REVIEW

Chemical and physical characterizations of the $n = 1$ Ruddlesden–Popper phases: $Nd_2 = \sqrt{5}r_{\nu}Ni_1 = \sqrt{2}C_2O_4 + \delta$ $(y = 1 \text{ and } 0.1 \leq x \leq 0.9)$

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Abstract The structural stability and physical properties of NdSrNi₁ – $_{x}Co_{x}O_{4} \pm \delta$ (0.1 ≤ x ≤ 0.9) mixed oxides, elaborated by conventional sol–gel process, have been investigated and obtained results show that substitution of nickel by cobalt at $x = 0.5$ enhances conductivity at room temperature; $\sigma = 17.24 \Omega^{-1}$ cm⁻¹ coinciding with minimum activation energy ($E_a = 0.05$ eV). Rietveld refinements of X-ray powder diffraction patterns at room temperature indicate that all compositions crystallize in a tetragonal system with I4/mmm space group and exhibit K2NiF4-type structure. Variations of a and c parameters display various behavior with increasing cobalt content. Changes in cell parameters are discussed in terms of crystal field theory. In addition, transition metal oxidation state is investigated on the basis of the Brown bond valence calculation. The deduced Global Instability Index (GII) value decreases when cobalt substance increases, indicating that the structure becomes more stable once cobalt is introduced. Oxygen stoichiometry of these compounds was determined from thermogravimetric analyses (TGA) followed by reduction in 5% H_2 in N₂ gas. Conductivity of NdSrNi₁ – $_xCo_xO_4 \pm \delta$ (0.1 ≤ x ≤ 0.9) oxides was measured by an ac four-probe method. Oxygen vacancies are the possible ionic charge carriers.

Specimens exhibit a semiconducting behavior in the whole range of temperature. The electrical transport mechanism agrees with an adiabatic small polaron hopping (ASPH) model.

Keywords Functional oxides . Sol–gel preparation . Rietveld analysis K_2N iF₄-type structure \cdot Thermogravimetric analysis . Electrical properties

Introduction

The Ruddlesden–Popper series displaying the K_2NiF_4 type structure are the central focus of researchers especially chemists and physicists as they present a wide range of chemical and physical properties. Indeed, they permeate diverse fields as superconductivity [\[1](#page-10-0)], magneto resistance [\[2,](#page-10-0) [3](#page-10-0)], catalysis [[4\]](#page-10-0), and mixed ionic–electronic conductivity [[5](#page-10-0)].

In particular, several authors have been attracted by the possible application of Ln_2MO_4 compounds (Ln = rare earth; $M = Co$, Ni) such as oxygen membrane or electrode in solid oxide fuel cells (SOFCs) in addition to solid oxide electrolysis cell (SOEC) devices [[6](#page-10-0)–[8](#page-10-0)]. Furthermore, doping these compounds with alkaline earth and transition metal can modulate not only the oxygen content in the system but also the average oxidation state of the transition metal ion [[9](#page-10-0), [10](#page-10-0)].

The majority of investigations are devoted to lanthanum compounds, only few were interested in neodymium rare earth. In fact, Moritomo et al. had explored La_{2-x} Sr_xCoO₄ (0.4 ≤ x ≤ 1.0) solid solution with a mixed valence of cobalt ion (Co^{2+}/Co^{3+}) and found steep decreases in electrical resistivity with an increase in the amount of Sr beyond the value 0.7 [\[11](#page-10-0), [12](#page-10-0)]. Next to these

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Table 1 Lattice parameters, volume, and tolerance factor for all compositions

Composition (x)	$a = b$ (Å)	$c(\AA)$	Volume (\AA^3)	Tolerance factor
0.0	3.7963(1)	12.2971(2)	177.224(4)	0.951
0.1	3.7875(2)	12.3456(4)	177.099 (2)	0.945
0.3	3.7834(9)	12.3528(6)	176.819(4)	0.937
0.5	3.7824(3)	12.3315(2)	176.421(2)	0.928
0.7	3.7804(4)	12.2798(8)	175.495(3)	0.920
0.9	3.7799(1)	12.2652(5)	175.240 (1)	0.912

attempts, Kharton et al. reported that the conductivity of layered La₂Ni₁ – _xCo_xO_{4 + δ} is predominantly p-type electronic within the stability domain $(x = 0.1 - 0.2)$ phase and occurs via small polaron mechanism, which is indicated by temperature-activated hole mobility and $p(O_2)$ dependencies of electrical properties [\[13\]](#page-10-0).

Over the last decade, great interest has been paid to neodymium oxides $Nd_{2-x}Sr_{x}CoO_{4}$ with the K₂NiF₄-type structure for the discovery of magneto resistance and superconductivity [[14](#page-10-0), [15\]](#page-10-0). D. Grandjean et al. [\[16](#page-10-0)] have studied the effect of cobalt substitution for copper in the parent compound NdSrCoO4, and highlighted that this doping produces quite dissimilar oxide coordination environments for the dopant and copper ion which may change spin state.

A previous study performed by H. Chaker et al. [\[17\]](#page-10-0) focused on the effect of copper substitution in $NdSrNiO₄$ which showed a metal-semiconductor transition at ~190 K [\[18\]](#page-11-0). Structural and conductivity results for the composition $NdSrNi_{0.8}Cu_{0.2}O₄$ reveal that this phase is a possible candidate for integrated devices like, for example, bottom electrode in parallel plate ferroelectric capacitors. At the second stage of this study, we used this electrode to epitaxial growth for a functional oxide, namely the Aurivillius phase ferroelectric [[19](#page-11-0)].

As we got more and more interested in our lab in mixed oxides based on neodymium, erbium, and dysprosium rare earths, we have already reported in previous works the obtained results after substitution of nickel for copper and chromium [\[20](#page-11-0)–[26\]](#page-11-0).

Within this framework, our present work is an extension to our previous studies on neodymium nickelatebased mixed oxide. It aims to investigate the effect of cobalt incorporation in the parent compound NdSrNiO4 on the structural features and the ability to accommodate oxygen non-stoichiometry. We suspect that these new functional K_2N i F_4 -type oxides should form a very interesting family of materials: their physical properties and the possibility to combine them with a ferroelectric oxide in heterostructures by epitaxy open a way for fabricating innovating and high-performance components for applications in microelectronic and as electrode in SOFCs operating at intermediate temperatures.

For this reason, we have synthesized new compounds in the NdSrNi₁ – $_xCo_xO_4 \pm \delta$ (0.1 ≤ x ≤ 0.9) system by sol–gel method. We report here some preliminary results of powder X-ray diffraction, thermogravimetric analysis, bond valence calculation, and resistivity measurements.

Fig. 2 Evolution of unit cell volume with x in the NdSrNi_{1 − x}Co_xO_{4 ± δ} $(0.0 \le x \le 0.9)$ solid solution

Materials and methods

Polycrystalline samples with particular compositions $(x = 0.1, 0.3, 0.5, 0.7, 0.7)$ and 0.9) have been synthesized in the solid solutions NdSrNi₁ – $_{x}Co_{x}O_{4}$ $_{\pm 0}$ from powders $Nd₂O₃$ (Aldrich, 99.99%), SrCO₃, NiO, and Co₃O₄ (Aldrich, 99.99%) by the sol–gel method in the desired mole ratio in the reaction:

$$
1/2Nd_2O_3 + SrCO_3 + (1-X)NiO + x
$$

× 1/3Co_3O_4 → NdSrNi_{1-x}Co_xO_{4±δ}+CO₂

Neodimium oxide, strontium carbonate, nickel, and cobalt oxides were initially independently dissolved in a minimum quantity of concentrated chlorid acid with magnetic stirring. After intimate mixture of these solutions with distilled water, citric acid $C_6H_8O_7 \cdot H_2O$ and ethylene glycol $C_2H_6O_2$ were then added. The obtained sol was left on a hot plate with constant stirring for approximately 5 h at 80 °C until it gelled and changed the color from transparent purple to green. After that, the formed pale green gel was treated by further heating, first at 200 °C for approximately 24 h then at 450 $^{\circ}$ C in a muffle furnace. The resulting powder was ground before it was returned to the furnace in air at 600 °C for 48 h. The calcined mixtures were subsequently pressed into pellets (13-mm diameter and 2-mm thickness under 10 $t/cm²$) and annealed in a tube furnace at 1150 °C in a flow of oxygen and cooled to room temperature at a rate of 100°/h for 3 days with intermediate regrinding and repelleting until no further reaction was established by powder X-ray diffraction.

Powder X-ray diffraction data for the six samples $(x = 0.1, 0.3, 0.5, 0.7, 0.09)$ were collected at room temperature using a Bruker D8 Advance, with a monochromatic CuK α_1 radiation ($\lambda = 1.54056$ Å). Data were

Table 2 The refined composition of samples from EDX analysis

\boldsymbol{x}	Element	Atomic	Refined	Nominal
composition		%	composition	composition
0.1	\mathcal{O} Co Ni	59.46 1.42 12.53	4.163 0.099 0.877	$\overline{4}$ 0.1 0.9
	Sr Nd	14.03 12.56	0.982 0.879	$\mathbf{1}$ $\mathbf{1}$
0.3	Ω	58.70	4.109	$\overline{4}$
	Co	4.42	0.309	0.3
	Ni	9.71	0.679	0.7
	Sr	13.72	0.960	$\mathbf{1}$
	Nd	13.45	0.941	$\mathbf{1}$
0.5	Ω	57.85	4.049	$\overline{4}$
	Co	6.71	0.469	0.5
	Ni	6.85	0.479	0.5
	Sr	14.1	0.987	1
	Nd	14.49	1.014	$\mathbf{1}$
0.7	Ω	58.42	4.089	$\overline{4}$
	Co	9.85	0.689	0.7
	Ni	4.14	0.289	0.3
	Sr	14.57	1.019	1
	Nd	13.02	0.911	$\mathbf{1}$
0.9	Ω	60.28	4.219	$\overline{4}$
	Co	11.71	0.819	0.9
	Ni	1.28	0.089	0.1
	Sr	14.14	0.989	$\mathbf{1}$
	Nd	12.59	0.881	1

collected with a 0.0105° step 2 θ width and \approx 2-s counting time per point over a 2θ range from 10 to 120°. The whole pattern profile refinements were carried out with the FullProf program [[27](#page-11-0)].

The morphological features and grain size of the particles were studied by scanning electron microscopy (SEM) on a JSM-6400 apparatus working at 20 kV at room temperature. For this reason, the samples were ground, deposited, and then metalized. Compositional analyses were performed by recording the energy dispersive X-ray analysis (EDX).

The total oxygen content of the samples was determined by hydrogen reduction in a TGDTA 92- SETARAM thermogravimetric analyzer (TGA). The material was heated in an alumina pan under flowing dried 5% H_2 in N₂ gas (1.5 l/h) from room temperature up to 1100 °C at a rate of 10 °C/min.

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The temperature dependence of dc resistivity $\rho(T)$ was measured by the conventional four-probe method under high vacuum conditions using the van der Pauw method [\[28\]](#page-11-0). The latter was implemented through cooling down the samples from 390 to 20 K with a rate of 1 K/min using the Cryodyne Industries Refrigerator, model 22.

Results and discussion

Structural study

Departing from the superposition of the patterns of all prepared compositions in the NdSrNi₁ – $_{x}Co_{x}O_{4}$ $_{\pm}$ $_{\delta}$ $(0.1 \le x \le 0.9)$ system, it is deduced that the corresponding compounds are structurally identical and no evidence for a refinement for $NdSrNi_{0.5}Co_{0.5}O_{4 - δ}$

structural phase transition was detected when varying the composition of the cobalt. In fact, auto-indexing of the whole pattern profile of the X-ray peaks for all compositions led to only one solution with all lines indexed and a good figure of merit. Profile analysis has been performed on the basis of a tetragonal system with space group I4/mmm. The peak shape was described by a pseudo-Voigt function (Thomson–Cox– Hasting (TCH)).

For each diffraction pattern, a zero-point shift and unit cell parameters were refined in addition to profile

Fig. 4 Observed (dotted) and calculated (solid) lines of the final Rietveld refinement. The lower line is relative to the difference pattern $(Y_{obs} - Y_{cal})$, whereas vertical tics correspond to the 2θ Bragg positions

parameters. The background was fitted with a linear interpolation between 24 chosen points. The lattice parameters deduced from the whole pattern profile refinement ("pattern matching" mode in FullProf) of the X-ray data are presented in Table [1](#page-1-0).

Figure [1](#page-1-0) portrays the variation of the lattice parameters as a function of x in NdSrNi₁ – $_{x}Co_{x}O_{4\pm\delta}$ (0.1 ≤ x ≤ 0.9) compounds (results for $x = 0$ have been reported from a preceding research [\[17\]](#page-10-0)). As it can be observed in this figure, the c parameter, at the first time, increases clearly with doping cobalt but subsequently,

NdSrNi0.5Co0.5O4

Fig. 5 Tetragonal K_2N i F_4 -type structure of NdSrNi₁ – $_{x}Co_{x}O_{4}$ $_{\pm \delta}$ $(0.0 \le x \le 0.9)$ compounds

it decreases remarkably. On the other side, a parameter decreases slightly then remains unchangeable.

To explain these evolutions, we need to take into account the Jahn–Teller effect in $(Ni/Co)O_6$ octahedron. In fact, the fractional substitution of Ni³⁺ adopting $t_{2g}^6 d_{x^2-y^2}^1 d_{z^2}^0$ configu-ration in the NdSrNiO₄ parent compound [[17](#page-10-0)] by Co^{2+} , which usually exhibits $t_{2g}^6 d_{z^2}^1 d_{x^2-y^2}^0$ configuration in low spin state as declared in tetragonal symmetry, causes an increase only of c parameter owing to the increase in electron quantity in the \vec{c} direction. Moreover, the slight decrease in a parameter is due to the deficiency of one electron in \vec{a} and \vec{b} directions compared to Ni³⁺ (change in e_g^1 configuration: d_{z^2} orbital is occupied instead of $d_{x^2-y^2}$).

Furthermore, for a higher percentage of cobalt than nickel, parameter c decreases dramatically. This results from the fact that Co^{2+} (0.65 Å) oxidizes to smaller Co^{3+} (0.545 Å) ion for low spin state in sixfold coordination [\[29\]](#page-11-0). This finding is in

Nd/Sr—O (2) Å \times 4 2.6991 (1)

good agreement with the statement of (+III) oxidation for nickel and the mixed valence (+II/+III) for cobalt. The unit cell volume, shown in Fig. [2,](#page-2-0) decreases with x in the series NdSrNi₁ – $_{x}Co_{x}O_{4} \neq \delta$ (0.0 ≤ x ≤ 0.9), but the reduction rate proves to be significant for $x \ge 0.5$.

In order to ascertain the cobalt content in the prepared materials, semiquantitative analysis using EDX was carried

Fig. 6 (Nd/Sr) and (Ni/Co) environments for NdSrNi_{0.5}Co_{0.5}O_{4 – δ}

BVS bond valence sum, GII Global Instability Index

out. These measurements prove the presence of all the chemical elements with their stoichiometric proportions. The results of the analysis are largely summarized in Table [2.](#page-2-0)

A general tendency was observed through SEM. In fact, the particle size extended and the porosity, as a matter of fact, diminished when we have progressively introduced cobalt instead of nickel. Consequently, the powder became more and more dense and crystallized with homogenous distribution. Figure [3](#page-3-0) displays SEM micrographs of the six samples with the same scale.

Identifying tolerance factor (t) for layered oxide having K₂NiF₄ structure is established by $t = \frac{A-O}{\sqrt{2}(B-1)}$ $\frac{A-O}{\sqrt{2}(B-O)}$, where A–O and B–O are the metal–oxygen bond lengths of A-site and B-site cation, respectively. Referring to Goldschmidt [\[30](#page-11-0)], K₂NiF₄-type structure can be formed when $(0.85 < t < 1.05)$. On this basis, the tolerance factor was calculated using tabulated values of ionic radii [\[29\]](#page-11-0). These are presented in Table [1](#page-1-0) for NdSrNi_{1 – x}Co_xO_{4 ± δ} (0.0 ≤ x ≤ 0.9) phases. Within this context, the tolerance factor values of these materials reside in the range $0.912 \le t \le 0.945$ for which the tetragonal distortion is favored [\[31](#page-11-0)].

The structural refinement of each composition in NdSrNi_{1 – x}Co_xO_{4 ± δ} family was performed in the tetragonal system (I4/mmm space group) beginning with atomic posi-tions taken from NdSrNiO₄ [[32](#page-11-0)], which involves Nd³⁺ and $Sr²⁺$ ions that were disorganized over the ninefold coordinate 4e sites with their ratio according to the nominal composition (Table [2](#page-2-0)). The Ni and Co ions are situated at the 2a site, and the O1, O2 atoms are placed at $(0, \frac{1}{2}, 0)$ in site 4c and at $(0, 0, 0)$ z) in site 4e, respectively. At the final stage, good reliability factors are obtained as a result to refinement of all parameters, i.e., atomic coordinates, global atomic thermal displacement, and preferential orientation along the [001]. To illustrate these notions, the corresponding refined parameters of the NdSrNi_{0.5}Co_{0.5}O₄ − δ compound are presented in Table [3](#page-4-0) and the final fit is displayed in Fig. [4](#page-4-0). Projection of the NdSrNi_{0.5}Co_{0.5}O₄ – δ structure is portrayed in Fig. [5,](#page-5-0) and some corresponding inter-atomic distances are given in Table [4.](#page-5-0) Consequently, Ni/Co atoms shape a distorted octahedral where the axial (Ni/Co)–O1 band in the ab plane is shorter than the apical (Ni/Co)–O2 band along the c direction (Fig. [6\)](#page-5-0). The coordination environments for strontium cation are highlighted also in Fig. [6](#page-5-0).

Under this situation, metal–oxygen bonds are strained and the extension of these strains can be deduced using BVM (bond valence method) expressed by $S_{ij} = \exp[(R_0 - R_{ij}/B)]$, where R_0 and B (constant = 0.37) are the experimentally given parameters and R_{ij} is the bond length of the cation–anion pair. Referring to Brown [[33](#page-11-0)], the sum of the bond valence around an ion has to be equal to the formal valence V_i of these ions

 $\left(V_i = \sum_{i} S_{ij}\right)$. It is called VSR (valence sum rule). The deij

viation observed between these two values can be attributed to instabilities of the structure. The root mean square of

Fig. 7 The relative weight loss curve for NdSrNi_{0.9}Co_{0.1}O_{4 − δ} and NdSrNi $_{0.7}Co_{0.3}O_4 - \delta$

Fig. 8 The relative weight loss curve for $NdSrNi_{0.3}Co_{0.7}O_{4-δ}$ and $NdSrNi_{0.1}Co_{0.9}O_4 - \delta$

deviations for all atoms is named global instability index and

expressed as $[34]$ $[34]$ $[34]$ GII = $\int (1-\lambda^2)$ ∑ N $\sum_{i=1}^N \left\{ \sum_i$ ∑ S_{ij} − V_i $\sqrt{\sum\limits_{i=1}^N\left\{\left(\sum\limits_{\mathbf{j}}S_{ij}\hspace{-0.2mm}-\hspace{-0.2mm}V_i\right)^2\right\}}$ N where N is the number of atoms in the asymmetric unit. Sánchez-Andújar M. and Señarís-Rodríguez MA have investigated the Global Instability Index (GII) of K_2N iF₄-type LnSrCoO₄ (Ln = La, Nd, and Gd) [[35\]](#page-11-0) and demonstrated that GII decreases with the increase of the ionic radius of the rare earth ion and/or the doping level.

In the present work, the following R_0 values are used for subsequent calculations: 1.637 Å for $Co^{3+}-O^{2-}$ bond, 1.686 Å for Ni³⁺–O^{2−} bond, 2.105 Å for Nd³⁺–O^{2−} bond, and 2.118 Å for $Sr^{2+}-O^{2-}$ bond [[36](#page-11-0)–[38](#page-11-0)]. Computed bond valence sum by FullProf for cobalt atom concerning all compositions resting on Brown theory for oxygen deficiency on the $z = 0$ plane is recapitulated in Table [5.](#page-6-0) The GII of the NdSrNi₁ – $_{x}Co_{x}O_{4} \neq \delta$ (0.1 ≤ $x \leq 0.9$) compounds illustrated also in Table [5](#page-6-0) denotes that the crystal structure becomes more and more stable when we substitute nickel for cobalt.

Thermogravimetric analysis and oxygen stoichiometry determination

In order to investigate the variation in the oxygen deficiency as a function of cobalt proportion, we have performed thermogravimetric analysis in hydrogen/nitrogen atmosphere. A common aspect in the reduction curves for all compositions was the observation of a step-like behavior. These types of wide plateaus in the weight loss curves which appear throughout the reduction process are generally related to the stabilization of a number of phases with different oxygen stoichiometry [[39](#page-11-0)]. Departing from the weight loss values, oxygen contents are deduced for all compositions taking into account $Nd₂O₃$, SrO, and metallic Ni and Co as final products. Figure [7](#page-6-0) shows the relative weight loss curve for NdSrNi_{0.9}Co_{0.1}O_{4 − δ} and NdSrNi_{0.7}Co_{0.3}O_{4 − δ} which exhibit similar behavior. As it can be seen, the reduction phenomenon comprises three steps. The first process corresponds to the reduction of Ni^{3+} to Ni^{2+} . The second one stands for the reduction of Co^{3+} to Co^{2+} . The third plateau in this weight loss curve corresponds to reduction of Ni^{2+} and Co^{2+} into metallic nickel and cobalt without any distinction between them. Within the same framework, for both cases where $x = 0.7$ and $x = 0.9$, close results are recorded and presented in Fig. 8. In this specimen, the reduction phenomenon occurs only with two steps bringing about two plateaus. The first one for ≈400 ≤ $T(C^{\circ})$ ≤ 550, which stands for reduction of the majority of cobalt cation from Co^{3+} to Co^{2+} and minority of nickel from Ni^{3+} to Ni^{2+} . A second wide plateau is observed at $T(C^{\circ}) \ge 750$ and corresponds to the complete reduction of cobalt and nickel ions to metal. Referring to Table 6, oxygen

con

 \mathbf{s} ol

Fig. 9 Temperature dependence of ρ for x= 0.1, 0.3, 0.5, 0.7 and 0.9 in NdSrNi Co O. The inset (a) show $ln(\rho/T)$ versus 1/T plots in the hightemperature region. full line is the fit to equation: $ρ(T) = BT \exp(Ea/$ kBT). The temperature dependence of the conductivity σ in scales of ln(σ) as a function of $(1/T1/4)$ is also shown in the inset (b) with their fits (full lines) using the variable range hopping (VRH) model

content for the synthesized phases proves to be consistently under 4.00 atoms per formula unit and seems to be essentially dependent on the cobalt content. In fact, oxygen content diminishes up to $x = 0.3$ then rises with the increase of cobalt content. This behavior confirms the presence of Co^{3+} in compounds having a percentage of cobalt more or equal to 50%.

Electrical transport properties

The electrical resistivity measurements of NdSrNi_{1 – x}Co_xO_{4 ± δ} $(0.1 \le x \le 0.9)$ demonstrate that all the members of the solid solution are semiconducting between 20 and 390 K. In this common feature, the conductivity may be accounted for on the basis of two different models.

Firstly, Emin–Holstein theory of adiabatic small polaron hopping model (ASPH) [[40](#page-11-0)], which is expressed as

$$
\rho(T) = \text{BTexp}\Big(E_a \Big/k_B T\Big) \tag{1}
$$

where E_a is the activation energy for hopping conduction and B is the residual resistivity. Figure 9 portrays the corresponding results for the ASPH model, where the straight line is a fit to Eq. (1). The activation energy E_a is deduced from the fit of ln (ρ/T) versus $1/T$ curve.

Secondly, Mott et al. [\[41](#page-11-0)] have established the variable range hopping (VRH) mechanism, expressed by the following equation:

$$
\rho(T) = \rho_0 \exp\left(\frac{T_0}{T}\right)^{1/4} \tag{2}
$$

where ρ_0 is considered as a constant which depends on electron–phonon interaction, though it is slightly affected by temperature [[42\]](#page-11-0). The characteristic VRH temperature is noted $T_0 = 16 \cdot \alpha^3 / k_\beta N(E_F)$ where $N(E_F)$ is the density of states at the Fermi level. As a matter of fact, the T_0 value is calculated from the slope of the plot ln(σ) versus $T^{-1/4}$ (Fig. 9). Referring to [\[40\]](#page-11-0), constant α was taken as 2.22 nm⁻¹. Table [7](#page-9-0) illustrates

Table 7 Characteristic temperature T_0 , density of state at the Fermi level $N(E_F)$, activation energy E_a , and correlation factors R for NdSrNi₁ – $_{x}Co_{x}O_{4}$ – $_{\delta}$ $(0.1 \le x \le 0.9)$ using ASPH (Eq. [1](#page-8-0)) and VRH (Eq. [2](#page-8-0)) models

the computed values of T_0 , $N(E_F)$, and E_a based on these two models. At this stage of analysis, it is worth noting that our results obtained for the semiconducting phase can be described quite well according to both models but the first one (ASPH) provides the best squared linear correlation coefficients (R^2) . To conclude, we would assert that the transport behavior is dominated by the adiabatic small polaron hopping mechanism.

In this work and in related references [[43](#page-11-0)–[45](#page-11-0)], it is inferred that the obtained values of activation energy for NdSrNi₁ – $_{x}Co_{x}O_{4} \pm \delta$ (x = 0.1, 0.3, and 0.5) compounds (Table 7) are lower than that for NdSrCoO₄ [\[43\]](#page-11-0). Summing up these results, it can be concluded that after the incorporation of nickel and cobalt in the same site, we obtain a semiconductor material with conductivity in the range 0.3–17.24 Ω^{-1} cm⁻¹ and low activation energy. The same result is found by several authors, for instance for the NdSrNi_{0.5}Cr_{0.5}O_{4 ± δ} compound, the $\sigma_{\text{RT}} = 8.86 \Omega^{-1} \text{ m}^{-1}; E_a = 0.09 \text{ eV} [25], \text{ for}$ $Nd_{1.8}Sr_{0.2}Ni_{0.6}Cu_{0.4}O_{4 \tImes}$ the $\sigma_{640^{\circ}C} = 42.7$ (3) S cm⁻¹;

 $E_a = 0.07$ eV [[46\]](#page-11-0) and for NdSrNi_{0.8}Cu_{0.2}O_{4 ± δ} the $\sigma_{\rm RT} = 1.13 \Omega^{-1} \text{ cm}^{-1}$; $E_a = 0.06 \text{ eV}$ [\[17\]](#page-10-0).

Based on the outcome of our investigation, it is possible to conclude that NdSrNi₁ – $_{x}Co_{x}O_{4}$ $_{\pm 0}$ compounds having low activation energy (in the range 0.04–0.09) can be used as mixed ionic–electronic conductors (MIECs) [\[46\]](#page-11-0).

The room temperature resistivity (ρ_{RT}) values as a function of cobalt content x are presented in Fig. 10. As it can be observed, the ρ_{RT} of the samples is constant up to $x \approx 0.5$, then increases rapidly with x . It is significant to note that room temperature resistivity exhibits an increase each time x increases, which is different from the evolution in the formal valence of cobalt. Referring to Brian et al. [\[18\]](#page-11-0), the half-filled $\sigma_{x^2-y^2}^*$ band is responsible for the metallic behavior of NdSrNiO4 compound. In the present research, when substituting Ni³⁺ for Co²⁺/Co³⁺ (electronic configuration $t_{2g}^6 d_{z^2}^1 d_{x^2-y^2}^0$ / $t_{2g}^6 d_{z}^0 d_{x^2-y^2}^0$), the half-filled $\sigma_{x^2-y^2}^*$ band changes gradually to a half-filled $\sigma_{z^2}^*$ band due to the presence of Co^{2+} cation. On the

other side, when we have compositions corresponding to a higher percentage of cobalt than that of nickel $(x > 0.5)$, the resistivity increases remarkably with the incorporation of the diamagnetic Co³⁺ (t_{2g}^6) cation.

On the basis of the electrical resistivity ($\rho_{300K} = 0.058 \Omega$ cm in bulk form) and the structural and microstructural characteristics of the NdSrNi_{0.5}Co_{0.5}O_{4 ± δ} material, and at least, it can be successfully deposited in thin film and used as a bottom electrode suitable for the epitaxial regrowth of functional ferroelectricbased parallel plate capacitors. Our material has great potential for other applications as an electrode in solid oxide fuel cells (SOFCs) operating at intermediate temperatures.

Conclusion

Cobalt can be successfully substituted for nickel in the parent compound NdSrNiO4. Refinements of X-ray diffraction data for all compositions $(x = 0.1, 0.3, 0.5, 0.7,$ 0.9) show that these phases crystallize in the tetragonal K₂NiF₄-type structure (I4/mmm space group with $z = 2$). Oxygen content deduced from thermogravimetric analysis (TGA) under reducing atmosphere displays substantial deviations from the ideal values after the incorporation of cobalt in the structure. These oxygen vacancies are the possible ionic charge carriers. The a and c unit cell parameters values exhibit certain variations with increasing cobalt content. This behavior is deduced from Jahn–Teller distortions as well as the presence of the mixed valence of Co^{2+}/Co^{3+} ions. Bond valence calculations show that cobalt has an oxidation state between two and three. The GII of NdSrNi₁ – $_{x}Co_{x}O_{4}$ – $_{\delta}$ (0.1 $\leq x \leq$ 0.9) compounds indicates that the crystal structure becomes more stable when we substitute nickel for cobalt. The transport mechanism in NdSrNi₁ – $_{x}Co_{x}O_{4}$ – $_{\delta}$ (0.1 ≤ x ≤ 0.9) solid solution has been investigated and shows an adiabatic small polaron hopping model. Departing the activation energy that has been undertaken, it is possible to conclude that NdSrNi_{1 – x}Co_xO_{4 ± δ} (x = 0.1, 0.3, and 0.5) compounds are semiconductors with low activation energy. This paper has clearly demonstrated that the composition NdSrNi_{0.5}Co_{0.5}O_{4 ± δ} exhibits the most interesting physical properties: maximum conductivity at room temperature $\sigma = 17.24 \Omega^{-1}$ cm⁻¹ is coinciding with minimum activation energy ($E_a = 0.05$ eV). These findings are of direct practical relevance in SOFCs and plate capacitors (bottom electrode).

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