

September 2020 Vol. 63 No. 9: 1656–1674 https://doi.org/10.1007/s11431-019-1463-6

### Research progress based on observations of the New Vacuum Solar Telescope

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Received June 14, 2019; accepted October 14, 2019; published online April 24, 2020

The purpose of this paper is to introduce the main scientific results made by the one-meter New Vacuum Solar Telescope (NVST), which was put into commission on 2010. NVST is one of the large aperture solar telescopes in the world, located on the shore of Fuxian lake of Yunnan province in China, aiming at serving solar physicists by providing them with high resolution photospheric and chromospheric observational data. Based on the data from NVST and complementary observations from space (e.g., Hinode, SDO and IRIS, etc), dozens of scientific papers have been published with a wide range of topics concentrating mainly on dynamics and activities of fine-scale magnetic structures and their roles in the eruptions of active-region filaments and flares. The achievements include dynamic characteristics of photospheric bright points, umbral dots, penumbral waves, and sunspot/light bridge oscillation, observational evidence of small-scale magnetic reconnection, and fine-scale dynamic structure of prominences. All these new results will shed light on the better understanding of solar eruptive activities. Data release, observation proposals, and future research subjects are introduced and discussed.

sunspots, solar flares, solar filaments/prominences, jets, Ellerman bombs, magnetic reconnection

Citation: Yan X L, Liu Z, Zhang J, et al. Research progress based on observations of the New Vacuum Solar Telescope. Sci China Tech Sci, 2020, 63: 1656–1674, https://doi.org/10.1007/s11431-019-1463-6

#### 1 Introduction of observational systems

NVST is a vacuum solar telescope with a 985 mm clear aperture located at the Fuxian Solar Observatory of the Yunnan Observatories [1]. The main instruments include the highresolution multi-channel imaging system [2] and the high dispersion multi-wavelength spectrograph [3, 4]. The two systems have identical focus and are arranged to be perpendicular to each other. The multi-wavelength spectrograph can observe two chromospheric lines H $\alpha$  656.3 nm and Ca II 854.2 nm, and one photospheric line Fe I 5324, simultaneously. In this paper, we only focus on the imaging system. The working wavelengths and the main properties of each channel are

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summarized below.

#### 1.1 High resolution multi-channel imaging system

The imaging system comprises five channels, i.e., two broadband interference filter channels and three narrow-band Lyot filter channels. The working wavelengths and the main properties of each channel are summarized as follows.

(1) H $\alpha$  channel, to monitor magnetic structures in the chromosphere. The central wavelength can be tunable in the range of (656.28 ± 0.4) nm and the full bandpass width is 0.025 nm (e.g., giving a FWHM of an order of 11 km/s).

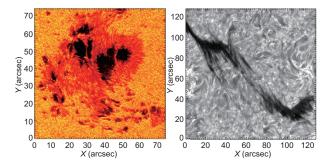
(2) Ca II H/K channel, to monitor magnetic structures in the lower chromosphere. H and K lines are alternative. The central wavelength can be tunable in the range of (393.3 (3) He I 1083 nm, a near infrared proxy of coronal holes, to explore magnetic fields in the low corona. The central wavelength can be tunable in the range of  $(1083 \pm 0.5)$  nm and the full bandpass width is 0.05 nm.

(4) G-band, to indicate small magnetic structures in the deep photosphere. The central wavelength is 430 nm and the full bandpass width is 1 nm.

(5) TiO band, sensitive to the temperature, to allow easy detection of BPs in granular lanes or umbral dots. It is centered at 705.8 nm and the full bandpass width is 1 nm.

Although only three channels including  $H\alpha$ , TiO and Gband have been used, the observations provide impressive evidence of the high-resolution capability of the NVST. Particularly, the off-band observation in H $\alpha$  is very impressive [5]. The post-processing with the speckle masking method produces the results with the resolution close to the diffraction limit of the TiO band. There are 3 levels of data products. Level-0 represents the raw status. Level-1 is achieved based on a luck image selection algorithm, followed by the dark current and flat field modification. Level-1.5 is reconstructed by a speckle masking method [5-7]. To do so, we have to take a large number of short-exposure images (H $\alpha$  is 20 ms, TiO is about 1 ms). Using a large number of frames (more than 100 frames) and a statistical algorithm, we are able to reconstruct one frame with the whole field of view ( $\approx$  3 arcmin). H $\alpha$  images have a pixel size of 0."163 and a cadence of 12 s. In addition, TiO images have a pixel size of 0."04. The time cadence of Level 1 and 1.5 data products is about 12 s in H $\alpha$  channel and 30 s in TiO channel.

Figure 1(a) shows an example of a complex active region (AR) named NOAA 12673 observed on September 5, 2017. The penumbra threads and fine structures in a light bridge are well resolved in the high resolution TiO image. Fine structures of the chromospheric fibrils, umbra dots, and penumbra



**Figure 1** (Color online) High resolution observation of photosphere and chromosphere. (a) A complex active-region NOAA 12673 in TiO image observed by the NVST at 03:45:18 UT on September 5, 2017. This active-region includes several sunspots and pores. The large sunspots are hosting light bridge. (b) Fine structure of a quiescent filament observed at H $\alpha$  line center by the NVST at 05:31:01 UT on July 23, 2018.

wave can be clearly found in the observation of a sunspot at the H $\alpha$  center images. Zhu et al. (2016) [8] found that the magnetic field lines from non-linear force-free field (NLFFF) extrapolations are consistent with the patterns of H $\alpha$  fibrils observed by the NVST. In particular, the fine structure of solar filaments observed by the NVST is very complicate. Figure 1(b) shows an example of a quiescent filament observed at 05:29 UT on July 23, 2018. Many thin threads perpendicular to the spine are found. These data can be used to study the structure and evolution of different atmospheres from the photosphere to the chromosphere.

#### 1.2 High dispersion multi-wavelength spectrographs

The NVST actually consists of two vertical spectrographs, which are alternative and co-existing in space by dispersing the light at two different directions. They share the same entrance slit, but the slit direction has to be changed by rotation. The length along the slit is about 136 arcseconds. Several slits with different widths (60-150 micrometer) can be selected. One spectrograph works at the optical wavelength band. The focus length of the collimator is about 6 m. A 1200 mm<sup>-1</sup> grating is used to obtain two chromospheric lines H $\alpha$  and Ca II 854.2 nm and one photospheric line Fe I 5324, simultaneously. In addition, the grating can be rotated to switch the observation window to other lines. The main properties of this spectrograph are listed in Table 1. The other spectrograph is for the near-infrared band. The focus length is about 9 m. A 316 mm<sup>-1</sup> echelle grating is used to obtain the high order spectrum of the He I 10830 nm (the 5th order) and Fe I 1.56 micrometer (the 3rd order). The precise reduction of these spectra observed by the NVST can be seen from Wang et al. (2013) [3], and Cai et al. (2017, 2018) [3,4,9].

#### 2 Data release and proposed observations

The NVST has been taking the daily observations since October, 2012. Quick-look movies (Level 1 or Level 1+) of all observation entries can be browsed at http://fso.ynao.ac.cn. Data requirement can be submitted through this website. In addition to the daily observations, the NVST also takes observations proposed by scientists around the world. Applications for observation proposals usually start around the middle of May and end in the middle of June. Application form can be downloaded from http://fso.ynao.ac.cn/proposal.aspx. The contents related to terminals may change every year according to the operation condition. Scientists affiliated to the CAS can access the application form by China Virtual Observatory (China-VO). The NVST scientific committees will review and evaluate all proposals. The accepted proposals will be carried out at the NVST from September to Novem-

 Table 1
 Main properties of the multi-wavelength optical-band spectrograph

Item	Parameter 156 mm×220 mm		
Grating size			
Grooves per mm	1200 g/mm		
Blazed angle (°)	36.8 (1st order, blazed wavelength 10000 Å)		
Present slit width (microns)	100 (0.45 arcsec)		
Linear dispersion (mm/Å)	0.75 @ Hα, 0.82 @ Ca II 8542, 2 @ Fe I 5324		
Raw spectra frame size (pixels)	2672×4008		
Pixel size (microns)	9×9		
Spectral sample (mÅ/pixel)	12 @ Hα, 11 @ Ca II 8542, 4.5 @ Fe I 5324		
Spatial sample (arcsec /pixel)	0.041		

ber. The principle investigator (PI) has the priority to keep the observation data for three months. After three months, the data will be open on our website.

Since 2015, the NVST has carried out many proposed observations from National Astronomical Observatories of China, Purple Mountain Observatories, Peking University, Shandong University, Xinjiang Observatories, National Space Science Center, Kunming University of Science and Technology, Hida Observatory of Japan, and Indian Institute of Astrophysics. In addition, the NVST has also taken coordinate observations with other space or ground-based telescopes, such as GST at BBSO, Domeless Telescope at Hida, Hinode, IRIS, and GREGOR.

## 3 Main scientific results based on the NVST data

High Resolution Multi-channel Imaging System can image the photosphere and the chromosphere simultaneously. Based on these high resolution data, many features in the photosphere and the chromosphere were studied. These results are divided into five parts according to their different properties.

# **3.1** Dynamic characteristics of photospheric bright points, umbral dots, penumbral wave, and sunspot/light bridge oscillation

Since the NVST started observations, many quiet regions and sunspots have been observed. In quiet sun, many dynamical features can be seen, such as the bright points, granules, granule dark lane, and so on. The light bridges and umbral dots are often seen in the umbra of sunspots. In the upper atmosphere of sunspots, especially in the chromosphere, penumbral wave and umbral oscillations are the main features. Using the high resolution data of the NVST, the following studies were carried out.

#### 3.1.1 Bright points

In high-resolution observation of photosphere, bright points are prominent features located between the granules. Using six high-resolution TiO-band image sequences taken by the NVST from 2012 to 2014, Ji et al. (2016) [10] investigated the morphologic, photometric dynamic properties of BPs in terms of equivalent diameter, the intensity contrast, lifetime, horizontal velocity, diffusion index, motion range, and motion type. In this study, the Laplacian and morphological dilation algorithm (LMD) [11] and three-dimensional segmentation algorithms were used to detect and trace the evolution of the BPs. Moreover, the statistical properties of the magnetic fields of BPs are also explored and compared by using the vector magnetograms.

Many useful parameters of BPs are obtained as follows. The areas of BPs range from 0.2% to 2%. The mean equivalent diameters range from (168±29) km to (195±36) km. The mean ratios of intensity contrast range from  $(0.99\pm0.04)$ to  $(1.06\pm0.05)$ . The mean lifetimes range from 104 to 141 s. The mean horizontal velocities range from  $(1.04\pm0.54)$ km s<sup>-1</sup> to  $(1.35\pm0.71)$  km s<sup>-1</sup>. The mean diffusion indices range from  $(0.86\pm0.39)$  to  $(1.31\pm0.54)$ ; The mean ratio of motion range values range from  $(0.96\pm0.67)$  to  $(1.30\pm0.80)$ , and the mean index of motion type values range from 0.58 to 0.69. Note that the rate of motion range is defined as  $m_r =$  $\sqrt{X_{\text{max}} - X_{\text{min}}^2 + Y_{\text{max}} - Y_{\text{min}}^2}/r$ , where  $X_{\text{max}}$  and  $X_{\text{min}}$  are the maximum and minimum coordinates of the path of a single BP in the x-axis, and  $Y_{\text{max}}$  and  $Y_{\text{min}}$  are the ones in the y-axis, r is the radius of the circle which corresponds to the maximum size of the BP during its lifetime. The index of motion type is defined as  $m_t = d/L$ , where  $d = \sqrt{X_n - X_1^2 + Y_n - Y_1^2}$  and L is the whole path length. In addition, the relationship between the properties of BPs and their embedded magnetic environments has been investigated. It is found that the stronger magnetic field corresponds to the bigger and brighter BPs, while the weaker magnetic field corresponds to the smaller and weaker, faster diffusion, and faster movement of the BPs.

The detailed parameters are shown in Table 2.

Liu et al (2018) [12] studied the properties of isolated BPs [13] and non-isolated BPs in an active region observed at TiO (705.8 nm) wavelength by the NVST with the seeing parameter  $r_0$  of (10.628 ± 0.220) cm. They presented a novel algorithm of identifying and tracking BPs by using the Laplace kernel and the second order derivatives in x- and y-directions to obtain the edge information, and by assuming that two BPs obtained in single frames in consecutive images is the same one appearing at different time if one has overlapped with the other. Based on these ideas, they identified and tracked the photospheric BPs, and further classified them into the isolated (see Figure 2(a)-(e)) and non-isolated ones (Figure 2(f)-(j)). The isolated BPs are the ones that did not experience split and mergence in their lifetimes. However, the nonisolated ones are those that experienced split and mergence for at least once in their lifetimes. In the algorithm, an evolving BP is considered as a non-isolated BP if it appeared in one frame as a single blob in a frame but more than one blob appeared in previous or next consecutive images. Otherwise, the evolving BP is considered as an isolated one.

Once the BPs were identified, tracked, and classified, they started to study the isolated and the non-isolated BPs in three regions with different background magnetic fields. Namely, Region 1 contains a small pore with the strongest background fields; Region 2 has no pores but contains elongated BPs with weaker background fields; Region 3 contains the least BPs with the weakest background fields. In these regions, they found that the number density of BPs obtained in single frame appeared most in Region 1, and least in Region 3. Therefore, the number density of snapshot BPs was considered to be strongly dependent on the strength of the background fields as well as the area coverage of BPs. On the other hand, the size/intensity contrast distributions of both the isolated and the non-isolated BPs were found to be independent of the strength of the background fields. However, their lifetimes are slightly shorter in the stronger background filed regions.

BPs are located in the dark inter-granular between the granule cells. By using the data of GST, Ji et al. (2012) [14] found that there are many ultrafine channels rooted in the dark inter-granular connecting from the photosphere to the corona, in which the energy may be transferred from the lower atmosphere to the higher one. Whether BPs play an important role in the energy transfer or not is deserved for further investigation.

#### 3.1.2 Umbral dot, penumbral wave

Umbral dots (UDs) are small isolated brightenings in sunspot umbrae under high resolution observations. Their appearance indicates the existence of convection. According to their locations inside an umbra, UDs can be divided into two types: central UDs (CUDs) and peripheral UDs (PUDs). To investigate UD properties, the difference between UDs and PUDs, and the relationship between their dynamic properties and magnetic field strengths, Ji et al. (2016) [15] selected highresolution TiO images of four ARs taken by the NVST. Six sunspots were selected from the four ARs (NOAA 11598, 11801, 12158, and 12178). A total of 1220 CUDs and 603 PUDs were identified by using the phase congruency technique [16]. Meanwhile, the line-of-sight magnetograms of the sunspots taken with the Helioseismic and Magnetic Imager (HMI) on-board the Solar Dynamics Observatory (SDO) were used to determine umbral magnetic field strengths. The main results are as follows: the diameters and lifetimes of

Table 2 Properties of the BPs of the six data sets. Table reproduced from Ji et al. (2016) [10], copyright by RAA

Data set	1	2	3	4	5	6
Area coverage	0.20%	0.99%	1.55%	1.53%	1.75%	1.99%
Equivalent diameter (km)	181±22	$168 \pm 29$	178±29	195±36	184±38	194±36
[min,max]	[111, 245]	[103, 402]	[109, 440]	[106, 445]	[122, 447]	[112, 473]
Intensity contrast	$0.99 \pm 0.04$	$1.01 \pm 0.04$	$1.03 \pm 0.04$	$1.05 \pm 0.06$	$1.05 \pm 0.04$	$1.06 \pm 0.05$
[min,max]	[0.91, 1.12]	[0.90, 1.19]	[0.90, 1.31]	[0.92, 1.30]	[0.92, 1.24]	[0.89, 1.28]
Lifetime (s)	104±104	133±133	114±114	$141 \pm 141$	121±121	124±124
[min,max]	[103, 582]	[102, 826]	[120, 723]	[119, 735]	[120, 572]	[114, 580]
Velocity (km s <sup>-1</sup> )	$1.35 \pm 0.71$	$1.23 \pm 0.64$	$1.06 \pm 0.55$	$1.04 \pm 0.54$	$1.06 \pm 0.55$	$1.05 \pm 0.55$
[min,max]	[0.01, 5.27]	[0, 6.80]	[0.08, 5.43]	[0.02, 5.32]	[0.06, 5.21]	[0.03, 5.75]
Diffusion index	1.31±0.65	1.21±0.78	0.91±0.43	$1.05 \pm 0.67$	$0.86 \pm 0.49$	0.93±0.77
[min,max]	[-3.64, 3.93]	[-4.91, 5.39]	[-4.17, 4.28]	[-5.21, 6.51]	[-7.00, 4.21]	[-5.70, 4.43]
Ratio of motion range	$1.30 \pm 0.80$	1.18±0.76	$1.11 \pm 0.77$	$1.02 \pm 0.62$	$0.96 \pm 0.67$	1.03±0.69
[min,max]	[0.31, 5.06]	[0.15, 6.39]	[0.04, 6.28]	[0.17, 5.42]	[0.13, 4.73]	[0.13, 4.79]
Motion type	$0.69 \pm 0.69$	$0.69 \pm 0.69$	$0.58 \pm 0.58$	$0.59 \pm 0.59$	$0.59 \pm 0.59$	$0.62 \pm 0.62$
[min,max]	[0.08, 0.99]	[0.04, 1.00]	[0.23, 0.99]	[0.03, 0.99]	[0, 1.00]	[0.04, 0.98]

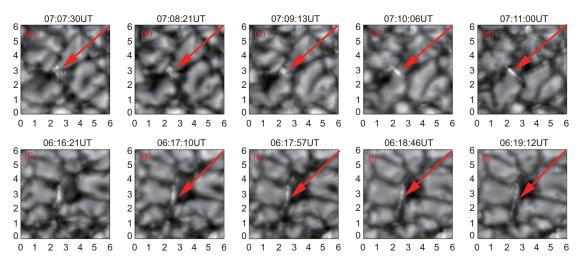


Figure 2 (Color online) Two different bright points observed by the NVST. (a)–(e) Evolution of the isolated bright points (BPs); (f)–(j) evolution of the non-isolated BPs. Image reproduced with permission from Liu et al. (2018), copyright by AAS.

UDs increased with the brightness, but velocities did not. Moreover, diameters, intensities, lifetimes and velocities depend on the surrounding magnetic fields. The CUD diameter was larger, brighter, and has longer lifetime and slower motion in a weak umbral magnetic field environment than that in a strong one. By using Goode Solar Telescope (GST) data, Su et al. (2016) [17] also found that the wavefronts rotate the counter-clockwise and clockwise rotation alternately in the sunspot umbra.

Typical 5-minute and 3-minute oscillations were found in the photosphere and chromosphere of the Sun for dozens of years. High resolution GST observation also revealed that the periods of the running umbral waves are from 2.2 to 2.6 min [17, 18]. Using high spatial and temporal resolution observations of the NVST and Atmospheric Imaging Assembly (AIA) on board the SDO, Wang et al. (2018) [19] discovered one minute oscillation at different heights above sunspot umbra. They used a novel time-frequency analysis method named the synchrosqueezing transform (SST) to represent their power spectra and to reconstruct the high-frequency signals at different solar atmospheric layers. The method SST is capable of resolving weak signals even when their strength is comparable to the high-frequency noise. A significant enhancement between 10 and 14 mHz (labeled as 12 mHz) at different atmospheric layers was found by analyzing the spectrum in the umbra. Moreover, the 12 mHz component exists only inside the umbra. Besides, Zhou et al. (2017) [20, 21] found that the period of the running waves within the umbra is half of that within the penumbra. The former is a typical 3-min oscillation with a period of 156 s and the latter is a typical 5-min oscillation with a period of 312 s.

#### 3.1.3 Light wall on the sunspot light bridge

Using high resolution photospheric and chromospheric data

from the NVST, Yang et al. (2015) [22] found a special structure rooted in the light bridge of the sunspot and this structure exhibited upward and downward movements. They named this structure as light-wall. The deprojected mean height, amplitude, oscillation velocity, and the dominant period are determined to be 3.6 Mm, 0.9 Mm, 15.4 km s<sup>-1</sup>, and 3.9 min, respectively. They interpreted the oscillations of the light wall as the leakage of p-modes from below the photosphere. In the following, Hou et al. (2016) [23] reported that the light wall can also appear near the umbral-penumbral boundary and along a neutral line between two small sunspots. Moreover, the wall body consists of multi-layer and multi-thermal structures that occur along magnetic neutral lines in active regions. In the following, Tian et al. (2018) [24] found that the surge-like activity previously reported above light bridges in the chromosphere has two components. One component is the ever-present short surges likely to be related to the upward leakage of magneto-acoustic waves from the photosphere. The other is the occasionally occurring long and fast surges that are obviously caused by the intermittent reconnection jets.

### **3.2** Formation, fine-scale structures, and eruption mechanism of solar filaments

Solar filaments/prominences are very common features in the solar atmosphere and relatively dense and cool plasma structures embedded in the low-density and hot corona. Due to the different locations of their appearances, they are given different names. When seen in emission above the limb, they are called as prominences, and they are called filaments when seen in absorption line (H $\alpha$ ) against the solar disk. According to their locations on the solar disk, they can be divided into three types, i.e., quiescent filaments, active-region fila-

ments, and intermediate filaments. Quiescent filaments are long-lasting structures with lifetimes of several days, weeks, or even much longer. Compared with quiescent filaments, active-region filaments are relatively short-lived structures lasting for several hours or days. The eruption of solar filaments is often associated with solar flares and coronal mass ejections (CMEs). A typical coronal mass ejection (CME) includes three parts, i.e., bright front, dark cavity, and bright core. It is generally believed that the bright core of a CME is a filament [25]. Therefore, studying solar filaments is very important for understanding the productions of CMEs.

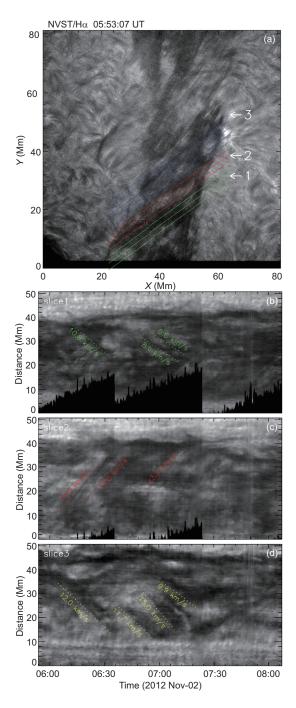
#### 3.2.1 Fine structure of solar prominences

The complex structures of solar prominences have always attracted people's attention. In general, quiescent prominences have two linked categories of structures: "spine" and "barb". A"spine" is the longest part of a quiescent prominence that forms a horizontal axis composed of many resolvable threads, while"barb" is connected to the spine and terminates in the chromosphere [26].

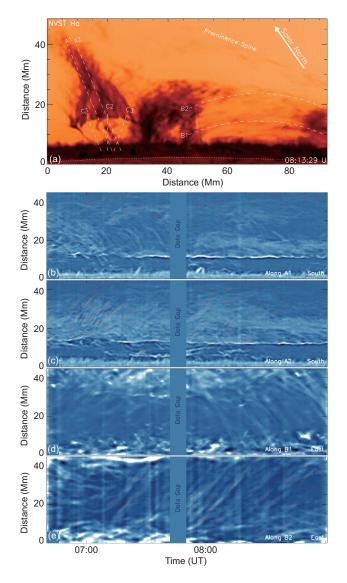
Due to the good seeing in Fuxian lake, a lot of quiescent prominences were observed by the NVST. The ubiquitous counter-streaming flows are the most prominent features in the prominences. Yan et al. (2015) [27] found that the largescale counter-streaming flows with a certain width in a solar prominence are co-aligned well with the prominence threads on November 2, 2012 (see Figure 3(a)). The velocity of these material flows ranged from 5.6 to 15.0 km s<sup>-1</sup> (see Figure 3(b)–(d)). In another quiescent prominence, the formation of several parallel tube-shaped structures in the barb is found and the width of the structures ranged from about 2.3 to 3.3 Mm during the evolution of the prominence on September 29, 2013. The parallel tube-shaped structures merged together accompanying with material flows from the spine to the barb. Li et al. (2018) [28] found that the material in the middle section of a prominence is mainly restricted to flow back and forth at a certain region. Shen et al. (2015) [29] have analyzed the evolution of a bubble structure and the counterstreaming flows in a prominence on May 20, 2014, in detail (see Figure 4(a)). They found that the collapse and oscillation of the bubble boundary were tightly associated with a flarelike feature located at the bottom of the bubble. To obtain the velocities of the upward and downward flow in the prominence, the time-distance diagrams were made along the curve lines A1, A2, B1, and B2 in Figure 4(a). The average speeds of the upward and downward flows along the curve lines A1 and A2 are  $(13.7 \pm 5.2)$  km s<sup>-1</sup> and  $(12.5 \pm 5.2)$  km s<sup>-1</sup>, respectively (see Figure 4(b) and (c)). The average speeds of the counter-streaming flows in the spine of the prominence along the curve lines B1 and B2 are  $(20.1 \pm 4.5)$  km s<sup>-1</sup> and

 $(15.8 \pm 3.8)$  km s<sup>-1</sup>, respectively (see Figure 4(d) and (e)).

Except for the counter-streaming flows, the Kelvin-Helmholtz instability was also identified in a quiescent prominence by Yang et al. (2018) [30] and Li et al. (2018) [31]. Using the NVST red-wing H $\alpha$  data on June 4, 2016 and



**Figure 3** (Color online) High-resolution H $\alpha$  image of the quiescent filament observed by NVST at 05:53:07 UT on November 2, 2012. The lines in (b) and (d) mark the areas that have material flows with the same di-rection and the lines in (c) denote material flows in opposite directions. The lines in (a) marked by the numbers indicate the positions of the time slices shown in Figure 5(b)–(d). Image reproduced from Yan et al. (2015), copyright by RAA.



**Figure 4** (Color online) A limb prominence and the counter-streaming. (a) A quiescent prominence observed by the NVST; (b)–(e) the time-distance diagrams showing the mass dynamical motion along the slit A1, A2, B1, and B2. Image reproduced with permission from Shen et al. (2015), copyright by AAS.

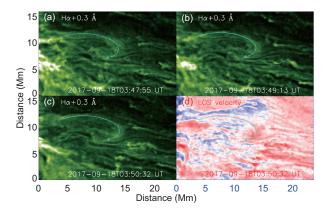
September 18, 2017, the vortex-like structures were detected on the interface/surface of prominence plumes by Yang et al. [30] and Li et al. [31]. Yang et al. [30] found 2'' vortex before the formation of 5'' size vortex in the prominence. Moreover, the 255 s oscillatory period was identified in the prominence. Li et al. [31] measured the speeds of the mass flows in the prominence and found that they are larger than the local sound speed. They suggested that the Kelvin-Helmholtz instability may occur in the interface/surface of prominence plumes. The vortex-like structure was also found during the evolution of the prominence (see H $\alpha$  off-band observation in Figure 5(a)–(c)). The line-of-sight velocity of the vortex-like structure exhibited red shift (see Figure 5(d)). The thread oscillation is another characteristic of prominences. Due to the high spatial and temporal resolution data of the NVST, two types of the small-scale oscillations of prominence threads are found by Li et al. (2018). In one type, the oscillation period firstly decreased and then increased. In the other type, the oscillation period grew quickly at the beginning and then decreased. Note that the direction of these small-scale oscillations mentioned above is perpendicular to that of the prominence threads.

The nature of the complex magnetic structures and the mass flows in the prominences is still unclear. The formation mechanisms and driven forces of the mass flows are deserved for further investigation. In addition, it is very difficult to explain the complex magnetic structures of solar prominences observed by the NVST. Therefore, there is still a long way to completely understand the magnetic structures and mass flow of solar prominences.

#### 3.2.2 Filament formation

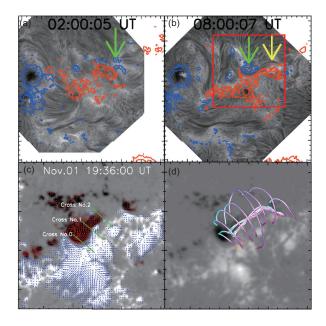
Due to the low time resolution data in the past, the formation process of solar filaments is difficult to capture. Up to now, formation and eruption mechanisms of solar filaments are still controversial. Especially, the studies on the formation of filaments are very rare. Thanks to the high temporal and spatial resolution observation of the NVST, several cases of filament formation were captured by the NVST. Based on these observations, several new formation mechanisms of solar filaments were proposed.

Yan et al. (2015) [32] studied the formation and magnetic structures of two successive active-region filaments in AR NOAA 11884. From October 31 to November 3, 2013, there are four-day observations with nice seeing. During this period, there are two active-region filaments formed



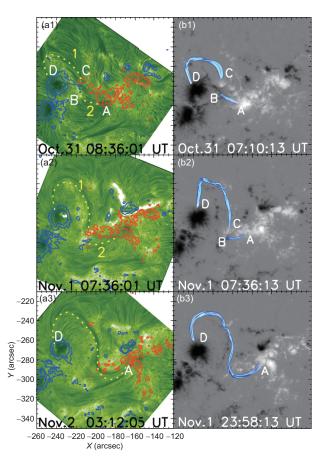
**Figure 5** (Color online) Evolution of a quiescent prominence observed by the NVST in off-band H $\alpha$  image. (a)–(c) H $\alpha$  +0.3 Å images; (d) line-of-sight velocity. The curved arrows indicate the vortex flow in the quiescent prominence. Image reproduced with permission from Li et al. (2018), copyright by AAS.

successively at the same position. Before the filament formation, there are several chromospheric fibrils connected a small sunspot with negative polarity (see Figure 6(a)). The arrow points to the small sunspot in Figure 6(a) and (b). Through studying the vector magnetograms and TiO observations, the small sunspot was found to rotate around its center in the photosphere. Due to the small sunspot rotation and the shearing motion between the two opposite polarities, the first filament formed at the polarity inversion line (PIL) (see Figure 6(b)). The right arrow in Figure 6(b) indicates the newly formed filament. Using the non-linear force free field extrapolation (NLFFF) based on the vector magnetograms observed by SDO/HMI, the twisted magnetic structure was found at the location of the filament (see Figure 6(c) and (d)). After the formation of the first filament, it experienced a failed eruption and disappeared. The second filament formed quickly at the same location after the vanish of the first one. Interestingly, one foot-point of two active-region filaments is always rooted in the small rotating sunspot. According to these observations, Yan et al. (2015) suggested that shearing motion and small pore rotation play an important role in the formation of two active-region filaments.



**Figure 6** (Color online) The formation of the active-region filaments observed by the NVST and its magnetic structure. (a), (b) H $\alpha$  images observed at 02:00:05 UT on October 31, 2013 and at 08:00:07 UT on November 1 by NVST superimposed on the corresponding line-of-sight magnetogram observed by the SDO/HMI. The contours indicate the positive and negative polarities. The contour levels are ±500 G and ±1000 G. The arrows indicate the small sunspot and the active-region filament. (c) The vector magnetogram observed by SDO/HMI at 19:36:00 UT on November 1. (d) The extrapolation of the filament structure and the surrounding magnetic fields of the filament superimposed on the longitudinal magnetic fields. Image reproduced from Yan et al. (2015), copyright by AAS.

Besides the case mentioned above, there is another case on the formation of an active region-filament reported by Yan et al. (2016) [33] and summed up by Guo et al. (2017) [34]. An inverse S-shaped filament formed in active region NOAA 11884 from October 31 to November 2, 2013. The clockwise rotation of a small positive sunspot around the main negative trailing sunspot dragged the chromospheric fibrils and resulted in the formation of a curved filament (see the dashed line 1 named CD in Figure 7(a1)–(a2)). Then the small positive sunspot cancelled with the negative magnetic flux to create a long inverse S-shaped active-region filament (see the dashed line named AD in Figure 7(a3)). In the chromospheric H $\alpha$  observations, the formation of the inverse Sshaped active-region filament was due to the magnetic reconnection between the curved filament 1 (the dashed line CD) and the chromospheric fibrils 2 (the dashed line AB) (see Figure 7(a1)–(a2)). Using the non-linear force free field extrapolation (NLFFF) based on the vector magnetograms observed by SDO/HMI, the formation process of the magnetic structure supporting the filament was reproducted (see Figure



**Figure 7** (Color online) Formation of an active-region filament observed by the NVST and the extrapolation magnetic fields by using NLFFF model. (a1–a3) H $\alpha$  images observed by the NVST; (b1–b3) the extrapolations based on the vector magnetograms observed by SDO/HMI. Image reproduced from Yan et al. (2016), copyright by AAS.

7(b1)–(b3)). In this case, the movement of the small sunspot and the magnetic cancellation play a key role in the formation of this filament.

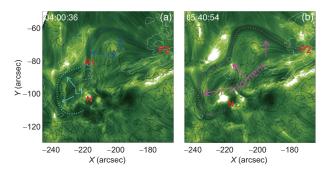
Different from these cases mentioned above, Xue et al. (2017) [35] presented a case of an active-region filament formation due to the tether-cutting magnetic reconnection between two groups of chromospheric fibrils (see Figure 8(a)). The two groups of the chromospheric fibrils L1 and L2 experienced tether-cutting reconnection at their joint and the active-region filament with the twisted structures formed (see Figure 8(b)). Wang and Liu (2019) [36] also mentioned this observational result in their review paper. The similar phenomenon were also reported by Yang et al. (2016) [37] and Chen et al. (2016) [38] by using SDO data.

In summary, several formation mechanisms of solar filaments are found based on NVST H $\alpha$  observations. These results are very important for us to understand the nature of solar filaments.

#### 3.2.3 Magnetic structures of active-region filaments

How the magnetic structures support the cool plasma against gravity in solar filaments is an important issue in solar physics. In order to explain the observations of the magnetic field configuration of solar filaments, there are two popular views proposed by some authors. One view is the sheared arcade model [39, 40], and the other is the flux rope model [41–44].

Based on the NVST observations, several active-region filaments were observed during their formation and activation stages. Yang et al. (2014) [45] studied a filament activation by using H $\alpha$  data. They found that a filament was activated due to magnetic field cancellation. After the rise of the activated filament, the twisted flux rope structure was exhibited by tracking the flows along helical threads. Using the photospheric vector magnetograms, Yan et al. (2015) [32] and Yan



**Figure 8** (Color online) Formation of an active-region filament observed by the NVST. L1 and L2 indicate the left and the right part of the chromospheric fibrils marked by curved lines in (a). The curved line in (b) indicates the newly formed filament. Image reproduced with permission from Xue et al. (2017), copyright by AAS.

et al. (2016) [33] extrapolated the magnetic structures of two active-region filaments by using non-linear force free field model and found that there is a flux rope at the locations of the active-region filaments. According to the connections of the filament threads, Xue et al. (2016) [46] also found the threads of the filament winded each other and formed a flux rope. By tracking two rotating features during a filament eruption, Chen et al. (2019) [47] found that a total twist of about 1.3  $\pi$  is estimated to be stored in the filament before its eruption. These observations are inclined to support that active-region filaments may have twisted magnetic structures before their eruptions. Awasthi, Liu and Wang (2019) [48] investigated the Doppler shifts and longitudinal oscillation in an active-region filament and found that the composite motions of the filament material suggest a double-decker host structure with mixed signs of helicity, comprising a flux rope atop a sheared-arcade system.

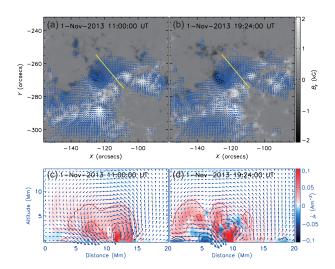
The exact magnetic structures of solar active-region filaments are still unclear. The measurement of chromospheric and coronal magnetic fields may help to solve this problem.

#### 3.2.4 Filament eruption

An eruptive filament is often taken as the bright core of a typical CME. The mechanism of filament eruption has been studied for many years. Several models have been proposed, e.g., kink instability [49, 50], torus instability [51]), break-out model [52], tether-cutting model [53], catastrophic loss of equilibrium [54, 55], and so on. Up to now, the trigger mechanism for a specified observed filament eruption remains unclear.

By using NVST and SDO data, Wang et al. (2016) [56] reported that the opposite horizontal electric currents under the filament appeared before onset of the filament eruptions. They suggested that these opposite electric currents were carried by the new flux emerging from below the photosphere, which might be the trigger mechanism of the filament eruption. Figure 8 presents the observation of vector magnetograms and the electric current perpendicular to the axis of the filament. The lines in Figure 9(a) and (b) indicate the position of the cross-section of the electric current (Figure 9(c) and (d)). Compared with Figure 9(a) and (b), it is obvious that the opposite current appeared before the filament eruption.

Sometimes, solar filaments consist of two parts, which have different twisted magnetic structures. Bi et al. (2015) [57] found a filament composed of two twisted flux ropes winding around each other. One has higher twist than the other. The one with higher twist erupted and the other with the lower twist part did not erupt. They suggested that the highly twisted flux rope became kink unstable when the



**Figure 9** (Color online) Vector magnetograms and electric currents along the axis of the filaments. (a), (b) Vector magnetograms observed by the SDO/HMI; (c), (d) the electric currents along the axis of the filaments derived from the extrapolated 3-D magnetic fields by using NLFFF model. The lines indicate the positions of the cross-section of the electric currents perpendicular to the filaments. Image reproduced with permission from Wang et al. (2016), copyright by AAS.

instability threshold declined with the expansion of the flux rope. Cheng et al. (2018) [58] also found that some filaments just experienced partial eruption. Except for the individual filament eruption, two adjacent filaments can be observed to erupt in succession. This phenomenon is called sympathetic events. Li et al. (2017) [59] reported that the eruption of one of the two adjacent filaments is likely to be triggered by kink instability, while the weakening of overlying magnetic fields due to the magnetic reconnection at an X-point between the two filament systems might result in the onset of the other filament eruption. The two filament eruptions are confined due to the strong constraints of the overlying large-scale quiescent filament. Sometimes, the change of the overlying surrounding magnetic structure can also work as the trigger of the filament eruption. Zhou et al. (2017) [60] found that the rising of a flux rope resulted in magnetic reconnection between two external magnetic loops. The external magnetic reconnection weakens the constraining effect of the overlying field and leads to the further rising of the filament. Finally, the filament erupted due to torus instability.

Yang and Chen (2019) [61] presented state-of-the-art highresolution H $\alpha$  and SDO/AIA multi-wavelength images tracking the evolution of an active region before, during, and after a blowout filament eruption that produced a CME. They found unambiguous observational evidences indicating that multiple interactions occurred between the emerging magnetic fields and the coronal loops overlying the filament. The interactions resulted in the formation of a sigmoid structure and the eruption of the filament. After the filament eruption, the emerging magnetic fields continued to reconnect with the remaining filament channel, leading to the reformation of the filament. In the photosphere, flux cancellation associated with  $\delta$  sunspots formation occurred during the observation. These observations provide compelling evidence that tethercutting reconnection accompanying magnetic flux cancellation first triggered the blowout eruption and then continued to restructure the magnetic field for the reformation of the filament. The observations demonstrate previously proposed mechanisms for the triggering of CME/flare eruptions and for the building of sheared fields where filaments form.

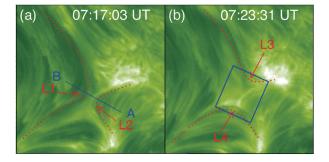
The aforementioned studies suggest that the trigger mechanism of solar filament eruptions may be diverse. Prediction of solar filament eruptions is also a very interesting topic.

### **3.3** Observational evidence and the role of small-scale magnetic reconnection in the solar eruptions

Magnetic reconnection is a very important process of energy release in Astrophysics. It is widely accepted that the free energy stored in the non-potential fields is released by magnetic reconnection during solar eruptions. There are several models of solar eruptions involving in magnetic reconnection. However, observational evidences have been relatively rare and typical only partial [62–66]. The role of magnetic reconnection in solar eruptions is still unclear.

#### 3.3.1 Observational evidence of small-scale magnetic reconnection

Observational features of large-scale magnetic reconnection were found in some solar eruptions. However, small-scale magnetic reconnection in solar activities is rarely reported. Using the high spatial and temporal H $\alpha$  line center images observed by the NVST, the magnetic reconnection between two groups of chromospheric fibrils was discovered by Yang et al. (2015) [67] for the first time. L1 and L2 denote the two groups of the chromospheric fibrils (see Figure 10(a)). After the magnetic reconnection, the two newly formed groups of the chromospheric fibrils were labelled as L3 and L4 (see Figure 10(b)). The thickness and length of the reconnection region are measured to be about 420 km and 1.4 Mm, respectively (see the box in Figure 10(b)). Meanwhile, it is found that there were two stages of the magnetic reconnection. The first stage is a slow phase with a duration of more than several tens of minutes at the beginning of the magnetic reconnection. The second stage is a rapid step lasting for only about three minutes. The inflows along the slit AB in Figure 10(a)range from 1.9 to 15.4 km s<sup>-1</sup> and the outflows range from 2.4 to 25.8 km s<sup>-1</sup>. The down-flow along the newly formed cusp structures ranges from 40 to 80 km s<sup>-1</sup>.



**Figure 10** (Color online) A small-scale magnetic reconnection observed by the NVST. L1 and L2 indicate the two groups of the chromospheric fibrils before the magnetic reconnection. L3 and L4 indicate the newly formed chromospheric fibrils after the magnetic reconnection. The box indicates the reconnection region. Image reproduced with permission from Yang et al. (2014), copyright by AAS.

Yang et al. (2016) [68] found another case of magnetic reconnection between two groups of loops with an X-shaped configuration in the chromosphere. They found that the antidirected loops approached each other and then began to reconnect. New loops formed and stacked together. They also found that a set of loops suddenly retracted toward the balance position and performed an overshoot movement, leading to a convergent oscillation with the mean period of about 45 s. Due to the sudden retraction, other lower loops were pushed outward and performed an oscillation with the period of about 25 s.

#### 3.3.2 Release of twist in a filament by magnetic reconnection

Except for the small-scale magnetic reconnection mentioned above, another magnetic reconnection event was also found during a filament eruption. Using high temporal and spatial resolution H $\alpha$  observations from the NVST, and supplemented by UV, EUV images and vector magnetograms observed by SDO, the X-ray data from Hinode and GOES, the magnetic reconnection occurring in the active region 12178 on October 3, 2014 was studied by Xue et al. (2016) [65] in detail. The magnetic reconnection occurs between the filament threads and the ambient chromospheric fibrils (see Figure 11(a) and (b)). For the first time, it was observed that the twist in the filament is released quickly by the magnetic reconnection. In this event, an unprecedented comprehensive set of observational evidences are detected, including reconnection inflows and outflows, current sheet, two hot cuspshaped structures, newly formed magnetic loops, shrinking of the loops, and so on (see Figure 11(c)-(f)). The estimated reconnection rate between 0.08 and 0.6 is consistent with the values of the theoretical expectations. In addition, the change of magnetic structures is supported by the change in the structure of the extrapolated nonlinear force-free field and by a data-constrained magnetohydrodynamic simulation (see Figure 11(g)–(i)). These results further confirm the occurrence of the magnetic reconnection during the filament eruption and the release process of the twist in the filament by the magnetic reconnection.

### 3.3.3 Interchange magnetic reconnection between a filament and nearby open fields

Interchange reconnection is defined as the interaction between closed and open fluxes [69]. It is considered to be a primary mechanism to transfer the dense mass, heat, and momentum to solar wind [70, 71]. Though this conception was proposed for many years, the observation evidences of interchange reconnection are very few (Kong et al. 2018) [72]. Using high resolution H $\alpha$  data from the NVST, Zheng et al. (2017) [73] studied a confined erupting filament and found evidence of interchange magnetic reconnection between the filament and adjacent open fields on December 24, 2014. During the filament eruption, the untwisting and rising motion of the filament bring their embedded closed field lines to meet with the adjacent oppositely open fields, and some of the closed field of the filament were opened. Moreover, the newly formed open field lines were observed to be rooted at negative polarities around one footpiont of the filament and the newly formed closed loops.

#### 3.3.4 Magnetic reconnection between a twisted arch filament system and coronal loops

It is well accepted that magnetic reconnection plays an important role in the solar activities. Similar to the results of Xue et al. (2016) [65] and Zheng et al. (2017) [73], Huang et al. (2018) [74] found a reconnection event occurred between the threads of a twisted arch filament system (AFS) and coronal loops. Their observations reveal that the relaxation of the twisted AFS drives some of its threads to encounter the coronal loops, providing inflows of the reconnection. Magnetic reconnection can transfer the free magnetic energy in the twisted magnetic system into the upper solar atmosphere.

#### 3.3.5 Magnetic reconnection between two active-region filaments

Using high spatial and temporal resolution  $H\alpha$  data from the NVST and simultaneous observations from the SDO, Yang et al. (2017) [75] investigated a rare event of the interaction between two filaments (F1 and F2) in AR NOAA 11967 on January 31, 2014 (see Figure 12(a)). The adjacent two filaments were almost perpendicular to each other. The interaction was driven by the movement of F1 and began when the two filaments collided with each other. During their interaction, the threads of F1 continuously slipped from the northeast to the

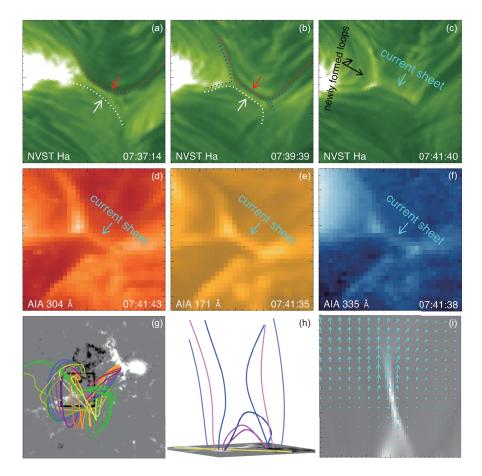
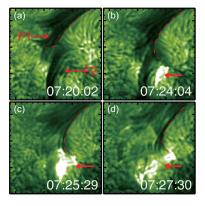


Figure 11 (Color online) The reconnection process of an active-region filament and the chromospheric fibrils observed by the NVST. The filament threads and chromospheric fibrils are indicated by arrows respectively in panels (a) and (b). The current sheet forming in magnetic reconnection is marked by the arrows in panels (c)–(f). The results of a data-constrained magnetohydrodynamic simulation reproducing the magnetic reconnection during the filament eruption are presented in panels (g)–(i). Image reproduced with permission from Xue et al. (2016), copyright by NPG.

southwest, and were accompanied by the brightenings at the junction of two filaments and the northeast footpoint of F2 (Figure 12(b)–(c)). At the same time, bright material initiated from the junction of two filaments was observed to move along F1. The magnetic connectivities of F1 were found to change after their interaction (see Figure 12(d)). These observations suggest that magnetic reconnection was involved in the interaction of two filaments and resulted in the eruption of F2.



#### 3.3.6 Oscillatory magnetic reconnection

Xue et al. (2019) [76] found that a small-scale oscillatory magnetic reconnection occurred in active region 11800 from July 24 to 25, 2013. It consists of four relatively independent magnetic reconnections which last for about 48, 158, 275, and 340 min, respectively. Moreover, the inflow/outflow regions of the previous magnetic reconnection are transformed into the outflow/inflow regions of the following magnetic reconnection, and a current sheet was observed in each magnetic reconnection. Four current sheets along two nearly

**Figure 12** (Color online) Interaction between two active-region filaments observed by the NVST. F1 and F2 indicate the two active-region filaments. The dashed lines in (b), (c) outline the newly formed filament threads and the red arrows indicate the brightening caused by the interaction of the two filaments. Image reproduced with permission from Yang et al. (2017), copyright by AAS.

perpendicular directions are formed alternately. Therefore, the oscillatory reconnection experiences two cycles, and the periods of the two oscillations are 206 and 615 min, respectively, which are much longer than those in the previous results. In addition, a small-scale flux rope is observed to form and disappear, which is caused by the oscillatory magnetic reconnection.

More and more evidences of magnetic reconnection are provided by the NVST observation. It is a fact that magnetic reconnection plays an important role in solar eruptions [55, 77]. However, how the particles are accelerated in the current sheet is still obscure. The simulations showed that there are many magnetic islands in current sheet [78–80]. So far, there is no reliable observational evidence that magnetic islands exist in current sheets. Higher resolution observational data are required to reveal this phenomenon.

#### 3.4 Small-scale eruptive activities

Following the improvement of the observation technology, many small-scale eruptive activities in the solar atmosphere are observed clearly. For example, there are various kinds of jet-like structures like spicules, chromospheric jets, X-ray/UV/EUV jets, etc, and sudden brightening in the extended H $\alpha$  wings named Ellerman bombs in the different solar atmosphere [81]. Studying on the triggering mechanisms of these phenomena is one of the important subjects in solar physics.

#### 3.4.1 Ellerman bombs

Ellerman bombs (EBs) [81] are defined as the two emission bumps in both wings of chromospheric lines, such as  $H\alpha$  6563 Å and Ca II 8542 Å lines. They are smallscale brightening events in the solar lower atmosphere. Recent high-resolution observations by the Interface Region Imaging Spectrograph (IRIS) [82] revealed that absorption lines (singly ionized or neutral) superimposed on the greatly broadened transition region lines [83] are typically found in small-scale compact bright regions from slit-jaw images (SJIs) taken with the 1400 Å and 1330 Å filters. This phenomenon indicates that there is local heating of the photosphere or lower chromosphere to  $8 \times 10^4$  K under the assumption of collisional ionization equilibrium named UV bursts. Is there a connection between UV bursts and EBs? Tian et al. (2016) [84] identified 10 UV bursts by using the IRIS observations. They found that three UV bursts are unambiguously and three UV bursts are possibly connected to EBs, which show intense brightening in the extended  $H\alpha$  wings without leaving an obvious signature in the H $\alpha$  core (see Figure 13). Their findings support the formation of these bombs in the photosphere and EBs can be heated much more efficiently than previously thought. They also found that the Mg II k and h lines can be used to identify EBs similarly to  $H\alpha$ , which opens a promising new window for EB studies. The remaining four UV bursts have no connection to EBs. These UV bursts were suggested to form in a higher layer, possibly in the chromosphere.

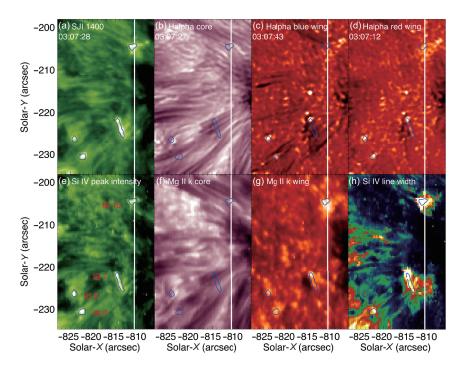
Chen et al. (2017) [85] have carried out a statistical study on the visibility of EBs in the AIA 1700 Å images by using NVST H $\alpha$  core and off-band images. There are 145 EBs that were identified from H $\alpha$  wing images, in which 74 EBs (about 51%) can be clearly identified as bright points (BPs) in the AIA 1700 Å images. On the contrary, they have identified 125 BPs in the 1700 Å images, in which 66 BPs (53%) corresponds to EBs in the H $\alpha$  wing images. These results show that most of the small-scale, compact, and transient brightenings in the AIA 1700 Å images can be identified as EBs.

Up to now, the heating mechanism of EBs is still controversial. Based on the connection between UV bursts and EBs, Tian et al. (2016) concluded that some EBs could be heated to  $1 \times 10^4 - 8 \times 10^4$ , which is much higher than previously thought. To address this issue, Ni et al. (2016) [86] carried out MHD simulations and found that reconnection around the temperature minimum region (TMR) can indeed heat some of the materials to such high temperature if the plasma beta is low. Fang et al. (2017) [87] have performed non-LTE calculations  $H\alpha$  and Ca II 8542 Å line profile and found that the heating of EBs cannot reach 10000 K. Consequently, the observations and theoretical studies are not consistent with each other. The heating of EBs is an important problem for the understanding of magnetic reconnection in the partially ionized lower solar atmosphere. In the future, observations and theoretical studies are required.

#### 3.4.2 Mini-filaments and jets

Due to the high resolution H $\alpha$  observation, Hong et al. (2017) [88] studied a mini-filament eruption in active region NOAA 12259 on January 14, 2015. The mini-filament experienced twice eruptions (a partial and a full eruption). The first eruption was associated with a fan-spine jet and the second eruption was associated with a blowout jet. It is worth mentioning that the blowout jet is accompanied by an interplanetary type-III radio burst. This is the first time to report that the mini-filament eruption is correlated with the interplanetary type-III radio burst. Tian et al. (2017) [89] reported two successive two-sided loop jets caused by magnetic reconnection between two adjacent filamentary threads, which were observed by the NVST at the southern periphery of active region NOAA 12035 on April 16, 2014.

Previous modeling and observations of jets support the occurrence of jets due to magnetic reconnection between emerging bi-polar and their overlying horizontal magnetic fields. The observations of Tian et al. (2017) [89] extend



**Figure 13** (Color online) IRIS bomb and EB bomb observed by the NVST and IRIS. IRIS/SJI 1400 Å image, and NVST H $\alpha$  core and wing (–1 Å and +1 Å) images taken around 03:05:38 UT. The dark filamentary structures in the H $\alpha$  wings, especially in the blue wing, are chromospheric spicules which could affect the detection of EBs. (e)–(h) Images of the Si IV 1393.755 Å intensity, Mg II k core and wing (sum of –1.33 Å and +1.33 Å), and Si IV 1393.755 Å line width. The locations of the IBs are marked by overplotting contours of the Si IV 1393.755 Å peak intensity. The white line in each panel indicates the slit location at 03:07:28 UT. Image reproduced with permission from Tian et al. (2016), copyright by AAS.

the jet model, i.e., a "tether-cutting model", which is often used to explain the formation of large-scale filaments or flux ropes. Similar to the observations of Hong et al. (2017) [88], Shen et al. (2017) [90] also studied a blowout jet that experienced two distinct ejection stages at the southern periphery region of active region NOAA 12035 on April 16, 2014 (see Figure 14). In the first stage, a small jet is triggered by the magnetic reconnection between the rising confining loops and the overlying open fields. A blowout jet is triggered by magnetic reconnection between the rising confining loops and the overlying open field in the second stage. They also observed periodic metric radio type III bursts at the every beginning of the two stages. In addition, two-side-loop jets were found by Zheng et al. (2018) [91] by using the observation of the NVST in active region NOAA 12681 on September 23, 2017. These jets were triggered by magnetic reconnection between the emerging loops and the overlying horizontal filament threads.

Using the data from SDO and NVST, a circular filament eruption associated with the formation of jet under a nonaxisymmetric fan spine configuration was studied. It is found that the non-axisymmetric fan spine configuration can lead to a non-radial eruption of filament toward magnetic null point. In this process, the eruptive filament interacted with the overlying magnetic loops, resulting in the disintegration of the filament, which caused the formation of a blowout jet [92].

#### 3.5 Formation and trigger mechanism of solar flares

The flares associated without CMEs are termed "confined flares" [93]. Using the H $\alpha$  observations from the NVST, Yang et al. (2014) [94] investigated the fine structures of three confined flares on October 12, 2013 (see Figure 15). All the three confined flares took place successively at the same location in active region NOAA 11861. They have similar morphologies and are termed as homologous confined flares. Supplemented by the SDO observations, many large-scale coronal loops above the confined flares are clearly found in multiwavelength observations. A reconnection between the H $\alpha$ fibrils of two emerging dipoles resulted in the change of the  $H\alpha$  fibril configuration. The reconnection also occurs between a set of emerging H $\alpha$  fibrils and a set of pre-existing large-scale loops, forming a typical three-legged structure. Liu et al. (2018) [95] also found that a C-class flare was triggered by magnetic reconnection between an emerging magnetic bi-pole region and the super-penumbral fibrils. These observations confirmed that the magnetic reconnection between the emerging loops and the pre-existing loops triggers solar flares and the large-scale overlying loops prevent the flares from erupting.

Based on multi-wavelength observations from the NVST and SDO, Xu et al. (2019) [96] found that the sunspot structures and their magnetic fields suffered significant transfor-

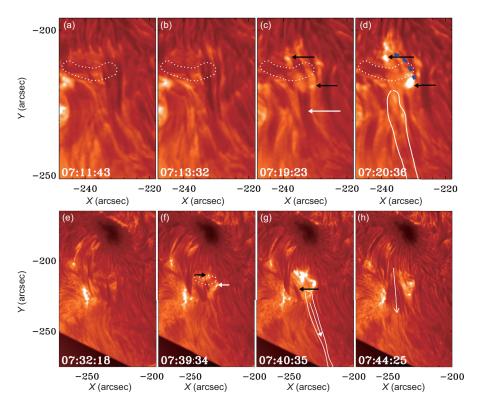
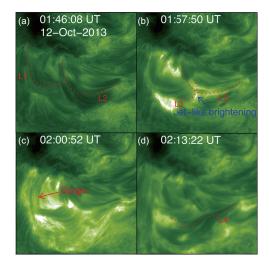


Figure 14 (Color online) Two blowout jets observed by the NVST. (a)–(d) Evolution of the first ejection stage of the blowout jet; (e)–(h) evolution of the second ejection stage of the blowout jet. Image reproduced with permission from Shen et al. (2018), copyright by AAS.

mations during an M-class flare. One penumbral segment decayed with partial fields transforming into umbra fields when swept by one flare ribbon. At the same time, an adjacent penumbral segment expanded permeating the granular area with the horizontal fields enhancing along the flaring PIL. Their observations support the magnetic implosion idea that the flare eruptions in the upper atmosphere could result in the magnetic field rearrangement on the surface of the sun.



**Figure 15** (Color online) Process of a confined flare observed by the NVST. The dashed lines display the loops identified from H $\alpha$  images. Image reproduced with permission from Yang et al. (2014), copyright by AAS.

In recent years, circular-ribbon flares have caught a lot of attentions due to challenging the traditional model of tworibbon flares. Using the NVST H $\alpha$  data, Li et al. (2018) [97] studied a mini-filament erupting in association with a circular ribbon flare on March 17, 2014. The circular flare ribbon was observed around the filament at the onset of the eruption. After the filament activation, its eruption looked like the formation of a surge, which ejected along one end of a large-scale closed coronal loops with a curtain-like shape. They carried out a potential field extrapolation and found that a null point was located above the mini-filament. They suggested that the null point reconnection may facilitate the eruption of the filament to produce the circular-ribbon flare. Xu et al. (2017) [98] reported two homologous circular-ribbon flares associated with two filament eruptions observed by the NVST in active region NOAA 11991 on March 5, 2014. Two small-scale filaments erupted in sequence associated with two homologous circular-ribbon flares and displayed an apparent writhing motion. Moreover, the post-flare loops and subsequent helical structures are clearly seen during their eruption. These observations imply that the small-scale filaments have twisted magnetic structures. In addition, Song et al. (2018) [99, 100] found that the eruptions of the mini-filaments are often associated with white-light flares.

Recent investigations reveal that the active regions with rotating sunspots have relatively high solar flare productivity [101]. For example, active regions, e.g., NOAA 10486, NOAA 10930, NOAA 11158, produced many flares and coronal mass ejections (CMEs). However, the relationship between evolution of rotating sunspots and solar activities is still unclear.

The triggering process of solar flares is not very clear. In order to address this issue, Yan et al. (2018) [102] used HMI data to study the occurrence of two successive X-class flares and two coronal mass ejections (CMEs) triggered by shearing motion and sunspot rotation in active region NOAA 12673 on September 6, 2017. From the evolution in the continuum intensity images and the line-of-sight magnetograms, it is very obvious that there was a shearing motion between the two sunspots from September 5 to the onset of the second Xclass flare on September 6. Moreover, the sunspot with negative polarity rotated around its umbral center and the other sunspot with positive polarity also exhibited a slow rotation (see Figure 16(a) and (b)). The structure of the sunspot with negative polarity is very complex seen from high resolution TiO images in Figure 16(c) and (d). Before the occurrence of the first X-class flare, it is found that there is a flux rope that forms along the polarity inversion line between the two sunspots (see Figure 17). The S-shaped structure seen from the EUV observation corresponds to this flux rope. Moreover, the first X-class flare is associated with the eruption of the flux rope.

The sunspot with negative polarity at the northwest of active region also began to rotate counter-clockwise before the onset of the first X-class flare, which is related to the formation of the second S-shaped structure in the EUV observation.

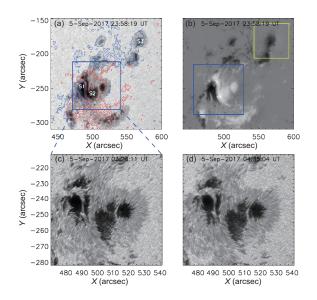


Figure 16 (Color online) Active-region NOAA 12673 in the continuum intensity image, line-of-sigh magnetogram, and TiO images. (a) The continuum intensity image; (b) line-of-sigh magnetogram; (c)–(d) TiO images. Image reproduced from Yan et al. (2018), copyright by AAS.

The eruption of the second S-shaped structure produced the second X-class flare and a CME. The successive formation and eruption of two S-shaped structures were closely related to the counter-clockwise rotation of three sunspots. The existence of a flux rope is found prior to the onset of two flares by using non-linear force free field extrapolation based on the vector magnetograms observed by SDO/HMI. These results suggest that shearing motion and sunspot rotation play an important role in the buildup of the free energy and the formation of flux ropes in the corona.

#### 4 Summary and prospect

The NVST is a very powerful solar telescope to observe the photosphere and chromosphere of the sun. High spatial and temporal observational data from the NVST are very suitable for researchers to study small-scale solar activities. Based on the these data, many new results were obtained.

By using TiO images, the photospheric BPs were studied in detail. The area, equivalent diameter, lifetime, horizontal velocity, diffusion index, motion range, and motion type of BPs were determined by using a new method named the Laplacian and morphological dilation algorithm [11,103]. Moreover, the relationships between the magnetic field intensity and the parameters of BPs were investigated. The diameters, intensities, lifetimes and velocities of UDs were found to depend on the surrounding magnetic field. In addition to the typical 5-min and 3-min oscillations in the photosphere and chromosphere of the sun, one minute oscillation at different heights above sunspot umbra was discovered.

Due to the high resolution chromospheric data (H $\alpha$  line center and off-band images), the fine structures and formation mechanisms of solar prominences or filaments were studied in detail. In addition to the counter-streaming flows, the kelvin-Helmholtz instability was identified in quiescent prominences. Moreover, the parallel tube-shaped structures in the barb were found during the evolution of the prominence. Magnetic reconnection between two groups of chromospheric fibrils driven by sunspot movement, shearing motion and sunspot rotation, and tether-cutting magnetic reconnection were found to be the formation mechanisms of solar active-region filaments. Note that there are two mechanisms on the filament channel formation, i.e., surface effects that reconfigure pre-existing coronal fields and magnetic emergence of a horizontal flux rope (see the review of Mackay et al. (2010) [104]. The results on the filament formation obtained using NVST data belong to the former mechanism. Though many works reported that solar eruptions are closely related to the shearing motion and sunspot rotation [105-107], the filament formation related to a small rotating pore was re-

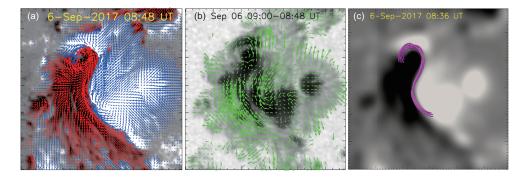


Figure 17 (Color online) A part of the sunspot groups in active-region NOAA 12673. (a) Vector magnetogram; (b) velocity field derived from vector magnetogram by using DAVE method; (c) a flux rope. Note that the field lines are the selected field lines extrapolated by using NLFFF method. Image reproduced from Yan et al. (2018), copyright by AAS.

ported for the first time.

The observational evidences of small-scale magnetic reconnection in solar eruptions are confirmed. Small-scale magnetic reconnection between the two groups of the small chromospheric fibrils and the release of twist in a filament by magnetic reconnection were discovered for the first time. Interchange magnetic reconnection between a filament and nearby open fields, magnetic reconnection between a twisted arch filament system and coronal loops, oscillatory magnetic reconnection, and magnetic reconnection between two active-region filaments were observed. The relatively large field-of-view of the NVST provides the opportunity to observe these different types of magnetic reconnection between the different magnetic structures. In comparison to the previous studies, all of the magnetic reconnection events observed by the NVST are on-disc and small-scale events.

Through analyzing the relationship between UV bursts and EBs, the results are inclined to support that the formation of EBs was in the photosphere and EBs can be heated much more efficiently than previously thought. Moreover, the UV bursts were suggested to form in a higher layer, possibly in the chromosphere.

The physical connections between the eruptions of minifilaments and solar flares/jets were studied in detail. Using NVST data, the mini-filament eruption correlated with the interplanetary type-III radio burst was reported for the first time. Besides, two successive two-sided loop jets caused by magnetic reconnection between two adjacent filamentary threads were observed.

By using  $H\alpha$  line center images, a series of confined flares were studied in detail. It is evidenced that magnetic reconnection between two groups of chromospheric fibrils led to the release of free energy in a series of successive confined flares. Part of magnetic fields in the penumbral segment transforming into umbra fields were found during another M-class confined flare. This result supports that the eruptions in the upper atmosphere can result in the magnetic field rearrangement in the lower atmosphere.

In the future, magnetograms with high spatial and temporal resolution will be provided at Fuxian Lake Observatory. High resolution magnetograms will be taken to study the magnetic field evolution of small-scale features in the photosphere such as BPs, foot-points of mini-filaments, EBs, and so on. Moreover, the spectrometers will start routine observations. According to the requirement of researchers, several observational models will be set up for users to choose. Note that these works are in progress. It is expected that more new results can be obtained due to the high resolution photospheric and chromospheric observations in combination with high resolution photospheric magnetograms.

We would like to thank the NVST, SDO/AIA, SDO/HMI teams for highcadence data support. This work was supported by the National Natural Science Foundation of China (Grant Nos. 11873087, 11633008), the Youth Innovation Promotion Association CAS (Grant No. 2011056), the Yunnan Talent Science Foundation of China (Grant No. 2018FA001), the Project Supported by the Specialized Research Fund for State Key Laboratories, and the grant associated with project of the Group for Innovation of Yunnan Province.

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