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# **Thermodynamic Analysis of a Modified Ejector-Expansion Refrigeration Cycle with Hot Vapor Bypass**

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**Abstract:** In this study, a modified ejector-expansion refrigeration cycle (MERC) is proposed for applications in small refrigeration units. A vapor bypass circuit is introduced into the standard ejector expansion refrigeration cycle (ERC) for increasing the ejector pressure lift ratio, thereby lowering the compressor pressure ratio in the MERC. A mathematical model has been established to evaluate the performances of MERC. Analysis results indicate that since a two phase vapor-liquid stream is used to drive the ejector in the MERC, a larger ejector pressure lift ratio can be achieved. Thus, the compressor pressure ratio decreases by 21.1% and the discharge temperature reduces from 93.6°C to 82.1°C at the evaporating temperature of -55°C when the vapor quality of two phase vapor-liquid stream increases from 0 to 0.2. In addition, the results show that the higher ejector component efficiencies are effective to reduce the compressor pressure ratio and the discharge temperature. Actually, the discharge temperature reduces from  $91.4^{\circ}$ C to  $82.1^{\circ}$ C with the ejector component efficiencies increasing from 0.75 to 0.85 at the two phase stream vapor quality of 0.2. Overall, the proposed cycle is found to be feasible in lower evaporating temperature cases.

#### **Keywords: compressor pressure ratio, discharge temperature, ejector, refrigeration cycle, vapor bypass**

## **1. Introduction**

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Recently, the use of two-phase ejector in single-stage vapor compression refrigeration and heat pump systems has drawn many attentions because the two-phase ejector can contribute to the system performance [1-4]. Principally, a two-phase ejector as expansion device applied in a single-stage vapor compression system can efficiently recover the loss of throttling process by generating isentropic expansion process to improve the system performance. Hence, a large number of theoretical and experimental studies have been conducted on the performance of the single-stage vapor

compression cycle with an ejector.

Lawrence et al. [5] proposed three different two-phase ejector refrigeration cycles, and the result showed that these three cycles show the improvement in *COP* compared to the conventional refrigeration cycle. Ersoy and Sag [6] carried out preliminary experiments on an R134a refrigeration system with a two-phase ejector, and found that this refrigeration system exhibits higher *COP* than that of the system without ejector by 6.2%-14.5%. Li et al. [7] evaluated performance of the R1234yf ejector expansion refrigeration cycle, and confirmed that the cycle is superior to the standard cycle, especially at extreme operating conditions. Hassanain et al. [8]

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conducted the ejector design and performance evaluation for an ejector expansion refrigeration system with refrigerant R134a. Also, Sag et al. [9] conducted an experimentally study on the ejector expansion refrigeration system in terms of energetic and exergetic aspects. This system applied R134a as refrigerant. The results showed that the total irreversibility of this ejector expansion refrigeration system is lower than that of the basic system. He et al. [10] carried out performance assessment of a transcritical  $CO<sub>2</sub>$  refrigeration system using an adjustable ejector, and revealed that the system has a better dynamic performance characteristic. To sum up, the ejector is an ideal expansion device, and it can be used in the vapor compression cycle system to improve the system performance.

Actually, the ejector is also a vapor compression device, i.e. the so-called thermo-compressor, which is widely used in heat-driven ejector refrigeration systems [11,12]. This means that an ejector can be used to mainly generate compression function similar to conventional compressors. Thus, ejectors could constitute an attractive alternative for conventional compressors in vapor compression refrigeration cycles. As well known, conventional single-stage compression cycles can only be operated in a limited temperature range. Their performances may degrade when they operate at low refrigeration temperatures, such as low coefficient of performance (*COP*), low cooling capacity, high discharge temperature and large pressure ratio of compressor. For this case, these drawbacks could be overcome by

utilizing two-stage compression cycles [13-15]. Typical two-stage vapor compression cycle usually uses two individual compressors or a compound compressor with vapor injection. Compared with the compressor, the ejector has no moving parts, lower cost, simple structure and lower maintenance requirements. This fact makes the use of ejector that is an alternative to the commonly used compressor in a two-stage compression cycle may have some technical advantages, particularly in terms of low cost and high reliability. However, this idea has not been founded in the opening literature. Obviously, the idea of using an ejector as an additional compressor could be of interest for new ejector applications in the single-stage vapor compression cycle, which is favorable for reliable operation in a wide range of refrigeration temperature similar to a two-stage vapor compression cycle. Hence, it is meaningful to propose a modified ejector-expansion refrigeration cycle from this idea, and the investigations on performances of this cycle are much needed.

In this study, a modified ejector-expansion refrigeration cycle (MERC) is presented based on the standard ejector expansion refrigeration cycle (ERC). In the MERC, a compressor discharge vapor bypass circuit is configured for bypassing the condenser. In this case, the two phase vapor-liquid stream can be used to drive the ejector in the MERC, thus allowing the ejector pressure lift ratio to be increased, resulting in decreased compressor pressure ratio. Thus, the MERC may employ a compressor to operate at wide temperature and pressure ranges which could be attractive for small refrigeration apparatus applications. A mathematical model has been established to simulate the performances of MERC. In the simulations, the main parameters are analyzed and discussed including the compressor pressure ratio and discharge temperature, the performances of ejector and the cooling capacity per unit volumetric displacement. The study mainly concentrates on a theoretical investigation of the performance of this proposed cycle for developing the ejector expansion refrigeration technologies application in small refrigeration units.

## **2. Cycle Description and Modeling**

The schematic diagram of the MERC system is shown in Fig. 1. The system components for MERC include a compressor, a bypass circuit with the flow regulating valve 1, a condenser circuit with the flow regulating valve 2, an ejector, a separator, a throttling device (expansion valve) and an evaporator. The schematic *P-h* diagram for the MERC is shown in Fig. 2. From Fig. 1, it is noticed that the MERC is very similar to the ERC. In the ERC, the ejector is mainly used as an expansion device to recover the energy loss in the throttling process, and its pressure lift is limited. However, the ejector in the MERC is applied to mainly provide the large compression effect in associate with expansion work recovery, resulting in the increase of the compressor suction pressure, for that a larger compression effect can be realized by the two phase flow driven ejector in the MERC. In this case, two flow regulating valves in both the bypass circuit and condenser circuit can be used to adjust relevant flow rates for obtaining the expected dryness quality of the mixed fluid (primary fluid of the ejector). The working process of MERC cycle is described as follows: a greater portion of the compressed refrigerant vapor from the compressor flows through the flow regulating valve 2 and enters the condenser. This refrigerant condenses into saturated or subcooled liquid at the condenser outlet by rejecting heat to the surroundings (process 2-2″-3′); the other portion of the vapor flows through the flow regulating valve 1 and mixes with the liquid refrigerant from the condenser (process 2-2′ and 2′, 3′-3); After that, this two-phase fluid flows into the nozzle of the ejector as the primary fluid to drive the ejector; Meanwhile, the secondary fluid, which is the saturated or superheated refrigerant vapor from the outlet of evaporator, is entrained to the ejector; In the ejector, the process 3-3″, (3″, 7)-4′ and 4′-4 are the expansion process of primary fluid, the mixing process between primary fluid and secondary fluid as well as compression process of mixed fluid, respectively; The two-phase refrigerant at the ejector outlet enters separator and separates into the saturated vapor (process 4-1) and the saturated liquid (process 4-5); At last, the saturated

vapor refrigerant flows into the compressor where it is continually compressed to the condensing pressure (process 1-2); The saturated liquid refrigerant flows through the throttling device and finally enters the evaporator where the working fluid becomes saturated or superheated (process 5-6-7).



**Fig. 1** The schematic diagram for the MERC system



**Fig. 2** The schematic *p-h* diagram for the MERC

The performances of MERC are theoretically assessed based on a mathematical model, and some basic assumptions made before establishing the model are as follows [16]:

(1) The steady-state and steady-flow process is carried out for all components;

(2) A variable isentropic efficiency is used to represent irreversible loss in the compressor;

(3) An isenthalpic process is used to represent throttling process in throttling device;

(4) The vapor and liquid discharged from the separator are saturated;

(5) The kinetic energies of the primary fluid and

secondary fluid inlets in ejector are neglected;

(6) The pressure drops and heat losses of refrigerant in the cycle are neglected;

(7) Ejector component efficiencies are introduced to represent friction losses in ejector.

On the basis of assumptions above, the mathematic model of MERC can be established in terms of the mass, momentum and energy conservations. Actually, the entrainment ratio  $\mu$  and the pressure lift ratio  $\gamma_{\rm eie}$  are important parameters, which can be used to evaluate the performance of ejector. The  $\mu$  and the  $\gamma_{\text{eie}}$  are defined as,

$$
\mu = \dot{m}_{\rm s}/\dot{m}_{\rm p} \tag{1}
$$

$$
\gamma_{\rm eje} = p_4/p_7 \tag{2}
$$

where  $\dot{m}_s$  is the mass flow rate of the ejector secondary fluid;  $\dot{m}_p$  is the mass flow rate of the ejector primary fluid;  $p_4$  is the ejector exit pressure;  $p_7$  is the pressure of the ejector secondary fluid inlet.

In the nozzle, the primary fluid is accelerated, which creates a low pressure zone at the nozzle exit plane. Thus, the secondary fluid can be entrained to the low pressure zone of the ejector. On the basis of the energy balance of the primary fluid between the inlet and outlet of the nozzle, the velocity of primary fluid at the nozzle outlet is defined as,

$$
w_{p2} = \sqrt{2\eta_n \left(h_{p1} - h_{p2s}\right)}
$$
 (3)

where  $h_{p1}$  is specific enthalpy of the primary fluid at the nozzle inlet; and  $h_{p2s}$  is the specific enthalpy of primary fluid at the nozzle outlet under an isentropic process;  $\eta_n$ is the nozzle efficiency, which is defined as,

$$
\eta_{n} = (h_{p2} - h_{p1})/(h_{p2s} - h_{p1})
$$
\n(4)

where  $h_{p2}$  is the practical specific enthalpy of primary fluid at the nozzle outlet

Keenan et al. [17] proposed two common mathematical models, the constant-area mixing (CAM) model and the constant-pressure mixing (CPM) model, to describe the mixing mechanism of fluids in the mixing chamber. The CPM is widely adopted in ejector design due to the superior performance in companion with the CAM [18-20]. And the CPM model is applied in the modeling i.e. the primary fluid mixes with the secondary fluid in the ejector at constant pressure. Then the momentum conservation equation for the mixing process of fluids in the mixing chamber can be obtained,

$$
\left(\dot{m}_{\rm s} + \dot{m}_{\rm p}\right) w_{\rm m2, is} = \dot{m}_{\rm p} w_{\rm p2} \tag{5}
$$

where  $w_{m2,s}$  is the ideal velocity of the mixed fluid at the mixing chamber outlet.

Furthermore, the mixing efficiency  $\eta_m$  is defined as,

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$$
\eta_{\rm m} = w_{\rm m2}^2 / w_{\rm m2, is}^2 \tag{6}
$$

where  $w_{m2}$  is the practical velocity of the mixed fluid at the mixing chamber outlet.

And then the  $w_{m2}$  can be obtained by combining Eq. (1) with Eq.  $(5)$  and Eq.  $(6)$ ,

$$
w_{\text{m2}} = w_{\text{p2}} \sqrt{\eta_{\text{m}}} / (1 + \mu) \tag{7}
$$

The energy conservation equation of the mixing process can be expressed,

$$
(m_s + m_p)(h_{m2} + w_{m2}^2/2) = m_s h_{s1} + m_p h_{p1}
$$
 (8)

where  $h_{s1}$  is the specific enthalpy of the secondary fluid inlet.

And then the specific enthalpy  $h_{m2}$  of the mixed fluid at the outlet of the mixing chamber is obtained as,

$$
h_{\text{m2}} = \frac{h_{\text{p1}} + \mu h_{\text{s1}}}{1 + \mu} - \frac{w_{\text{m2}}^2}{2} \tag{9}
$$

In the diffuser chamber, the kinetic energy of mixed fluid is converted into pressure energy. At the diffuser exit, the velocity is reduced and the pressure is raised enough to cause discharge. Applying the energy conservation to the compression progress of the diffuser, the velocity of mixed fluid at the ejector outlet can be obtained as,

$$
w_{d2} = \sqrt{w_{m2}^2 - 2(h_{d2} - h_{m2})}
$$
 (10)

where  $h_{d2}$  is the specific enthalpy of mixed fluid at the ejector outlet, which can be obtained by,

$$
h_{\rm d2} = h_{\rm m2} + \left(h_{\rm d2, is} - h_{\rm m2}\right) / \eta_{\rm d} \tag{11}
$$

where  $h_{d2, is}$  is the specific enthalpy of mixed fluid at the ejector outlet under an isentropic process for the same exit pressure;  $\eta_d$  is the diffuser efficiency.

As assumed that the kinetic energy of mixed fluid at the ejector out is neglected, the  $\mu$  can be obtained by combining Eq.  $(1)$  with Eqs.  $(3-11)$ ,

$$
\mu = \sqrt{\eta_n \eta_m \eta_d \frac{h_{\rm p1} - h_{\rm p2, is}}{h_{\rm d2, is} - h_{\rm m2}}} - 1
$$
 (12)

The allocation of mass flow rate  $\omega$  between the bypass pipe and condenser is written as,

$$
\dot{m}_{2'} = \omega \dot{m}_2, \dot{m}_{3'} = (1 - \omega)\dot{m}_2 \tag{13}
$$

Mass balances for the ejector, flash tanks and condenser can be written as,

$$
\dot{m}_3 + \dot{m}_7 = \dot{m}_4 = \dot{m}_1 + \dot{m}_5 \tag{14}
$$

$$
\dot{m}_2 = \dot{m}_{2'} + \dot{m}_{3'} = \dot{m}_3 \tag{15}
$$

Using Eq. (1) and Eqs. (13-15), the fundamental equation between the entrainment ratio and refrigerant quality at the ejector outlet can be expressed as,

$$
1 + \mu = 1/x_4 \tag{16}
$$

For the compressor, the input power can be written as

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$$
\dot{W} = \dot{m}_{\rm p} \left( h_2 - h_1 \right) = \frac{\dot{m}_{\rm p} \left( h_{2,\rm is} - h_1 \right)}{\eta_{\rm is}} \tag{17}
$$

$$
\dot{m}_{\rm p} = \frac{\eta_{\rm v} \dot{V}}{v_{\rm l}} \tag{18}
$$

where  $h_1$  is the refrigerant specific enthalpy at the compressor inlet;  $h_2$  is the refrigerant specific enthalpy at the compressor outlet;  $h_{2,is}$  is the refrigerant specific enthalpy at the compressor outlet under an isentropic process;  $\vec{V}$  is the volumetric displacement of the compressor;  $v_1$  is the refrigerant specific volume at the inlet of the compressor;  $\eta_{is}$  is the compressor isentropic efficiency;  $\eta_v$  is the compressor volumetric efficiency, which can be obtained [7,21],

$$
\eta_{\text{is}} = 0.874 - 0.0135 \frac{p_2}{p_1} = 0.874 - 0.0135 \gamma_{\text{com}} \quad (19)
$$

$$
\eta_{\rm v} = 0.959 - 0.00642 \frac{p_2}{p_1} \tag{20}
$$

where  $\gamma_{\text{com}}$  is the compressor pressure ratio,  $p_1$  is the refrigerant pressure at the compressor inlet;  $p_2$  is the refrigerant pressure at the compressor outlet.

The cycle cooling capacity is

$$
\dot{Q}_{c} = \dot{m}_{s} (h_{7} - h_{6}) = \frac{\mu \eta_{v} \dot{V} (h_{7} - h_{6})}{v_{1}}
$$
(21)

where  $h_6$  is the refrigerant specific enthalpy at evaporator inlet;  $h_7$  is the refrigerant specific enthalpy at the evaporatoroutlet. In addition, an actual cooling capacity per unit volumetric displacement  $q<sub>v</sub>$  can be obtained by Eq. (21),

$$
q_{\rm v} = \mu \eta_{\rm v} \left( h_7 - h_6 \right) / v_1 \tag{22}
$$

The cycle cooling coefficient of performance (*COP*) can be expressed by

$$
COP = \dot{Q}_{\rm c}/\dot{W} \tag{23}
$$

Finally, the performance of the MERC is simulated using the above model. The simulation code is written in Fortran. The refrigerant properties are calculated by the NIST database and subroutines [22]. The flow diagram of the calculation process is shown in Fig. 3.



**Fig. 3** The flow diagram of the calculation process for MERC

# **3. Results and Discussion**

The R290 has excellent thermodynamic properties, the zero ODP and the low GWP, thus it is widely used in refrigeration systems [23]. Therefore, R290 is selected as the working fluid to evaluate the performance of MERC in this research. The correlative conditions are given as follows: the condensing temperature  $t_c$  changes from 30 to 40 $^{\circ}$ C; the evaporating temperature  $t_e$  changes from -55 to -45°C [24]; the ejector component efficiencies  $\eta_{\text{n}}$ ,  $\eta_{\text{m}}$ and  $\eta_d$  are assumed to be varied from 0.75 to 0.85 based on the study of literatures [25-28]. It should be noted that although in the MERC the two streams flowing through the regulating valves and condenser undergo a pressure drop Δ*p* in practice, the effect of Δ*p* on MERC

**Table 1** The performances of MERC at the  $t_e$  of -55 $^{\circ}$ C,  $t_e$  of 40 $^{\circ}$ C and ejector nozzle isentropic efficiencies of 0.85

	$\Delta p = 0 \times 10^4$ Pa			$\Delta p = 1 \times 10^4$ Pa			$\Delta p = 2 \times 10^4$ Pa		
	$\gamma_{\rm com}$	$t_{\rm d}$ /°C	$q_v/J \cdot m^{-3}$	$\gamma_{\rm com}$	$t_{\rm d}/^{\circ}C$	$q_v/\text{J}\cdot\text{m}^{-3}$	$\gamma_{\rm com}$	$t_{\rm d}$ /°C	$q_v/\text{J}\cdot\text{m}^{-3}$
$x_3=0$	16.1	93.6	0.378	16.2	94.2	0.378	16.3	94.8	0.377
$x_3 = 0.05$	15.2	90.6	0.375	15.3	91.1	0.375	15.4	91.7	0.374
$x_3 = 0.1$	14.3	87.6	0.371	14.4	88.2	0.370	14.5	88.8	0.370
$x_3 = 0.15$	13.5	84.8	0.366	13.6	85.4	0.365	13.7	85.9	0.365
$x_3 = 0.2$	12.7	82.1	0.359	12.8	82.7	0.359	12.9	83.2	0.359

performances is relatively small at the conventional range of pressure drop, as shown in Table 1. Hence, the Δ*p* is ignored when effects of the  $x_3$  on the performance of MERC are studied in Section 3.1~3.3.

## **3.1 Effects of the** *x***3 on the performance of MERC at the different** *t***<sup>e</sup>**

The mass flow rates of the bypass circuit and the condenser can be varied by controlling the regulating valves, resulting in variable vapor quality  $x_3$  at the ejector nozzle inlet. Actually, the ejector performances (the pressure lift ratio  $\gamma_{\text{eie}}$  and the entrainment ratio  $\mu$ ) are varied with the  $x_3$ . Therefore, the compressor pressure ratio  $\gamma_{\text{com}}$  and the discharge temperature  $t_d$  will be affected by the  $x_3$  due to the interaction of the  $\gamma_{\text{com}}$  and the  $\gamma_{\text{eie}}$ . In this section, effects of the  $x_3$  on the performance of MERC are investigated at different *t*e. The  $t_c$  is assumed to be 40 $\degree$ C, the nozzle isentropic efficiency  $\eta_n$ , the mixing efficiency  $\eta_m$  and the diffuser isentropic efficiency  $\eta_d$  are fixed at 0.85, respectively.

Fig. 4 shows that the changes of the  $\gamma_{\text{com}}$  and  $t_d$  with  $x_3$ at different  $t_e$  for MERC. It can be seen that both the  $\gamma_{\text{com}}$ and  $t<sub>d</sub>$  are reduced with the increase of  $x<sub>3</sub>$ , and as the reduction in  $t_e$ , the drops of  $\gamma_{com}$  and  $t_d$  with  $x_3$  are raised. For instance, when the  $x_3$  increases from 0 to 0.2, the  $\gamma_{\text{com}}$ decreases by 18.0% at the  $t_e$  of -45°C, while the  $\gamma_{\text{com}}$ decreases by 21.1% at the  $t_e$  of -55<sup>o</sup>C. Furthermore, the  $t_d$ reduces from 93.6 to 82.1°C and from 76.9 to 70.6°C at the  $t_e$  of -55°C and -45°C when the  $x_3$  increase from 0 to 0.2. The reason is that the decrease of  $t<sub>e</sub>$  could result in a serious rise of *γ*com in a ERC, and the rise of *γ*com can impair the  $\eta_{\text{is}}$ , thus the  $t_{\text{d}}$  could be increased hugely. However, the  $\gamma_{\text{eje}}$  increases with the rise of  $x_3$  especially at a lower  $t_e$  in MERC, as shown in Fig. 5. Therefore, the  $\gamma_{\text{com}}$  reduces with the  $x_3$  and the reduction rate rises with the decrease of  $t_e$ . Hence, the  $t_d$  is dropped sharply with the rise of  $x_3$  especially at the lower  $t_e$ . The result indicates that a lower evaporating temperature can be achieved when the MERC is applied in a practical



**Fig. 4** The changes of the  $\gamma_{\text{com}}$  and  $t_d$  with  $x_3$  at different  $t_e$ 

refrigeration unit, since the MERC is able to reduce the compressor pressure ratio and discharge temperature remarkably.

Fig. 5 shows that the changes of the  $\gamma_{eje}$  and  $\mu$  with  $x_3$ at different  $t<sub>e</sub>$  for MERC. It can be seen in Fig. 5 that when the  $x_3$  increases from 0 to 0.2, the  $\gamma_{\text{eie}}$  increases by  $21.9\% \sim 26.7\%$  at the  $t_e$  in the range of -45~ -55°C. This case may be explained as that compared with the liquid stream, the change in specific enthalpy of the primary fluid through the nozzle is larger when the two-phase vapor-liquid stream is used as primary stream. And thus the velocity of primary fluid at the nozzle outlet is higher for this two-phase vapor-liquid stream from energy conservation. This results in an increase of the kinetic energy of the mixed stream in the ejector mixing chamber because of the momentum conservation for the mixing process. And then more pressure energy of the exit fluid would be obtained because the higher kinetic energy can be converted to pressure energy in the diffuser. Hence the higher pressure lifting ratio of ejector *γ*eje can be obtained when two-phase primary fluid is used. Furthermore, Fig. 5 shows that the entrainment ratio *μ* decreases by 22.7%  $\sim$  23.5% at the  $t_e$  in the range of -45 $\sim$ -55 $\degree$ C when the  $x_3$  increases from 0 to 0.3. The reason is that as the rises of the  $x_3$  and the  $\gamma_{\text{eje}}$ , the ejector outlet vapor quality  $x_4$  increases. Hence, the reduction in  $\mu$  is obtained because of the fundamental relation between the entrainment ratio and refrigerant quality at the ejector outlet, which is shown in Eq. (16).



**Fig. 5** The changes of the  $\gamma_{\text{eie}}$  and  $\mu$  with  $x_3$  at different  $t_e$ 

Fig. 6 shows that the changes of the  $q_v$  and  $\omega$  with  $x_3$ at different  $t<sub>e</sub>$  for MERC. It can be seen in Fig. 6 that the  $q_v$  decreases with the rise of  $x_3$  when the  $t_e$  is fixed. The reason is that the  $\mu$  is reduced with the increase of  $x_3$ , therefore the  $\dot{m}$ , through the evaporator decreases, thus the amount of  $q_v$  is reduced with the increase of  $x_3$ . However, the  $\gamma_{\text{eie}}$  rises with the increase of  $x_3$ , therefore the  $v_1$  decreases. Furthermore, the  $\eta_v$  increases due to the reduction in  $\gamma_{\text{com}}$ , which results in the increase of the  $\dot{m}_{\text{n}}$ . On the basis of these reasons, the effect of  $x_3$  on  $q_y$  is

limited in the given range of  $x_3$  actually. For instance, the  $q_v$  decreases by 5.0%  $\sim$  8.2% at the  $t_e$  in the range of -55 $\sim$  $-45^{\circ}$ C when the  $x_3$  increases from 0 to 0.2. Besides, it can be seen in Fig. 6 that the allocation of mass flow rate  $\omega$ increases with the rise of  $x_3$ , and it increases in a continuously accelerated speed as the  $t<sub>e</sub>$  goes up, i.e. when the  $t_e$  is -55°C, the  $\omega$  increases from 0 to 0.154; however, when the  $t_e$  is -45<sup>o</sup>C, the  $\omega$  increases from 0 to 0.164 with the  $x_3$  changing from 0 to 0.2. The result means that the mass flow rate of the refrigerant through the bypass circuit should be decreased at a lower  $t<sub>e</sub>$  to maintain the same  $x_3$ . Note, the *COP* decreases from 1.027 to 0.897, i.e. decreases by 12.7%, at the *t*e of -55°C when the  $x_3$  changes from 0 to 0.2. It means that the reduction in *COP* is the cost of obtaining a lower  $t_e$  and keeping an applicable  $t_d$ .



**Fig. 6** The changes of the  $q_v$  and  $\omega$  with  $x_3$  at different  $t_e$ 

## **3.2. Effects of the** *x***3 on the performance of MERC at the different** *t***<sup>c</sup>**

The changes of the performances of the ejector and the MERC with the  $x_3$  at different  $t_c$  will be discussed in this section. The  $t_e$  is assumed to be -55°C, the  $\eta_n$ ,  $\eta_m$  and  $\eta_d$ are fixed at 0.85, respectively. It can be seen in Fig. 7 that when the  $x_3$  changes from 0.1 to 0.2, the  $\gamma_{\text{com}}$  decreases by 19.8%  $\sim$  21.1% at the *t*<sub>c</sub> in the range of 30 $\sim$  40°C. Furthermore, when the  $t_c$  is 30, 35 and 40<sup>o</sup>C, the  $t_d$  is reduced to 70.2, 76.3 and 82.1 $\degree$ C with the  $x_3$  increasing to 0.2. The results show the  $\gamma_{\text{com}}$  and the  $t_d$  decrease obviously as the  $x_3$  rises, and decrease rates accelerate continuously as the  $t_c$  goes up. The reason is similar to Fig. 4, so that there is no more detailed description here. Fig. 8 shows that when the  $x_3$  changes from 0.1 to 0.2, the  $\gamma_{\text{eie}}$  increases by 24.7% ~ 26.7% at the  $t_c$  in the range of 30~ 40°C. The explanation on the variation of *γ*eje with the  $x_3$  is the same as what is analyzed in Fig. 5. Furthermore, the simulation results show the *γ*eje increases with the  $t_c$  at a constant of the  $x_3$ . This is due to the fact that the primary fluid in the motive nozzle with higher pressure energy would result in the mixed fluid with higher kinetic energy in the mixing chamber, so that

the higher pressure lifting ratio can be obtained at the ejector exit. In addition, it can be seen in Fig. 8 that when the  $x_3$  changes from 0.1 to 0.2, the  $\mu$  decreases from 0.568 to 0.438, 0.540 to 0.415 and 0.513 to 0.392 for the  $t_c$  of 30, 35 and 40°C, respectively. It means that the  $\dot{m}_s$ through the evaporator decreases with the rise of  $x<sub>3</sub>$  at a constant  $\dot{m}_p$ , which can lead to the change of the  $q_v$ . It can be seen in Fig. 9 that when the  $x_3$  changes from 0.1 to 0.2, the  $q_v$  is decreased by 5.8%  $\sim$  5.0% at the  $t_c$  in the range of 30~40 $^{\circ}$ C. As same as the Fig. 6, the effect of  $x_3$ on  $q_v$  is limited at different  $t_c$  for MERC. Furthermore, Fig. 9 shows when the  $x_3$  changes from 0.1 to 0.2, the  $\omega$ 



**Fig. 7** The changes of the  $\gamma_{\text{com}}$  and  $t_d$  with  $x_3$  at different  $t_c$ 



**Fig. 8** The changes of the  $\gamma_{\text{eie}}$  and  $\mu$  with  $x_3$  at different  $t_c$ 



**Fig. 9** The changes of the  $q_v$  and  $\omega$  with  $x_3$  at different  $t_c$ 

increases from 0 to 0.160, 0 to 0.157 and 0 to 0.154 for the  $t_c$  of 30, 35 and 40 $\degree$ C, respectively. It means that the refrigerant vapor through the bypass circuit should be decreased to maintain the same  $x_3$  when the MERC systems are applied in the higher ambient temperatures.

# **3.3 Effects of the** *x***3 on the performance of MERC at the different**  $\eta_n$ ,  $\eta_m$  and  $\eta_d$

The working fluid flow in the ejector is very complicated, which causes large amount of the irreversibility losses, decreasing the ejector performance, and thus the irreversibility losses of the working process in ejector are usually considered by introducing the ejector component efficiencies [29]. In this section, three different ejector component efficiencies (0.75, 0.80 and 0.85) have been selected from ranges of  $0.7~1.0$  [25-28] to discuss the variations of the cycle performance and the ejector performance. The  $t_e$  and the  $t_c$  is assumed to be -55 and 40°C, respectively.

Fig. 10 shows the changes of the  $\gamma_{\text{com}}$  and  $t_d$  with  $x_3$  at different ejector component efficiencies for MERC. It can be seen in Fig. 10 that when the  $x_3$  changes from 0.1 to 0.2, the *γ*<sub>com</sub> decreases by 15.7%, 18.3% and 21.1% for the ejector component efficiencies of 0.75, 0.85 and 0.95, respectively. The result is related to the fact that high ejector component efficiencies can result in the low irreversible loss in the nozzle, the mixing chamber and the diffuser, thus a higher pressure of refrigerant at the ejector outlet could be obtained. It results in the increase of the *γ*eje , which can be shown in Fig. 11, therefore the  $\gamma_{\text{com}}$  decreases. Furthermore, when  $x_3$  is 0.2, the  $t_d$  is reduced to 91.4, 86.7 and 82.1°C at the ejector component efficiencies of 0.75, 0.80 and 0.85. In a short conclusion, lifting the ejector component efficiencies is effective to reduce  $\gamma_{\text{com}}$  and  $t_d$ . Fig. 11 shows that when the  $x_3$  changes from 0.1 to 0.2, the  $\gamma_{\text{eie}}$  increases by 18.6%, 22.4% and 26.7% for ejector component efficiencies of 0.75, 0.80 and 0.85, respectively, i.e. *γ*eje is more sensitive to the variation of the  $x_3$  when ejector component efficiencies are relatively high. In addition, the *μ* decreases with  $x_3$ , i.e. when  $x_3$  changes from 0.1 to 0.2 the *μ* decreases from 0.497 to 0.370, 0.504 to 0.380 and 0.513 to 0.392 for ejector component efficiencies of 0.75, 0.80 and 0.85, respectively.

It can be seen in Fig. 12 that when the  $x_3$  changes from 0.1 to 0.2, the *q*v is decreased by 12.6%, 9.0% and 5.0% at the ejector component efficiencies of 0.75, 0.80 and 0.85. The result is related to the fact that the increase of ejector component efficiencies can observably reduce the irreversible loss in the ejector, thus the high  $\mu$  could be obtained. Therefore, the reduction rate in  $q_v$  due to the deceasing  $\mu$  can be decreased with the ejector component efficiencies increase. Furthermore, Fig. 12 shows when the  $x_3$  changes from 0.1 to 0.2, the  $\omega$  increases from 0 to

0.147, 0 to 0.150 and 0 to 0.154 for the ejector component efficiencies of 0.75, 0.80 and 0.85, respectively. It means that the allocation of mass flow rate between the bypass pipe and condenser needs to increase to maintain the same  $x_3$  at the higher ejector component efficiencies.



**Fig. 10** The changes of the  $\gamma_{\text{com}}$  and  $t_d$  with  $x_3$  at different  $\eta_n$ ,  $\eta_{\rm m}$  and  $\eta_{\rm d}$ 



**Fig. 11** The changes of the  $\gamma_{\text{eie}}$  and  $\mu$  with  $x_3$  at different  $\eta_n$ ,  $\eta_m$ and  $\eta_d$ 



**Fig. 12** The changes of the  $q_v$  and  $\omega$  with  $x_3$  at different  $\eta_n$ ,  $\eta_m$ and  $n_d$ 

#### **4. Conclusions**

A modified ejector-expansion refrigeration cycle (MERC) is presented in this paper. In the MERC, a bypass circuit with a flow regulating valve to regulate the  $x_3$  is integrated in the ERC for increasing the  $\gamma_{\text{eie}}$  and lowering the *γ*<sub>com</sub>. Furthermore, the theoretical investigation on the MERC performances (the  $\gamma_{\text{com}}$ , the  $t_{\text{d}}$ and the  $q_v$  etc.) is conducted. The main conclusions include:

1) The  $\gamma_{\text{com}}$  and  $t_d$  can be decreased when the  $x_3$  is raised, i.e. the  $\gamma_{\text{com}}$  decreases by  $18.0\% \sim 21.1\%$  with the  $x_3$  increasing from 0 to 0.2 at the  $t_e$  in the range of -45~-55°C. The  $t<sub>d</sub>$  reduces from 93.6 to 82.1°C at the  $t<sub>e</sub>$  of  $-55^{\circ}$ C when the  $x_3$  increases from 0 to 0.2.

2) A small reduction in  $q_v$  is conducted with the rise of  $x_3$ , for instance, the  $q_v$  decreases by 5.0%  $\sim$  8.2% at the  $t_e$ in the range of -55~-45°C when the  $x_3$  increases from 0 to 0.2.

3) The higher ejector component efficiencies are effective to reduce  $\gamma_{\text{com}}$  and  $t_{d}$ , for instance, when the  $x_3$ changes from 0.1 to 0.2, the  $\gamma_{\text{com}}$  decreases by 15.7%, 18.3% and 21.1% for the ejector component efficiencies of 0.75, 0.80 and 0.85, respectively.

In general, the key point of this article is lowering the  $\gamma_{\text{com}}$  and  $t_d$  by the MERC, aiming to obtain a lower  $t_e$ . The proposed MERC could broaden the ejector expansion refrigeration technologies, which could work in a wider temperature range than the ERC system. Certainly, further experimental studies on the MERC system operating performance are required in the next step to confirm the practicability of this cycle system.

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