

# Effect of  $(Ca_{0.8}Sr_{0.2})_{0.6}La_{0.267}TiO<sub>3</sub>$  on Phase, Microstructure, and Microwave Dielectric Properties of  $Mg_{0.95}Zn_{0.05}TiO<sub>3</sub>$ Synthesized by Polymeric Precursor Method

## ABDUL NAEEM, $^1$  ASAD ULLAH, $^{1,4}$  TAHIRA MAHMOOD, $^1$ YASEEN IQBAL, $^2$  ASAD MAHMOOD, $^1$  and MUHAMMAD HAMAYUN $^3$

1.—National Centre of Excellence in Physical Chemistry, University of Peshawar, Peshawar 25120, Pakistan. 2.—Department of Physics, Material Research Centre, University of Peshawar, Peshawar 25120, Pakistan. 3.—Department of Chemistry, Institute of Natural Sciences, University of Gujrat, Gujrat, Pakistan. 4.—e-mail: asadncepc@upesh.edu.pk

A number of compounds in the  $(1-x) {\rm Mg}_{0.95} {\rm Zn}_{0.05} {\rm TiO}_3\text{-}x {\rm (Ca}_{0.8} {\rm Sr}_{0.2})_{0.6}$  $La_{0.267}TiO<sub>3</sub>$  (x = 0 to 0.25) composition series have been obtained via a polymeric precursor route to investigate the effect of increasing  $(Ca_{0.8}Sr_{0.2})_{0.6}$ .  $La_{0.267}TiO_3$  proportion on the phase, microstructure, and microwave dielectric properties of the sintered ceramics. Composite powders having nanometric particles were obtained by calcining the precursors at 700°C. Refinement results revealed that these samples comprised a mixture of  $Mg_{0.95}Zn_{0.05}TiO_3$  and  $(Ca_{0.8}Sr_{0.2})_{0.6}La_{0.267}TiO<sub>3</sub>$  ceramics. A combination of optimum microwave dielectric properties, i.e., dielectric constant of 25.17, quality factor of  $58,\!754$  GHz, and temperature coefficient of resonant frequency of  $-5.8$  ppm/  $\rm{^{\circ}C}$ , was achieved for the  $x = 0.2$  composition sintered at 1200 $\rm{^{\circ}C}$ .

Key words: Microwave dielectrics, electroceramics, quality factor, polymeric

## INTRODUCTION

Microwave dielectrics have played an important role in the rapid development of advanced communication systems, due to their wide range of applications in terrestrial and satellite communication technologies such as the global positioning system (GPS), software radio, and direct broadcast satellite television (DBS TV). $^{1-4}$  The development of highfrequency wireless communication technology and the consequent growth of the related industry over recent decades have substantially increased demand for microwave dielectric materials. For practical applications in microwave devices, a dielectric material must have a reasonably moderate dielectric constant  $(\varepsilon_r > 20)$  for size miniaturization, higher quality factor ( $Q \times f > 20,000$ ) for target frequency selectivity, and near-zero temperature coefficient of resonant frequency  $(\tau_f)$  for temperature stability of the system. $5,6$ 

Magnesium titanate ( $MgTiO<sub>3</sub>$ ) with (R3) symmetry is a promising microwave dielectric material, having  $\varepsilon_{\rm r} = 17$ ,  $Q \times f = 160,000$  GHz, and  $\tau_f = -51$  ppm/°C.<sup>[7](#page-7-0)</sup> Huang et al.<sup>[8](#page-7-0)</sup> reported enhanced  $Q \times f$  for  ${ {\rm Mg}_{0.95} Z_{\rm N0.05} {\rm TiO}_3}$  ceramic via partial replacement of  $Mg^{2+}$  with  $Zn^{2+}$  in MgTiO<sub>3</sub>, but a high negative  $\tau_f$  value rendered it unsuitable for commercial applications. Fortunately, it is possible to adjust the high negative  $\tau_f$  of  $Mg_{0.95}Zn_{0.05}TiO_3$  via forming a composite with another material having high positive  $\tau_f$  value.<sup>[9](#page-7-0)</sup> One such candidate material to compensate high negative  $\tau_f$  values is CaTiO<sub>3</sub> with orthorhombic (Pbnm) structure, exhibiting  $\varepsilon_{\rm r} = 170$ ,  $Q \times f = 6000$  GHz, and  $\tau_{\rm f} = +800$  ppm/°C.<sup>10</sup> Partial substitution of La<sup>+3</sup> for  $\text{Ca}^{+2}$  in  $\text{CaTiO}_3$  has been reported to improve the low  $Q \times f$  of CaTiO<sub>3</sub>.<sup>[11](#page-7-0)</sup> Similarly,  $\varepsilon_r$  can be enhanced by partial substitution of  $Sr^{+2}$  for  $Ca^{+2}$ .<sup>[12](#page-7-0)</sup> The authors also reported that  $(Ca_{0.8}Sr_{0.2})_{0.6}La_{0.267}TiO_3$  ceramic synthesized by a polymeric precursor method exhibited reasonable microwave dielectric properties, such as  $\varepsilon_r = 116.2$ ,  $Q \times f = 15,730 \text{ GHz}, \text{ and } \tau_f = +231 \text{ ppm}^{\circ} \text{C}^{13} \text{ Con-}$ (Received October 11, 2015; accepted April 28, 2016;  $\mathcal{L} \wedge \mathcal{L} = 13,130$  Virial, and  $t_f = +231$  ppm  $\vee$ . Conservationally, the solid-state sintering route is applied milionally, the solid-state sintering route is a

published online May 18, 2016)

for ceramics processing, but a common problem associated with this method is the requirement for high calcination and sintering temperatures, causing irregular grain growth during sintering. The main advantage of fabricating ceramics by chemical routes is the formation of highly stoichiometric, pure, and homogeneous nanosized ceramic powder, ideal to achieve dense ceramics at relatively lower temperatures. The preparation and characterization of nanostructured materials have attracted increased research attention in recent years, not only for academic research but also for industrial applications, due to their unusual properties compared with materials synthesized by conventional methods.<sup>14,15</sup> The polymeric precursor route is simple and can be employed to produce nanosized multicomponent ceramics.<sup>[3,16,17](#page-7-0)</sup> This method consists of formation of a polymeric complex starting from a polyhydroxy alcohol and an alpha-hydroxycarboxylic acid, with metallic cations homogeneously distributed throughout the matrix. The advantage of the polymeric precursor method is reduced segregation of different metal ions in the metal-chelate complexes. Immobilization of metalchelate complexes in such a rigid organic polymer net can reduce segregation of individual metals during pyrolysis at high temperatures.

In the present study, the effect of  $(Ca_{0.8}Sr_{0.2})_{0.6}$  $La<sub>0.267</sub>TiO<sub>3</sub>$  on the phase constitution, microstructure, and microwave dielectric properties of  $Mg_{0.95}Zn_{0.05}TiO<sub>3</sub>$  fabricated by the polymeric precursor route was investigated.

#### EXPERIMENTAL PROCEDURES

## Synthesis

High-purity  $(299%)$  magnesium nitrate [Mg(NO<sub>3</sub>)<sub>2</sub>·6H<sub>2</sub>O], zinc nitrate [Zn(NO<sub>3</sub>)<sub>2</sub>·6H<sub>2</sub>O], calcium nitrate  $\text{[Ca}(\text{NO}_3)_2 \cdot 4\text{H}_2\text{O}]$ , strontium nitrate  $[Sr(NO<sub>3</sub>)<sub>2</sub>]$ , lanthanum nitrate  $[La(NO<sub>3</sub>)<sub>3</sub>·6H<sub>2</sub>O]$  $(\geq 99\%$ , Scharlau), ethylene glycol  $(C_2H_6O_2)$ , and citric acid  $(C_6H_8O_7H_2O)$  were acquired from Scharlau (Barcelona, Spain), and  $\geq 97\%$  pure titanium isopropoxide  $[Ti(C_3H_7O)_4]$  was acquired from Sigma Aldrich (Dorset, UK). Titanium citrate solution was prepared by first dissolving stoichiometric amount of Ti $(C_3H_7O)_4$  in  $C_2H_6O_2$  followed by slow addition of  $C_6H_8O_7H_2O$  in a beaker with stirring for 30 min. Citrate solution of the other metals was obtained by dissolving metal nitrates in deionized water followed by addition of citric acid. Clear solution was obtained by mixing both citrate solutions slowly while stirring at  $60^{\circ}$ C. To promote polymerization, the clear solution was kept overnight in an oven at  $120^{\circ}$ C. The solution became viscous with loss of excess solvent, and finally brown resin was obtained. To remove excess organic residue, the resin was heated at  $250^{\circ}$ C for 3 h. The resulting precursor was calcined at different temperatures

 $(500^{\circ}\text{C}, 600^{\circ}\text{C}, 700^{\circ}\text{C}, \text{ and } 800^{\circ}\text{C})$  to obtain  $(1 - x)M_{0.95}Zn_{0.05}T-x(Ca_{0.8}Sr_{0.2})_{0.6}La_{0.267}TiO_3$  composite ceramic powders, hereinafter referred to as  $(1 - x)$ MZT-xCSLT, where x is the molar fraction of CSLT. The calcined powders were ball-milled for 8 h prior to pressing into 11-mm-diameter  $\sim$ 5-mm-high pellets at 2000  $\text{kg/cm}^2$ . The pellets were sintered at  $1200\textdegree C$  at heating/cooling rate of  $3\textdegree C/min$ .

#### Characterization

A Netzsch STA 409PC/PG thermal analyzer was employed for thermal analysis of the polymeric precursor at air flow rate of 50 mL/min with heating rate of  $15^{\circ}$ C/min up to  $800^{\circ}$ C. The mass and dimension method was used to calculate the densities of the sintered pellets. For structural analysis of calcined powders and sintered pellets, a Philips (PW1830 APD system) x-ray diffractometer (XRD) was used with step size of  $0.05^{\circ}$  at 5 s/step in the range of  $2\theta = 10^{\circ}$  to  $80^{\circ}$ . Samples for XRD measurements were prepared by grinding sintered pellets with SiC papers (400 grade). The microstructure of calcined powders and sintered ceramics was examined using a field-emission scanning electron microscope (FE-SEM; Philips XL30). SEM samples were prepared by grinding pellets with various grades (400 to 1200) of SiC papers, and polished with diamond paste  $(6 \mu m)$  and an oxide polishing suspension, then etched at  $1100^{\circ}$ C for 1 h. A Renishaw microscope system (New Mills, Wotton-Under-Eagle, UK) with a 514-nm excitation laser diode, was employed to record the Raman spectra from sintered pellets at room temperature. The resonance cavity method employing an Agilent R3767CH vector network analyzer in the temperature range of 20°C to 75°C was used to measure the microwave dielectric properties of the sintered pellets.<sup>[18](#page-7-0)</sup>

#### RESULTS AND DISCUSSION

#### Thermogravimetry

Thermogravimetry/differential scanning calorime- $\mathrm{try}\,(\mathrm{TG}/\mathrm{DSC})\,\mathrm{curves\,of}\,(1-x) \mathrm{MZT\text{-}xCSLT\,polymeric}$ precursors are shown in Fig. [1.](#page-2-0) The first weight loss  $(\sim8\%)$  was observed from room temperature to  $\sim$ 250 $\degree$ C, accompanied by an associated small endotherm on the DSC curve at 108°C. This weight loss may be attributed to removal of adsorbed water from atmosphere, decomposition of  $\mathrm{NO_3^-}$  absorbed on the surface of the precursor, and burning of residual organics. A drastic weight loss  $(\sim 66\%)$  was observed in the temperature range of  $250^{\circ}$ C to  $550^{\circ}$ C, which may attributed to decomposition of citrate complex and nucleation of metal ions into metal oxides.<sup>[19](#page-7-0)</sup> The DSC profile for the second weight loss appeared split into three peaks  $(400.4^{\circ}C, 471.7^{\circ}C,$  and  $513.8^{\circ}C)$ , corresponding to combustion of citrate complex and phase formation. Thermal decomposition of citrate complex occurs as follows<sup>20</sup>:

<span id="page-2-0"></span>

$$
(Mg/Zn/Ca/Sr/La)Ti(C_6H_5O_7)_2
$$
  
\n
$$
\rightarrow (Mg/Zn/Ca/Sr/La)TiO(C_2O_4)_2
$$
  
\n
$$
\rightarrow (Mg/Zn/Ca/Sr/La)Ti(OH)_2C_2O_4
$$
  
\n(1)

The small weight loss  $(\sim 1.2\%)$  observed at temperatures up to  $700^{\circ}\mathrm{C}$  could be assigned to solid-state reaction of residual metal carbonates and  $TiO<sub>2</sub>$ associated with a very small endothermic peak on  $\mathrm{DSC}$  at 635°C.

#### Phase and Microstructure

 $\bold{XRD}$  patterns of the  $(1-x)\bold{MZT-xCSLT}$  samples calcined in air atmosphere at different temperatures of 500°C, 600°C, 700°C, and 800°C are shown in Fig. 2. No XRD peaks indicative of crystallization could be seen in the XRD pattern of the sample calcined at  $500^{\circ}$ C. The emergence of all the peaks matching the relevant PDF files for the  $(1 - x)$ MZT $x$ CSLT sample calcined at 600°C indicated crystallization of the final phases, as identified by the TGA and DSC profiles. The XRD pattern of the sample calcined at 600°C revealed formation of MZT as major phase along with minor CSLT phase and impurity  $MgTi<sub>2</sub>O<sub>5</sub>$  (MTO5) phase. When the calcination temperature was increased to 800°C, the observed increase in peak intensity for MZT and CSLT indicates an increase in the crystallinity of these phases, while the decrease in the peak intensity for MTO5 indicates a decrease in the amount of this phase. Figure 3 shows the XRD patterns of  $(1-x)MZT-xCSLT$   $(x = 0$  to 0.25) ceramics sintered at  $1200^{\circ}$ C. As expected, all the XRD peaks and relevant intensities recorded for the  $x = 0$  sample matched PDF# 00-006-0494 for MZT with rhombohedral  $(R3)$  symmetry as major phase. With an increase in  $x$  to 0.1, peaks due to CSLT with orthorhombic (Pbnm) symmetry (PDF # 00-022- 0153) emerged and persisted throughout the employed composition range. A single low-intensity



Fig. 2. XRD patterns of  $(1 - x)$ MZT-xCSLT powders calcined at different temperatures.



Fig. 3. XRD patterns of  $(1 - x)$ MZT-xCSLT ceramics sintered at 1200°C.

peak matching PDF# 00-035-0792 for the impurity phase MTO5 was observed in the XRD patterns for all samples except for  $x = 0$ .

Figure [4](#page-3-0) shows the Rietveld refinement results for the  $x = 0$  and 0.2 compositions. Refinement of these XRD data was carried out using MAUD software.<sup>[21](#page-7-0)</sup> These refinements produced good agreement factors:  $R_{wp} = 11.08$  to 8.69,  $R_p = 8.28$  to 6.40,  $R_{\rm exp}$  = 7.19 to 6.63, and goodness of fit of 1.49 to 1.23. The refined crystallographic data for The refined crystallographic data for  $(1-x)$ MZT- $x$ CSLT  $(x = 0, 0.1, 0.15, 0.2, 0.25)$ ceramics are listed in Table [I](#page-4-0). As is evident from Table [I](#page-4-0), there was no significant change in the lattice parameters with increasing CSLT content,

<span id="page-3-0"></span>

Fig. 4. Rietveld refinement of ceramics sintered at 1200°C: (a) MZT and (b) 0.8MZT-0.2CSLT.

demonstrating the formation of composite ceramics without any diffusion of CSLT into MZT. An appropriate increase in the XRD peak intensity of CSLT phase was observed, consistent with the increase in the amount of CSLT, from  $x = 0.1$  to 0.25. Furthermore, the slight shift of diffraction peak positions of MZT and CSLT towards lower angles than for  $MgTiO<sub>3</sub><sup>22</sup>$  $MgTiO<sub>3</sub><sup>22</sup>$  $MgTiO<sub>3</sub><sup>22</sup>$  and  $CaTiO<sub>3</sub><sup>23</sup>$  $CaTiO<sub>3</sub><sup>23</sup>$  $CaTiO<sub>3</sub><sup>23</sup>$  respectively, might be due to the larger ionic radius of the constituent  $\text{Zn}^{2}$  compared with Mg<sup>2+</sup> for the same coordination number (6). Similarly, the ionic radii of  $Sr^{+2}$  and La<sup>+3</sup> are larger than that of Ca<sup>+2</sup>.<sup>[24](#page-7-0)</sup>

Figure [5](#page-5-0)a and b shows FE-SEM images of  $(1 - x)$ MZT- $x$ CSLT $(x = 0$  and 0.2) powders calcined at 700°C. Close inspection of these images demonstrates agglomeration of relatively smaller  $(<100$  nm) particles. Figure [5c](#page-5-0) and d shows FE-SEM images of sintered  $(1 - x)$ MZT- $x$ CSLT ceramics. The microstructure of the single-phase MZT ceramic sample (i.e.,  $x = 0$  sintered composition) comprised densely packed grains of size ranging from 3  $\mu$ m to 5  $\mu$ m with no variation in contrast. On the other hand, the microstructure of the sintered composition with added CSLT comprised relatively smaller grains with further smaller grains at triple junctions, demonstrating the multiphasic nature of the sample.

#### Raman Spectroscopy

Raman spectra of the sintered  $(1-x)\mathrm{MZT\text{-}}x\mathrm{CSLT}$ samples are shown in Fig. [6.](#page-5-0) The ilmenite structure

of MgTiO<sub>3</sub> is known to exhibit  $5A_g + 5E_g$  Raman-active vibrational modes.<sup>[25](#page-7-0)</sup> The vibrations of Mg/Zn and Ti  $(A_g \text{ modes})$  along z-axis appeared at  $223.97$   $\text{cm}^{-1}$  and  $304.84$   $\text{cm}^{-1}$ . The other  $A_{\text{g}}$  modes associated with vibrations of O atoms were observed at 396.36  $\rm cm^{-1}$ , 496.31  $\rm cm^{-1}$ , and 714.54  $\rm cm^{-1}$ . The Raman mode observed at 279.3  $\text{cm}^{-1}$  may be due to the  $E<sub>g</sub>$  mode, an antisymmetric vibration of oxygen octahedron. The Raman modes observed at  $325.06 \text{ cm}^{-1}$  and  $350.60 \text{ cm}^{-1}$  may be due to twisting vibrations of oxygen octahedron with Mg/Zn and Ti atoms, while the mode observed at  $483.62 \text{ cm}^{-1}$ may be associated with antisymmetric vibration of Mg/Zn and Ti atoms with O octahedra, and the  $E_{\rm g}$ mode at 640.05  $\text{cm}^{-1}$  may arise from antisymmetric vibration of Ti–O stretching region.

There are 24 Raman-active modes for orthorhombic  $(Pbnm)$  symmetry, i.e.,  $7A_g + 5B_{1g} + 7B_{2g} +$  $5B_{3g}$ .  $26.27$  Three weak  $A_g$  Raman modes at  $161.09 \text{ cm}^{-1}$ ,  $205.9 \text{ cm}^{-1}$ , and  $256.06 \text{ cm}^{-1}$  related to bending vibration of O–Ti–O can be seen for the CSLT phase. It is possible that the other bands predicted for CSLT may be overlapped by other intense bands, or many modes have very low polarizability so as to prevent their being seen in the spectrum, as reported for  $CaTiO<sub>3</sub>$ .  $^{28,29}$  $^{28,29}$  $^{28,29}$ 

#### Density and Microwave Dielectric Properties

The mixing model reported in Ref. [30](#page-8-0) was used to calculate the theoretical density of the composite ceramic system comprising two phases, using Eq. 2:

$$
\frac{1}{\rho} = \frac{(1-x)}{\rho_1} + \frac{x}{\rho_2},\tag{2}
$$

where x is the molar fraction of CSLT, and  $\rho_1$  and  $\rho_2$ are the theoretical densities of MZT and CSLT phase, respectively. The values of the theoretical density for both phases were calculated from XRD data. The dielectric constant and quality factor are closely related to the bulk density, and both are known to increase with an increase in density. The apparent densities, dielectric constants, and quality factors of the 0.8MZT-0.2CSLT ceramics sintered at different temperatures are shown in Fig. [7](#page-5-0). The 0.8MZT-0.2CSLT composition sintered at 1200°C showed the highest relative density, i.e., 93.1% of theoretical density, and exhibited the optimum microwave dielectric properties:  $\varepsilon_r = 25.17$  and  $Q \times f = 58,754 \text{ GHz}$ . Figure [8a](#page-6-0) shows that the dielectric constant increased from 16.18 to 27.26 with increase in  $x$  from 0 to 0.25. The dielectric constant of ceramic composites can be predicted using different mixing models, given mathematically by Eqs.  $3-5$ :

Parallel mixing model $31$ :

$$
\varepsilon_{\mathbf{r}} = (1-x)\varepsilon_{\mathbf{r}_1} + y\varepsilon_{\mathbf{r}_2};\tag{3}
$$

Series mixing model $32$ :

<span id="page-4-0"></span>

<span id="page-5-0"></span>

Fig. 5. FE-SEM images of powder samples calcined at 700°C: (a) MZT and (b) 0.8MZT-0.2CSLT ceramics, and sintered at 1200°C: (c) MZT and (d) 0.8MZT-0.2CSLT.



Fig. 6. Room-temperature Raman spectra of  $(1 - x)$ MZT-xCSLT ceramics.

$$
\frac{1}{\varepsilon_{\rm r}} = \frac{(1-x)}{\varepsilon_{\rm r_1}} + \frac{x}{\varepsilon_{\rm r_2}}; \tag{4}
$$





Lichtenecker model $^{33}$  $^{33}$  $^{33}$ :

$$
\ln \varepsilon_{\rm r} = (1 - y) \ln \varepsilon_{\rm r_1} + y \ln \varepsilon_{\rm r_2}, \tag{5}
$$

where x is the molar fraction of CSLT phase, and  $\varepsilon_{\rm r1}$ and  $\varepsilon_{r2}$  are the dielectric constants of MZT and CSLT phase, respectively. Figure [8](#page-6-0)b shows a comparison of the calculated and experimental values of

<span id="page-6-0"></span>

Fig. 8. (a) Density, (b) dielectric constant, (c) quality factor, and (d) temperature coefficient of resonant frequency for  $(1-x)MZT-xCSLT$ ceramics sintered at 1200°C.





 $\varepsilon_{\rm r}$  for the  $(1-x){\rm MZT\text{-}xCSLT}$  composite ceramics for different  $x$  values. In the present case, it can be seen from Fig. 8a that the results obtained through the Lichtenecker theoretical model were comparatively closer to the experimental values than the other two models (Eqs. [4](#page-3-0) and [5](#page-5-0)). In microwave dielectrics having two components, the  $Q$  is predicted by Eq.  $6^{34}$ :

$$
\frac{1}{Q} = \frac{1-x}{Q_1} + \frac{x}{Q_2},
$$
\n(6)

where  $Q_1$  and  $Q_2$  are the quality factors of component phase 1 and 2, respectively. In the present study, the measured Q values were lower than the predicted ones. The deviation of the measured values from the predicted values may be due to

<span id="page-7-0"></span>the relatively lower density, porosity, and dissimilar microstructures of the constituent phases. Figure [8c](#page-6-0) shows the quality factor values of the  $(1 - x)$ MZTxCSLT ceramics as a function of x. The  $Q \times f$  values of the current samples decreased from 163,480 GHz to 46,465 GHz with increase in the CSLT content from  $x = 0.00$  to 0.25, due to the much lower  $Q \times f$ value of the CSLT phase. The dielectric loss at microwave frequencies is governed by the density, porosity, second phases, grain boundaries, and inclusions in real homogeneous ceramics.<sup>[35,36](#page-8-0)</sup> Moreover, the higher-dielectric-constant CSLT phase exhibits higher polarizability, resulting in higher intrinsic losses, which may be a possible reason for the observed substantial decrease in  $Q \times f$ .

Figure [8](#page-6-0)d shows the temperature coefficient of resonant frequency  $(\tau_{\text{f}})$  of the sintered  $(1-x)\text{MZT}$ xCSLT ceramics. There are several factors which affect  $\tau_f$ , such as the composition of the ceramics, the type and amount of additives, and second phase(s). The  $\tau_f$  of the composites was predicted using Eq.  $7^{37}$  $7^{37}$  $7^{37}$ :

$$
\tau_f = (1 - x)\tau_{f1} + x\tau_{f2}, \tag{7}
$$

where x is the molar fraction of CSLT phase, and  $\tau_{\text{f1}}$ and  $\tau_{f2}$  are the temperature coefficients of resonant frequency of the MZT and CSLT phase, respectively. The  $\tau_{\rm f}$  value was observed to change from  $-60$  ppm/  $\rm{^{\circ}C}$  to 7.2 ppm/ $\rm{^{\circ}C}$  as the molar ratio of CSLT was increased from 0 to 0.25. Figure [8d](#page-6-0) also shows that the measured values of  $\tau_f$  deviated slightly from the values predicted using Eq. 7. A near-zero  $\tau_f$  value  $(i.e., -5.8~ppm/°C)$  was achieved for the 0.8MZT-0.2CSLT ceramic sintered at 1200°C (Table [II\)](#page-6-0). Although the dielectric constant values of these  $(1-x)$ MZT- $x$ CSLT ceramics are comparable to those reported for  $(1 - y)Mg_{0.95}Co_{0.05}TiO_3 \cdot y(Ca_{0.8}Sr_{0.2})_{0.6}$  $\text{La}_{0.267}\text{TiO}_3$ , i.e.,  $(1-y)\text{MCoT-yCSLT}^{20}$ , the  $Q \times f$ values (163,480 GHz to 46,465 GHz for  $x = 0$  to 0.25) of the  $(1 - x)$ MZT- $x$ CSLT ceramics achieved in the present study are higher than those (138,766 GHz to 39,180 GHz at  $y = 0$  to 0.25) reported for  $(1 - y)MCoT-yCSLT<sup>20</sup>$  This observed improvement may be due to the fact that the ionic polarizability of Zn  $(2.04 \text{ A})$  is higher than that of Co  $(1.65 \text{ A})$ , resulting in stronger atomic vibrations causing a decrease in the dielectric loss.<sup>3</sup>

## **CONCLUSIONS**

 $(1 - x)$ MZT- $x$ CSLT powders with submicron particle size were successfully synthesized by a chemical synthesis route. XRD refinement revealed that, in the sintered composites, CSLT phase (with Pbnm symmetry) coexisted with MZT phase (with R3 symmetry) along with a small proportion of MTO5 impurity phase. As expected, the dielectric constant and  $\tau_f$  values increased while the  $Q \times f$  value decreased with increase in the amount of CSLT. However, the  $Q \times f$  value decreased with increasing CSLT amount. The 0.8MZT-0.2CSLT composition

sintered at  $1200^{\circ}$ C exhibited  $\varepsilon_r = 25.17$ ,  $Q \times f = 58,754 \text{ GHz, and } \tau_f = -5.8 \text{ ppm} / \text{°C}.$ 

## ACKNOWLEDGEMENTS

One of the authors acknowledges the laboratory support extended by Prof. Robert Freer, School of Materials, University of Manchester, Manchester M13 9PL, UK, for partial characterization of the materials. We acknowledge financial support extended by the Directorate of S&T, KP via ADP No. 130314 and Pilot Research Studies for the enhancement of research facilities in Materials Research Laboratory, and Mineral Upgradation Pilot Plant, University of Peshawar.

#### REFERENCES

- 1. F. Liu, X.-Y. Liu, C.-L. Yuan, J.-J. Qu, G.-H. Chen, C.-R. Zhou, and F. Liu, J. Eur. Ceram. Soc. 35, 2091 (2015).
- 2. X. Guo, B. Tang, J. Liu, H. Chen, and S. Zhang, J. Alloys Compd. 646, 512 (2015).
- 3. A. Ullah, Y. Iqbal, T. Mahmood, A. Mahmood, A. Naeem, and M. Hamayun, Ceram. Int. 41, 15089 (2015).
- 4. A. Naeem, A. Mahmood, Y. Iqbal, A. Ullah, T. Mahmood, and M. Humayun, J. Alloys Compd. 645, 290 (2015).
- 5. J. Qu, F. Liu, C. Yuan, X. Liu, and G. Chen, Mater. Sci. Eng. B 191, 15 (2015).
- H. Li, B. Tang, Y. Li, Z. Qing, and S. Zhang, *Mater. Lett.* 145, 30 (2015).
- 7. K. Wakino, Ferroelectrics 91, 69 (1989).
- 8. C.-L. Huang, Y.-B. Chen, and C.-F. Tasi, J. Alloys Compd. 460, 675 (2008).
- 9. L. Li, J. Ye, S. Zhang, Z. Gao, and S. Li, J. Alloys Compd. 648, 184 (2015).
- 10. R. Kell, A. Greenham, and G. Olds, J. Am. Ceram. Soc. 56, 352 (1973).
- 11. I.-S. Kim, W.-H. Jung, Y. Inaguma, T. Nakamura, and M. Itoh, Mater. Res. Bull. 30, 307 (1995).
- 12. P. Wise, I. Reaney, W. Lee, T. Price, D. Iddles, and D. Cannell, J. Eur. Ceram. Soc. 21, 1723 (2001).
- 13. A. Naeem, A. Ullah, Y. Iqbal, T. Mahmood, and A. Mahmood, J. Alloys Compd. 672, 298 (2016).
- 14. W.K. Leutwyler, S.L. Bürgi, and H. Burgl, Science 271, 933 (1996).
- 15. A.L. Rogach, D.V. Talapin, E.V. Shevchenko, A. Kornowski, M. Haase, and H. Weller, Adv. Funct. Mater. 12, 653 (2002).
- 16. H. Wu, Z. Feng, Q. Mei, J. Guo, F. Hou, P. Li, and X. Jiang, Ceram. Int. 41, 7645 (2015).
- 17. Y. Liu, Y. Pu, Z. Sun, and Q. Jin, Mater. Res. Bull. 70, 195 (2015).
- 18. T. Nishikawa, K. Wakino, H. Tamura, H. Tanaka, and Y. Ishikawa, 1987 IEEE MTT-S International Microwave Symposium Digest (1987), p. 277.
- 19. C. Mao, X. Dong, and T. Zeng, *Mater. Lett.* 61, 1633 (2007). 20. A. Ullah, A. Naeem, Y. Iqbal, T. Mahmood, and A.
- A. Ullah, A. Naeem, Y. Iqbal, T. Mahmood, and A. Mahmood, J. Mater. Sci. 27, 3506 (2016).
- 21. L. Lutterotti and C. Maud, Newsletter, (IUCr) No. 24 (December 2000).
- 22. M. Zhang, L. Li, W. Xia, and Q. Liao, J. Alloys Compd. 537, 76 (2012).
- 23. H. Kay and P. Bailey, Acta Crystallogr. 10, 219 (1957).
- 24. R.D. Shannon, Acta Crystallogr. Sect. A 32, 751 (1976).
- 25. C.-H. Wang, X.-P. Jing, W. Feng, and J. Lu, J. Appl. Phys. 104, 034112 (2008).
- 26. P. McMillan and N. Ross, *Phys. Chem. Miner.* 16, 21 (1988). 27. T. Hirata, K. Ishioka, and M. Kitajima, *J. Solid State* 27. T. Hirata, K. Ishioka, and M. Kitajima, J. Solid State
- Chem. 124, 353 (1996).
- 28. S. Keshri and S.S. Rajput, Ceram. Int. 40, 4257 (2014).
- 29. U. Balachandran and N. Eror, Solid State Commun. 44, 815 (1982).
- <span id="page-8-0"></span>30. H. Rocha, F.N.A. Freire, M.R.P. Santos, J.M. Sasaki, T. Cordaro, and A.S.B. Sombra, Phys. B 403, 2902 (2008).
- 31. P. Sarah and S. Suryanarayana, Bull. Mater. Sci. 26, 745 (2003).
- 32. D. Barrow, T. Petroff, R. Tandon, and M. Sayer, J. Appl. Phys. 81, 876 (1997).
- 33. Z. Hashin and S. Shtrikman, J. Appl. Phys. 33, 3125 (1962).
- 34. S.H. Yoon, G.-K. Choi, D.-W. Kim, S.-Y. Cho, and K.S. Hong, J. Eur. Ceram. Soc. 27, 3087 (2007).
- 35. D.A. Sagala and S. Nambu, J. Am. Ceram. Soc. 75, 2573 (1992).
- 36. B. Silverman, Phys. Rev. 125, 1921 (1962).
- 37. A. Paladino, J. Am. Ceram. Soc. 54, 168 (1971).
- 38. H.-J. Hagemann, J. Phys. C 11, 3333 (1978).