



Structural evolution and electronic properties of medium-sized $\text{CrSi}_n^{-/0}$ ($n = 19\text{--}25$) clusters

Kai Wang¹ · Chaoyong Wang¹ · Wei Li¹ · Le Liu¹ · Yarui Wang¹ · Jiaye Chen¹ · Jun Zhao¹ · Junji Guo¹

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Abstract

Herein, the structural evolution, electronic and magnetic properties of $\text{CrSi}_n^{-/0}$ for $n = 19\text{--}25$ have been investigated at density functional theory (DFT) level. The global minimum structures of these clusters have been fully searched through a self-developed genetic algorithm code combined with DFT calculations. Both of anionic and neutral CrSi_n for $n = 19\text{--}25$ share the same configurations. All these clusters prefer to adopt endohedral structures (Cr@Si_{14} for sizes $n = 19\text{--}21$, and Cr@Si_{13} for larger ones) as the structural motif with the remaining Si atoms attached on the surface. Among all these clusters, CrSi_{22} and CrSi_{23} are the most prominent through the analysis of HOMO–LUMO gaps, average binding energies, and the second order of energy differences. All of these medium-sized cluster anions possess the same total magnetic moment of $1 \mu_B$, but with very different contributions from that of small sizes ($n \leq 18$), except for size 22.

Keywords Chromium-doped silicon clusters · Structural evolution · Density functional theory · Genetic algorithm

Introduction

Silicon clusters have potential applications in micro-nano devices as building blocks to continue the miniaturization trend of Moore's Law. However, the pure silicon clusters are always chemically reactive, limiting their applications [1, 2]. Fortunately, the introduction of transition metals into silicon clusters can significantly improve their stabilities [3–5], and may also introduce novel electronic and magnetic properties [6, 7], possibly allowing them to be used as building block for cluster-assembled materials. For example, the caged clusters of Si_{16} encapsulating a group-IV metal atom form superatoms with large HOMO–LUMO gaps [8–10]. The wheel-like $\text{V}_3\text{Si}_{12}^-$ cluster exhibits a ferromagnetic state with a total magnetic moments of $4\mu_B$ [7]. Therefore, more and more experimental and theoretical efforts have been devoted to this type silicon clusters in recent thirty years [8–50]. In the early stages, people mainly studied the stability of these clusters through mass spectrometry technology.

For example, Beck studied the stability of TMSi_n ($\text{TM} = \text{Cr}, \text{Mo}, \text{W}, \text{Cu}$) clusters by mass spectrometry experiments, revealing several stable clusters including $\text{TMSi}_{15,16}$ ($\text{TM} = \text{Cr}, \text{Mo}^{0/+}, \text{W}$) [11, 12]. Neukermans et al. also studies a series of similar systems including TMSi_n ($\text{TM} = \text{Cr}, \text{Mn}, \text{Cu}, \text{Zn}$) clusters through the similar mass spectrometry experiments [13]. Subsequently, more advanced technology such as photoelectron spectroscopy (PES) was used to investigate the electronic properties of a series of dopants ($\text{Sc}, \text{Ti}, \text{V}, \text{Y}, \text{Zr}, \text{Nb}, \text{Lu}, \text{Tb}, \text{Ho}, \text{Hf}, \text{Ta}, \text{Mo},$ and W) doped Si_n^- clusters with sizes $n \leq 20$ [14–18]. The demand for understanding the structure information of these type clusters urges the research on their structures. The structures of MSi_{8-16} ($\text{M} = \text{Sc}\text{--}\text{Zn}, \text{Zr}, \text{Hf}, \text{Ta}; n \leq 18$) [19–28], $\text{M}_2\text{Si}_{1-8}$ ($\text{M} = \text{Cr}, \text{Mn}, \text{Fe}, \text{Co}, \text{Ni}$) [29, 30], $\text{Mo}_2\text{Si}_{9-16}$ [31], $\text{M}_2\text{Si}_{11-18}$ ($\text{M}_2 = \text{Mo}_2, \text{Nb}_2, \text{Ta}_2, \text{W}_2, \text{NbMo}, \text{TaW}$) [32], $\text{Ag}/\text{AuSi}_{10,12,14}$ [51], and $\text{AuSi}_{1-12}^{0/+}$ [52] clusters were predicted based on DFT calculations, and their structural evolutions is clarified accordingly. Considering the experimental photoelectron spectrum as the fingerprint of a cluster or molecule [53–55], the comparison of experimental and theoretical spectra have been used as criteria to identify the ground-state structures of TMSi_n^- ($\text{TM} = \text{Sc}_{1-2}$ [33, 34], V_{1-2} [35], Cr_{1-2} [36, 37], Co [38], Nb_{1-2} [39, 40], Ag [41], La [42], Ta [43], Au_{1-2} [44, 45], $n \leq 14$; $\text{TM} = \text{Ti}_{1-2}$ [10], V_{1-3} [46–48], Cr_{1-2} [47, 49]; $14 \leq n \leq 20$) clusters in recent years. Both experimental

✉ Kai Wang
kkwangkai@huuc.edu.cn

¹ Henan Engineering Research Centre of Building-Photovoltaics, School of Mathematics and Physics, Henan University of Urban Construction, Pingdingshan 467036, China

and theoretical results show that TM atoms prefer to be surrounded by Si atoms, and forming endohedral structures with the growing number of Si atoms, except for TM = Cu, Ag and Au for small sizes with TM atoms absorbed on the surface of the bare Si_n clusters [41, 44, 50].

Among all 3d transition metals, chromium ($3d^54s^1$) is the only one has 6 unpaired electrons in the valence shell, making it an interesting dopant for the modulation of magnetic properties. Therefore, Cr-Si clusters have received widespread attention [36, 37, 47, 49, 56, 57] due to their potential applications in spintronic materials [58] and contact materials [59]. Recent studies have shown that small-sized CrSi_n exhibits strange structures and magnetic properties [36, 37]. For example, $\text{Cr}_3\text{Si}_{12}^-$ possess a wheel-like structure [37]. Both of $\text{Cr}_2\text{Si}_{13}^-$ and $\text{Cr}_3\text{Si}_{12}^-$ exhibit large magnetic moments of $3\mu_B$ and $7\mu_B$, respectively [37]. Our recent research on CrSi_{14-18} cluster anions shows that Cr atoms do not always contribute positive magnetic moments [47]. It is worth exploring the structural information and magnetic behaviors of Cr atoms in larger-sized clusters. Based on this, we conducted a systematically theoretical investigation on the structural evolution, electronic and magnetic properties of anionic and neutral CrSi_n clusters for $n = 19-25$.

Computational method

The critical step in theoretical study of a cluster is to obtain its reliable structure. Therefore, low-lying structures of $\text{CrSi}_n^{-/0}$ ($n = 19-25$) clusters were globally searched through a self-developed genetic algorithm code incorporated with the ORCA 5.0.4 software [60, 61] for DFT calculations. For each size, more than 3000 configurations were generated by the genetic algorithm code to fully search on the potential energy surface. The validity and efficiency of our global optimization algorithm code have been well demonstrated by several of our recent works about TMSi_{14-20}^- (TM = V₁₋₃, Cr₁₋₂) [46–49], TMGe_{8-17}^- (TM = Ti-Ta) [62], and $\text{In}_{3-16}\text{X}_{0,1}^-$ (X = Si, Ge) [63], $\text{Cr}_2\text{Ge}_{15-20}^-$ [64], CsSi_{5-16}^- [65], and Ge_{4-30}^- [66] clusters. Considering that the scheme of BP86 functional with Karlsruhe-type basis sets has proven to be suitable for the description of structural evolution and electronic properties of $\text{Cr}_{1-2}\text{Si}_{14-20}^-$ and $\text{Cr}_2\text{Ge}_{15-20}^-$ clusters in our previous works [47, 49, 64], therefore, this scheme was also adopted for the calculations of $\text{CrSi}_n^{-/0}$ ($n = 19-25$) clusters. Briefly, the def2-SVP basis set [67, 68] and the BP86 functional [69, 70] were adopted for DFT calculations during the global search. Then the higher-quality def2-TZVP basis set [67, 68] was employed to further optimize the top 10–20 candidate isomers to get more accurate geometric structures, and the diffuse def2-TZVP basis set (def2-TZVPD) [71] was adopted to obtain more accurate energies. All the structures have positive vibrational frequencies

from vibration analyses. Zero-point-energy corrections were considered for the energy calculations. Various spin multiplicities (SM) were considered in order to obtain the lowest-energy spin state. The vertical detachment energy (VDE) was obtained from the energy differences between the anionic and neutral clusters at the relaxed structure of the anionic state. The adiabatic detachment energy (ADE) were calculated as the energy differences between the anionic cluster and the relaxed neutral clusters using the anionic structure as initial configuration.

The stabilities of anionic and neutral CrSi_n cluster are evaluated by the average binding energies (E_b), defined as Eqs. (1) and (2) [72, 73]:

$$E_b(\text{CrSi}_n^-) = [(n-1)E(\text{Si}) + E(\text{Si}^-) + E(\text{Cr}) - E(\text{CrSi}_n^-)]/(n+1) \quad (1)$$

$$E_b(\text{CrSi}_n) = [nE(\text{Si}) + E(\text{Cr}) - E(\text{CrSi}_n)]/(n+1) \quad (2)$$

where $E(\text{CrSi}_n^-)$ and $E(\text{CrSi}_n)$ is the energies of the anionic and neutral CrSi_n clusters, respectively; $E(\text{Si})$, $E(\text{Si}^-)$ and $E(\text{Cr})$ are the total energies of the neutral Si, anionic Si and neutral Cr atoms, respectively. The second order of energy difference (Δ_2E) of a $\text{CrSi}_n^{-/0}$ was calculated using the formula (3):

$$\Delta_2E(\text{CrSi}_n^{-/0}) = E(\text{CrSi}_n^{-/0}) + E(\text{CrSi}_n^{-/0}) - 2E(\text{CrSi}_n^{-/0}) \quad (3)$$

All the graphs of the structures of clusters involved in this work were rendered by using the molecular graphics program VMD 1.9.3 [74]. The average bond length, Wiberg bond order, and electron spin density were calculated by using the post processing program of Multiwfn 3.8 (dev) [75].

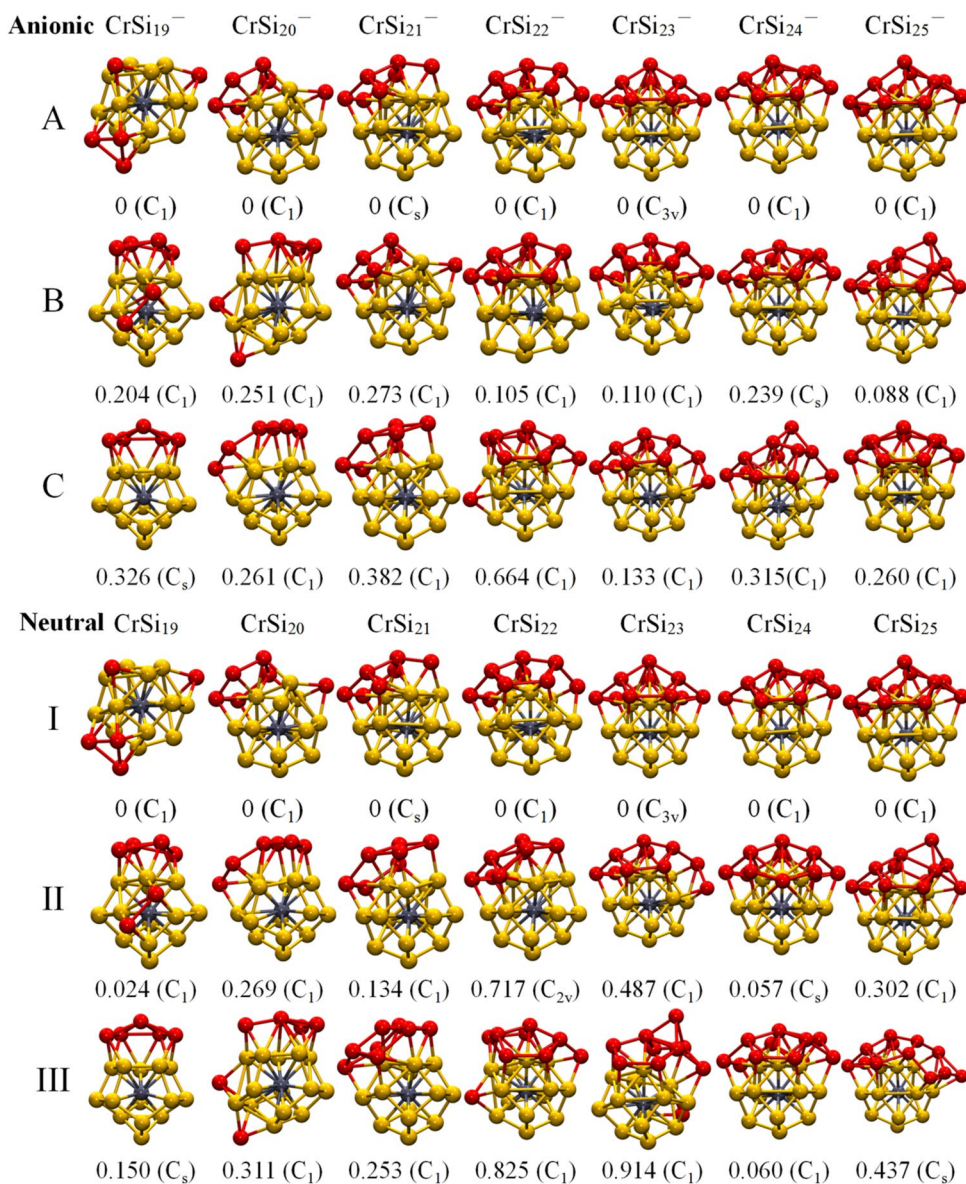
Computational results

Structures of $\text{CrSi}_n^{-/0}$ ($n = 19-25$) clusters

The optimized structures of first three lowest-lying isomers of anionic and neutral CrSi_{19-25} at the BP86/def2-TZVP level are displayed in Fig. 1. The calculations show that there are four typical endohedral structures (see Fig. 2) as the motifs in these low-lying isomers of $\text{CrSi}_{19-25}^{-/0}$ clusters. These motifs are also adopted for TMSi_{14} (TM = Mn, Fe, Co) [76], TiSi_{17-20}^- [10], and VSi_{16-20}^- [46], and CrSi_{16-18}^- [49] clusters.

For anionic CrSi_n^- clusters, the lowest-energy isomers ($n\text{As}$) prefer to adopt endohedral structures (motif I (Cr@Si₁₄) for sizes $n = 19-21$, and motif IV (Cr@Si₁₃) for larger ones) as the structural motif with the extra Si atoms attached on the surface. The lowest-energy structures of CrSi_{19}^- and CrSi_{20}^- also are adopted for VSi_{19}^- and VSi_{20}^- clusters, respectively [46]. Motif I is adopted for isomers 21B, 21C,

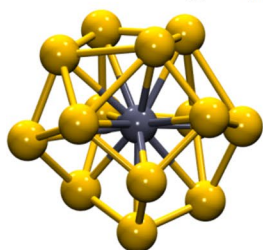
Fig. 1 Geometric structures of low-energy isomers for $\text{CrSi}_{19-25}^{-/0}$ clusters. For each size, the total energy difference (eV) with respect to the lowest-lying isomer is provided below the structure, and the symmetry is given inside the parentheses. Gray balls represent chromium atoms. The silicon atoms formed a cage and adsorbed on the surface are highlighted as yellow and red, respectively



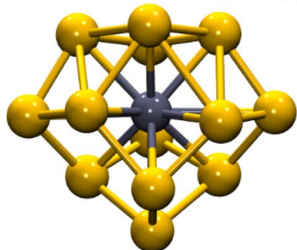
and 22B as the core structural motif with excess Si atoms attached to the surface, showing the strong dominance in size range of $n=19-21$. Starting from size 22, the dominance of

motif I begins to weaken, instead, motif IV begins to show high competitiveness, which is adopted for all the first three isomers for sizes $n=23-25$. The motif of isomer 19B is also

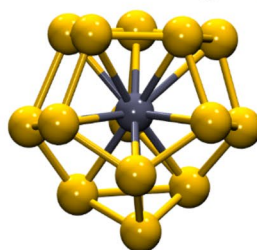
Motif I: $\text{Cr}@Si_{14}$



Motif II: $\text{Cr}@Si_{14}$



Motif III: $\text{Cr}@Si_{13}$



Motif IV: $\text{Cr}@Si_{13}$

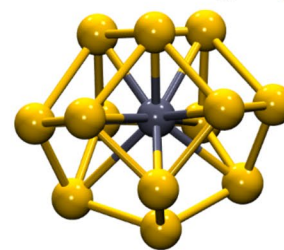


Fig. 2 Four typical structural motifs present in isomers of CrSi_n clusters with $n=19-25$ taken from isomers 19A, 22A, 19B, and 20C, respectively. Motifs II and III can be generated from motifs I and IV by removing one of the top Si atoms, respectively

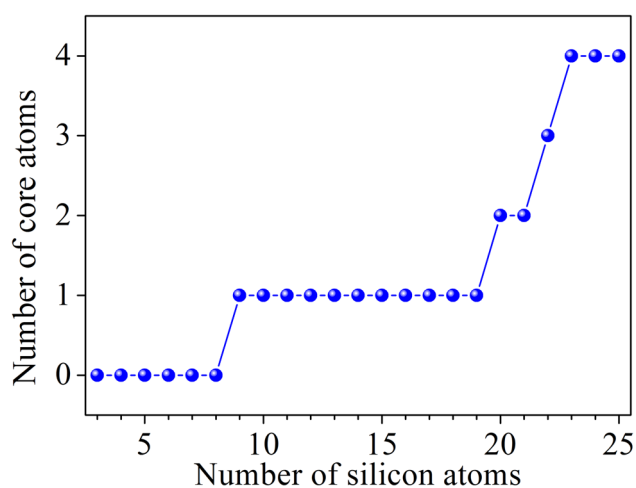


Fig. 3 A plot showing the core evolution in the lowest-energy isomers of CrSi_n ($n=3-25$) clusters. Gray and yellow balls represent chromium and silicon atoms, respectively. The results for sizes $n=3-18$ are from previous works [36, 49, 78]

an endohedral Cr@Si_{13} structure (motif II), which is also adopted for TMSi_{14} (TM = Mn, Fe, Co) clusters [76]. The motif III always represents a building block (Cr@Si_{14}) for the ground-state TMSi_n (TM = Ti^- , V^- and Cr^- for $n=17, 18$ [10, 44, 49]) clusters, but is only adopted by isomers 19C, 20B and 20C among all these isomers for $n=19-21$, indicating its weak competitiveness in large size clusters.

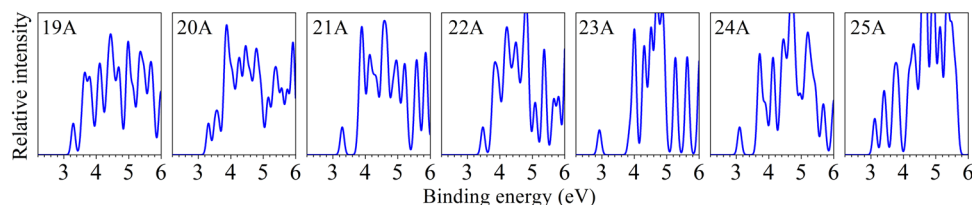
For neutral CrSi_n ($n=19-25$) clusters, their structures are very similar to that of anionic states. All the lowest-energy isomers of neutral CrSi_n ($n=19-25$) share the same configurations as that of anionic state. The calculated results (see Table S3 of Supplementary materials) show that all the root mean square deviations (RMSD) of anionic and neutral states are small with less than 0.13 Å, except for sizes 20 and 22 with values of 0.221 Å and 0.215 Å, respectively. This result indicates that the reduction of one electron in an anionic system has little effect on its geometric structure. Isomers 19II and 19III of neutral CrSi_{19} adopt the same configurations as 19B and 19C of anionic state, respectively. Isomers 20II, 20III, 21II, 22III, 23II, 24III, and 25II, adopt the same configurations as 20C, 20B, 21C, 22C, 23C, 24B, and 25B, respectively. Isomers 22II and 23III adopt the motif I, while 24II and 25III are based on the motif IV with excess Si atoms adsorbed on the surface.

For sizes $n=21-25$ for CrSi_n^{-0} clusters, each successive cluster size is composed of its predecessor with an extra Si atom adsorbed onto the surface, giving evidence of a stepwise growth. Moreover, the core also increases with the cluster size (see Fig. 3). Single Cr atom as core first observed in CrSi_9 still exists in CrSi_{19} , but for larger sizes, the clusters prefer to grow by core expansion. Both of the global minimum CrSi_{20} and CrSi_{21} possess a two-atom core, while CrSi_{22} has a three-atom core. The four-atom core of CrSi_{23} is a tetrahedral bipyramid and is persistent in the global minima of sizes $n=23-25$. Significantly, the structure of CrSi_{23} is very similar to the smallest core-shell silicon cluster Si_{27}^- [77], indicating that Cr atom accelerate the formation of cage-like structure for pure silicon cluster. From another point of view, the influence of Cr atom on the structure of pure silicon cluster starts to weaken.

Photoelectron spectra of CrSi_n^- ($n=19-25$) clusters

Figure 4 shows the simulated PES for the lowest-energy isomers of CrSi_n^- ($n=19-25$) cluster anions. For CrSi_{19}^- , the spectrum possesses six distinguishable peaks at 3.29, 3.65, 3.79, 4.11, 4.44 and 4.69 eV. In the PES of CrSi_{20}^- , there are two small peaks at 3.31 and 3.57 eV, followed by four prominent peaks located at 3.87, 4.26, 4.46, and 4.79 eV. For CrSi_{21}^- , the first peak locates at 3.29 eV, followed by three high-intensity peaks centered at 3.89, 4.13, and 4.59 eV. Similar spectrum is also observed in CrSi_{22}^- , the PES also possesses a clear small peak at 3.48 eV, followed by three prominent peaks located at 3.85, 4.21, and 4.49 eV. In the case of CrSi_{23}^- , there is a large separation between the highest occupied state and (2.92 eV) the next lower-lying state (3.87 eV), indicating a large gap between the highest occupied state (HOMO) and the lowest unoccupied state (LUMO) of the corresponding neutral cluster. This result corresponds to the large HOMO–LUMO gap of 1.50 eV for neutral CrSi_{23} . Similar to CrSi_{23}^- , the simulated spectrum of CrSi_{24}^- also has a small peak at 3.10 eV and three prominent peaks situated at 3.73, 4.13, and 4.45 eV. For CrSi_{25}^- , four discrete peaks are located at 3.12, 3.41, 3.77, and 4.30 eV, respectively. In short, for the simulated PES of CrSi_n^- ($n=19-25$) clusters, the spectral characteristics and peak positions are obviously different, indicating their significant differences in electronic properties. The

Fig. 4 Theoretical photoelectron spectra of the lowest-lying isomers of CrSi_n^- ($n=19-25$). Gaussian broadening with a width of 0.06 eV is used



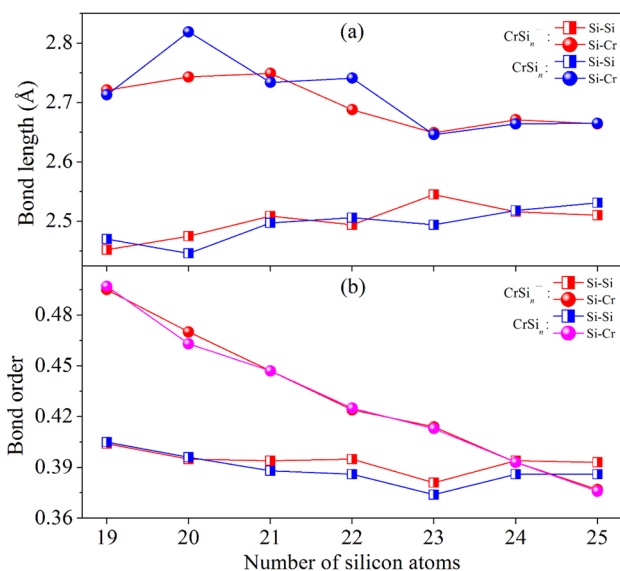


Fig. 5 Average bond lengths (a) and Wiberg bond orders (b) of Si-Si and Si-Cr of the lowest-energy structures of $\text{CrSi}_n^{-/0}$ ($n=19-25$) clusters

simulated PESs are expected to provide some guidance for future PES measurements.

Bonding and electronic properties of $\text{CrSi}_n^{-/0}$ ($n=19-25$) clusters

To further explore the bonding and electronic properties of the $\text{CrSi}_n^{-/0}$ ($n=19-25$) clusters, the average bond lengths, Wiberg bond orders, vertical detachment energies (VDEs), adiabatic detachment energies (ADE), HOMO–LUMO gaps, average binding energies (E_b), and the second order of energy difference (Δ_2E) for the lowest-energy structures were calculated and presented in Figs. 5 and 6 as well as summarized in Table 1.

From Fig. 5, the average bond length and bond order of Si–Si and Si–Cr bonds in anionic state are close to those in the neutral state. From Fig. 5(a), the average Si–Cr bond lengths (2.649 Å–2.819 Å) have significantly larger values than that of Si–Si bonds in the range of 2.446 Å–2.545 Å. The Wiberg bond orders (Fig. 5b) of Si–Cr bonds show a gradual decrease trend for $\text{CrSi}_n^{-/0}$ ($n=19-25$) clusters in

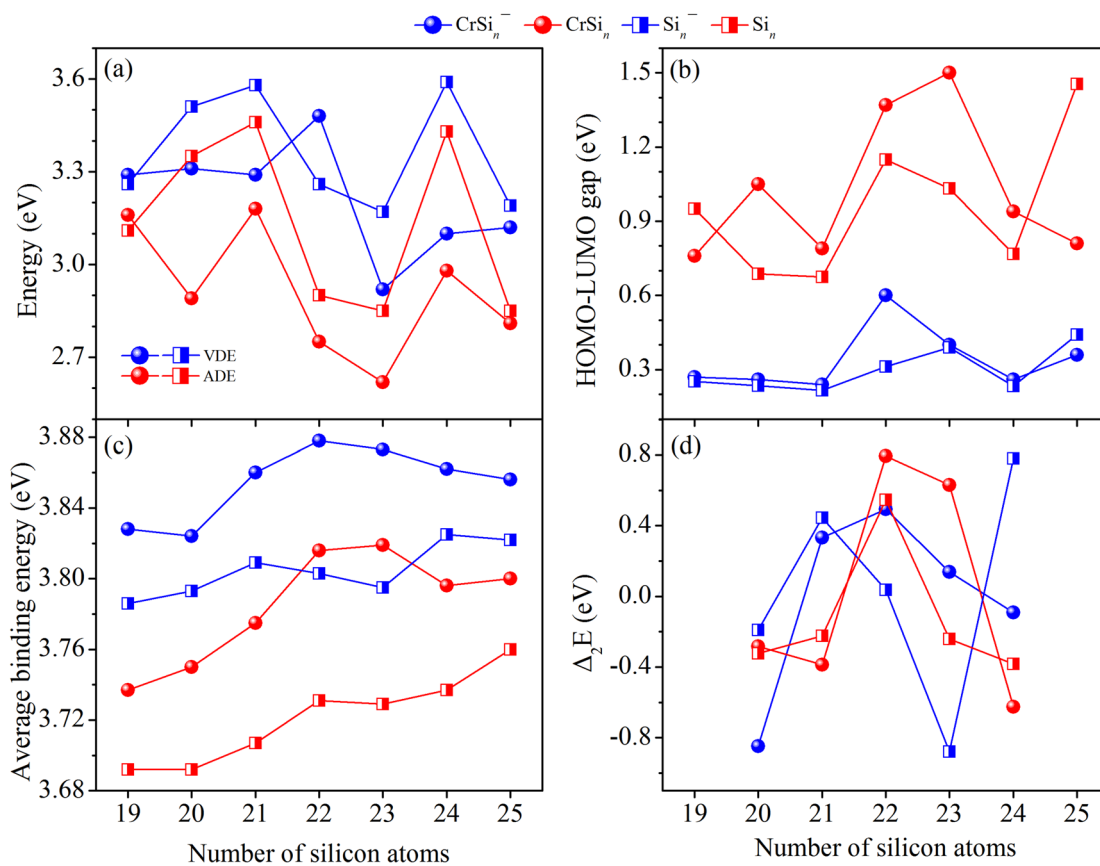


Fig. 6 Size-dependent vertical detachment energies (VDE), adiabatic detachment energies (ADE), HOMO–LUMO gaps, binding energies (E_b), and second order of energy difference (Δ_2E) of the lowest-lying

structures of $\text{CrSi}_{19-25}^{-/0}$ clusters with the results of $\text{Si}_{19-25}^{-/0}$ for comparison. The ground-state structures of Si_{19-25} clusters are from previous works [77, 79]

Table 1 Spin multiplicities (SM), average bond lengths (BL, in Å), bond orders (BO), vertical detachment energies (VDE, in eV), adiabatic detachment energies (ADE, in eV), HOMO–LUMO gaps (E_{HL} , in eV), average binding energies (E_b , in eV), and the second order of energy difference (Δ_2E , in eV) of the lowest-energy structures of $\text{CrSi}_n^{-/0}$ ($n=19\text{--}25$) clusters

Size	SM	BL		BO		VDE	ADE	E_{HL}	E_b	Δ_2E
		Si–Si	Si–Cr	Si–Si	Si–Cr					
CrSi_{19}^-	2	2.452	2.721	0.404	0.495	3.29		0.27	3.828	
CrSi_{20}^-	2	2.475	2.743	0.395	0.470	3.31		0.26	3.824	-0.848
CrSi_{21}^-	2	2.509	2.749	0.394	0.447	3.29		0.24	3.860	0.333
CrSi_{22}^-	2	2.494	2.688	0.395	0.424	3.48		0.60	3.878	0.494
CrSi_{23}^-	2	2.545	2.649	0.381	0.414	2.92		0.40	3.873	0.138
CrSi_{24}^-	2	2.516	2.671	0.394	0.393	3.10		0.26	3.862	-0.092
CrSi_{25}^-	2	2.510	2.664	0.393	0.377	3.12		0.36	3.856	
CrSi_{19}	1	2.470	2.713	0.405	0.497		3.16	0.76	3.737	
CrSi_{20}	1	2.446	2.819	0.396	0.463		2.89	1.05	3.750	-0.283
CrSi_{21}	1	2.497	2.734	0.388	0.447		3.18	0.79	3.775	-0.387
CrSi_{22}	1	2.506	2.741	0.386	0.425		2.75	1.37	3.816	0.794
CrSi_{23}	1	2.494	2.646	0.374	0.413		2.62	1.50	3.819	0.631
CrSi_{24}	1	2.518	2.664	0.386	0.393		2.98	0.94	3.796	-0.625
CrSi_{25}	1	2.531	2.665	0.386	0.376		2.81	0.81	3.800	

range of 0.495–0.376, while that of Si–Si bond fluctuate in a small size range of 0.374–0.405.

It shows that the VDEs (Fig. 6a) of CrSi_n^- are very close for $n=19$ to 21, rise to 3.48 eV at $n=22$, and then abruptly drop to 2.92 eV at $n=23$, gradually rising up again from $n=24$ to 25. The trend of ADEs of CrSi_n^- is roughly consistent with that of VDE, except for sizes 20 and 22, which corresponds to the large root mean square deviations (RMSD, see Fig. S3) values of anionic and neutral structures. The trend of VDE of Si_n^- is roughly consistent with that of ADE for Si_n^- . All the ADE and VDE values for neutral and anionic Si_n^- are bigger than that of CrSi_n^- clusters, except for sizes 19 and 22, respectively. This result indicates that Cr atoms can usually reduce the detachment energy of silicon cluster.

The trend of the HOMO–LUMO gaps for CrSi_n^- (in Fig. 6b) is similar to that of the VDE, except for size 23. The HOMO–LUMO gap of neutral CrSi_n^- and Si_n^- clusters are significantly larger than that of anionic states. This is because there are unpaired electrons in the anionic states. It is worth noting that the neutral CrSi_{22} and CrSi_{23} have large HOMO–LUMO gaps of 1.37 eV and 1.50 eV, respectively. All the HOMO–LUMO gap values for neutral and anionic CrSi_n^- are bigger than that of Si_n^- clusters, except for sizes

19 and 25. It is worth noting that the influence of Cr atoms on the HOMO–LUMO gap value of Si_n^- is generally very small, but for the neutral state.

From Fig. 6c, the average binding energies (E_b) exhibit similar trends in both anionic and neutral CrSi_n^- and Si_n^- clusters. E_b of CrSi_n^- for sizes 19 and 20 are significantly lower than that of larger sizes, increase gradually until coming to the maximum at $n=22$, and then decrease monotonously. For neutral CrSi_n^- clusters, the largest value of E_b occurs in size 23. Compared with pure Si_n^- clusters, CrSi_n^- clusters have a larger average binding energy, which is consistent with our expected results that doping transition metals can improve the stability of pure silicon clusters.

Interestingly, the trends of the Δ_2E (Fig. 6d) of CrSi_n^- and Si_n^- clusters are consistent with that of average binding energy curves. CrSi_{22}^- and CrSi_{23} also possess the largest Δ_2E values, indicating the largest relative stabilities. For anionic and neutral Si clusters, their local maximum values occur at sizes 21 and 22, respectively. Significantly, the CrSi_{22} and CrSi_{23} clusters have the considerable HOMO–LUMO gap, E_b , and Δ_2E values in the neutral and anionic states, respectively. The results indicate that these two clusters have high stabilities among these clusters.

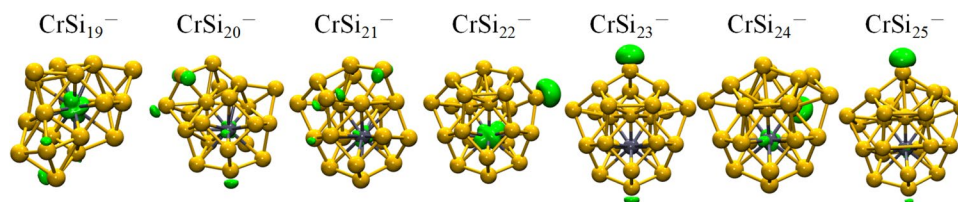


Fig. 7 Isosurface maps of the electron spin density of the lowest-energy $\text{CrSi}_{19\text{--}25}^-$ clusters. The isosurface is set to $\pm 0.006 e/\text{\AA}^3$. Green and red isosurfaces indicate the positive and negative electron spin density, respectively

Table 2 The total magnetic moments (μ_T) of the CrSi_{19-25}^- clusters were obtained by Hirshfeld population analysis, along with the local magnetic moment on the Cr atom (μ_{Cr}) and the Si atom with the maximum value ($\mu_{\text{Si-Max}}$). All magnetic moments are in μ_B

Cluster	μ_T	μ_{Cr}	$\mu_{\text{Si-Max}}$
CrSi_{19}^-	1	0.31	0.13
CrSi_{20}^-	1	0.07	0.18
CrSi_{21}^-	1	0.08	0.16
CrSi_{22}^-	1	0.55	0.40
CrSi_{23}^-	1	-0.04	0.36
CrSi_{24}^-	1	0.12	0.26
CrSi_{25}^-	1	0.01	0.34

Magnetic properties of CrSi_n^- ($n = 19 - 25$) clusters

Our results (see Tables S1 and S2 in Supplementary materials) show that all these of anionic and neutral CrSi_{19-25} clusters are in their lowest spin states, therefore, we just explore the magnetic properties of the CrSi_{19-25}^- cluster anions by the calculations of spin electron density and Hirshfeld population. The results are displayed in Fig. 7 and listed in Table 2, respectively. From Fig. 7, one can see that the green zones (the excess alpha electrons) appear around the Cr atom and some Si atoms with low coordination number, while the red regions (the excess beta electrons) are hardly distributed. Different from the results that Cr atoms in small-sized CrSi_n^- ($n \leq 18$) clusters [49] contribute most of the total magnetic moments of $1 \mu_B$, their contributions in larger clusters are not significant, even less than that of some Si atoms, except for size 22 with $\mu_{\text{Cr}} = 0.55 \mu_B$. This result indicates that the source of magnetic moment changes obviously for larger sized clusters studied here, which may be related to the saturation of Cr atoms.

Conclusions

The structural evolution, electronic and magnetic properties of $\text{CrSi}_{19-25}^{-/0}$ clusters have been computationally investigated. Global research for of the minimum structures of these clusters have been performed on a self-developed genetic algorithm code combined with DFT calculations. All the anionic CrSi_n^- clusters share the same configurations as that of their neutral states for $n = 19-25$. All these $\text{CrSi}_{19-25}^{-/0}$ clusters prefer to adopt endohedral structures ($\text{Cr}@\text{Si}_{14}$ for sizes $n = 19-21$, and $\text{Cr}@\text{Si}_{13}$ for $n = 22-25$) as the structural motif with the remaining Si atoms attached on the surface. The simulated photoelectron spectra show that these clusters have obvious size dependence. From the average binding energy results, the doping Cr atoms can significantly improve the stabilities of the pure silicon clusters. Among all these clusters, both of CrSi_{22} and CrSi_{23} possess large HOMO – LUMO gaps, average binding energies, and the second order of energy differences, indicating their high stabilities. All of the clusters possess total magnetic moment

of $1 \mu_B$, but with very different contributions from that of small sizes ($n \leq 18$), except for size 22.

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Author contributions K.W. wrote the manuscript. K. W., C. W., L. L., Y. W., J. C., J. Z. and J. G. analyzed the data and discussed the results. K. W. conceived the project, designed and supervised the whole research. K.W. revised and edited the manuscript. All authors contributed to the manuscript preparation.

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Availability of data and materials See the supplementary material for relative energies (Tables S1 and S2) of the $\text{CrSi}_{19-25}^{-/0}$ clusters anion for three different spin multiplicities (SM). The root mean square deviations of the lowest-energy isomers of anionic and neutral CrSi_{18-25} clusters were summarized in Table S3. The theoretical photoelectron spectra of the top three lowest-lying isomers of CrSi_n^- ($n = 19-25$) were shown in Fig. S1. Mulliken population analysis results of magnetic moments for CrSi_n^- ($n = 19-25$) were summarized in Table S4. The natural electron configuration of Cr atoms for CrSi_{19-25}^- clusters were shown in Table S5. The comparison of photoelectron spectra for CrSi_{19-25}^- and Si_{19-25}^- clusters were shown in Fig. S3. The calculated data of electronic properties for Si_{19-25}^- and Si_{19-25} were summarized in Table S6. Relative energies of the top three lowest-energy structures of $\text{CrSi}_n^{-/0}$ ($n = 19-25$) for the three different functionals were listed in Table S7.

Declarations

Ethical approval This declaration is not applicable.

Competing interests The authors declare no competing interests.

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