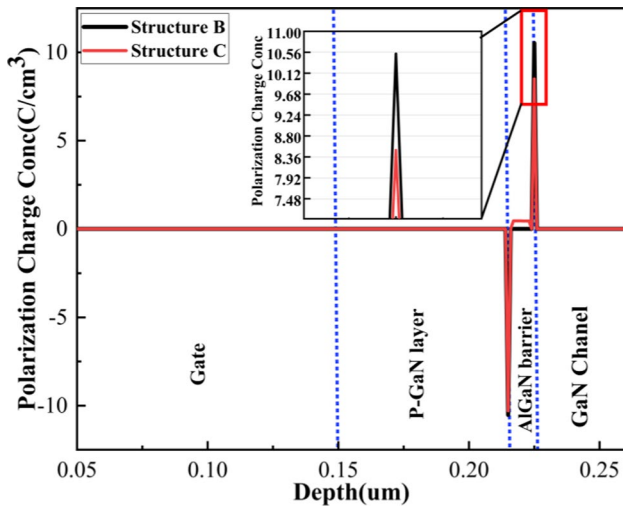


Fig. 3 Polarization-induced charge profile along the vertical axis.

(P_{Tot}) is obtained by adding the spontaneous polarization strength (P_{SP}) and the piezoelectric polarization strength (P_{PE}). The 2DEG at the interface of the AlGaN/GaN HEMT heterojunction can be described by the polarization charge formula, Eq. 1:

$$\sigma = P_{spont (AlGaN)}(x) - P_{spont (GaN)}(0) + P_{pie (AlGaN)}(x) \quad (1)$$

Here, the spontaneous polarization strength of the GaN layer is $P_{spont(GaN)} = -0.034$, and the spontaneous polarization of the AlGaN layer is $P_{spont(AlGaN)} = -0.052x - 0.034$.

The piezoelectric polarization intensity acting on the AlGaN barrier layer due to the lattice mismatch in the AlGaN/GaN heterojunction can be expressed by Eq. 2:

$$P_{pie (AlGaN)} = 2 \frac{a_s - a_0}{a_0} \left(E_{31} - \frac{C_{13}}{C_{33}} E_{33} \right) \quad (2)$$

Here, x represents the mole fraction of Al in the AlGaN layer, E_{31} and E_{33} are piezoelectric constants, C_{13} and C_{33} are elastic constants, a_0 and a_s are the lattice constant in the strained layer and the substrate lattice constant, respectively.

$$V_{th} = \varphi_B - \Delta E_c - \frac{qN_{D_{AlGaN}} d_{AlGaN}^2}{2\epsilon_0 \epsilon_{AlGaN}} - \frac{\sigma}{\epsilon_{AlGaN}} d_{AlGaN} \quad (3)$$

In Eq. 3, φ_B is the Schottky barrier height of the gate metal, ΔE_c is the conduction band offset at the AlGaN/GaN interface, q is the electron charge, σ is the net polarization charge density at the interface, d is the thickness of the AlGaN barrier, and ϵ_0 and ϵ_{AlGaN} are the vacuum permittivity and AlGaN permittivity, respectively.

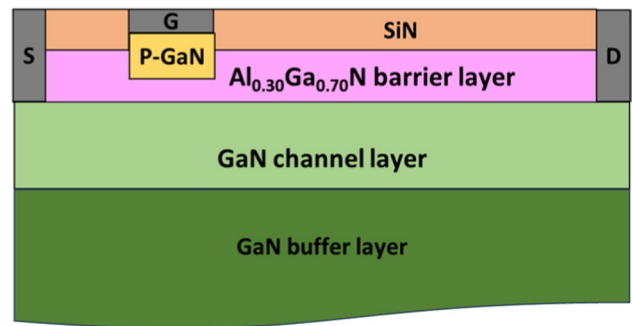


Fig. 4 Schematic of structure B.

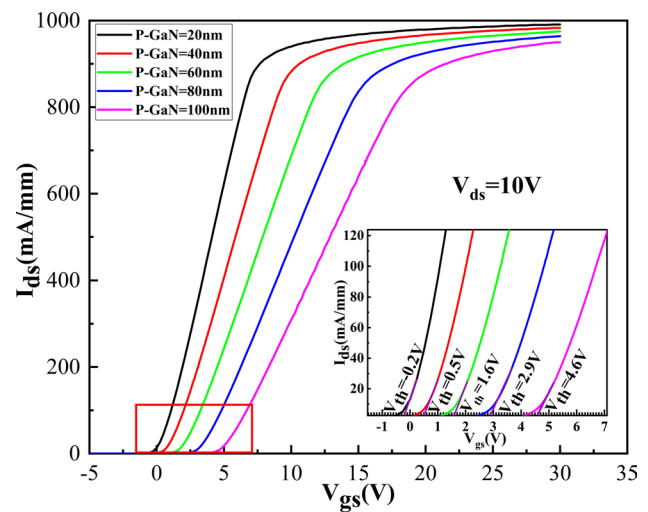


Fig. 5 Transfer characteristic curves of structure B with different thicknesses of P-GaN cap.

DC characteristics of the recessed gate HEMT covered with a P-GaN layer HEMT

We investigated the effect of the thickness of P-GaN on the transfer characteristics of the device when the P-GaN cap was doped with $2 \times 10^{18} \text{ cm}^{-3}$ P-type concentration. Figure 4 shows a schematic of structure B.

As shown in Figs. 5 and 6, we consider that when the etching thickness of the AlGaN barrier layer is fixed to 10 nm, the thickness of P-GaN is increased to 20 nm, 40 nm, 60 nm, 80 nm, and 100 nm, respectively. With an increase in the thickness of the P-GaN cap, the distance between the gate and the GaN channel increases, and so does the V_{th} . However, the transconductance decreases significantly. This is because the increased P-GaN thickness weakens the gate control over the channel. Figure 7 presents the relationship between the threshold voltage and transconductance for

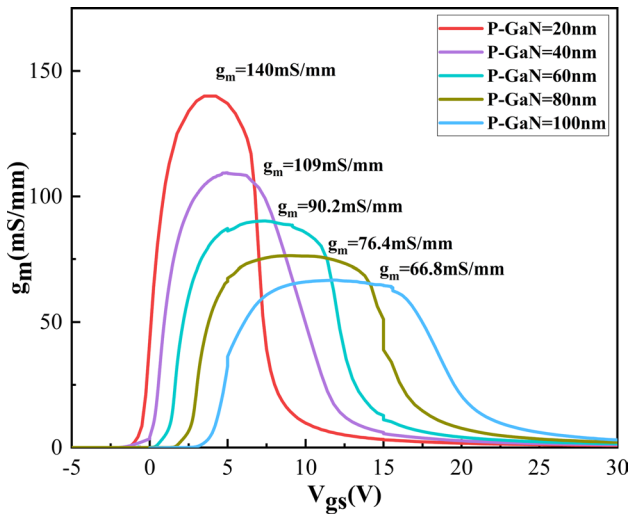


Fig. 6 Transconductance curves of structure B with different P-GaN caps.

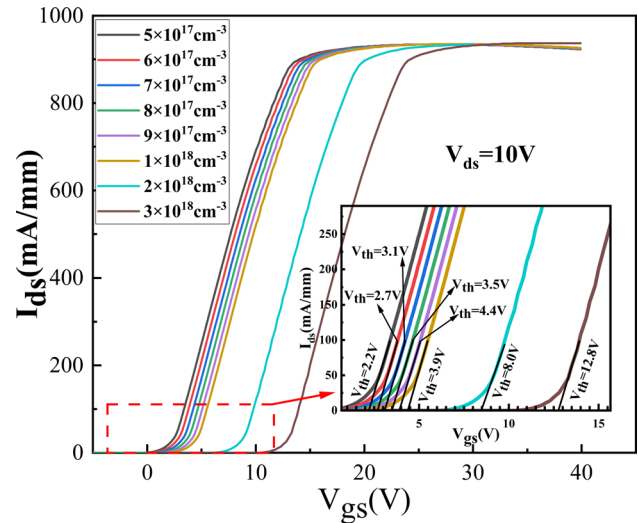


Fig. 8 Transfer characteristics of P-GaN cap with different P-doping concentrations.

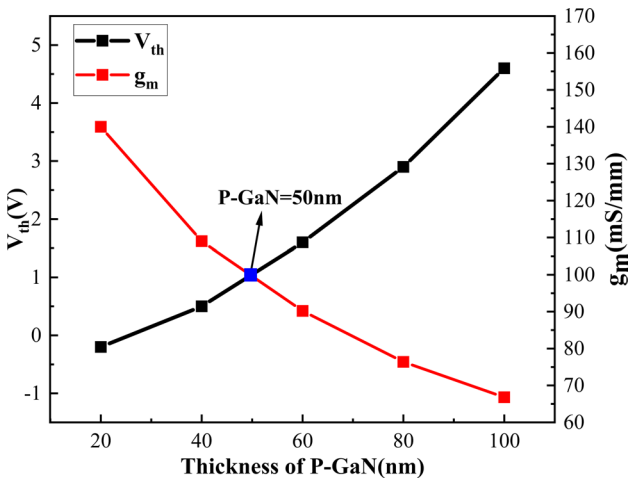


Fig. 7 Relationship between threshold voltage and transconductance for different thicknesses of GaN cap.

different thicknesses of the GaN cap layer. From Fig. 7, we can conclude that a P-GaN thickness of 50 nm is optimal.

After determining a P-GaN cap thickness of 50 nm, we further discuss the effect of *P*-type concentrations on the device. Figure 8 illustrates the transfer characteristics for different *P*-type concentrations in the P-GaN cap layer. For *P*-type concentrations below $2 \times 10^{18} \text{ cm}^{-3}$, although the threshold voltage increases with the doping concentration, the increase in V_{th} is not significant. However, when the *P*-type concentration is $2 \times 10^{18} \text{ cm}^{-3}$ and $3 \times 10^{18} \text{ cm}^{-3}$, the threshold voltage increases significantly to 8 V and 12.8 V, respectively. This is because the hole concentration can deplete the 2DEG under the gate, leading to an increase in the conduction band and the threshold voltage.²⁷ When

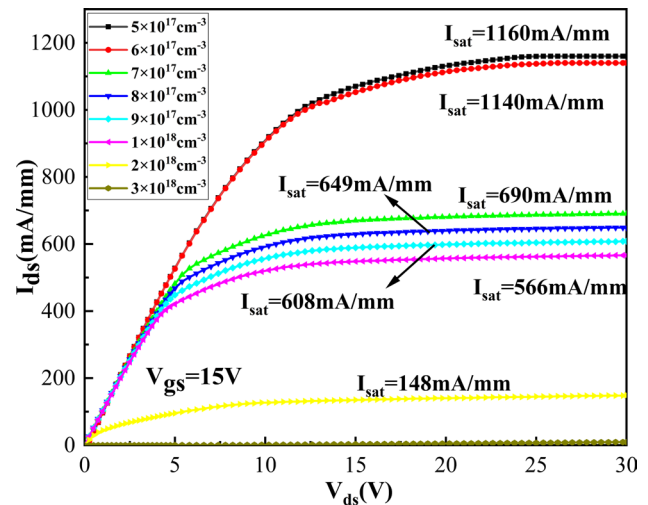


Fig. 9 Output characteristics of P-GaN cap with different P-doping concentrations.

the *P*-type concentration is $1 \times 10^{18} \text{ cm}^{-3}$, the 2DEG in the channel is entirely depleted. Thus, at hole concentrations of $2 \times 10^{18} \text{ cm}^{-3}$ and $3 \times 10^{18} \text{ cm}^{-3}$, P-GaN increases the energy gap between the conduction band and the Fermi level of the AlGaIn/GaN heterojunction. As a result, the gate bias voltage needs to provide additional voltage to reduce the energy gap between the conduction band and the Fermi level, resulting in a significant increase in the threshold voltage. Figure 9 shows the output characteristics for different *P*-type concentrations in the P-GaN cap. Although the device exhibits the highest threshold voltage at a hole concentration of $3 \times 10^{18} \text{ cm}^{-3}$, the source–drain current is practically zero. This is because the high concentration of holes diffuses into

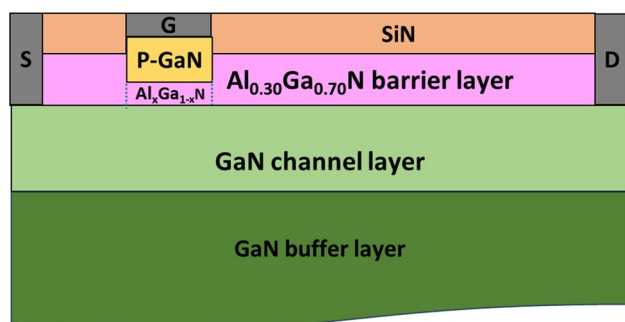


Fig. 10 Schematic of structure C.

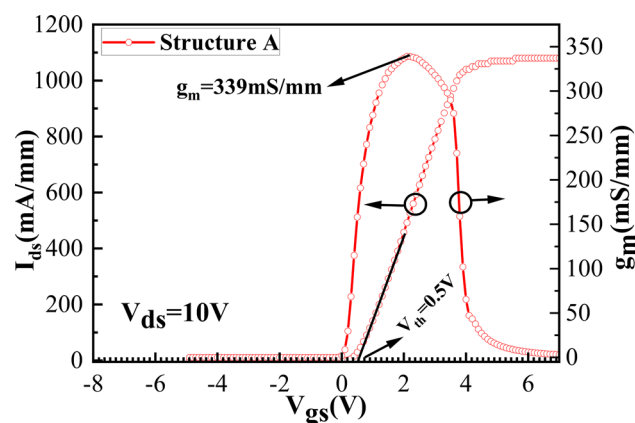


Fig. 11 Transfer characteristics and transconductance curves of structure A.

the GaN channel layer, leading to the recombination of the 2DEG with the holes, resulting in only a few electrons participating in conduction in the channel under the on state, and consequently leading to a very low source-drain current. Additionally, the diffusion of holes into the GaN channel, combined with the lower hole mobility compared with electron mobility, also contributes to the low source-drain current. To strike a balance between the threshold voltage and the source-drain current, we chose a hole concentration of $2 \times 10^{18} \text{ cm}^{-3}$ in this study.

To examine the performance of structure C, we propose three different structures: structure A, structure B, and structure C. Figure 10 shows a schematic of structure C. Figures 11 and 12a and b depict the transfer characteristics and transconductance (g_m) curves of the three structures. For structure A, the V_{th} and g_m can reach 0.5 V and 339 mS/mm, respectively. In the case of structure A, when the thickness of the AlGa_xN barrier layer beneath the gate is etched to 3 nm, the device has a V_{th} of only 0.5 V, indicating the limitation in etching the AlGa_xN barrier layer to improve the V_{th} . For structure B, the holes in the P-GaN cap can effectively deplete the electrons in the GaN channel layer, significantly

increasing the V_{th} to 8.0 V. However, the activation energy of Mg in the P-GaN layer is mostly high (0.15–0.21 eV), and the compensation effect of H in P-GaN prevents the activation concentration of P-GaN from reaching high levels. This is also a critical challenge faced by GaN technology. Therefore, according to Eq. 3, we can reduce the 2DEG beneath the recessed gate caused by polarization charges by reducing the lattice mismatch between the AlGa_xN barrier layer and the GaN channel layer. The V_{th} of structure C is 8.6 V, which is 0.6 V higher than that of structure B. The transconductance (95.7 mS/mm) of structure C is higher than that of structure B (90.4 mS/mm). This is attributed to the gradual change in the Al composition in the AlGa_xN barrier layer from 0.3 at the P-GaN/AlGa_xN interface to 0.24 at the AlGa_xN/GaN interface. The lower Al composition at the AlGa_xN/GaN interface reduces the scattering probability for electrons in the GaN channel and thus increases the electron mobility in the GaN channel. Deng et al. reported that increasing the threshold voltage of the device would decrease the density or mobility of 2DEG in the GaN channel, resulting in a decrease in the maximum saturation current output.²⁸ However, according to Fig. 12a and b, for the transfer and output characteristics of structure C and structure B, not only does the threshold voltage of structure C increase by 0.5 V, but the saturation current of the device is slightly larger than that of structure B. Moreover, structure C also has higher transconductance, which characterizes the controllability of the gate to the channel,²⁹ and shows that the controllability of structure C to the channel is higher than that of structure B.

To investigate the impact of the gradual change in the Al composition in the AlGa_xN barrier layer from the P-GaN/AlGa_xN interface to the AlGa_xN/GaN interface on the GaN channel, we set up six groups of barrier layers with different Al composition grades: Al: 0.3 → 0.20; 0.3 → 0.22; 0.3 → 0.24; 0.3 → 0.26; 0.3 → 0.28; 0.3 → 0.30. Figure 13a illustrates the transfer characteristics of the six sets of barrier layers with varying Al fractions. The threshold voltage of the device exhibits an increasing trend as the Al tapers from 0.3 to 0.2, attributed to the reduction in the 2DEG concentration resulting from the weakening of the polarization electric field. Figure 13b shows the gate leakage currents of the AlGa_xN barrier layers with six different Al compositions. The gate leakage current exhibits an increasing trend as the Al content decreases. Notably, the gate leakage current reaches its maximum severity at Al values of 0.20 and 0.22. Conversely, when Al is set to 0.3 and 0.28, the gate leakage current is almost zero. When Al is 0.24 and 0.26, the gate leakage is also considered small for P-GaN gate HEMT devices with V_{gs} greater than 15 V.^{30–34} As the AlGa_xN barrier increases, the gate leakage current is significantly reduced. However, an excessive Al fraction leads to a decrease in the device's threshold voltage. Based on the

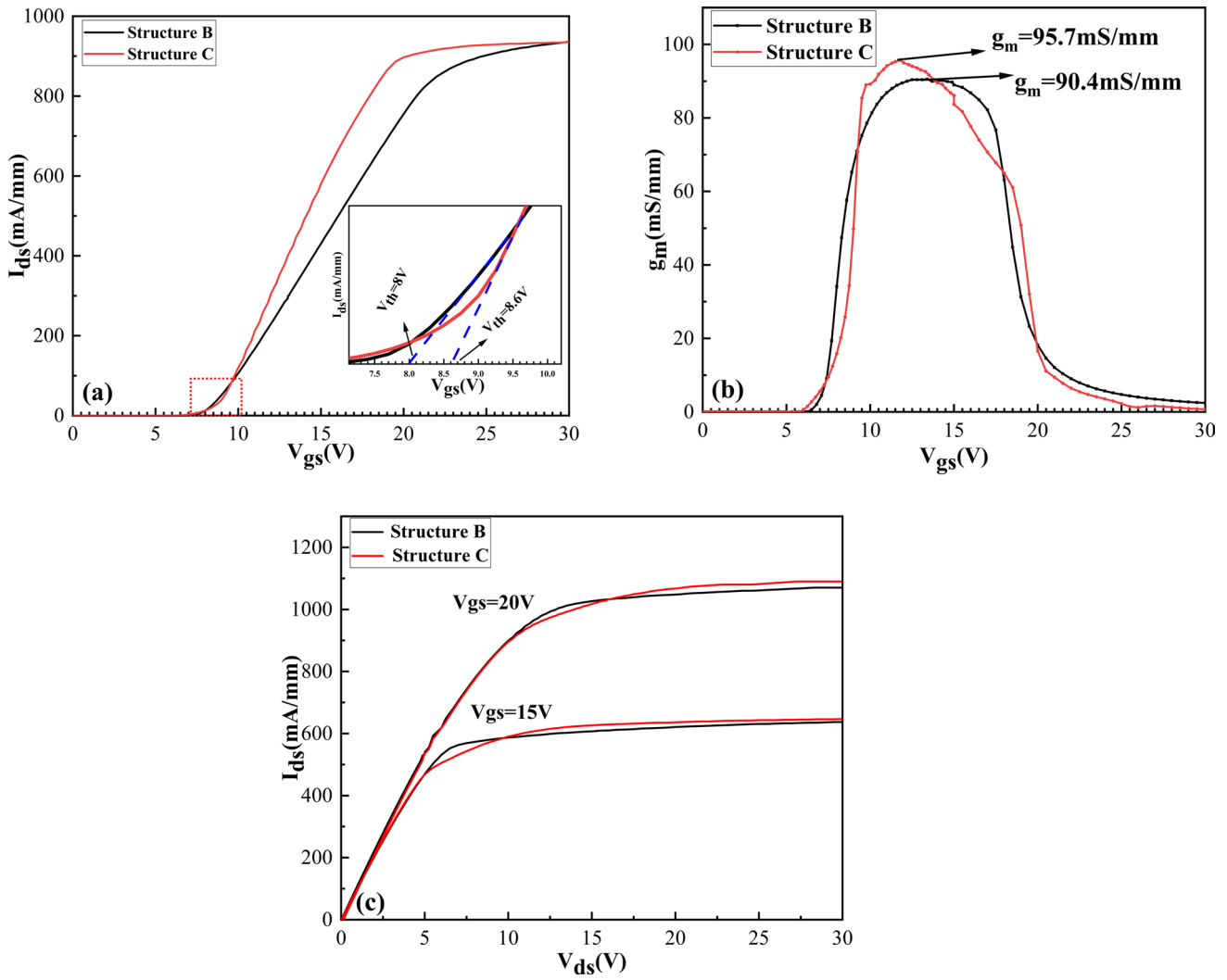


Fig. 12 (a) Transfer characteristic, (b) transconductance and (c) output characteristic curves of structure B and structure C.

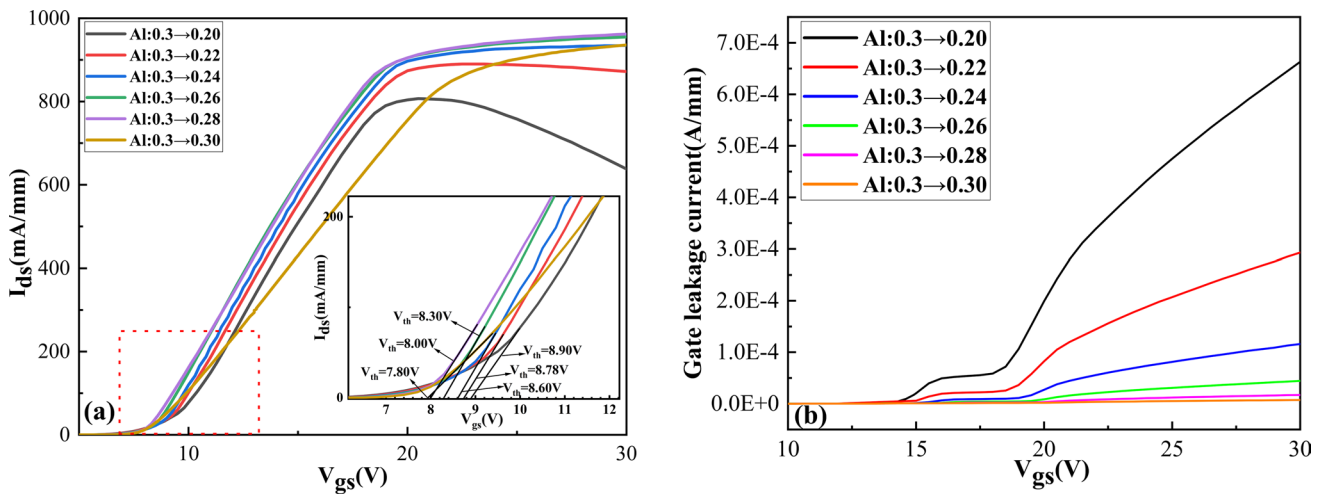


Fig. 13 (a) Transfer characteristic curves of different Al compositions in the graded AlGaN barrier layer, (b) gate leakage current of different Al compositions in the graded AlGaN barrier layer at $V_{ds} = 10$ V and $V_{gs} = 0-30$ V.

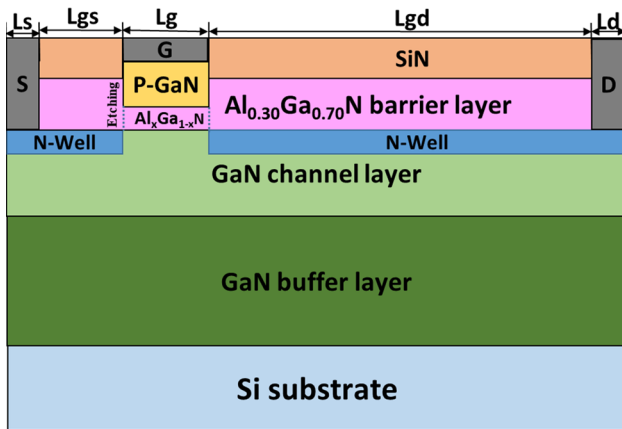


Fig. 14 Structure D: P-GaN cap and recessed gate HEMT with graded AlGaN layer and added N-well in GaN channel.

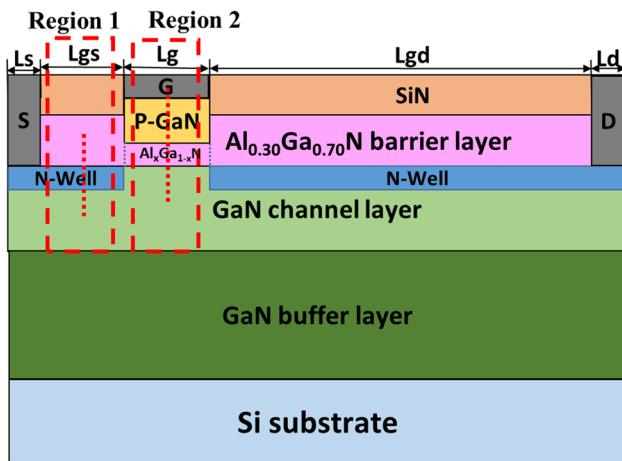


Fig. 15 Schematic of structure D in the simulation software, showing outline at two locations perpendicular to the x -axis.

aforementioned analysis, we determined that an Al composition gradient of 0.24 represents the optimal value for the AlGaN barrier layer.

Figure 14 shows the addition of two 20-nm-thick N-wells with N-doping in the GaN channel of structure D (Fig. 1c). The remaining structural parameters remain unchanged. Figure 15 presents a schematic of structure D with two regions perpendicular to the x -axis, designated as region 1 and region 2, in the simulation software. In Fig. 16a and b, the conduction band profiles along the cutline in region 1 and region 2 are shown when N-doping in the N-well is set to $4 \times 10^{18} \text{ cm}^{-3}$. In region 1, the triangular potential well widens below the Fermi level due to the influence of the N-well potential. In region 2, without the presence of an N-well, the rise in the conduction band of the P-GaN cap causes the triangular potential well at the AlGaN/GaN interface to be higher than the Fermi level, leading to the depletion

of electrons in the channel. If the N-well is confined within region 1, the doping distribution in the N-well will not affect the electron concentration in the channel below the gate, and thus it will not affect the threshold voltage.

Figure 17 shows the transfer characteristics of $I_{ds}(V_{gs})$ under different N-doping distributions in the N-well. Although the doping concentration in the N-well does not affect the threshold voltage of the device, the current density of the device increases with the doping concentration in the N-well. Figure 18 illustrates the output characteristics of structure D devices under different N-type doping concentrations in the N-well, in which the inset shows the effect of different doping profiles on the N-well drain current at $V_{gs} = 15 \text{ V}$. The output characteristics of the structure D devices increase with increasing N-doping concentration. This indicates that the addition of the N-well expands the triangular potential well below the Fermi level, thereby increasing the density of carriers participating in the conduction current in the channel. When the N-doping concentration exceeds $8 \times 10^{18} \text{ cm}^{-3}$, the saturation current of the device tends to level off. To achieve a better doping effect, we chose an N-well doping concentration of $8 \times 10^{18} \text{ cm}^{-3}$. Compared with the saturation current of the enhanced structure C device (646 mA/mm) (represented by the light purple dashed line in the inset), the saturation current of the structure D device is significantly higher. When the N-doping concentration is $8 \times 10^{18} \text{ cm}^{-3}$, the saturation current of the structure D device is 718 mA/mm. Although the ionized impurity scattering produced by doping N-type impurities in the GaN channel layer affects the carrier mobility in the GaN channel layer, the charge compensation effect of the N-well plays a major role.³⁵ As a result, the electrons after the ionization of N-type impurities participate in the conductivity, which leads to a significant increase in the saturation current of the device. Additionally, Fig. 19 summarizes the threshold voltage and saturation current performance of the proposed composite trench-gate HEMT compared with other reported GaN-based HEMTs.^{25,34,36–40} It can be observed that the HEMT of structure D exhibits a significant improvement in output performance at high threshold voltages compared with other GaN-based HEMTs. This points to a wide range of applications of this approach in high-power electronic devices.

Conclusion

In this paper, we propose a HEMT structure combining a P-GaN cap with a recessed gate and an AlGaN barrier layer with a graded Al component. Optimization of the P-GaN cap thickness and P-doping was performed using the Silvaco TCAD simulation tool to achieve an enhanced HEMT device. Additionally, the influence of

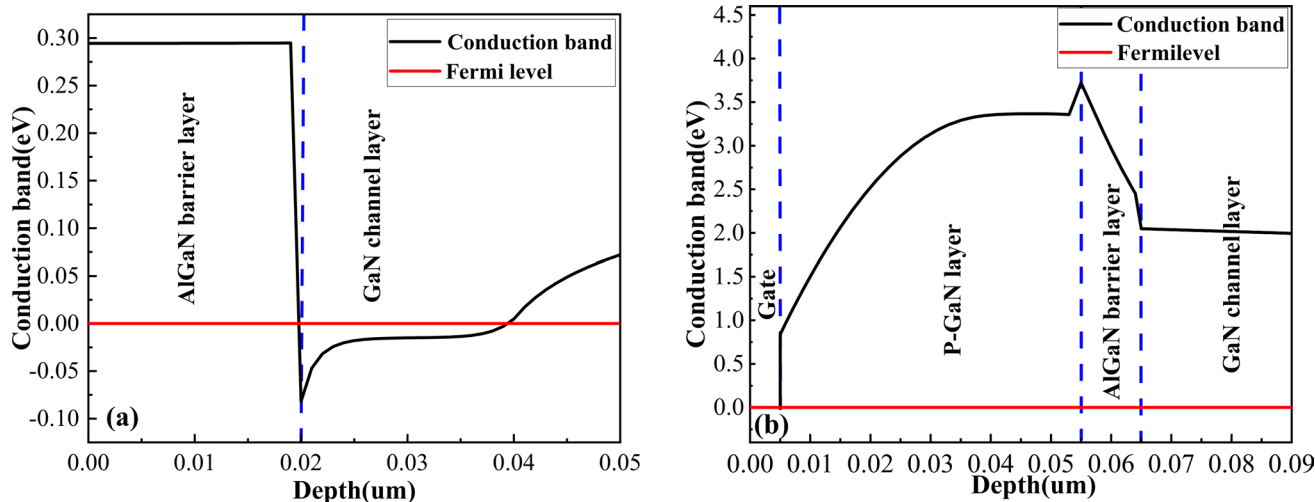


Fig. 16 Conduction band profile of structure D along (a) region 1 and (b) region 2 below the gate, showing triangular potential wells below the Fermi level at zero gate voltage in Fig. 15.

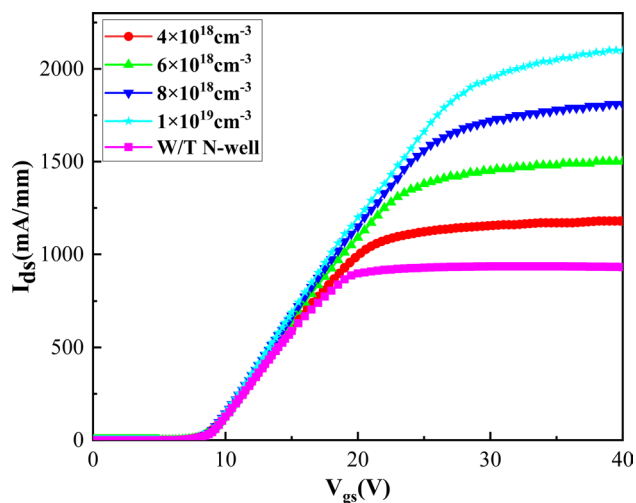


Fig. 17 Transfer characteristic curves for different N-doping concentrations in the N-well.

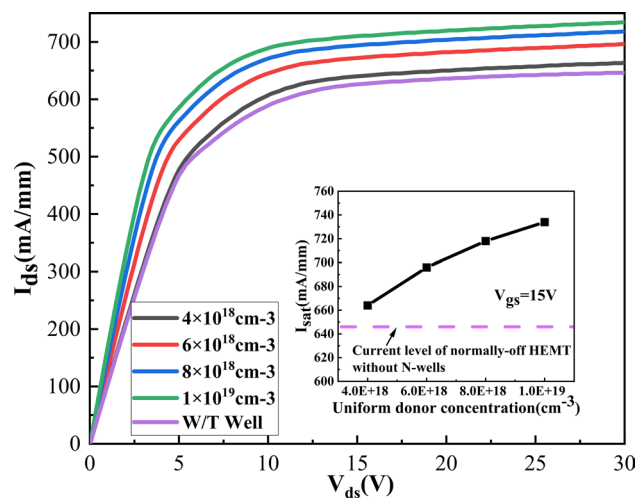


Fig. 18 Output characteristic curves for different N-doping concentrations in the N-well. The inset shows the effect of different doping profiles on the N-well drain current at $V_{gs} = 15$ V.

piezoelectric polarization in the AlGaN barrier layer under the gate electrode was investigated using polarization engineering techniques. Through comparative analysis, it was found that when the Al composition is graded from 0.3 to 0.24, the device exhibits a threshold voltage of $V_{th} = 8.6$ V and a transconductance of $g_m = 97.5$ mS/mm. Most design schemes for enhanced-mode HEMTs significantly impact the source–drain current of the device.

Therefore, in the optimized structure mentioned above, an N-well is added to the GaN channel layer, and the effect of the N-doping concentration on the device is discussed. Simulation results demonstrate that the participation of electrons from the N-well in conduction compensates for the channel degradation, thereby improving the device’s current density. In particular, the saturation current of the device shows the most significant increase when the N-doping concentration is $8 \times 10^{18} \text{ cm}^{-3}$ ($I_{sat} = 718$ mA/mm). Therefore, this structure holds great potential for applications in the field of high-power devices.

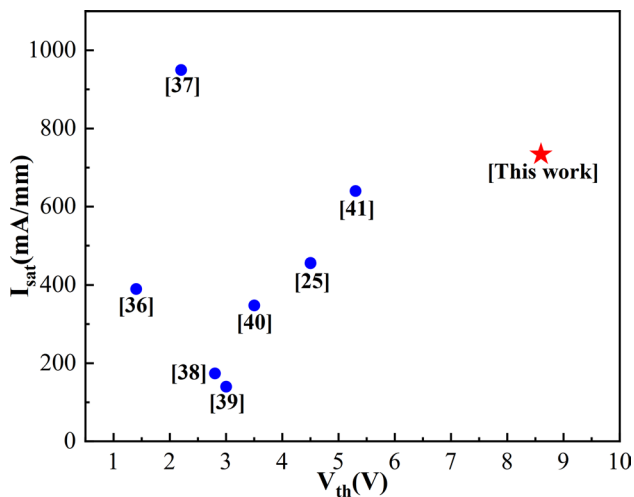


Fig. 19 Comparison of output and transfer characteristics of GaN-based HEMTs. The red star represents this work (HEMT devices of structure D).

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Data Availability The data that support the findings of this study are available from the corresponding authors upon reasonable request.

Conflict of interest The authors declare that they have no conflict of interest.

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