



# Application of Moiré Interferometry to the Characterization of Orthotropic Materials in the Antisymmetric Configuration using the Virtual Fields Method

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## Abstract

Moiré interferometry is an effective full-field deformation measurement technique and has been utilized for mechanical behavior analysis of materials and structures. For isotropic materials, Moiré patterns can be obtained by performing standard tests, such as, tensile and bending tests, to calculate the displacement and strain. Then, the mechanical properties can be characterized. However, standard tests are not sufficient to characterize the mechanical parameters of anisotropic materials due to the complexity of their material properties. Thus, in this work, Moiré interferometry was combined with the Virtual Fields Method to obtain the four in-plane elastic constants ( $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$ ) of orthotropic materials in the form of a diametrically compressed disk. Firstly, according to finite element method simulation results, optimized parameters can be achieved when the principal direction of the material does not coincide with the loading direction, making the loading configuration antisymmetric. Therefore, Moiré interferometry experiment was simulated to demonstrate the feasibility of measurement in the antisymmetric configuration. Finally, the  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$  and  $Q_{66}$  values of a unidirectional carbon fiber composite were measured in a real Moiré interferometry experiment using the proposed method, yielding results that agreed closely with those obtained using the strain gauges.

**Keywords** Moiré interferometry · Antisymmetric configuration · Virtual fields method · Material characterization · Orthotropic material

## Introduction

Moiré interferometry [1] is a widely used optical interference technique for full-field measurements. It is well known for its high accuracy (wavelength scale) and insensitivity to out-of-plane displacements, and has been successfully applied in many fields, such as, material characterization [2–5], fracture mechanics [6], residual stress [7] and so on [8]. In terms of material characterization, Moiré interferometry has the advantages of full field measurement and high accuracy, and efforts have been made in many fields, such as, welding [2], composites [3, 4], concrete materials [5] and so on.

To characterize the mechanical properties of anisotropic materials, traditional methods involving standard tests, such as, tensile and bending tests, are not suitable because of the

complexity of the mechanical properties. Therefore, in this study, the virtual fields method (VFM) [9–11] was combined with Moiré interferometry to measure the four in-plane elastic constants ( $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$ ) of an orthotropic material. The VFM is able to inverse all of the in-plane elastic constants from the full-field strain simultaneously and requires less computational time than other methods of solving the inverse problem, such as, the finite element model updating technique [12], the constitutive equation gap method [13], and deformation-pattern-based digital speckle correlation [14].

Many efforts have been made to characterize the mechanical properties of orthotropic materials by combining the VFM and full-field measurement techniques [4, 15–20]. According to some research [17–19], the off-axis angle (the angle between the principal axis of the material and the coordinate axis) significantly influences the accuracy of the inversed constants. Therefore, a diametrically compressed disk can be used as the loading configuration together with the VFM to facilitate adjustment of the off-axis angle to obtain optimal values in the characterization.

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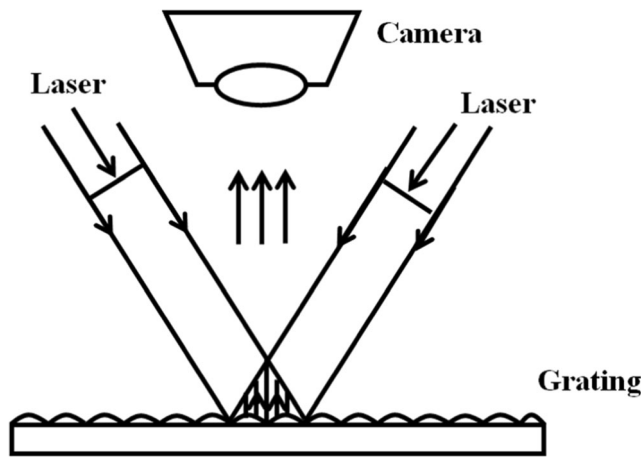


Fig. 1 Schematic diagram of Moiré interferometry

Notably, when the principal direction of the material does not coincide with the loading direction, the deformation field will be antisymmetric. However, in Moiré interferometry, the fringe patterns generated by real deformations and rigid-body rotations are coupled when the displacement fields are antisymmetric. Consequently, further study is necessary to discuss the feasibility of Moiré interferometry application in such situations.

In this work, firstly, the strain field of the disk configuration obtained from the finite element method (FEM) simulation was used for the VFM program to inverse the four elastic constants ( $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$ ) with high accuracy. Then, a Moiré interferometry experiment was simulated to demonstrate the feasibility of combining VFM and Moiré interferometry. The results of the simulated experiments illustrate that Moiré interferometry can be employed to measure strain fields for the VFM with sufficient accuracy when the displacement field is antisymmetric. Finally, the values of  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$  and  $Q_{66}$  for a unidirectional carbon fiber composite were measured in a real experiment conducted using a Moiré interferometer, which yielded results very similar to those obtained using the strain gauges.

## Theory

### Moiré Interferometry

In Moiré interferometry, high-frequency cross-line gratings are fabricated on the surface of the specimen.

Upon loading, the gratings deform with the specimen. Then, two laser beams are incident on the deformed gratings at symmetric angles to form virtual reference gratings. The specimen gratings subsequently interfere with the virtual reference gratings to generate Moiré fringe patterns, which can be recorded by a camera, as shown in Fig. 1. [1]

The relationships between the number of the fringes and the displacement are as follows [1]:

$$\begin{cases} u(x, y) = \frac{N_x}{2f} \\ v(x, y) = \frac{N_y}{2f} \end{cases} \quad (1)$$

where  $u(x, y)$  and  $v(x, y)$  are the displacements along the  $x$  and  $y$  axes, respectively;  $N_x$  and  $N_y$  are the number of fringes along the  $x$  and  $y$  axes, respectively; and  $f$  is the initial frequency of the specimen gratings.

In this work, gratings with a frequency of 1200 lines/mm were used. Then one fringe indicates the displacement of 417 nm.

### Constitutive Equation of Orthotropic Materials

A brief introduction and some notations used in the constitutive equation of orthotropic materials are given in this section. For orthotropic materials, there are five in-plane engineering constants:  $E_1$ ,  $E_2$ ,  $\nu_{12}$ ,  $\nu_{21}$ , and  $G_{12}$  [21]. Considering that  $\nu_{12}/E_1 = \nu_{21}/E_2$ , there are four independent elastic constants in total. Then, to describe the relation between stress and strain, four elastic constants denoted as  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  are commonly used and the constitutive equation is as follows:

$$\begin{bmatrix} \sigma_1 \\ \sigma_2 \\ \tau_{12} \end{bmatrix} = \begin{bmatrix} Q_{11} & Q_{12} & 0 \\ Q_{12} & Q_{22} & 0 \\ 0 & 0 & Q_{66} \end{bmatrix} \begin{bmatrix} \varepsilon_1 \\ \varepsilon_2 \\ \gamma_{12} \end{bmatrix} \quad (2)$$

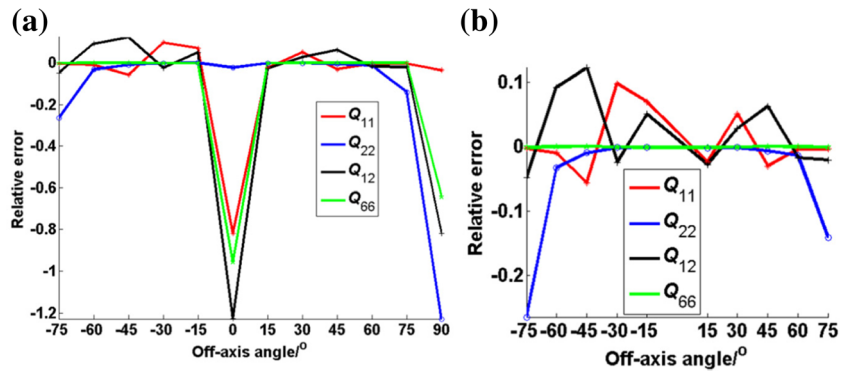
where  $(\sigma_1, \sigma_2, \tau_{12})$  is the stress tensor and  $(\varepsilon_1, \varepsilon_2, \gamma_{12})$  is the engineering strain tensor. The subscript '12' indicates the principal axis of the material.

The relationships between the four elastic constants and the engineering constants are then:

Table 1 ABAQUS model parameters

Model 1							
Diameter/m	Material	$E_1$ /Pa	$E_2$ /Pa	$\nu_{12}$	$G_{12}$ /Pa	Load/N	Element type
15	lamina	$70 \times 10^9$	$20 \times 10^9$	0.3	$5 \times 10^9$	50,000	CPS3
Model 2							
Diameter/m	Material	$E_1$ /Pa	$E_2$ /Pa	$\nu_{12}$	$G_{12}$ /Pa	Load/N	Element type
15	lamina	$130 \times 10^9$	$10 \times 10^9$	0.3	$5 \times 10^9$	80,000	CPS3

**Fig. 2** Relative errors of  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  in Model 1 with off-axis angles of **a**  $-75^\circ, -60^\circ, -45^\circ, -30^\circ, -15^\circ, 0^\circ, 15^\circ, 30^\circ, 45^\circ, 60^\circ, 75^\circ$ , and  $90^\circ$  **b**  $-75^\circ, -60^\circ, -45^\circ, -30^\circ, -15^\circ, 15^\circ, 30^\circ, 45^\circ, 60^\circ$ , and  $75^\circ$



$$\begin{cases} Q_{11} = \frac{E_1}{1-\nu_{12}\nu_{21}} \\ Q_{22} = \frac{E_2}{1-\nu_{12}\nu_{21}} \\ Q_{12} = \nu_{12}Q_{22} = \nu_{21}Q_{11} \\ Q_{66} = G_{12} \end{cases} \quad (3)$$

**Virtual Fields Method**

The virtual fields method is based on the principle of virtual work (PVW). When static and in-plane problems are considered and the body force is ignored, the simplified PVW equation can be written as follows [22]:

$$\int_S \sigma : \varepsilon^* dS = \int_{\partial S} \bar{\mathbf{T}} \cdot \mathbf{u}^* dl, \forall \mathbf{u}^* \text{KA} \quad (4)$$

where  $\sigma$  is the stress tensor,  $\varepsilon^*$  is the virtual strain tensor,  $S$  is the external surface,  $\bar{\mathbf{T}}$  is the traction on the boundary, and  $\mathbf{u}^*$

is the virtual displacement tensor. “KA” stands for “kinematically admissible”, which means that  $\mathbf{u}^*$  must satisfy the boundary condition.

For the in-plane orthotropic elastic model, introducing equation (3) into equation (4) yields:

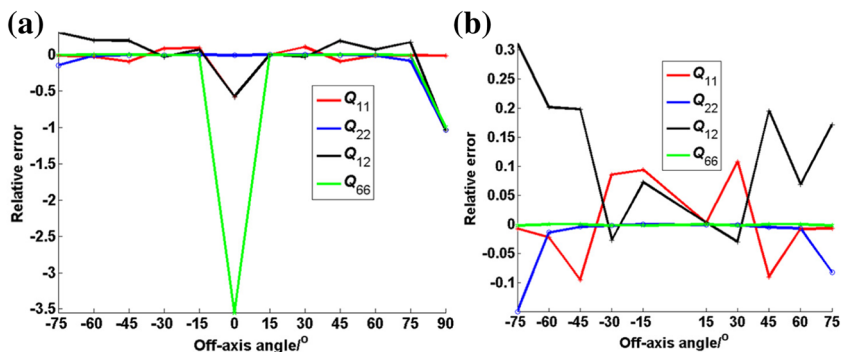
$$\begin{aligned} & Q_{11} \int_S \varepsilon_1 \varepsilon_1^* dS + Q_{22} \int_S \varepsilon_2 \varepsilon_2^* dS \\ & + Q_{12} \left( \int_S \varepsilon_1 \varepsilon_2^* dS + \int_S \varepsilon_2 \varepsilon_1^* dS \right) + Q_{66} \int_S \tau_{12} \gamma_{12}^* dS \\ & = \int_{\partial S} \bar{\mathbf{T}} \cdot \mathbf{u}^* dl \end{aligned} \quad (5)$$

where  $(\varepsilon_1^*, \varepsilon_2^*, \gamma_{12}^*)$  is the virtual strain tensor.

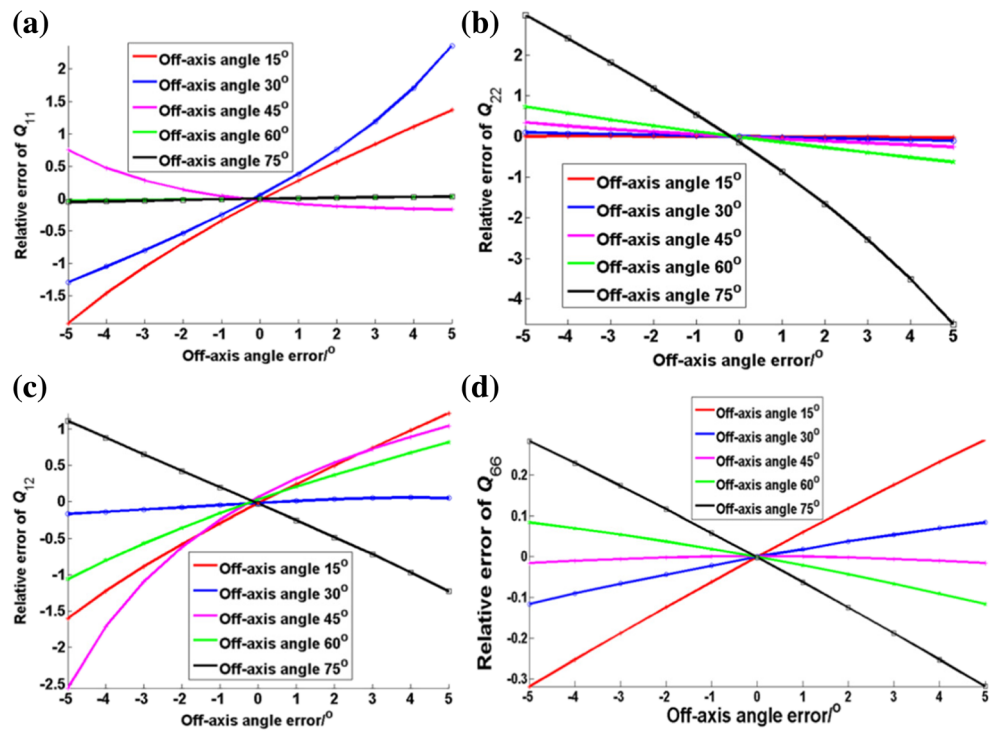
In this work, polynomials were used to define the virtual strain field  $(\varepsilon_1^*, \varepsilon_2^*, \gamma_{12}^*)$  in equation (5) [22]. Since there are four unknown parameters to be calculated, four groups of the virtual fields are required. Then equation (4) becomes:

$$\begin{bmatrix} \int_S \varepsilon_1 \varepsilon_1^{*(1)} dS \int_S \varepsilon_2 \varepsilon_2^{*(1)} dS \int_S \varepsilon_1 \varepsilon_2^{*(1)} dS + \varepsilon_2 \varepsilon_1^{*(1)} dS \int_S \gamma_{12} \gamma_{12}^{*(1)} dS \\ \int_S \varepsilon_1 \varepsilon_1^{*(2)} dS \int_S \varepsilon_2 \varepsilon_2^{*(2)} dS \int_S \varepsilon_1 \varepsilon_2^{*(2)} dS + \varepsilon_2 \varepsilon_1^{*(2)} dS \int_S \gamma_{12} \gamma_{12}^{*(2)} dS \\ \int_S \varepsilon_1 \varepsilon_1^{*(3)} dS \int_S \varepsilon_2 \varepsilon_2^{*(3)} dS \int_S \varepsilon_1 \varepsilon_2^{*(3)} dS + \varepsilon_2 \varepsilon_1^{*(3)} dS \int_S \gamma_{12} \gamma_{12}^{*(3)} dS \\ \int_S \varepsilon_1 \varepsilon_1^{*(4)} dS \int_S \varepsilon_2 \varepsilon_2^{*(4)} dS \int_S \varepsilon_1 \varepsilon_2^{*(4)} dS + \varepsilon_2 \varepsilon_1^{*(4)} dS \int_S \gamma_{12} \gamma_{12}^{*(4)} dS \end{bmatrix} \begin{bmatrix} Q_{11} \\ Q_{22} \\ Q_{12} \\ Q_{66} \end{bmatrix} = \begin{bmatrix} \int_{\partial S} \bar{\mathbf{T}} \cdot \mathbf{u}^{*(1)} dl \\ \int_{\partial S} \bar{\mathbf{T}} \cdot \mathbf{u}^{*(2)} dl \\ \int_{\partial S} \bar{\mathbf{T}} \cdot \mathbf{u}^{*(3)} dl \\ \int_{\partial S} \bar{\mathbf{T}} \cdot \mathbf{u}^{*(4)} dl \end{bmatrix} \quad (6)$$

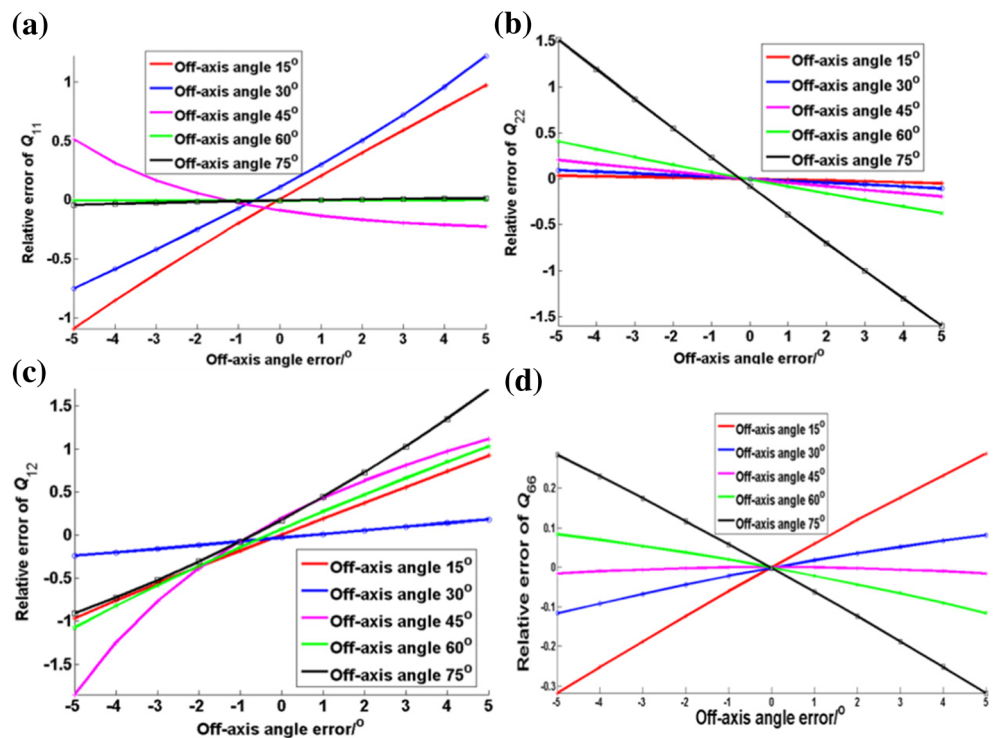
**Fig. 3** Relative errors of  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  in Model 2 with off-axis angles of **a**  $-75^\circ, -60^\circ, -45^\circ, -30^\circ, -15^\circ, 0^\circ, 15^\circ, 30^\circ, 45^\circ, 60^\circ, 75^\circ$ , and  $90^\circ$  **b**  $-75^\circ, -60^\circ, -45^\circ, -30^\circ, -15^\circ, 15^\circ, 30^\circ, 45^\circ, 60^\circ$ , and  $75^\circ$



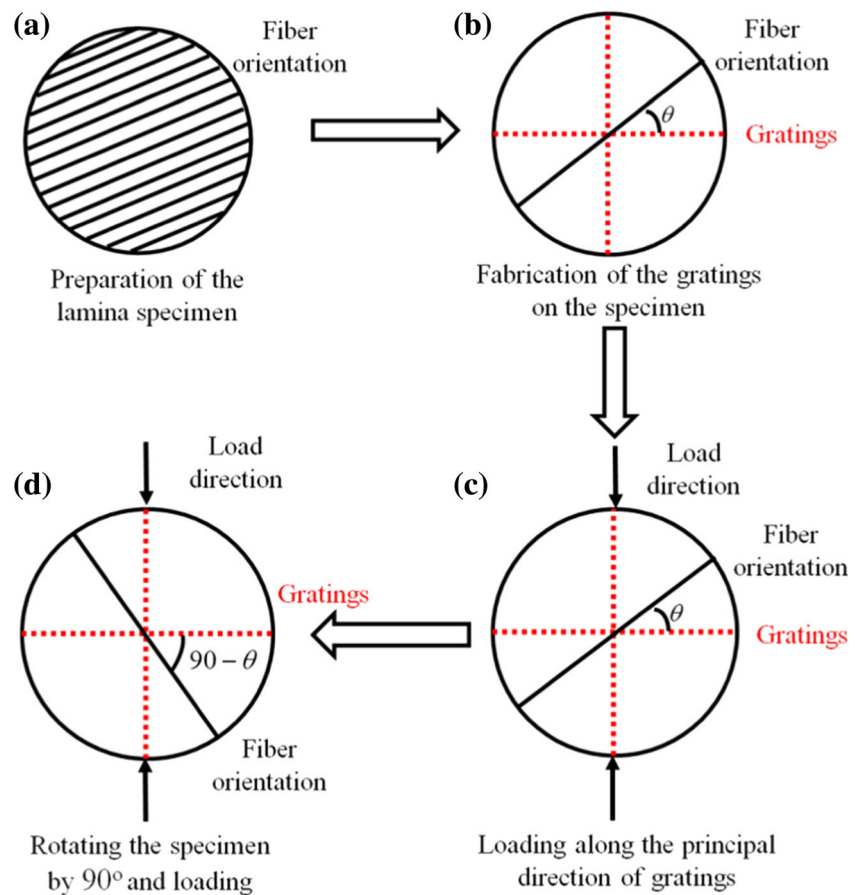
**Fig. 4** Relative errors of **a**  $Q_{11}$ , **b**  $Q_{22}$ , **c**  $Q_{12}$  and **d**  $Q_{66}$  in Model 1 with off-axis angles of  $15^\circ$ ,  $30^\circ$ ,  $45^\circ$ ,  $60^\circ$ , and  $75^\circ$  including off-axis angles errors ranging from  $-5^\circ$  to  $5^\circ$



**Fig. 5** Relative errors of **a**  $Q_{11}$ , **b**  $Q_{22}$ , **c**  $Q_{12}$  and **d**  $Q_{66}$  in Model 2 with off-axis angles of  $15^\circ$ ,  $30^\circ$ ,  $45^\circ$ ,  $60^\circ$ , and  $75^\circ$  including off-axis angles errors ranging from  $-5^\circ$  to  $5^\circ$



**Fig. 6** Flow chart of Moiré interferometry experiment

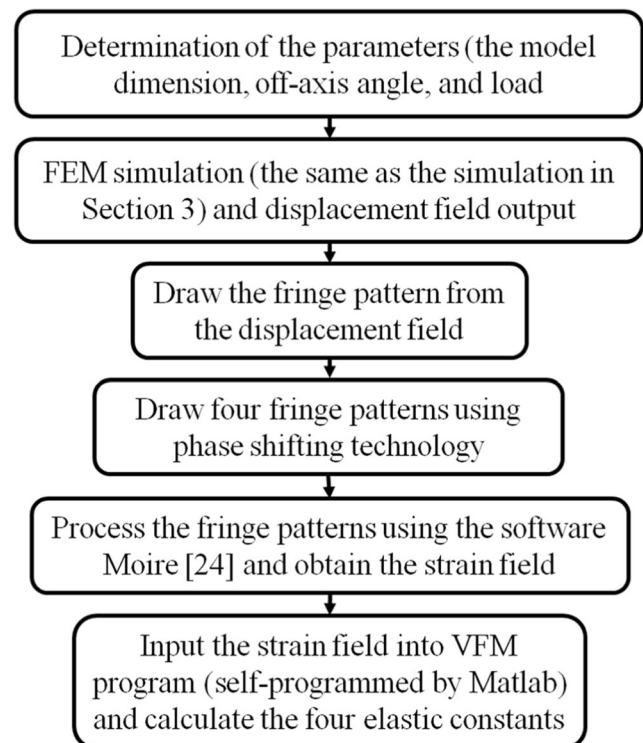


and  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  can be obtained directly by inverting equation (6).

### Numerical Simulation Analysis

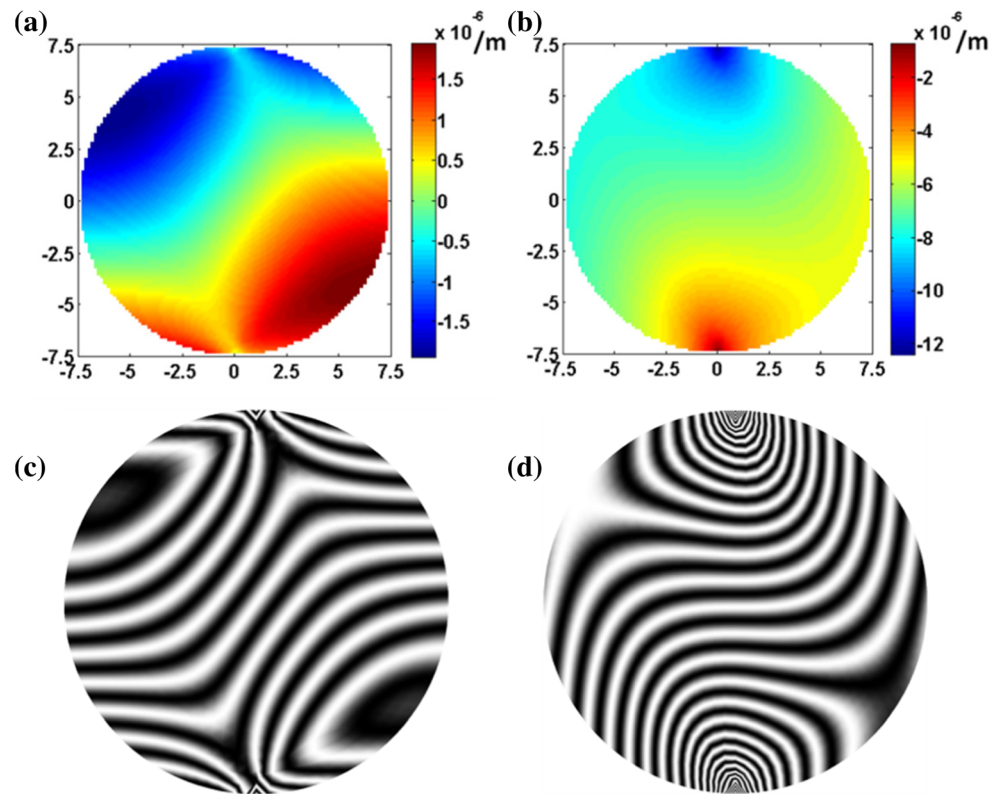
To verify the feasibility of the proposed method, numerical simulations were conducted using ABAQUS to obtain the strain field. The calculations were performed for two different models to make the results more convincing. Notably,  $E_1$  and  $E_2$  (defined in Section 2.2) differ significantly from one another for the unidirectional carbon fiber composite. Therefore,  $E_1$  was chosen to be larger than  $E_2$  in both models. The applied parameters are summarized in Table 1. In the FEM simulation, the top of the disk was subjected to a concentrated force and the bottom was fixed. A static step was applied for loading. CPS3 elements were used, and their total number was 1906.

Since the angle between the principal axis of the material and the coordinate axis (called the off-axis angle) significantly influences the measurement accuracy [17–19], cases with 12 different off-axis angles ( $-75^\circ$ ,  $-60^\circ$ ,  $-45^\circ$ ,  $-30^\circ$ ,  $-15^\circ$ ,  $0^\circ$ ,  $15^\circ$ ,  $30^\circ$ ,  $45^\circ$ ,  $60^\circ$ ,  $75^\circ$ , and  $90^\circ$ ) were simulated. The relative errors of the inversed parameters in Models 1 and 2 with those off-axis angles are illustrated in Figs. 2 and 3, respectively.

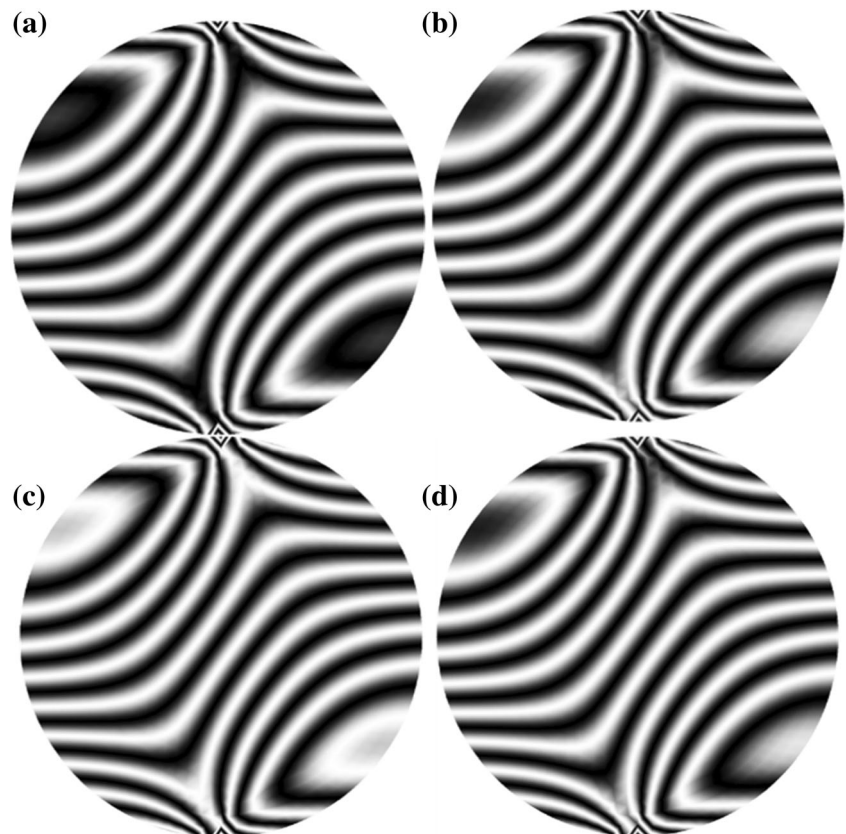


**Fig. 7** Flow chart of the simulated Moiré interferometry experiment

**Fig. 8** Displacement fields obtained from the FEM simulations: displacements along the **a**  $x$  axis and **b**  $y$  axis, and fringe patterns along the **c**  $x$  axis and **d**  $y$  axis generated from the displacement fields



**Fig. 9** Four fringe patterns generated using phase shifting technology



Figures 2 and 3 show the accuracies of the four elastic constants corresponding to different off-axis angles calculated by employing the VFM using the strain field in FEM simulation of both Models 1 and 2. As mentioned, the simulations were performed using 12 angles from  $-75^\circ$  to  $90^\circ$ , and the relative errors of  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  are presented. Obviously, the angles of  $0^\circ$  and  $90^\circ$  yielded highly inaccurate results. For clarity, curves without the  $0^\circ$  and  $90^\circ$  data are also provided.

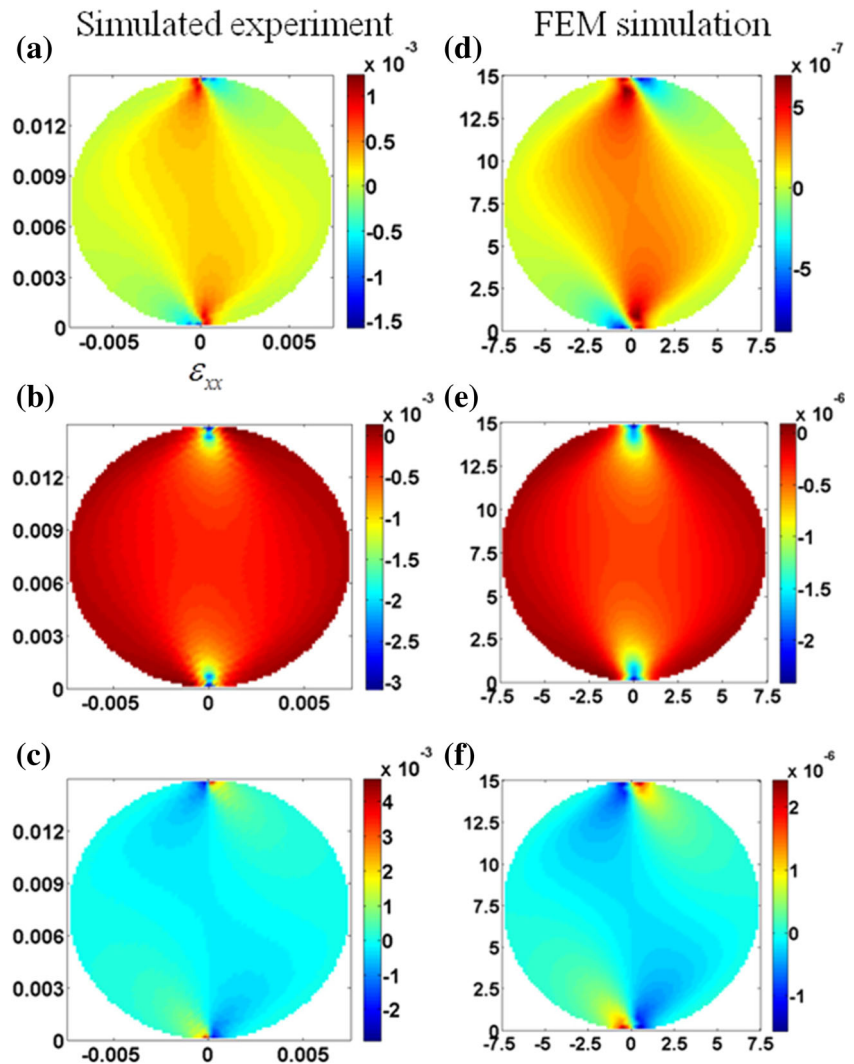
The results indicate that the different parameters have different optimal off-axis angles, which is of great importance for measurement accuracy. Therefore, to find the corresponding off-axis angles, further analysis was conducted by comparing the accuracy of each parameter with different off-axis angles. Practically, since the off-axis angle measurement error can hardly be avoided, the influence of this error was also considered. The results are presented in Figs. 4 and 5 for Models 1 and 2, respectively. Notably, for off-axis angles with the same magnitude and opposite signs, the inversed parameters are

almost the same. Thus, in the following part, only positive angles are discussed.

As shown in Figs. 4 and 5, for each parameter, the two models yielded similar optimal off-axis angles. For off-axis angles of  $60^\circ$  and  $75^\circ$ , the off-axis angle error has little influence on  $Q_{11}$ . The relative error of  $Q_{11}$  remains below 5% when the off-axis angle error is its maximum of  $5^\circ$ .  $Q_{22}$  exhibits similar behavior but has optimal off-axis angles of  $15^\circ$  and  $30^\circ$ . It is clear that the accuracy of  $Q_{66}$  is the highest, while that of  $Q_{12}$  is the lowest. When the off-axis angle is  $45^\circ$ , the relative error of  $Q_{66}$  remains below 2%. However, for  $Q_{12}$ , the minimum relative error is about 20% when the off-axis angle error is  $5^\circ$  and the off-axis angle itself is  $30^\circ$ . Even so, the measurement accuracy is still acceptable since it is also very difficult to obtain accurate  $Q_{12}$  value using traditional methods, such as, strain gauges according to the measured results presented in Section 5.

The analysis in this section demonstrates that the proposed approach using the VFM is feasible and has some advantages,

Fig. 10 Strain fields obtained from a-c the simulated experiment; d-f the FEM simulation



**Table 2**  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  values obtained from the simulated experiments

Model 1 (see Table 1)	$Q_{11}$ /GPa	$Q_{22}$ /GPa	$Q_{12}$ /GPa	$Q_{66}$ /GPa
off-axis angle: $30^\circ$	72.84	20.58	6.19	4.96
Model 1 (see Table 1)	$Q_{11}$ /GPa	$Q_{22}$ /GPa	$Q_{12}$ /GPa	$Q_{66}$ /GPa
off-axis angle: $-60^\circ$	71.60	20.26	6.54	4.97
Reference values for Model 1	71.85	20.53	6.16	5.00
Model 2 (see Table 1)	$Q_{11}$ /GPa	$Q_{22}$ /GPa	$Q_{12}$ /GPa	$Q_{66}$ /GPa
off-axis angle: $30^\circ$	-41.72	10.01	3.47	4.97
Model 2 (see Table 1)	$Q_{11}$ /GPa	$Q_{22}$ /GPa	$Q_{12}$ /GPa	$Q_{66}$ /GPa
off-axis angle: $-60^\circ$	129.52	10.01	3.36	4.97
Reference values for Model 2	130.91	10.07	3.02	5.00

such as, a simplified experimental process, low cost and reasonable accuracy. Furthermore, according to the FEM simulation results, it is obvious that the off-axis angle influenced the accuracies of  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  differently. For off-axis angles of  $60^\circ$  and  $75^\circ$ , highly accurate  $Q_{11}$  results were obtained and for  $Q_{22}$ , the corresponding off-axis angles were  $15^\circ$  and  $30^\circ$ . For  $Q_{12}$ ,  $30^\circ$  was the most suitable off-axis angle. Finally, for  $Q_{66}$ , the off-axis angle of  $45^\circ$  provided better results than the other angles, however, the accuracy of  $Q_{66}$  was also acceptable for other off-axis angles, such as,  $30^\circ$  and  $60^\circ$ , since  $Q_{66}$  was the least sensitive to errors as mentioned above.

## Moiré Interferometry Experiment Simulation

For the VFM, full field measurement is the key to improving the precision of the mechanical characterization. In this work, Moiré interferometry was combined with the VFM to perform a diametrical compression test on a disk. To demonstrate the feasibility and effectiveness of the proposed method, the Moiré interferometry experiment was simulated, and the detailed process and error analysis are described in this section.

## Simulated Experiments

The flow chart of the experiments is provided in Fig. 6. Firstly, the disk was prepared with orthotropic gratings (1200 lines/

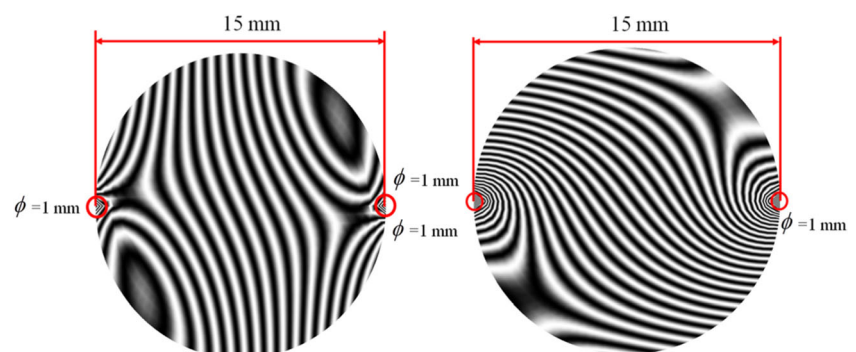
mm) on it. The angle between the principal direction of the gratings and the fiber orientation of the lamina (denoted by  $\theta$ ) was changeable because the gratings could be fabricated on the specimen along any direction. Then, since the displacement field can be measured only in the principal direction of the gratings in Moiré interferometry, the load should be applied in the principal direction of the gratings. In this case, the off-axis angle was  $\theta$ , as shown in Fig. 6c. After that, the specimen was rotated by  $90^\circ$ , and the off-axis angle became  $-(90-\theta)^\circ$ , as shown in Fig. 6d.

According to the analysis in Section 3,  $30^\circ$  was selected as  $\theta$  so that the displacement field corresponding to the off-axis angle of  $30^\circ$  and  $-60^\circ$  could be obtained. Highly accurate values of  $Q_{11}$  and  $Q_{22}$  were acquired using the off-axis angle of  $-60^\circ$  and  $30^\circ$ , respectively. For  $Q_{12}$ , an off-axis angle of  $30^\circ$  provided satisfying results, and for  $Q_{66}$ , both angles yielded little error.

To verify the feasibility of combining the VFM with Moiré interferometry, the above mentioned experiments were firstly conducted by performing numerical simulations.

Figure 7 provides a flow chart of the simulated experiments. For clarity, Model 2 (see Table 1) with an off-axis angle of  $30^\circ$  is used as an example. The displacement field and the corresponding fringe pattern of in this case are depicted in Fig. 8. Figure 8a and b show the displacements in the horizontal and vertical directions, respectively, and the fringe patterns in Fig. 8c and d were generated from the displacement fields in Fig. 8a and b, respectively. Since the

**Fig. 11** Areas of strain concentration in the fringe patterns





**Table 3**  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  values obtained using the data without the strain concentration areas

Model 1 (see Table 1) off-axis angle: $30^\circ$	$Q_{11}$ /GPa 115.62	$Q_{22}$ /GPa 21.60	$Q_{12}$ /GPa 6.99	$Q_{66}$ /GPa 5.04
Model 1 (see in Table 1) off-axis angle: $-60^\circ$	$Q_{11}$ /GPa 74.831	$Q_{22}$ /GPa 21.38	$Q_{12}$ /GPa 7.93	$Q_{66}$ /GPa 5.06
Reference values for Model 1	71.85	20.53	6.16	5.00
Model 2 (see Table 1) off-axis angle: $30^\circ$	$Q_{11}$ /GPa 125.33	$Q_{22}$ /GPa 10.56	$Q_{12}$ /GPa 3.49	$Q_{66}$ /GPa 5.06
Model 2 (see Table 1) off-axis angle: $-60^\circ$	$Q_{11}$ /GPa 134.82	$Q_{22}$ /GPa 10.66	$Q_{12}$ /GPa 4.10	$Q_{66}$ /GPa 5.10
Reference values for Model 2	130.91	10.07	3.02	5.00

grating frequency in the real experiment was 1200 lines/mm, each fringe in Fig. 8 represents the displacement of 0.4167  $\mu\text{m}$ , as explained in Section 2. In addition, phase shifting technology [23] was used to process the fringe patterns in this work. Therefore, four fringe patterns with different phase were generated. An example is provided in Fig. 9.

Then, the strain field presented in Fig. 10a was acquired from the fringe patterns using software Moiré [24]. The strain field obtained from the FEM simulation is also provided in Fig. 10b for comparison. The VFM program was then employed to calculate the four elastic constants. Both of the models described in Table 1 were used with off-axis angles of  $30^\circ$  and  $-60^\circ$ . The results of the simulated experiments are presented in Table 2.

From Table 2, it is evident that both of the models provided reasonably accurate results for  $Q_{22}$  and  $Q_{12}$  when the off-axis angle was  $30^\circ$ . In contrast, more accurate values of  $Q_{11}$  were achieved when the off-axis angle was  $-60^\circ$  than when it was  $30^\circ$ . For  $Q_{66}$ , the error is very small in each case. It is obvious that the results of the simulated experiments are in accordance with the FEM simulation results presented in Section 3. Therefore, it is feasible to use the same configuration and calculation method in real experiments.

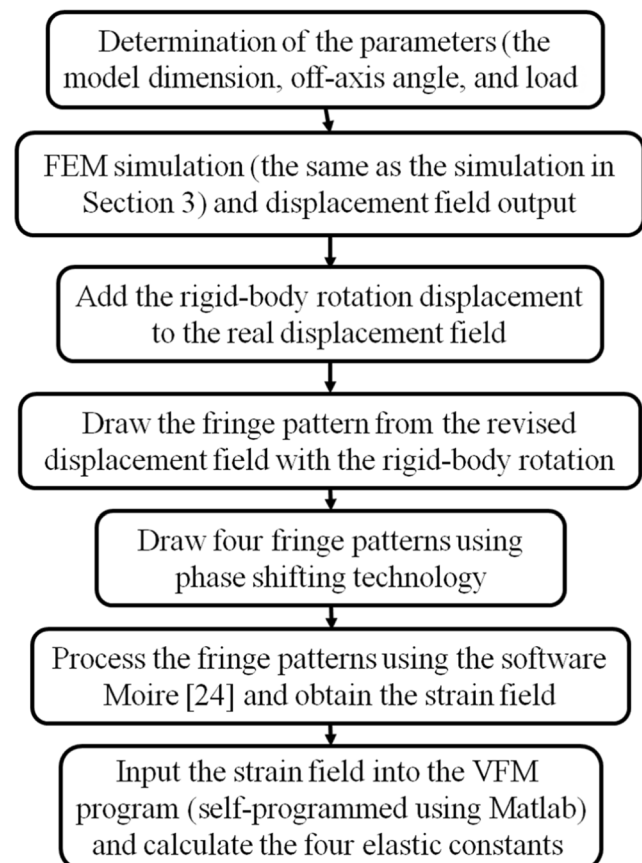
### Error Analysis

In real experiments, many factors induce errors in the measured deformation fields. In Moiré interferometry, since the displacement fields are extracted from the fringes, high quality gratings and null fields are the key to the accuracy of the measured displacements [25]. The gratings used in the real experiment as well as the null fields will be shown in Section 5. However, in this study, it was necessary to consider some additional factors because of the load configuration, as will be discussed in the following subsections.

### Influence of the strain concentration

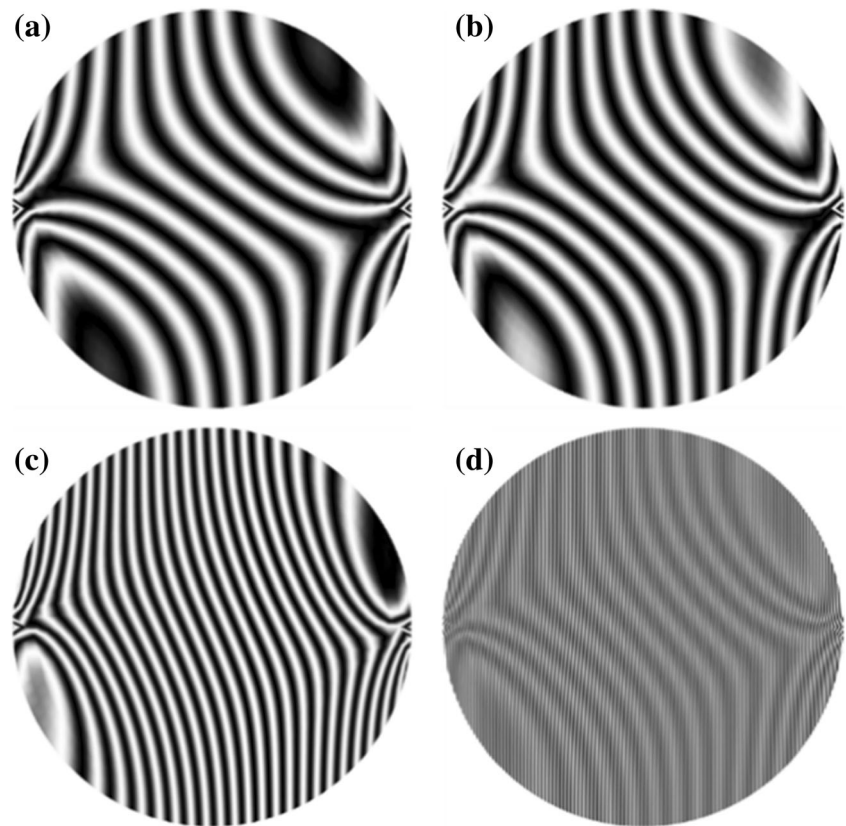
In this work, the concentrated force was applied to the top and bottom of the disk specimen. Therefore, strain concentration

around the loading points was inevitable, resulting in high fringe densities in these areas, as shown in Fig. 11. Two problems would result from the strain concentration. First, the deformation in these areas may be out of the range of elasticity. In addition, it is difficult to obtain accurate displacements from dense fringes in real experiments. To analyze how the strain concentration affects the inversed results, the data from the simulated experiments were used and the strain in the concentration areas was set to zero, as shown in Fig. 11. The results obtained using the data without strain concentration areas are presented in Table 3.



**Fig. 12** Flow chart of the simulated Moiré interferometry experiment with rigid-body rotation

**Fig. 13** Fringe patterns **a** without rigid-body rotation and with; **b**  $0.0025^\circ$ ; **c**  $0.025^\circ$ ; **d**  $0.25^\circ$  rigid-body rotation



The results presented in Table 3 indicate that removing the data from the strain concentration areas did not significantly influence on the results in general. The errors of  $Q_{11}$  in the  $-60^\circ$  case,  $Q_{22}$  in the  $30^\circ$  case, and  $Q_{66}$  in both cases are less than 5%. However, as mentioned above,  $Q_{12}$  is more sensitive to errors than the other parameters. Even so, the accuracy of  $Q_{12}$  is still acceptable based on the comparison to the strain gauge results, which will be presented in Section 5.

#### Influence of rigid-body rotation

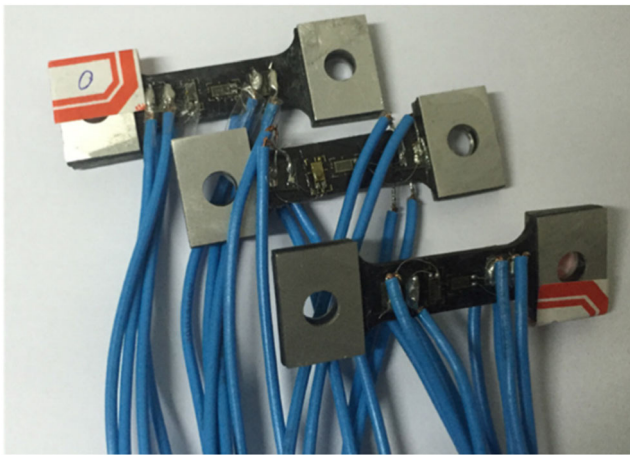
Moiré interferometry is usually applied to measure symmetric displacement field. For such fields, the rigid-body

displacement can be removed by adjusting the specimen position to make the fringe pattern symmetric. However, in the present work, the deformation field was antisymmetric, because the material was orthotropic. Therefore, in this part, the influence of rigid-body rotation on the accuracy of the obtained parameters is discussed, and the results indicate that this influence can be ignored.

The process is explained in detail Fig. 12. Firstly, the parameters to be used in the simulated experiments, such as, the model dimension and the off-axis angle, were determined. Then, the displacement field was obtained using ABAQUS, as described in Section 3. After that, rigid-body rotation was added to the real displacement as shown in equation (7) and

**Table 4**  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  values obtained using the displacement with  $0.025^\circ$  rigid-body rotation

		$Q_{11}/\text{GPa}$	$Q_{22}/\text{GPa}$	$Q_{12}/\text{GPa}$	$Q_{66}/\text{GPa}$
Model 1 (see Table 1) off-axis angle: $30^\circ$	Real displacement	72.84	20.58	6.19	4.96
	Displacement with rigid-body rotation	137.80	20.53	5.50	4.96
Model 1 (see Table 1) off-axis angle: $-60^\circ$	Real displacement	71.56	20.26	6.54	4.97
	Displacement with rigid-body rotation	71.47	20.23	6.42	4.97
Model 2 (see Table 1) off-axis angle: $30^\circ$	Real displacement	-41.72	10.01	3.47	4.97
	Displacement with rigid-body rotation	-42.18	10.09	3.49	4.97
Model 2 (see Table 1) off-axis angle: $-60^\circ$	Real displacement	129.52	10.01	3.36	4.97
	Displacement with rigid-body rotation	128.80	9.99	3.23	4.97



**Fig. 14** Unidirectional composite specimens measured using strain gauges

**Table 5** Elastic constants measured using strain gauges

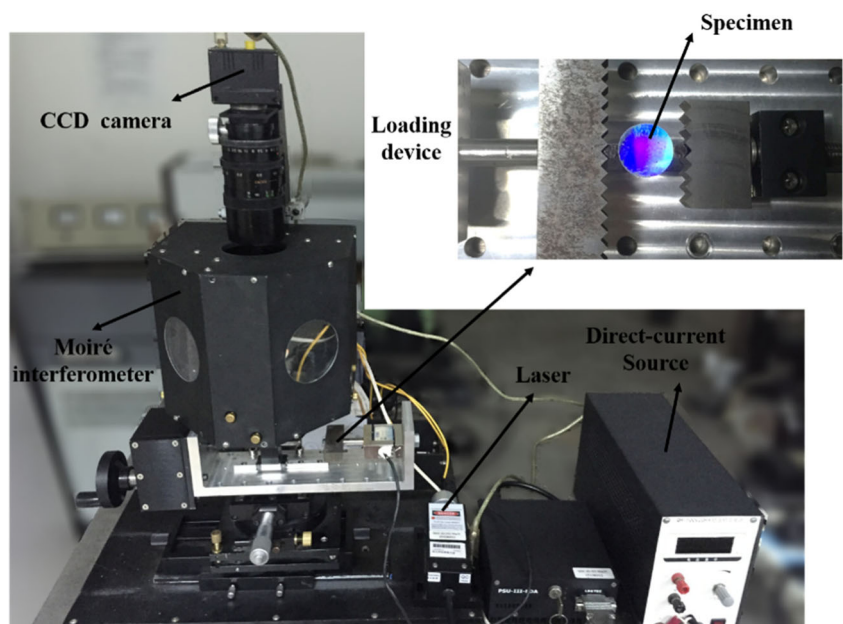
Material	$E_1/\text{GPa}$	$\nu_{12}$	$E_2/\text{GPa}$	$\nu_{21}$	$G_{12}/\text{GPa}$
Unidirectional composite	156.70	0.3043	8.74	0.0228	4.81

the fringe patterns were generated numerically. Finally, the VFM program was used to calculate the four elastic constants.

$$\begin{cases} u' = (\cos\theta - 1) \cdot x + \cos\theta \cdot u - \sin\theta \cdot y - \sin\theta \cdot v \\ v' = (\cos\theta - 1) \cdot y + \sin\theta \cdot x + \sin\theta \cdot u + \cos\theta \cdot v \end{cases} \quad (7)$$

Here,  $u'$  and  $v'$  are the revised displacements with rigid-body rotation,  $u$  and  $v$  are the real displacements, and  $\theta$  is the rotation angle.

**Fig. 15** Experimental setup



To determine a suitable value of  $\theta$ , fringe patterns with rigid-body rotation when  $\theta=0.0025^\circ$ ,  $0.025^\circ$ , and  $0.25^\circ$  were generated, as shown in Fig. 13. It is obvious that when  $\theta=0.025^\circ$ , the difference between the fringe patterns with and without rigid-body rotation displacement is sufficiently large to be recognized in real experiments. Therefore,  $0.025^\circ$  was chosen as the rotation angle for the analysis. Then, both of the models described in Table 1 with off-axis angles of  $30^\circ$  and  $-60^\circ$  were applied to see how the rigid-body rotation affected the inversed values of the four elastic constants. The results are presented in Table 4.

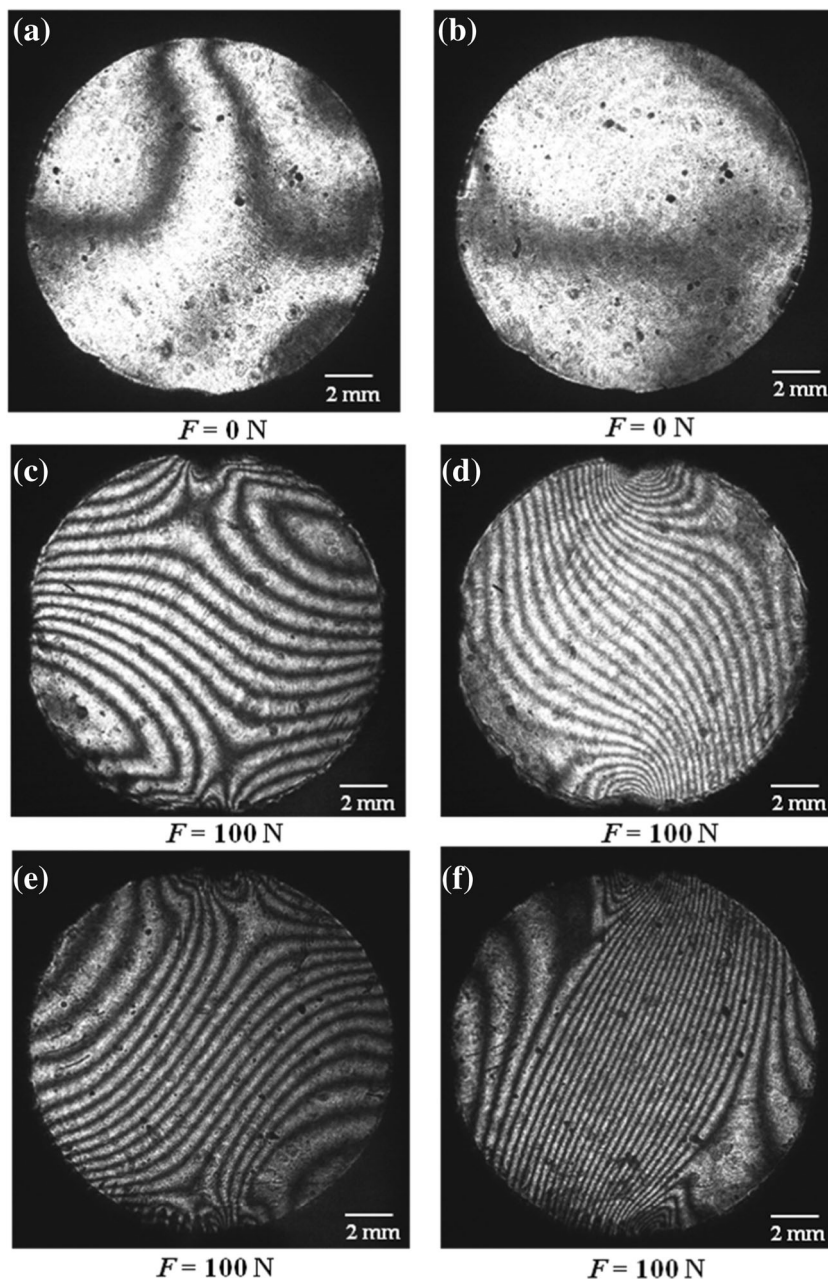
It is obvious that the elastic constants calculated using the displacements with rigid-body rotation are very close to those obtained using the real displacements. The results indicate that the strain fields measured using moiré interferometry are reasonably accurate for the VFM when the deformation field is antisymmetric, even though the rigid-body rotation cannot be eliminated.

## Experimental Validation and Discussion

### Experimental Validation

A real experiment was conducted using a unidirectional carbon fiber composite and Moiré interferometry to demonstrate the feasibility of the proposed method. Firstly, for comparison, strain gauges were used to measure the four elastic constants of the composite. Three types of specimens with off-axis angles of  $0^\circ$ ,  $45^\circ$ , and  $90^\circ$  were used, as shown in Fig. 14. From the normal strain of the  $0^\circ$  specimen,  $E_1$  and  $\nu_{12}$  were obtained. Similarly,  $E_2$  and  $\nu_{21}$  were calculated using the  $90^\circ$  specimen.

**Fig. 16** Null-field patterns along the **a**  $x$  and **b**  $y$  axes; fringe patterns with an off-axis angle of  $30^\circ$  and a load of 100 N along the **c**  $x$  and **d**  $y$  axes, and fringe patterns with an off-axis angle of  $-60^\circ$  and a load of 100 N along the **e**  $x$  and **f**  $y$  axes



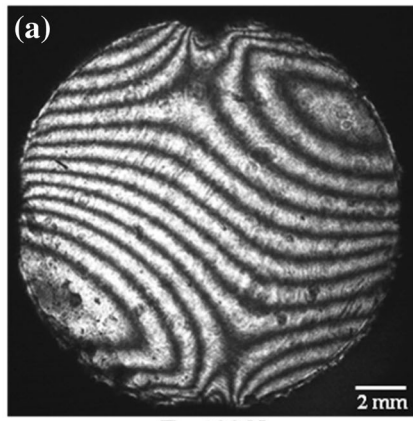
Then,  $G_{12}$  was obtained using the  $45^\circ$  specimen. The results are presented in Table 5.

Theoretically,  $E_1$ ,  $\nu_{12}$ ,  $E_2$ , and  $\nu_{21}$  should satisfy the relationship  $E_1\nu_{21} = E_2\nu_{12}$ . However,  $E_1\nu_{21}$  and  $E_2\nu_{12}$  were found to be 3.57 GPa and 2.66 GPa respectively, which indicates that the measurement error of  $Q_{12}$  by strain gauges is large. Here,  $Q_{12} = \nu_{12}Q_{22}$  was used to calculate  $Q_{12}$  as the reference value, and  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  were measured using the strain gauges to be 157.79 GPa, 8.80 GPa, 2.68 GPa, and 4.81 GPa, respectively.

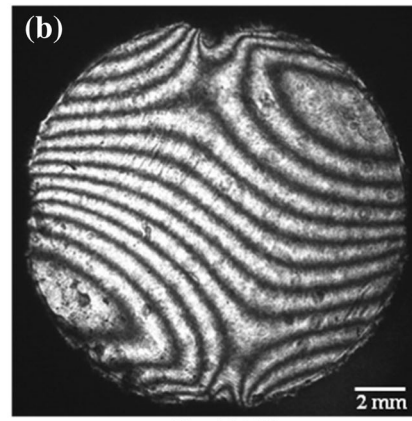
Then, the experiment was performed to determine  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  for the unidirectional carbon fiber composite using Moiré interferometry and the virtual fields method.

The experimental setup and the specimen with 1200 lines/mm gratings are shown in Fig. 15. The wavelength of the laser was 532 nm (TEM00, 20 mW), and the diameter and thickness of the specimen were 15 mm and 2 mm, respectively. The experimental process was explained in Section 4 and depicted in Fig. 6. According to this process, two tests were conducted with off-axis angle of  $30^\circ$  and  $-60^\circ$  and a load of 100 N. The null fields and fringe patterns of the two tests captured by a

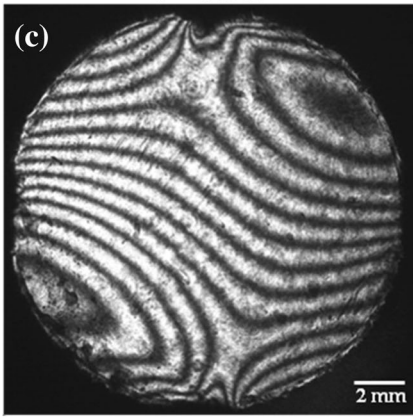
**Fig. 17** Four fringe patterns obtained using phase shifting technology with an off-axis angle of  $30^\circ$  and a load of 100 N along the **a-d**  $x$  and **e-f**  $y$  axes



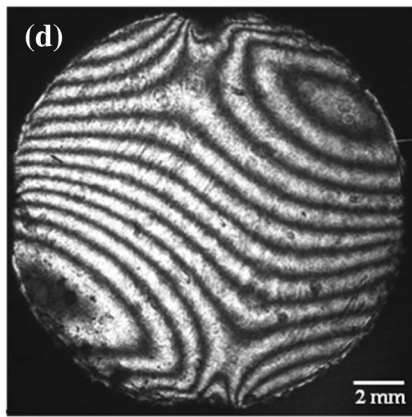
$F = 100 \text{ N}$



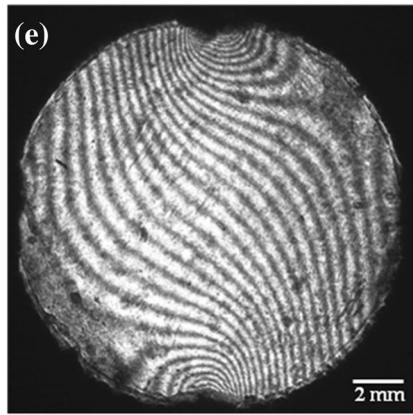
$F = 100 \text{ N}$



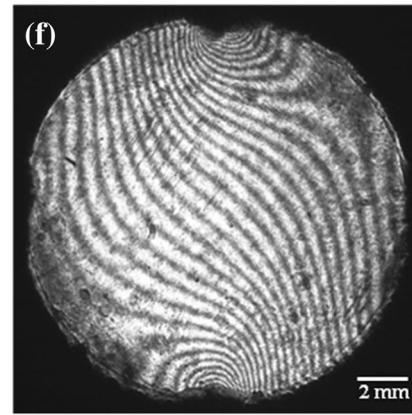
$F = 100 \text{ N}$



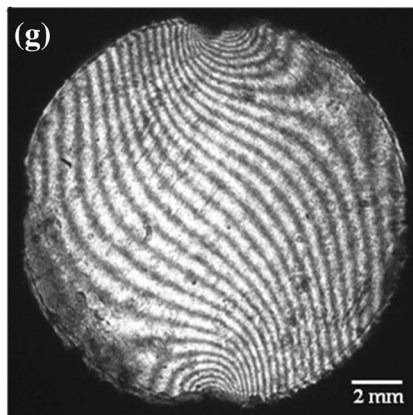
$F = 100 \text{ N}$



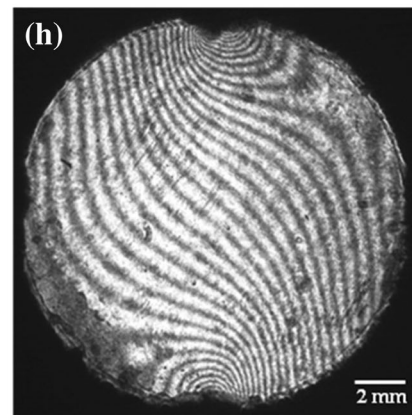
$F = 100 \text{ N}$



$F = 100 \text{ N}$

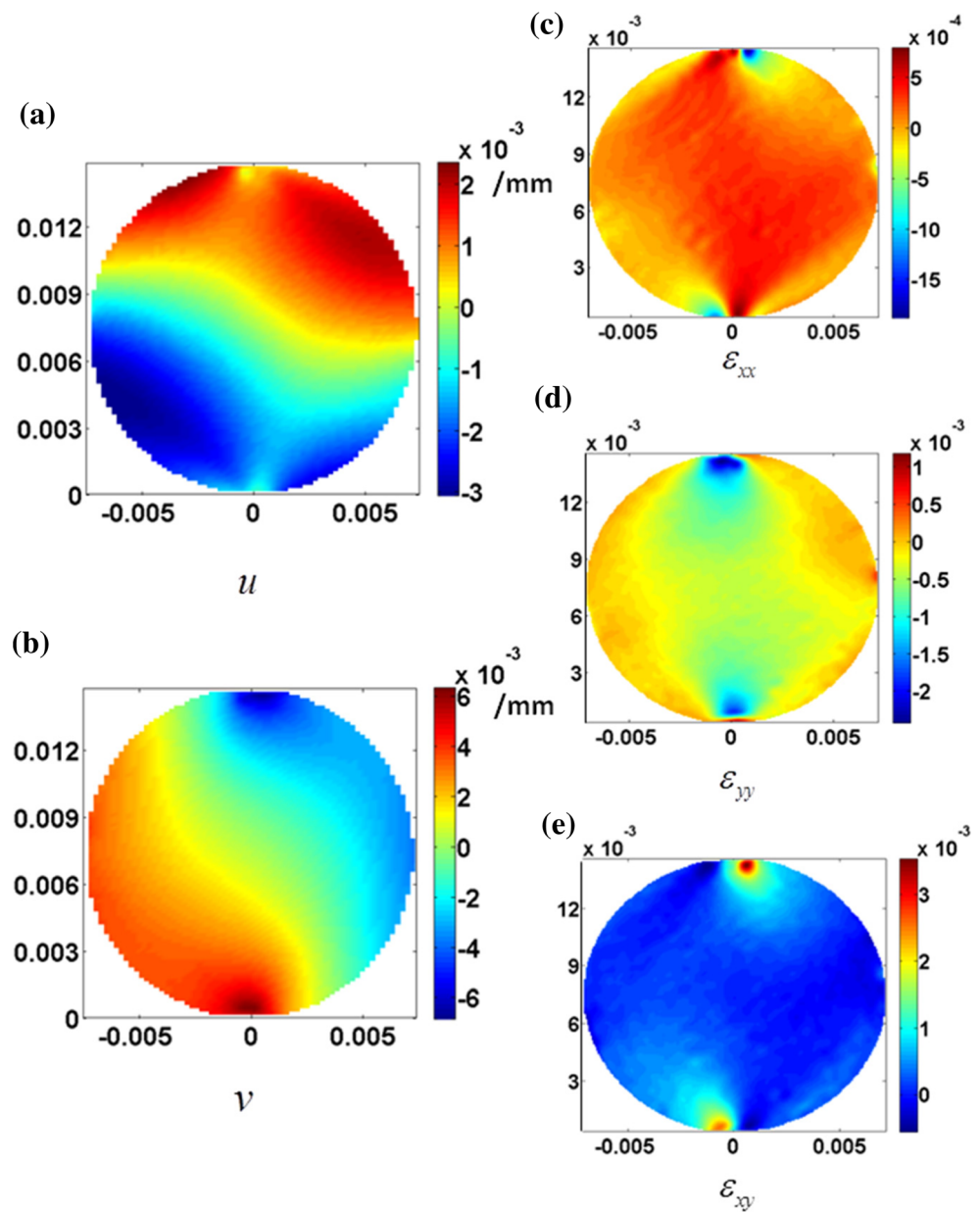


$F = 100 \text{ N}$



$F = 100 \text{ N}$

**Fig. 18** a–b Displacement and c–e strain fields obtained with an off-axis angle of  $30^\circ$  and a load of 100 N



camera (DH-HVI302UM) are shown in Fig. 16. In addition, since the phase shifting technology [23] was applied, the four fringe patterns obtained by the phase shifting with an off-axis angle of  $30^\circ$  are also provided in Fig. 17 as an example. The displacement and strain fields corresponding to the four fringe patterns are shown in Fig. 18.

After the strain fields were obtained, the VFM was used to calculate the four elastic constants. The results of the two tests are presented in Table 6.

## Discussion

From Table 6, it is obvious that the experimental values of  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  obtained in the  $30^\circ$  test are close to those

measured using strain gauges and that the  $Q_{11}$  and  $Q_{66}$  values obtained in the  $-60^\circ$  test exhibit better performance. These results are in accordance with the conclusions drawn based upon the simulated experiments described in Section 4.

To summarize, the main advantage of the VFM is that all of the elastic constants can be obtained from one test. In our

**Table 6**  $Q_{11}$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  values obtained in the real experiments

	$Q_{11}$ /GPa	$Q_{22}$ /GPa	$Q_{12}$ /GPa	$Q_{66}$ /GPa
Off-axis angle: $30^\circ$	-2.847	10.0387	2.035	4.941
Off-axis angle: $-60^\circ$	150.55	5.825	3.883	5.131
Strain gauges results	157.79	8.80	2.68	4.81

work, to improve the measurement accuracy, two tests were conducted. However, only one specimen was used because the disk was selected as the configuration. In contrast, the strain gauge-based approach requires three specimens, and a standard tensile test is conducted on each of them to obtain all the four in-plane elastic constants. In addition, the measurement error of  $Q_{12}$  in the strain-gauge based approach is large, as mentioned in Section 5.1. Therefore, according to the measurement results in Table 6, the accuracy of the proposed approach is acceptable in general.

## Conclusion

In this work, Moiré interferometry was combined with the virtual fields method to characterize the four elastic constants of an orthotropic material. The loading configuration was a diametrically compressed disk. Moiré interferometry was successfully applied in this antisymmetric configuration. FEM simulations, simulated experiments as well as real experiments were conducted, and all of them demonstrated the feasibility of the proposed approach. According to the results of the simulations and experiments, the following conclusions can be drawn:

- (1) In the FEM simulation, cases with off-axis angles of  $-75^\circ$ ,  $-60^\circ$ ,  $-45^\circ$ ,  $-30^\circ$ ,  $-15^\circ$ ,  $0^\circ$ ,  $15^\circ$ ,  $30^\circ$ ,  $45^\circ$ ,  $60^\circ$ ,  $75^\circ$ , and  $90^\circ$  were analyzed. The results show that when the off-axis angle was  $-30^\circ$  or  $30^\circ$ ,  $Q_{22}$ ,  $Q_{12}$ , and  $Q_{66}$  were less sensitive to errors than they were at the other angles and that when the off-axis angle was  $-60^\circ$  or  $60^\circ$ ,  $Q_{11}$  and  $Q_{66}$  exhibited better performance, which provided guidance for the experiment.
- (2) The errors of the four elastic constants obtained in the simulated Moiré interferometry experiment were less than 0.35 GPa. The influence of strain concentration was also analyzed, illustrating that the maximum error of the constants was around 3.9 GPa when the area with strain concentration was removed. In addition, to simulate the real condition of Moiré interferometry in antisymmetric measurements, rigid-body rotation was added to the displacement field. The results show that the influence of rigid-body rotation on the measurements was negligible.
- (3) The results obtained in the real experiment performing using unidirectional carbon fiber composite were compared with the values measured using strain gauges and were found to exhibit reasonable accuracy. Therefore, the strain measured using Moiré interferometry in the antisymmetric configuration is suitable for application in the VFM to characterize the mechanical properties of orthotropic materials.

Furthermore, in our former research [17], efforts were made to characterize fused deposition modeling-fabricated materials, whose principle axes were unknown in advance, using the virtual fields method and digital image correlation. However, the digital image correlation measurement accuracy was not sufficiently high for application to high-strength materials. Thus, the Moiré interferometry method used in the present work enables more accurate strain field determination. Therefore, in the future, unknown materials can be characterized using a combination of the methods employed in the present work and our former research.

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