

Fully Electrical Modeling of Thermoelectric Generators with Contact Thermal Resistance Under Different Operating Conditions

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The Seebeck effect is used in thermoelectric generators (TEGs) to supply electronic circuits by converting the waste thermal into electrical energy. This generated electrical power is directly proportional to the temperature difference between the TEG module's hot and cold sides. Depending on the applications, TEGs can be used either under constant temperature gradient between heat reservoirs or constant heat flow conditions. Moreover, the generated electrical power of a TEG depends not only on these operating conditions, but also on the contact thermal resistance. The influence of the contact thermal resistance on the generated electrical power have already been extensively reported in the literature. However, as reported in Park et al. (Energy Convers Manag 86:233, [2014\)](#page-9-0) and Montecucco and Knox (IEEE Trans Power Electron 30:828, [2015\)](#page-10-0), while designing TEG-powered circuit and systems, a TEG module is mostly modeled with a Thévenin equivalent circuit whose resistance is constant and voltage proportional to the temperature gradient applied to the TEG's terminals. This widely used simplified electrical TEG model is inaccurate and not suitable under constant heat flow conditions or when the contact thermal resistance is considered. Moreover, it does not provide realistic behaviour corresponding to the physical phenomena taking place in a TEG. Therefore, from the circuit designer's point of view, faithful and fully electrical TEG models under different operating conditions are needed. Such models are mainly necessary to design and evaluate the power conditioning electronic stages and the maximum power point tracking algorithms of a TEG power supply. In this study, these fully electrical models with the contact thermal resistance taken into account are presented and the analytical expressions of the Thévenin equivalent circuit parameters are provided.

Key words: Thermoelectric generator (TEG), constant temperature gradient, constant heat flow, fully electrical modeling, contact thermal resistance

INTRODUCTION

In recent years, energy harvesting technologies appeared as a solution to satisfy the power-supply requirements of a large variety of devices such as wireless sensors and biomedical devices.¹ Free

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availability of heat makes harvesting energy from it one of the most viable sources of electricity. $2,3$ Thermoelectric generators (TEGs) directly convert heat into electrical energy based on the Seebeck effect. $4-6$ TEGs have gained considerable attention^{[7](#page-9-0)} due to their noiseless operation, stability, and absence of moving parts, thus providing high reli-(Received May 30, 2016; accepted August 30, 2016; accepted May 30, 2016; accepted August 30

Therefore, TEGs are used to recover waste heat in a wide range of applications, ranging from low power (autonomous wireless sensors^{[10](#page-10-0)} and battery charg- ing^{11}) via medium power (automotive, 12 12 12 stoves, 13 13 13 and hybrid photovoltaic-TEG systems¹⁴) up to high power (geothermal power generation^{[15](#page-10-0)}).

A TEG module is made of two dissimilar semiconductors: p - and n -type, which are electrically connected in series, to generate a higher voltage and a higher electrical power, and thermally in parallel to keep each semiconductor element at the same temperature difference. These semicondutor elements are usually arranged in a planer array and fastened between two ceramic hot and cold plates.^{[16](#page-10-0)} The basic structure of a TEG formed by the association of three semiconductor $p-n$ junctions with the electrical load connected to its terminals is presented in Fig. 1. For a given temperature difference between the hot and the cold sides, heat flows through the TEG module as presented in Fig. 1 .^{[17](#page-10-0)} The circulation of heat flow through the semiconductor elements causes a displacement of electrons and holes in opposite directions and, therefore, the generation of an electrical current I passing through the connected load R_L .^{[18](#page-10-0)}

The physical effects that appear in a TEG module when subjected to a temperature difference are the thermal conduction and the Seebeck, Peltier, and Joule effects.^{[16](#page-10-0)} The Thomson effect is commonly neglected when all considered physical properties of a TEG module are temperature independent. $19-21$ The heat flow rates through the hot side Q_H and the cold side Q_C of a TEG module, are expressed respectively by:

$$
Q_H = \frac{T_H - T_C}{\theta_m} + \alpha T_H I - \frac{1}{2} R_E I^2. \tag{1}
$$

$$
Q_C = \frac{T_H - T_C}{\theta_m} + \alpha T_C I + \frac{1}{2} R_E I^2.
$$
 (2)

where R_E and θ_m are the electrical and thermal resistances of N pairs of p –n elements, respectively, and α is the Seebeck coefficient of the TEG module. T_H and T_C denote the internal hot and cold side

Fig. 1. General structure of a TEG module.

temperatures at the terminals of the p and n elements. T'_H and T'_C in Fig. 1 are the external temperatures applied on both hot and cold sides of the TEG module.

Depending on the applications, TEGs can be used either under constant temperature gradient or constant heat flow conditions. In the first case, the temperature difference $\Delta T'$ across a TEG module is considered constant over time, even though its value slowly evolves with variable environmental conditions. On the other hand, the constant heat flow conditions imply that the input heat flow Q_H is considered constant over time. Moreover, the generated electrical power of a TEG depends not only on the working conditions (constant temperature gradient or heat flow), but also on the contact thermal resistance. The influence of these contact thermal resistances on the TEG's generated electrical power and its efficiency in different operating conditions have already been discussed in the literature.^{[22](#page-10-0)–[25](#page-10-0)} Indeed, the thermal contact plates placed on the both sides of the TEG present an interface between the surrounding environment and its semiconductor elements and have a finite thermal resistance. If its value is high or not negligible, the semiconductor elements can perceive a different temperature gradient or a heat flow than the ones really applied on the both sides of the TEG module. However, while designing TEG-powered systems, the TEG module is mostly modeled with a simplified Thévenin equivalent circuit whose resistance is constant and equal to the TEG's electrical resistance and voltage is proportional to the temperature gradient applied to the TEG's terminals, where the coefficient of proportionality is the Seebeck coefficient. This widely used simplified electrical model of TEG provides approximated, inaccurate, and often unrealistic behaviour because the parameters of the thermal environment and those of TEG are not fully taken into account. On the other hand, the recent studies on TEG modelling provide accurate and realistic behaviours through electro-thermal models. $22,23$ Therefore, from the circuit designer's point of view, a faithful and fully electrical model under different working conditions (constant temperature gradient or heat flow) is needed. In this study, these fully electrical models of the TEG module with the contact thermal resistance taken into account are presented and the analytical expressions of the Thévenin equivalent circuit parameters are provided. By providing a realistic fully electrical TEG model reflecting different working and environmental conditions, the benefits are twofold: first, the design of TEG-powered circuits and systems is simplified and more realistic (system modeling and simulation), and second, the harvested thermal energy can be maximized by associating to the models the most adapted maximum power point tracking (MPPT) techniques, whose efficiency relies on the faithfulness of the used electrical models.

The main contributions of the present work can be summarized as follows:

- The unified approach allowing to model a TEG module with or without contact thermal resistance, regardless the working conditions, as a Thévenin equivalent circuit whose parameters are function of these working conditions and the TEG's module properties;
- The fully electrical model of the TEG module under constant temperature gradient conditions. Although some recent studies 22,23 22,23 22,23 have already addressed the electrical behaviour of the TEG module under constant temperature gradient presenting partially the equivalent electrical model, the model proposed in this work is fully electrical and fits to the unified approach described above;
- The fully electrical model of the TEG module under constant heat flow conditions. This is the major and most original contribution of this work. To the best of our knowledge, a Thévenin fully electrical TEG model under constant heat flow has never been presented before in the literature. Montecucco et al.^{[25](#page-10-0)} discovered experimentally the polynomial relationship between the load voltage and the TEG's terminals temperature gradient without identifying analytically this relationship in function of the TEG's parameters and working conditions.

This paper is organized as follows. The fully electrical TEG models under constant temperature gradient and constant heat flow conditions, without and with contact thermal resistance taken into account, are presented in ''TEG Electrical Modeling Under Constant Temperature Gradient Conditions'' and '['TEG Electrical Modeling Under Constant Heat](#page-5-0) [Flow Conditions](#page-5-0)" sections, respectively. In the "[Re](#page-8-0)[sults and Discussion](#page-8-0)'' section, the presented models are discussed and validated through simulation. Finally, the "Conclusion" section concludes this paper giving some perspectives for future works.

TEG ELECTRICAL MODELING UNDER CONSTANT TEMPERATURE GRADIENT **CONDITIONS**

In this section, an equivalent fully electrical Thévenin model of a TEG module without contact thermal resistance is first described. Next, considering the contact thermal resistance, the modified fully electrical TEG model is presented.

TEG Electrical Modeling Without Contact Thermal Resistance

As the contact thermal resistance is not considered, the thermal gradient $\Delta T'$ applied on both hot and cold sides of a TEG module is the same as the

thermal gradient ΔT at the terminals of its p and n elements (see Fig. [1](#page-1-0)), consequently:

$$
\Delta T' = (T'_H - T'_C) = \Delta T = (T_H - T_C). \tag{3}
$$

The open circuit voltage of the TEG module is proportional to the temperature gradient across it. The proportionality factor is the Seebeck coefficient a of the TEG module:

$$
V_{OC}\Big|_{\Delta T' = const, \theta_c = 0} = \alpha \Delta T.
$$
 (4)

where $\alpha = \alpha_p - \alpha_n$, α_p and α_n are the Seebeck coefficients of p and n semiconductor elements, respectively.

The output voltage of the TEG module is defined with:

$$
V_O\Bigg|_{\Delta T'=cnst,\theta_c=0} = \alpha \Delta T - R_E I.
$$
 (5)

The electrical power supplied to the load R_L is the difference between the input and output heat flow, which can be also defined as the product of the output voltage of the TEG module V_O and the load current I:

$$
P\Big|_{\Delta T' = \text{cnst}, \theta_c = 0} = R_L I^2 = Q_H - Q_C = \alpha \Delta T I - R_E I^2.
$$
\n(6)

Consequently, under constant temperature gradient conditions, the TEG module without contact thermal resistance θ_c can be electrically modeled by an equivalent constant voltage source V_{eq} in series with an equivalent internal resistance R_{eq} :

$$
V_{eq}\Big|_{\Delta T'=cnst,\theta_c=0} = \alpha \Delta T = \alpha \Delta T'.
$$
 (7)

$$
R_{eq}\Big|_{\Delta T'=cnst, \theta_c=0} = R_E. \tag{8}
$$

The equivalent Thévenin circuit with the parameters V_{eq} and R_{eq} presented with Eqs. 7 and 8, respectively, is the electrical TEG model commonly used in the literature in the TEG-powered electron-ics circuit and system designs.^{[1,](#page-9-0)[26](#page-10-0)} As seen from the previous demonstration, this electrical TEG model is only valid under constant temperature gradient conditions and does not consider contact thermal resistance, thus greatly approximating the realistic behaviour of the TEG module.

TEG Electrical Modeling with Contact Thermal Resistance

In prior research works on thermoelectric power generation systems, the temperature gradient across the thermoelectric module is considered

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Fig. 2. Thermal modeling of the TEG module with contact thermal resistance.

constant independently of electrical current levels drawn by the connected load R_L . In fact, the waste heat is transferred from the environment to the semiconductor elements of the TEG module though the contact plates, since it is at the origin of the thermal contact resistance θ_c . Because of this contact thermal resistance θ_c , the thermal gradient $\Delta T' = (T_H' - T_C')$ applied on both hot and cold sides of the TEG module is "seen" as $\Delta T = (T_H - T_C) < \Delta T'$ at the terminals of the p and n semiconductor elements.

Using the duality between electrical and thermal phenomena, a thermal modeling of the TEG module with contact thermal resistance can be represented as in Fig. 2.^{[19](#page-10-0)} This duality allows representing the temperature by an electrical voltage, the heat flow by an electrical current, and the thermal resistance by an electrical resistance.

The main goal of using the above mentioned electro-thermal duality is to model the TEG module as an equivalent Thévenin circuit (fully electrical model). To achieve this, the same approach is used throughout the paper: the electrical power of the TEG module is derived and presented as more as possible under the following form $P = V_{eq}I - R_{eq}I^2$, where V_{eq} is the equivalent Thévenin generator and R_{eq} is the corresponding equivalent resistance respectively. In the further discussion, θ_m is the thermal resistance of N pairs of p and n semiconductor elements, whereas θ_c is the contact thermal resistance of contact plates on the top and the bottom of the p and n elements. Each junction formed by p and n elements is connected to two metal contacts from the top and bottom (see Fig. [1\)](#page-1-0), which can be modeled by two series resistors with the value of θ_c/N .^{[27](#page-10-0)}

From Fig. 2, the expression of the heat flow through the metal contacts on the hot side of the TEG module can be written as:

$$
Q_H = \frac{T'_H - T_H}{\theta_c}.\tag{9}
$$

This heat flow is equal to the one flowing through the TEG module (see Eq. [1\)](#page-1-0), which allows writing:

$$
\frac{T_H + \alpha \theta_m T_H I - \frac{1}{2} R_E \theta_m I^2 - T_C}{\theta_m} = \frac{T'_H - T_H}{\theta_c}.
$$
 (10)

The difference between Q_H and Q_C (Eqs. [1](#page-1-0) and [2\)](#page-1-0) provides the relationship between the heat flow difference and the power transmitted to the load R_L :

$$
Q_H - Q_C = I \cdot V_O \Big|_{\theta_c \neq 0}.
$$
\n(11)

where $V_O|_{\rho \in \mathcal{A}^0}$ is the output voltage at the terminals of the load R_L .

When considering the contact thermal resistance, and following the same analysis to determine the expression of the heat flow through the metal contacts on the cold side of the TEG module, it can be deduced from Fig. 2:

$$
\frac{T_H'-T_H}{\theta_c} - \frac{T_C-T_C'}{\theta_c} = V_O\Big|_{\theta_c \neq 0} I.
$$
 (12)

Knowing that $T_H = \Delta T + T_C$, the temperature applied at the cold side of the p and n semiconductor elements is:

$$
T_C = -\frac{\theta_c V_O \Big|_{\theta_c \neq 0} I + \Delta T - T'_H - T'_C}{2}.
$$
 (13)

By replacing T_c by its expression in Eq. 10, the expression of the thermal gradient at the terminals of the p and n elements with contact thermal resistance can be given by:

$$
\Delta T\Big|_{\Delta T' = cnst, \theta_c \neq 0} = \left(\frac{1}{\alpha \theta_c \theta_m I + \theta_m + 2 \theta_c}\right) \times \left\{\theta_c \theta_m [(\alpha \theta_c V_O)\Big|_{\theta_c \neq 0} + R_E)I^2 + (V_O)\Big|_{\theta_c \neq 0} - \alpha (T'_H + T'_C))I] + \theta_m \Delta T'\right\}.
$$
\n(14)

It should be noted that $\Delta T|_{\Delta T' = const, \theta_c \neq 0}$ is load current-dependent.

The output voltage across the load R_L can be defined by the following expression:

$$
V_O\Big|_{\theta_c \neq 0} = \alpha \Delta T \Big|_{\Delta T' = \text{cnst}, \theta_c \neq 0} - R_E I. \tag{15}
$$

By substituting Eq. 15 in Eq. 14, the thermal gradient at the terminals of the p and n elements becomes:

$$
\Delta T\Big|_{\Delta T'=cnst,\theta_c\neq 0} = \left(\frac{1}{\alpha^2 \theta_c^2 \theta_m I^2 - \theta_m - 2 \theta_c}\right) \times \left\{\alpha R_E \theta_c^2 \theta_m I^3 + \alpha \theta_c \theta_m (T'_H + T'_C)I - \theta_m \Delta T'\right\}.
$$
\n(16)

Therefore, the open circuit voltage generated by the TEG module is expressed with:

$$
V_{OC}\Big|_{\Delta T'=cnst, \theta_c \neq 0} = \alpha \Delta T\Big|_{\Delta T'=cnst, \theta_c \neq 0} (I=0)
$$

=
$$
\frac{\theta_m}{\theta_m + 2 \theta_c} \alpha \Delta T'.
$$
 (17)

The output voltage across the load R_L can be expressed by:

$$
V_O\Biggl|_{\Delta T'=cnst,\theta_c\neq 0} = \frac{\alpha R_L \Delta T|_{\Delta T'=cnst,\theta_c\neq 0}}{R_L + R_E}
$$

=
$$
\frac{\alpha R_L \theta_m \left[\alpha \theta_c I(R_E \theta_c I^2 + T'_H + T'_C) - \Delta T' \right]}{(R_L + R_E)(\alpha^2 \theta_c^2 \theta_m I^2 - \theta_m - 2 \theta_c)}.
$$
(18)

with:

$$
R_L = \frac{V_O|_{\Delta T' = \text{cnst}, \theta_c \neq 0}}{I}.
$$
 (19)

By substituting Eq. 19 in Eq. [18,](#page-3-0) the resulting expression of the output voltage as a function of the load current I is derived:

$$
V_O \Big|_{\Delta T' = \text{cnst}, \theta_c \neq 0} = \left(\frac{1}{\alpha^2 \theta_c^2 \theta_m I^2 - \theta_m - 2 \theta_c} \right) \times \left\{ \left[\alpha^2 \theta_c \theta_m (T'_H + T'_C) + R_E (\theta_m + 2 \theta_c) \right] I - \alpha \theta_m \Delta T' \right\}.
$$
\n(20)

The electrical power generated by the TEG module is defined as the product of the output voltage V_O with the load current I , and can be expressed with the following expression:

$$
P\Big|_{\Delta T' = \text{cnst}, \theta_c \neq 0} = \left(\frac{1}{\alpha^2 \theta_c^2 \theta_m I^2 - \theta_m - 2 \theta_c}\right) \times \left\{ \left[\alpha^2 \theta_c \theta_m (T_H' + T_C') + R_E (\theta_m + 2 \theta_c) \right] I^2 - \alpha \theta_m \Delta T' I \right\}.
$$
\n(21)

As expected,

$$
\lim_{\theta_c \to 0} P\Big|_{\Delta T' = const, \theta_c \neq 0} = P\Big|_{\Delta T' = const, \theta_c = 0}.
$$
 (22)

Knowing that when the contact thermal resistance converges towards zero $(\theta_c \rightarrow 0)$, the thermal gradient $\Delta T'$ is equal to ΔT (Eq. [3\)](#page-2-0).

Equation 21 does not allow expressing the electrical power generated by the TEG module in the desired form with equivalent voltage and resistance, as was the case without contact thermal resistance $(Eq. 6)$ $(Eq. 6)$. For these reasons, an approximate expression of the electrical power can be proposed. This can be achieved by neglecting the term $\alpha^2 \theta_c^2 \theta_m I^2$ behind $\theta_m + 2 \theta_c$. The V–I characteristic of the TEG module with the contact thermal resistance is nonlinear, due to the term $\alpha^2 \theta_c^2 \theta_m I^2$. Analytically, this non-linearity occurs for large values of the load current greater few orders of magnitude than the short-current value I_{SC} , which explains why this term can be neglected. Moreover, the V–I characteristic is linear for positive values of the load current, $I > 0$. This approximation is also justified by the order of magnitude of $\alpha^2 \approx 10^{-4}$, whereas the order of magnitudes of other parameters is about unity. By applying this approximation to Eq. 21 , it becomes:

$$
P\Big|_{\Delta T'=cnst,\theta_c\neq 0} \approx \frac{\theta_m}{\theta_m+2\,\theta_c} \alpha \Delta T'I
$$

$$
-[R_E + \frac{\alpha^2 \,\theta_c \,\theta_m}{\theta_m+2\,\theta_c} (T'_H + T'_C)]I^2.
$$
(23)

If we apply the same approximation to Eq. 20, the output voltage V_O of the TEG module becomes:

$$
V_O\Big|_{\Delta T' = \text{cnst}, \theta_c \neq 0} \approx \frac{\theta_m}{\theta_m + 2 \theta_c} \alpha \Delta T'
$$

-
$$
[R_E + \frac{\alpha^2 \theta_c \theta_m}{\theta_m + 2 \theta_c} (T'_H + T'_C)]I.
$$
 (24)

Consequently, under constant temperature gradient conditions, the TEG module with contact thermal resistance θ_c can also be modeled by an equivalent constant voltage source V_{eq} in series with an equivalent internal resistance R_{eq} (Thévenin equivalent circuit):

$$
V_{eq}\Big|_{\Delta T'=cnst, \theta_c \neq 0} = \alpha \Delta T' \frac{\theta_m}{\theta_m + 2\theta_c}.
$$
 (25)

$$
R_{eq} \Big|_{\Delta T' = \text{cnst}, \theta_c \neq 0} \approx R_E + \frac{\alpha^2 \theta_c \theta_m (T'_H + T'_C)}{\theta_m + 2\theta_c}.
$$
 (26)

Conclusion

To conclude this section, one can summarize the results of the proposed unified approach to model the TEG module as a Thévenin equivalent circuit under constant temperature gradient conditions. The equivalent electrical model of the TEG module under constant temperature gradient conditions can be represented as in Fig. 3, where V_{eq} and R_{eq} are given by Eqs. [7](#page-2-0) and [8](#page-2-0) for the model without contact thermal resistance, and by Eqs. 25 and 26 for the model with contact thermal resistance.

Fig. 3. Equivalent electrical model of the TEG under constant temperature gradient conditions.

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TEG ELECTRICAL MODELING UNDER CONSTANT HEAT FLOW CONDITIONS

In most practical applications, as in the case of thermal energy recovery systems for automobile exhaust gas, the TEG module is subject to a constant thermal input flow instead of a constant temperature gradient. This is designated as ''constant heat'' operation conditions in the literature. Kumar et al.[28](#page-10-0) observed a significant difference between the electric power generated by a TEG module with an exhaust gas flow rate and with a temperature gradient. Even if the available input thermal heat flow changes over time, its rate of variation will be few order of magnitude slower than the dynamic electrical response of the TEG module.²⁹

In this section, by applying the same approach used in the previous sections for the TEG module under constant gradient conditions, the equivalent fully electrical models of the TEG module without and with contact thermal resistance under constant heat flow conditions are presented. In both cases, the input thermal power Q_H flowing through the TEG module and the cold source temperature T_{C}^{\prime} are assumed constant throughout this discussion. In most TEG systems, the cold side temperature remains almost constant with relatively small changes in thermal power flowing into the cold side. 25 25 25 Moreover, to the best of our knowledge, no fully electrical TEG model under constant heat flow conditions has been previously presented in the literature.

TEG Electrical Modeling Without Contact Thermal Resistance

Given that the contact thermal resistance θ_c is not taken into account, Eq. [3](#page-2-0) is verified. Moreover, from Eq. [1](#page-1-0) it can be obtained:

$$
\Delta T\Big|_{Q_H=cnst,\theta_c=0} = \frac{\frac{1}{2}R_E\theta_m I^2 - \alpha\theta_m T_c'I + \theta_m Q_H}{1 + \alpha\theta_m I}.
$$
 (27)

The open circuit voltage V_{OC} of the TEG module is expressed with:

$$
V_{OC}\Big|_{Q_H=cnst, \theta_c=0} = \alpha \Delta T\Big|_{Q_H=cnst, \theta_c=0, I=0}
$$
\n
$$
= \alpha \theta_m Q_H.
$$
\n(28)

The electrical power delivered by the TEG module is defined by:

$$
P\Big|_{Q_H=cnst,\theta_c=0} = \alpha \Delta T \Big|_{Q_H=cnst,\theta_c=0} I - R_E I^2
$$

=
$$
-\frac{1}{2} \alpha \theta_m R_E I^3 - (R_E + \alpha^2 \theta_m T_C')I^2 + \alpha \theta_m Q_H I}{1 + \alpha \theta_m I}.
$$
(29)

An approximate expression of this electrical power can be proposed by using the development of Maclaurin up to the 4th order:

$$
P\Big|_{Q_H = cnst, \theta_c = 0} = \alpha Q_H \theta_m I - \{R_E + \alpha^2 Q_H \theta_m^2 + \alpha^2 \theta_m T'_C - \left[\frac{\alpha R_E \theta_m}{2} + \alpha^3 \theta_m^2 T'_C + \alpha^3 \theta_m^3 Q_H\right] I \quad (30) + \left[\alpha^4 \theta_m^3 T'_C + \alpha^4 Q_H \theta_m^4 + \frac{\alpha^2 R_E \theta_m^2}{2}\right] I^2 \} I^2.
$$

Consequently, under constant heat flow conditions, the TEG module without contact thermal resistance can also be modeled by an equivalent constant voltage source V_{eq} in series with a variable equivalent internal resistance R_{eq} (Thévenin equivalent circuit):

$$
V_{eq}\Big|_{Q_H=cnst,\theta_c=0} = \alpha \theta_m Q_H.
$$
 (31)

$$
R_{eq} \Big|_{Q_H = const, \theta_c = 0} \approx R_E + \alpha^2 Q_H \theta_m^2 + \alpha^2 \theta_m T_C'
$$

$$
- \Big[\frac{\alpha R_E \theta_m}{2} + \alpha^3 \theta_m^2 T_C' + \alpha^3 \theta_m^3 Q_H \Big] I
$$

$$
+ \Big[\alpha^4 \theta_m^3 T_C' + \alpha^4 Q_H \theta_m^4 + \frac{\alpha^2 R_E \theta_m^2}{2} \Big] I^2.
$$
 (32)

TEG Electrical Modeling with Contact Thermal Resistance

Using the expression of the thermal gradient at the terminals of the p and n elements under constant heat flow conditions given by Eq. 27, and knowing that T_C' is different from T_C when considering contact thermal resistance, the thermal gradient at the terminals of the p and n elements with contact thermal resistance can be expressed as:

$$
\Delta T\Big|_{Q_H=cnst,\theta_c\neq 0} = \frac{\frac{1}{2}R_E\theta_m I^2 - \alpha \theta_m T_C I + Q_H \theta_m}{1 + \alpha \theta_m I}.
$$
 (33)

By substituting the expression of T_C given by Eq. [13](#page-3-0) in Eq. 33, we obtain:

$$
\Delta T\Big|_{Q_H=cnst,\theta_c\neq 0} = \left(\frac{1}{\alpha\theta_m I + 2}\right)
$$

$$
\times \left\{ (\theta_m (R_E + \alpha \theta_c V_O\Big|_{\theta_c\neq 0}) I^2 - \alpha\theta_m (T'_H + T'_C)I + 2Q_H\theta_m \right\}. \tag{34}
$$

However, this thermal gradient must be expressed only as a function of Q_H and T_C' . From Eq. [9](#page-3-0) we can write:

$$
T'_H = T_H + Q_H \theta_c = \Delta T + T_C + Q_H \theta_c. \tag{35}
$$

and from Eq. [12](#page-3-0):

$$
T_C = -\theta_c V_O \Big|_{\theta_c \neq 0} I + T'_C + Q_H \theta_c.
$$
 (36)

By substituting Eq. [36](#page-5-0) in Eq. [35](#page-5-0), and then by substituting the result in Eq. [34,](#page-5-0) the thermal gradient across the p and n elements becomes: -

$$
\Delta T\Big|_{Q_H = \text{cnst}, \theta_c \neq 0} = \left(\frac{1}{1 + \alpha \theta_m I}\right)
$$

\$\times \left\{\theta_m \left(\frac{1}{2} R_E + \alpha \theta_c V_O\right|_{\theta_c \neq 0})I^2 - \alpha \theta_m (T_C' + Q_H \theta_c)I\$
\$+ Q_H \theta_m\$}. \tag{37}

As expected, under constant heat flow conditions:

$$
\lim_{\theta_c \to 0} \Delta T \Big|_{Q_H = \text{cnst}, \theta_c \neq 0} = \Delta T \Big|_{Q_H = \text{cnst}, \theta_c = 0}.
$$
 (38)

The open circuit voltage of the TEG module is given by:

$$
V_{OC}\Big|_{Q_H=cnst,\theta_c\neq 0} = \alpha \Delta T\Big|_{Q_H=cnst,\theta_c\neq 0} (I=0)
$$
\n
$$
=\alpha \theta_m Q_H.
$$
\n(39)

By following the same analysis used for the determi-nation of the output voltage in "[TEG Electrical](#page-2-0) [Modeling with Contact Thermal Resistance](#page-2-0)'' section, the output voltage of the TEG module can be obtained:

$$
V_O\Big|_{Q_H = \text{cnst}, \theta_c \neq 0} = \left(\frac{1}{\alpha^2 \theta_c \theta_m I^2 - \alpha \theta_m I - 1}\right)
$$

\$\times \left\{\frac{1}{2} \alpha R_E \theta_m I^2 + (R_E + \alpha^2 \theta_m T_C' + \alpha^2 Q_H \theta_c \theta_m)I\$ - $\alpha Q_H \theta_m\right\}.$ (40)

By the same, the electrical power generated by the TEG module is given by:

$$
P\Big|_{Q_H = const, \theta_c \neq 0} = \left(\frac{1}{\alpha^2 \theta_c \theta_m I^2 - \alpha \theta_m I - 1}\right) \times \left\{\frac{1}{2} \alpha R_E \theta_m I^3 + (R_E + \alpha^2 \theta_m T_C' + \alpha^2 Q_H \theta_c \theta_m) I^2 - \alpha Q_H \theta_m I\right\}.
$$
\n(41)

Similarly, using the development of Maclaurin up to the 4th order, the expression of the electrical power can be approximated with:

$$
P\Big|_{Q_{H}=cnst,\theta_{c}\neq0} \approx \alpha Q_{H}\theta_{m}I - \Big\{R_{E} + \alpha^{2}\theta_{c}\theta_{m}Q_{H}\n+ \alpha^{2}\theta_{m}^{2}Q_{H} + \alpha^{2}\theta_{m}T_{C}' - \Big[\frac{\alpha R_{E}\theta_{m}}{2} + \alpha^{3}\theta_{m}^{2}T_{C}'\n+ \alpha^{3}\theta_{m}^{3}Q_{H} + 2\alpha^{3}\theta_{m}^{2}\theta_{c}Q_{H}\Big]I + \Big[\alpha^{4}\theta_{m}^{3}T_{C}' + \alpha^{4}Q_{H}\theta_{m}^{4}\n+ \frac{\alpha^{2}R_{E}\theta_{m}^{2}}{2} + \alpha^{2}R_{E}\theta_{m}\theta_{c} + \alpha^{4}\theta_{m}^{2}\theta_{c}Q_{H}(\theta_{c} + 3\theta_{m})\n+ \alpha^{4}\theta_{m}^{3}\theta_{c}^{2}T_{C}'\Big]I^{2}\Big\}I^{2}.
$$
\n(42)

The choice to develop only up to the 4th order using Maclaurin series can be justified by the results reported by Montecucco et al.^{[25](#page-10-0)} where a quadratic relationship between the TEG's output voltage and the temperature gradient has been observed experimentally. If this nonlinear voltage relationship is transposed to the power, the four order polynomial current dependence is obtained.

As expected, under constant heat flow conditions:

$$
\lim_{\theta_c \to 0} P \Big|_{Q_H = \text{cnst}, \theta_c \neq 0} = P \Big|_{Q_H = \text{cnst}, \theta_c = 0} . \tag{43}
$$

Applying the same development of Maclaurin to Eq. 40, the output voltage of the TEG module becomes:

$$
V_O\Big|_{Q_H=cnst,\theta_c\neq 0} \approx \alpha Q_H \theta_m - \left\{R_E + \alpha^2 \theta_c \theta_m Q_H\right.+ \alpha^2 \theta_m^2 Q_H + \alpha^2 \theta_m T'_C - \left[\frac{\alpha R_E \theta_m}{2} + \alpha^3 \theta_m^2 T'_C\right.+ \alpha^3 \theta_m^3 Q_H + 2\alpha^3 \theta_m^2 \theta_c Q_H\right]I + \left[\alpha^4 \theta_m^3 T'_C + \alpha^4 Q_H \theta_m^4\right.+ \frac{\alpha^2 R_E \theta_m^2}{2} + \alpha^2 R_E \theta_m \theta_c + \alpha^4 \theta_m^2 \theta_c Q_H (\theta_c + 3\theta_m)+ \alpha^4 \theta_m^3 \theta_c^2 T'_C\Big|I^2\Big]I.
$$
 (44)

Consequently, under constant heat flow conditions, the TEG module with contact thermal resistance can also be modeled with an equivalent constant voltage source V_{eq} in series with a variable equivalent internal resistance R_{eq} (Thévenin equivalent circuit):

$$
V_{eq}\Big|_{Q_H=cnst, \theta_c\neq 0} = \alpha \theta_m Q_H.
$$
 (45)

$$
R_{eq}\Big|_{Q_H=cnst,\theta_c\neq 0} \approx R_E + \alpha^2 \theta_c \theta_m Q_H + \alpha^2 \theta_m^2 Q_H
$$

+ $\alpha^2 \theta_m T'_C - \left[\frac{\alpha R_E \theta_m}{2} + \alpha^3 \theta_m^2 T'_C + \alpha^3 \theta_m^3 Q_H\right]$
+ $2\alpha^3 \theta_m^2 \theta_c Q_H |I + [\alpha^4 \theta_m^3 T'_C + \alpha^4 Q_H \theta_m^4 + \frac{\alpha^2 R_E \theta_m^2}{2}$
+ $\alpha^2 R_E \theta_m \theta_c + \alpha^4 \theta_m^2 \theta_c Q_H (\theta_c + 3\theta_m)$
+ $\alpha^4 \theta_m^3 \theta_c^2 T'_C |I^2$. (46)

Fig. 4. Equivalent electrical model of the TEG module under constant heat flow conditions.

	Constant $\Delta T'$		Constant Q_H	
	$\boldsymbol{V_{eq}}$	$\bm{R_{eq}}$	$\boldsymbol{V_{eq}}$	R_{eq}
$\theta_c=0$	$\alpha \Delta T'$	$R_E\,$	$\alpha\theta_m Q_H$	$R_E + \alpha^2 Q_H \theta_m^2 + \alpha^2 \theta_m T_C' - \frac{\alpha R_E \theta_m}{2} +$ $\alpha^3\theta_m^2T'_C+\alpha^3\theta_m^3Q_H]I+[\alpha^4\theta_m^3T'_C+$ $\alpha^4 Q_H \theta_m^4 + \frac{\alpha^2 R_E \theta_m^2}{2}]I^2 \ .$
$\theta_c\neq 0$	$\frac{\theta_m}{\theta_m+2\theta_c} \alpha \Delta T'$	$R_E+\frac{\alpha^2\theta_c\theta_m(T_H'+T_C')}{\theta_m+2\theta_c}$	$\alpha\theta_m Q_H$	$R_E + \alpha^2 \theta_c \theta_m Q_H + \alpha^2 \theta_m T_C' + \alpha^2 \theta_m^2 Q_H - [\frac{\alpha R_E \theta_m}{2}]$ $+\alpha^3\theta_m^2T'_C+\alpha^3\theta_m^3Q_H+2\alpha^3\theta_m^2\theta_cQ_H I+$ $[\alpha^4\theta_m^3T'_C+\alpha^4Q_H\theta_m^4+\frac{\alpha^2R_E\theta_m^2}{2}+\alpha^2R_E\theta_m\theta_c$ $+\alpha^4\theta_m^2\theta_cQ_H(\theta_c+3\theta_m)+\alpha^4\theta_m^3\theta_c^2T'_C]I^2$

Table I. Summary of electrical modeling of the TEG module under constant temperature gradient and heat conditions

Fig. 5. $P = f(I, \theta_c)$ under constant temperature gradient conditions, where I_{SC} is the short-circuit current.

Conclusion

To conclude this section, it is important to underline one of the major contributions of this paper: a

TEG module under constant heat flow conditions can also be represented as in Fig. [4](#page-6-0), by an equivalent Thévenin electrical model where V_{eq} and R_{eq} are given by Eqs. [31](#page-5-0) and [32](#page-5-0) for the TEG module without contact thermal resistance, and by Eqs. 45 and [46](#page-6-0) for the TEG module with contact thermal resistance.

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 $\beta=0$ $\beta = 6.7$ \mathbb{Z} $\beta = 13$. $V_{OC}(\theta_{c}=0)$ $\beta=30$ $\beta = 66.$ 5 $\overline{3}$ $\overline{2}$ $\overline{1}$ 100 10 20 30 60 80 90 $[\%]$ $I_{SC}(\theta_c)$

 $\cdot \beta = 13$

 $\beta = 30$ $\beta = 66$

 $\beta = 100$

Fig. 7. $V_O = f(I, \theta_c)$ under constant temperature gradient conditions. Fig. 8. $V_O = f(I, \theta_c)$ under constant heat flow conditions.

RESULTS AND DISCUSSION

Table [I](#page-7-0) summarizes the analytical expressions of the parameters for the fully electrical TEG models (Thévenin equivalent circuit) detailed previously. It includes the expressions of the equivalent open circuit voltage and the equivalent internal resistance of the TEG module under constant temperature gradient, and constant heat flow conditions with $\theta_c = 0$ and with $\theta_c \neq 0$.

Under constant temperature gradient conditions, as seen from Table [I,](#page-7-0) the equivalent open-circuit voltage of the TEG module with the contact thermal resistance taken into account is lower that the one in the absence of the latter. Apertet et al. observed in^{22} in^{22} in^{22} a decrease in the open-circuit voltage due to an additional term depending on the load current and TEG's parameters including the contact thermal resistance θ_c . However, from the electrical modeling's point of view, the equivalent Thévenin generator which represents the open-circuit voltage of the TEG module in the absence of the load current, should not be load current dependent, as was presented in^{[22](#page-10-0)} (Eq. 11 in²²). In the presented fully electrical TEG model, the equivalent voltage source is constant for a given temperature gradient and is valid regardless the value of the load current. On the other hand, the equivalent electrical resistance of the TEG module with the contact thermal resistance taken into account is larger than without it. The additional term presented in Table [I](#page-7-0) depends on the thermal resistance, Seebeck coefficient, and operating conditions of the TEG module. This term has also been identified by Apertet et al. 22 22 22 in terms of the figure of merit $Z\overline{T}$, electrical resistance R_E , and thermal conductances.

Under constant heat flow conditions, the results presented in Table [I](#page-7-0) show that there is no difference between the equivalent open-circuit voltage of the TEG module with and without the contact thermal resistance taken into account. However, the relation

In order to appreciate the faithfulness of the presented fully electrical TEG models under different operation conditions, the analytical expressions of Table [I](#page-7-0) were validated through simulation using Matlab software. The used physical parameters of the TEG module are presented in Table [II.](#page-7-0) The percentage ratio β between the contact thermal resistance θ_c and the thermal resistance θ_m have been varied from 0% (case of the TEG module with $\theta_c = 0$) to 100% (case of the TEG module with a value of θ_c equal to its thermal resistance θ_m). For all values of θ_c , the electrical power and the output voltage of the TEG module were observed.

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Figure [5](#page-7-0) shows the TEG's electrical power versus load current for different values of the contact thermal resistance θ_c under constant temperature gradient conditions. The parabolic curves in lines result from the analytical expression of the electrical power without approximation $(Eq. 21)$ $(Eq. 21)$, and the curves with markers result from the approximate one (Eq. [23](#page-4-0)). It can be seen from Fig. [5](#page-7-0) that the electric power of the TEG module obtained with the proposed approximated expression is close to the exact one. As expected from the physical modeling studies, the electrical power decreases as θ_c increases.[22,23](#page-10-0) This electrical power reduction results from the decrease of the effective thermal gradient applied to the TEG module, which moves from $\Delta T'$ to ΔT ($\Delta T < \Delta T'$), essentially caused by the presence of the contact thermal resistance θ_c .

Figure [6](#page-7-0) shows the TEG's electrical power versus load current for different values of the contact thermal resistance θ_c under constant heat flow conditions. Similarly as in the constant temperature gradient case, the parabolic curves in lines result from the exact analytical expression of the electrical power (Eq. [41\)](#page-6-0), whereas the curves with markers result from the approximate expression obtained by using the development of Maclaurin $(Eq. 42)$ $(Eq. 42)$ $(Eq. 42)$. As in the previous case, the electrical power decreases as θ_c increases. However, this decrease is smaller compared to the one obtained in the case of the TEG module under constant temperature gradient conditions. It does mean that the electrical modeling of the TEG under constant heat flow conditions is very little influenced by the presence of θ_c compared to the one under constant temperature gradient conditions.

Figure [7](#page-8-0) shows the TEG's output voltage V_O versus load current I for different values of the contact thermal resistance θ_c under constant temperature gradient conditions. It can be seen that the exact analytical expression of the output voltage of the TEG module (Eq. [20](#page-4-0)) presented with lines for different values of the contact thermal resistance are superposed to the curves with markers used to present the approximated expression given by Eq. [24](#page-4-0). This justifies the validity of the used approximation. The equivalent internal resistance of the TEG module R_{eq} determines the slope of the linear V–I characteristic. As expected from the physical modeling studies, Fig. [7](#page-8-0) illustrates that more θ_c increases, more the output voltage of the TEG module decreases and the slope of the V–I characteristic becomes steeper. $22,23$ This can be interpreted by the increase of the equivalent internal electrical resistance of the TEG module "seen" by the connected load R_L , as represented with Eq. [26](#page-4-0). Notice that the open circuit voltage of the TEG module decreases as θ_c increases.

Figure [8](#page-8-0) shows the TEG's output voltage V_O versus load current I for different values of the contact thermal resistance θ_c under constant heat flow conditions. Similarly, the curves with lines used to represent the exact analytical equation of the output voltage of the TEG module $(Eq. 40)$ $(Eq. 40)$, are superposed to

the curves with markers used to represent the approximated expression given by Eq. [44](#page-6-0). This result justifies the validity of the approximation made with the development of Maclaurin. Unlike the constant temperature gradient conditions, the TEG module under constant heat flow conditions has the same open circuit voltage V_{OC} , whatever the value of the contact thermal resistance θ_c . The slope of the V–I characteristic in this case is slightly modified with θ_c . This can explained by comparing the equivalent internal electrical resistance of the TEG module with $\theta_c = 0$ and $\theta_c \neq 0$ given by Eqs. [32](#page-5-0) and [46](#page-6-0) respectively, which are almost equal. The main difference lies in the additional terms present in Eq. [46](#page-6-0) whose contribution is low since they are multiplied either by α^2 or α^4 .

CONCLUSION

In this study, the fully electrical modeling of the TEG module taking into account its contact thermal resistance under different operation conditions is presented and discussed. An unified fully electrical TEG modeling approach based on a Thévenin equivalent circuit were proposed and the analytical expressions of the parameters V_{eq} and R_{eq} were identified without and with the contact thermal resistance: first under constant temperature gradient, and second under constant heat flow conditions. The electrical model of the TEG module under constant temperature gradient conditions is linear and depends only on the TEG's electrical and thermal parameters. On the other hand, the electrical model of the TEG module under constant heat flow conditions is non linear and load current dependent. The proposed fully electrical TEG models were validated through simulation. The obtained simulation results fit into the already reported physical behaviour of the TEG modules presented in recent research studies.^{[22,23,25](#page-10-0)} The proposed fully electrical models of the TEG module can be used in the design of the TEG powered electronics circuits and systems in simulation and modeling phases as well in the development of the most adapted MPPT techniques to maximize the harvesting of the available waste thermal energy. 1,26 1,26 1,26

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