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# Effects of vdW Interaction and Electric Field on Friction in  $MoS<sub>2</sub>$

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**Abstract** Sliding behaviors between two  $MoS<sub>2</sub>$  layers have been investigated using DFT calculations including vdW dispersion. Contribution of vdW interactions to energy corrugations and maximum lateral frictional forces of the sliding system has been calculated. Our investigations show that the smaller the external normal load is, the larger the contribution of vdW interaction to the friction is. The energy corrugation and lateral frictional force as a function of the electric field are derived, suggesting that friction can be reduced by an external electric field. The reduced friction is attributed to a weaker chemical interaction enabled by the charge depletion between the adjacent S planes. In-depth understanding of the relationship between friction and interlayer interaction shows that friction can be tuned by external electric fields.

**Keywords** First principles  $\cdot$  Friction  $\cdot$  Sliding  $\cdot$  MoS<sub>2</sub>  $\cdot$ Electric field - vdW interaction

## 1 Introduction

Thanks to recent developments in nanotechnology, the hope is high to build mechanical devices on the scale of the nanometer. For this purpose, it is important to determine

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the mechanical and, especially, frictional properties in such nanodevices. The understanding and control of frictional properties of matter at the nanoscale remain a great challenge for a wide spectrum of sciences, ranging from engineering studies of micromachine lubrication to atomistic simulations of nanoscale systems. In particular, microelectro-mechanical system (MEMS) and nanoelectromechanical system (NEMS) devices hold a tremendous potential in a variety of applications. One important challenge in applications of MEMS devices is that at these small scales, surface forces, e.g., adhesion and friction, become dominant and can lead to a fast degradation of material's performance. Therefore, fundamental understanding of friction at the atomic scale is essential for the successful design of MEMS and NEMS devices [[1,](#page-6-0) [2](#page-6-0)].

Among the various members of the transition metal dichalcogenides (TMDs) family of compounds,  $MoS<sub>2</sub>$ , usually used as a solid lubricant in MEMS, is well known for its lubricating behavior  $[3-5]$ . Its lubricating behavior stems from easy glide between  $MoS<sub>2</sub>$  sheets which is intrinsic to its crystal structure. Similar to graphite,  $MoS<sub>2</sub>$ crystallizes in the hexagonal structure where a sheet of molybdenum atoms is sandwiched between two hexagonally packed sulfur layers. The bonding within the S–Mo–S sandwich is covalent, while weak van der Waals forces hold the sandwich together resulting in interlamellar easy glide. Experimentally, pure  $MoS<sub>2</sub>$  can obtain a superlow friction coefficient of 0.002 (superlubricity) under high vacuum [\[6](#page-6-0)]. Lee et al. [[2\]](#page-6-0), using friction force microscopy, have investigated the nanoscale frictional characteristics of atomically thin  $MoS<sub>2</sub>$  sheets and found that friction monotonically increased as the number of layers decreased. The alloying  $MoS<sub>2</sub>$  into much harder matrices such as amorphous carbon greatly stimulates the production of efficient self-lubricating coatings [\[3](#page-6-0)]. Singer et al. [\[7](#page-6-0)] have reported the friction coefficient of thin  $MoS<sub>2</sub>$ coatings decreases as the load increases. Moreover, the Amontons–Coulomb law has been observed on the nanoscale, when a  $MoS<sub>2</sub>$  flake moves on a  $MoS<sub>2</sub>$  surface [\[8](#page-6-0)]. All these evidences imply the need for a clear understanding of its tribological properties at the atomistic level, and particular attention must also be paid to unraveling the nature of the interlayer forces, as an interplay of electrostatic interactions, van der Waals (vdW) forces, and Pauli repulsion  $[2, 9-12]$  $[2, 9-12]$ . In recent years, compared to tremendous attention paid to layered  $MoS<sub>2</sub>$  with exceptional electronic, optical, and catalytic properties (see review  $[13–15]$  $[13–15]$ ), less work  $[9, 16–21]$  $[9, 16–21]$  $[9, 16–21]$  $[9, 16–21]$  $[9, 16–21]$  $[9, 16–21]$  studied its tribological behavior. Previously, the static potential energy surfaces and atomic-scale energetic barriers encountered during sliding at  $MoS_2$  (0001) and  $MoO_3$  (001) surfaces and at the  $MoS_2$ /  $MoO<sub>3</sub>$  interface have been firstly calculated by using firstprinciples method [[16\]](#page-6-0). The results suggest that the lowest energy pathway is to slide  $MoO<sub>3</sub>$  over  $MoS<sub>2</sub>$ , and the highest energy pathway involves  $MoO<sub>3</sub>$  (001) interlayer sliding. Blumberg et al. [\[18](#page-6-0)] also determined the sliding potential of  $MoS<sub>2</sub>$  by a simple and intuitive explanation on the basis of the registry index (RI) concept. Onodera et al. [\[9](#page-6-0)] found that the predominant interaction between two  $MoS<sub>2</sub> sheets is Coulombic repulsion, which directly affects$ the  $MoS<sub>2</sub>$  lubrication using a computational chemistry method. Moreover, this group has investigated the friction anisotropy of two-layered  $MoS<sub>2</sub>$  at atomistic level adopting the same method [\[17](#page-6-0)].

The first-principles method is widely recognized as a powerful tool to examine electronic structures and the interactions between surfaces at the atomistic scale. Moreover, within the generalized gradient approximation (GGA) calculations including van der Waals corrections, the frictional figure of merit for a pair of layered  $MoS<sub>2</sub>$  nanostructures has been determined by Cahangirov et al. [\[19](#page-7-0)]. Very recently, by means of ab initio modeling of the static potential energy surface and charge distribution analysis, Levita et al. [\[20](#page-7-0)] demonstrated how electrostatic interactions, negligible in comparison with van der Waals and Pauli contributions at zero load, progressively affect the sliding motion at increasing loads. Our group first investigated the effect of strain on atomic-scale friction in layered  $MoS<sub>2</sub>$  [\[21](#page-7-0)]. The results indicate that one can change the interfacial interactions and then tune the interlayer friction by applying biaxial in-plane strains to the layered  $MoS<sub>2</sub>$ . However, all the above researches have not directly studied vdW interactions between the  $MoS<sub>2</sub>$  layers, especially at increasing loads. The past investigations showed that vdW interactions play an important role on the electronic structures [\[12](#page-6-0)], the interaction energy of layered materials [[22,](#page-7-0) [23](#page-7-0)], and negative friction coefficient at low load [\[24](#page-7-0), [25](#page-7-0)]. Thus, there are important practical significances to investigate vdW interactions between the lamellar nanostructures at increasing loads. Moreover, Park et al. [\[26](#page-7-0)] reported a method by the use of electric fields for controlling friction in a nanoscale contact. They worked with a silicon pn junctions. When a bias of  $+4$  V was applied to the sample, an increase in friction by up to a factor of 2 was observed in the p-doped region. The mechanism for this increase is not clear, but estimates of the contribution due to electronic friction (the drag force that results when charge carriers move in an electric field) are too low to explain the result. Previous investigations reveal that the friction can be tuned by modifying atomic and then electronic structures of the sliding surfaces [[21,](#page-7-0) [27–30](#page-7-0)]. Moreover, one can change the electronic structure of bilayer  $MoS<sub>2</sub>$  by vertical electric field [\[31](#page-7-0), [32](#page-7-0)]. In a word, it is necessary to study the effect of the external electric field on the sliding properties in the layered materials.

In this paper, for the purpose of better understanding the frictional behaviors between  $MoS<sub>2</sub>$  layers, particularly vdW interactions and effects of the external electric field, we investigate their influence on the interlayer sliding using density functional theory (DFT) calculations including vdW correction. This static approach will help providing a rationale of the microscopic processes which could lead to further insights into the tunability of the macroscopic frictional properties of  $MoS<sub>2</sub>$ . In-depth understanding of the relationship between the atomic-scale friction and the interfacial interaction shows that the interlayer friction in  $MoS<sub>2</sub>$  can be tuned by external electric fields. Moreover, as the load decreases, the contribution of the vdW interactions to the friction increases. These investigations may be helpful for friction control, as well as electronic device and lubricant design.

#### 2 Details of Calculation Methods

All of calculations in this work were performed using the Vienna ab initio simulation package (VASP) code [\[33](#page-7-0), [34](#page-7-0)], equipped with the projector augmented-wave (PAW) method [\[35](#page-7-0), [36\]](#page-7-0) for electro-ion interaction. The exchange– correlation functional is represented with the GGA in the form of the Perdew–Burke–Ernzerhof (PBE) [\[37](#page-7-0)]. The electronic wave functions are expanded in a plane-wave basis with the energy cutoff of 500 eV to converge the relevant quantities. Total energy and electronic structures of the systems are calculated using a  $15 \times 15 \times 1$  Monkhorst–Pack (MP) grid [[38\]](#page-7-0). The convergence for energy is chosen to be  $10^{-5}$  eV between two electronic relaxation steps, and the maximum force allowed on each atom is  $\leq 0.005$  eV/Å upon ionic relaxation. For an accurate description of the geometric and electronic structure of

<span id="page-2-0"></span>molybdenum disulfide, we carried out DFT calculations including van der Waals dispersion forces (DFT-D2) proposed by Grimme [[39\]](#page-7-0).

As shown in Fig. 1, a supercell used in the present study consists of two layers of  $MoS<sub>2</sub>$  sheets to simulate their relative sliding. To eliminate the interactions between adjacent supercells, we consider the supercell with a length of  $30 \text{ Å}$  in the z direction, leaving a vacuum region of more than 20  $\AA$ . The loading method is similar to that described in the previous work  $[21]$  $[21]$ . The potential energy surface (PES) experienced by the upper monolayer upon translation above the lower one is constructed by calculating the interlayer interaction energy for different relative lateral positions of the two layers by shifting the upper layer by a step of 0.2  $\AA$  along the x and y directions, respectively. In order to study the load effect on the friction along the considered path  $(x$  direction as shown in Fig. 1), we investigate the variation of the system energy corrugation (the difference between the maximum and minimum energies) and the maximum static lateral force under different loads. The load has been evaluated as the ratio  $L = F<sub>z</sub>/A$ , between the z component of the residual force on the cell  $F<sub>z</sub>$  and the cell area A. An external electric field with a tolerable magnitude in the range of  $0.1-0.3$  V/ $\AA$  has been applied perpendicularly to the sheets by adding an artificial



Fig. 1 Ball and stick model of  $MoS<sub>2</sub>$  used in the present study. Yellow and *aqua balls* represent S and Mo atoms, respectively. **a** Top view, b side view (Color figure online)

dipole sheet in the middle of the vacuum part in the periodic supercell [[40\]](#page-7-0).

According to the DFT-D2 method [\[39](#page-7-0)], the total energy of a system is given by

$$
E = E_c + E_{\text{vdW}} \tag{1}
$$

where  $E_c$  is the usual self-consistent Kohn–Sham energy and  $E_{\text{vdW}}$  is the vdW dispersion correction. The dispersion energy  $E_{\text{vdW}}$  is given by

$$
E_{\rm vdw} = -s_6 \sum_{i=1}^{N-1} \sum_{j=i+1}^{N} \frac{C_6^{ij}}{R_{ij}^6} \frac{1}{1 + e^{-d(R_{ij}/R_r - 1)}} \tag{2}
$$

where  $s<sub>6</sub>$  (0.75 for PBE potential) is a global scaling factor, N represents the number of atoms in the system,  $C_6^{ij}$  is the dispersion coefficient for atom pair *ij*, given by  $C_6^{ij} =$  $(C_6^i C_6^j)^{1/2}$ ,  $R_{ij}$  denotes the interatomic distance, and  $R_i$ represents the sum of atomic van der Waals radii, respectively. d is a damping parameter (20.0 for PBE potential). The dispersion coefficient  $C_6$  is 24.670 and 5.570 J nm<sup>6</sup>/ mol and vdW radii  $R_0$  is 1.639 and 1.683 Å for Mo and S atoms, respectively.

The corresponding lateral frictional force along the sliding direction  $f$  is given by

$$
f = \frac{\partial E}{\partial x} = \frac{\partial E_c}{\partial x} + \frac{\partial E_{\text{vdW}}}{\partial x} = f_c + f_{\text{vdW}}
$$
(3)

where  $f_c$  and  $f_{\text{vdW}}$  are the contribution of the Kohn–Sham energy and the vdW interaction to the lateral frictional force, respectively.

In the lamellar system, the energy corrugation along the sliding path is mainly composed of the change in the electrostatic and the vdW interaction between atoms of the two layers, especially in the situation of quasi-static or very low velocity sliding [[23\]](#page-7-0). That is,

$$
\Delta E = \Delta E_c + \Delta E_{\text{vdw}} \tag{4}
$$

where  $\Delta E$ ,  $\Delta E_c$ , and  $\Delta E_{\text{vdW}}$  are the energy corrugation, the variation of the electrostatic interaction energy, and the vdW interactions energy, respectively. Therefore, we can evaluate the contribution of vdW interaction to the friction between the two sheets as the ratio  $\Delta E_{\text{vdW}}/ \Delta E$  and  $\Delta f_{\text{vdW}}^{max}/ \Delta f$  max where  $\Delta f^{\text{max}}$  and  $\Delta f^{\text{max}}_{\text{vdW}}$  are the maximum lateral frictional force and the contribution of vdW interaction, to that, respectively.

## 3 Results and Discussion

### 3.1 vdW Interaction on Friction of  $MoS<sub>2</sub>$  Layers

To check the efficiency of the computational method, the MoS2 bilayer has been first optimized using the DFT-D2 method mentioned previously. Our calculations show that

the hexagonal lattice constant  $a_0$  is 3.198 Å, the length of the Mo–S bond is 2.417  $\AA$ , the interplanar binding energy is 0.149 meV, and the thickness of  $MoS<sub>2</sub>$  monolayer is 3.119  $\AA$ , in good agreement with previous calculations [\[20](#page-7-0)]. The interlayer distance has been calculated to be 6.22 Å, which lies within a 1.5 % margin from the experimental data  $(6.15 \text{ Å})$  [[2\]](#page-6-0).

Without an external electric field, the static potential energy surface (PES) at zero load mapped by the contour plot as a function of relative displacement of two  $MoS<sub>2</sub>$ layers in the  $x$  and  $y$  directions is shown in Fig. 2a. For convenience, the contour plot is expanded with intrinsic periodicity. The minimum potential energy is set to 0 eV. From this energy surface, we can find that the potential energy surface has the same sixfold symmetry as the surface structure of  $MoS<sub>2</sub>$  (0001). The energetically favored path (black solid line) during sliding is shown in Fig. 2a. Such a zigzag path in the potential energy surface has been predicted in previous work [[16–](#page-6-0)[20\]](#page-7-0) as well. The corresponding potential energy profile for the linear translation of a  $MoS<sub>2</sub>$  layer above a fixed one along the x direction is shown in Fig. 2b. Four characteristic potential energies

(Min1, Saddle, Min2, and Max) are also marked in this figure. According to these four potential energies, their geometric configurations are shown in Fig. 2c, d, e, f, respectively. For the minimum energy configuration, Min1, the arrangement of two-layer  $MoS<sub>2</sub>$  is similar to its bulk structure (Fig. 2c). Namely, the Mo (S) atom of the upper layer straightly falls on top of the S (Mo) atom of the bottom one. For the other stable configuration, Min2, the Mo atom of its upper layer is faced to the Mo atom of the other layer, while the S atom of one layer is located in the sixfold symmetry center of the other one (Fig. 2e). When the S atom of one layer is faced to that of the other (Fig. 2f), the system has the maximum potential energy, 0.053 eV/cell, in agreement with the previous calculation [\[20](#page-7-0)].

Next, we consider the vdW interaction on the friction of  $MoS<sub>2</sub>$  layers. When one layer slides above the other along the  $x$  direction, the interlayer distance, vdW interaction energy, and the corresponding potential energy at three certain normal loads (0, 5, 9 GPa) are plotted in Fig. [3](#page-4-0)a, b, c, respectively. For the vdW and potential energy, they have been set to 0 eV at the initial position  $(x = 0)$ . At the



Fig. 2 a Potential energy surface (PES) for the sliding motion of  $MoS<sub>2</sub> bilayer at zero load. Black solid line represents the minimum$ energy path. The color scale refers to the energy range of the PES corrugation (in eV); b the potential energy profile for the linear

translation of a  $MoS<sub>2</sub>$  layer above a fixed one along the x direction. c, d, e, and f Represents Min1, saddle, Min2, and Max geometric configuration, respectively. Yellow and aqua balls represent S and Mo atoms, respectively (Color figure online)

<span id="page-4-0"></span>

Fig. 3 Interlayer distance (a), vdW energy (b), and potential energy (c) as a function of the sliding distance along the  $x$  direction. The potential energy at the initial position  $(x = 0)$  has been set to 0 eV

same load, such as 5 GPa (red line in Fig. 3), they show a similar trend with the relative displacement of the two  $MoS<sub>2</sub>$  layers: The vdW interaction and potential energy increase or decrease with the increase or decrease in the interlayer distance. At a certain sliding distance (e.g., Min1, Saddle, and Max configuration), the interlayer distance and vdW interaction energy decrease with the increase in external normal load. However, due to the electrostatic repulsions between the two adjacent S planes, the potential energy raises. With respect to Min2 arrangement, its potential energy reduces with the increase in load. This may be the result of competition between van der Waals and electrostatic repulsion interactions.

Understanding of the frictional properties of  $MoS<sub>2</sub>$  layers requires comparing the energy corrugation and maximum lateral frictional force along the sliding path. The energy corrugations and maximum lateral frictional forces as a function of external normal loads are plotted in Fig. 4a. As the Amontons's friction law expected, both energy corrugations and maximum lateral frictional forces monotonically increase with normal loads. It shows that the



Fig. 4 a Energy corrugation and lateral frictional force as a function of a normal load; b contribution of vdW interaction to the energy corrugation and lateral frictional force as a function of a normal load

interlayer friction of  $MoS<sub>2</sub>$  on the atomic scale also obeys the basic friction law observed experimentally by Miura [\[8](#page-6-0)]. We have further studied the vdW interactions on friction of  $MoS<sub>2</sub>$  layers. As shown in Fig. 4b, the contributions of vdW interactions to energy corrugation and lateral frictional force as a function of a normal load have been exhibited. From the figure, we can find that vdW interactions increasingly contribute to the friction with the decreased load. Particularly, at a small load, vdW interactions dominate the interlayer friction of the sheet materials [\[24](#page-7-0), [25\]](#page-7-0). Of course, the above analysis is an oversimplification of the experimental situation in which layers cannot move as rigid blocks. Instead, in a physical system, it is possible that some regions will slip, while others will not. The regions that do slip will not settle fully into their local minima and will interact elastically with other regions that have not slipped. Thus, the predicted maximum frictional forces based on this approach can be expected to be generally larger than the experimental values. These limitations notwithstanding our results give a trend in the sliding behavior as a function of a normal load, in close agreement with previous experimental measurements [[8\]](#page-6-0).

## 3.2 Effects of Electric Field on Friction of  $MoS<sub>2</sub>$ Layers

Among several strategies currently being employed to change electronic structures of two-dimensional materials, tuning of electronic structures by external electric fields is a particularly interesting one. Recent theoretical studies also suggest the possibility of employing a similar strategy to manipulate the band gap in bilayer  $MoS<sub>2</sub>$  [[31,](#page-7-0) [32\]](#page-7-0). In this case, an external electric field reduces the fundamental band gap of the bilayer structure and a transition from

semiconductor to metal occurs at an appropriate situation. It is well known that the friction between two  $MoS<sub>2</sub>$  sheets originates from interlayer interactions. Moreover, electronic structures of the system are also sensitive to interlayer interactions. Therefore, modulation of electronic structures by electric fields may be practicable to tune the interlayer friction.

Effects of external electric fields on sliding properties of  $MoS<sub>2</sub>$  layers are also investigated in this paper. The electric field is perpendicularly imposed on the system by adopting an artificial dipole sheet in the middle of the vacuum (Fig. [1](#page-2-0)). Applied by an external electric field, the energy corrugation and maximum lateral frictional force of the relative sliding system at a fixed normal load have been calculated by using the above methods. Figure 5 shows the variation along the  $x$  direction of the system energy corrugation (Fig. 5a) and the maximum static lateral force (Fig. 5b) under external electric fields. Our results indicate that the energy corrugation and the corresponding maximum lateral frictional force take on a similar trend. That is, at a constant external electric field, higher energy corrugation and maximum lateral frictional force will be experienced when one layer slides relatively to another under higher load. And even it is more interesting that, compared with a zero electric field, the energy corrugation and maximum lateral frictional force at a fixed load slightly increase with the applied electric field from 0 to  $0.25$  eV/ $\AA$ , while they decrease with a more intensive electric field  $(0.25 \text{ eV/A})$ , in which the 0.25 eV/Å is a critical value.

In order to comprehensively understand the effects of electric fields on the friction of the two  $MoS<sub>2</sub>$  layers, electronic structures of this system at a fixed load need to be investigated. Firstly, energy band structures under different electric fields of the two  $MoS<sub>2</sub>$  layers have been calculated. In Fig. 6, we provide an overview of the band



Fig. 5 Energy corrugation (a) and maximum lateral frictional force (b) at a fixed normal load as a function of the external electric field



Fig. 6 Energy band structures of the  $MoS<sub>2</sub>$  bilayer applied with different electric fields. The normal load is 0.5 GPa. The fermi level is set to 0 eV

structures of the  $MoS<sub>2</sub>$  bilayer at a 0.5-GPa load considered in this work as a function of applied external field. It is apparent at a glance that the band gap in this case is driven continuously to zero with increasing external electric fields. We note that the fundamental band gap at zero electric field is an indirect gap between the valence-band (VB) maximum at  $\Gamma$  and the conduction-band (CB) minimum at K. The VB maxima at  $\Gamma$  and CB minima at  $K$  inexorably approach the Fermi energy value with increasing external electric fields. At the electric field of about  $0.25 \text{ eV/A}$ , similar to the critical value of friction transition (Fig. 5), the bilayer translates from semiconductor to metal. The band-gap modulation discussed above arises from the wellknown Stark effect, which has been observed in the previous studies as well [[31,](#page-7-0) [32](#page-7-0)]. External perpendicular electric field induces a potential difference between the two layers. As a result, the energy bands belonging to different MoS2 layers are separated from each other entirely. In  $MoS<sub>2</sub> bilayer, the stronger the electric field is, the larger$ the band splitting is, and thus the smaller the band gap is.

Next, the self-consistent difference of charge density is obtained by subtracting the charge density of the  $MoS<sub>2</sub>$ bilayer with an electric field from that of the  $MoS<sub>2</sub>$  bilayer without an electric field. The isosurfaces of the difference of charge density and the variation of its value averaged over  $(x, y)$  planes parallel to  $MoS<sub>2</sub>$  layers (called the difference of linear charge density along the z axis) are presented in Fig. [7.](#page-6-0) Generally, charge distribution of the sliding interface dominates the friction between the two relative sliding interfaces [[28\]](#page-7-0). The charge density accumulation for  $0.20 \text{ eV/A}$  (depletion for  $0.30 \text{ eV/A}$ ) between the two adjacent S planes is the key feature to explain the increase (decrease) in the energy corrugation and maximum lateral frictional force. This charge accumulation (depletion) enhances (weakens) the chemical interaction between the two layers and results in the increase (decreasing) in energy corrugation and lateral frictional force as shown in Fig. 5. Our theoretical investigations may

<span id="page-6-0"></span>

**Fig. 7** a Difference of charge density ( $\Delta \rho = 0.0005 \text{ e}/\text{\AA}^3$ ) of the MoS<sub>2</sub>. **b** Changes in the linear charge density distribution along the  $z$  direction. The normal load is 0.5 GPa. Red (blue) isosurface plots correspond to the charge density accumulation (depletion). For the definition of difference of charge density, see the text (Color figure online)

provide important information on how to reduce friction between  $MoS<sub>2</sub>$  interfaces. Of course, theoretical calculations need to be verified by experiment.

## 4 Conclusions

Sliding properties of two  $MoS<sub>2</sub>$  layers have been investigated by the use of density functional theory (DFT) calculations including van der Waals corrections. Static potential energy surface of the relative sliding system is plotted in this paper. Contribution of vdW interactions to energy corrugations and maximum lateral frictional forces of the sliding system has been calculated by the DFT-D2 method. Our results show that the smaller the external normal load is, the larger the contribution of vdW interaction to the friction is. Furthermore, the energy corrugation and maximum lateral frictional force as a function of the electric field are derived as well, suggesting that the friction between two  $MoS<sub>2</sub>$  layers can be reduced by imposing an electric field. The reduced friction is attributed to a weaker chemical interaction enabled by the charge depletion between the adjacent S atomic planes. In-depth understanding of the relationship between friction and interfacial interaction shows that friction can be tuned with an external electric field.

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