

Loss and Fractionation of Noble Gas Isotopes and Moderately Volatile Elements from Planetary Embryos and Early Venus, Earth and Mars

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Abstract Here we discuss the current state of knowledge on how atmospheric escape processes can fractionate noble gas isotopes and moderately volatile rock-forming elements that populate primordial atmospheres, magma ocean related environments, and catastrophically outgassed steam atmospheres. Variations of isotopes and volatile elements in different planetary reservoirs keep information about atmospheric escape, composition and even the source of accreting material. We summarize our knowledge on atmospheric isotope ratios and discuss the latest evidence that proto-Venus and Earth captured small H₂-dominated primordial atmospheres that were lost by hydrodynamic escape during and after the disk dispersed. All relevant thermal and non-thermal atmospheric escape processes that can fractionate various isotopes and volatile elements are discussed. Erosion of early atmospheres, crust and mantle by large planetary impactors are also addressed. Further, we discuss how moderately volatile elements such as the radioactive heat producing element 40 K and other

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rock-forming elements such as Na can also be outgassed and lost from magma oceans that originate on large planetary embryos and accreting planets. Outgassed elements escape from planetary embryos with masses that are $\leq M_{\text{Moon}}$ directly, or due to hydrodynamic drag of escaping H atoms originating from primordial- or steam atmospheres at more massive embryos. We discuss how these processes affect the final elemental composition and ratios such as K/U, Fe/Mg of early planets and their building blocks. Finally, we review modeling efforts that constrain the early evolution of Venus, Earth and Mars by reproducing their measured present day atmospheric ³⁶Ar/³⁸Ar, ²⁰Ne/²²Ne, noble gas isotope ratios and the role of isotopes on the loss of water and its connection to the redox state on early Mars.

Keywords Protoplanetary disk · Primordial atmospheres · Steam atmospheres · Atmospheric escape · Noble gases · Isotopes · Magma oceans · Rock-forming elements · Planetary evolution

1 Introduction

The earliest atmospheric records of Venus, Earth and Mars are stored in the noble gases and their isotopes. The reason is that these elements do not react with other atmospheric species or the surface. Noble gas elements such as Ne, Ar, Kr and Xe and their isotopes contain historical records related to the early evolutionary processes starting during the accretion phase of planetesimals to planetary embryos before and after disk dispersal (Ozima and Podosek 2002; Porcelli et al. 2002).

Depending on how fast planetary embryos grow to larger masses within the disk and at which orbital locations they originate or migrate, embryos will evolve to objects with different volatile and elemental abundances, which results in different elemental ratios compared to their initial ones (Lammer et al. 2020a; this issue; Benedikt et al. 2020). Heating by shortlived ²⁶Al and ⁶⁰Fe radioisotopes shape the thermal history and interior structure of growing planetesimals and planetary embryos during the early stages of planetary formation (e.g., Lichtenberg et al. 2016; Young et al. 2019; Benedikt et al. 2020). The subsequent thermomechanical evolution, such as magma ocean-related internal differentiation, rapid volatile degassing, and the possible accumulation of H_2 /He-envelopes has important implications for the final structure, composition and evolution of terrestrial protoplanets (Ikoma and Genda 2006; Elkins-Tanton 2012; Massol et al. 2016; Stökl et al. 2016; Lichtenberg et al. 2016; Benedikt et al. 2020). Under these extreme conditions, noble gases with low condensation temperatures such as Ar and Ne, significant amounts of major rock-forming elements (e.g., K, Si, Mg, Fe, Ca, Al, Na, S, P, Cl), and other moderately volatile elements like Rb or Zn with condensation temperatures of ≈ 1000 K and lower (Lodders et al. 2009) populate the hot accretionary atmospheres to significant amounts (Albarède and Blichert-Toft 2007; Schaefer and Fegley 2007; Fegley et al. 2016).

If one considers for instance a mixture of ordinary chondrites with enstatite-like material and/or carbonaceous chondrites with ureilites as recently suggested by Schiller et al. (2018) and as indicated by specific isotopic abundances to be the parent bodies of the Earth, then the terrestrial K/U ratio would be higher than the present-day ratio. This indicates that most likely some loss process during proto-Earth's accretion was necessary to deplete moderately volatile elements and fractionate K from U. Since radioactive decay of the heat producing element ⁴⁰K contributes to drive a long-lived magnetic dynamo and influences the long-term temperature evolution of a planet's interior (e.g. Turcotte and Schubert 2002; Murthy et al. 2003; Nimmo 2007; O'Neill et al. 2020; this issue), it is very important to investigate how atmospheric escape can modify the amount of ⁴⁰K and other moderately volatile elements such as Si, Mg, Fe, Ca, Al, Na, S, P, Cl, Rb, Zn until an Earth-like planet finishes its accretion.

It is currently believed that magma oceans played a large role in the early stages of terrestrial planet formation. Evidence that supports this hypothesis falls into four different categories:

- (i) Samples found on the Moon, Mars and Vesta imply widespread silicate melting followed by fractional crystallization (e.g. Wakita and Schmitt 1970; Foley et al. 2005).
- (ii) Ancient zircons have crystallization ages that fall into the time of planetary formation, which can be interpreted as accretion happening in a compressed timescale. As a result, heat production is more likely to cause melting instead of being radiated away into space (Harrison 2009).
- (iii) Short accretionary timescales are also supported by ¹⁸²W and ¹⁴²Nd contents entrapped in crustal and mantle rocks, since those require early core formation (e.g. Foley et al. 2005; Caro et al. 2006).
- (iv) Metallic core formation requires some kind of magma ocean (Terasaki et al. 2005).

According to magma ocean models, the slope of the adiabat in the pressure-temperature space is steeper than the solidus, they intersect at depth and the magma ocean starts to solidify upwards from the bottom (e.g., Elkins-Tanton 2012; Salvador et al. 2017). Cold plumes descend from the surface, resulting in a hotter return flow from the interior. In this way the solid region at the bottom grows as the magma ocean cools. The volatiles and trace gases contained in the magma ocean have low solid to liquid partition coefficients and as such do not crystallize (Ikoma et al. 2018). Hence, the liquid phase becomes more and more enriched in these elements as the magma ocean solidifies. Once the saturation limit for a given species is reached, it is outgassed catastrophically from the magma ocean into the atmosphere.

Although planetary embryos grow within the protoplanetary disk (Wade and Wood 2005), these bodies are differentiated and depleted in moderate volatile elements (Hin et al. 2017; Young et al. 2019) due to the loss of a magma ocean-related catastrophically outgassed steam atmospheres (Odert et al. 2018; Benedikt et al. 2020) as soon as they are not surrounded by the nebula gas anymore.

Hydrodynamic atmospheric escape, which is controlled by planetary parameters (e.g. Lammer et al. 2016; Owen and Wu 2016; Guo 2019) and the EUV flux of their young host stars (e.g. Koskinen et al. 2013; Erkaev et al., 2014, 2015, 2016; Lammer et al. 2014; Odert et al. 2018; Kubyshkina et al. 2018a; Guo 2019; Zahnle et al. 2019; Benedikt et al. 2020), can be very efficient if the mass of the protoplanetary body is small and the main upper atmospheric species is hydrogen. If hydrogen, which either originates from gas accretion in the disk or from dissociation of H_2O vapour of a magma ocean-related outgassed steam atmosphere, escapes efficiently, H atoms can drag along heavier species to space (Zahnle and Kasting 1986; Hunten et al. 1987; Zahnle et al. 1990, 2019; Chassefière 1996a, 1996b; Odert et al. 2018; Guo 2019; Benedikt et al. 2020). The conditions under which this so-called post-nebula atmospheric escape and volatilization processes take place, further lead to the modification of volatile composition and elemental abundances, as well as to the odd elemental and isotopic ratios that are observed in the Solar System planets (e.g. Gillmann et al. 2009; Marty 2012; Jellinek and Jackson 2015; Bonsor et al. 2015; Carter et al. 2015; Benedikt et al. 2020b).

The main aim of this work is to review the latest knowledge on various atmospheric evolution and escape processes that fractionate isotopes and elements during Venus', Earth's and Mars' atmosphere evolution since their accretion phase in the disk. Section 2 summarizes our knowledge on the main atmospheric isotopic species and the composition of terrestrial planets obtained from various space missions. In Sect. 3 we discuss the evidence that accreting proto-Venus and Earth captured primordial atmospheres during their growth within the disk. Here we also discuss briefly the new knowledge obtained from exoplanet research and the discovery of H₂-dominated low mass exoplanets. In Sect. 4 we briefly review various thermal and non-thermal atmospheric loss processes and discuss their role in isotopic fractionation of different atmospheric species. In Sect. 5 the role of impact erosion of primordial atmospheres and crustal material as well as the role of this process in the modification of volatile elements of accreting protoplanets are discussed. In Sect. 6 we address the loss and isotopic and elemental fractionation of rock-forming volatile elements and noble gases from Solidifying magma oceans from Moon- to Mars-like planetary embryos. In Sects. 7 and 8 we address the latest attempts for constraining the early evolution of Venus and Earth and Mars by reproducing their present noble gas and elemental ratios until these planets finished accretion. Section 9 concludes the review.

2 Atmospheric Isotopic Compositions of the Terrestrial Planets

Noble gases can be seen as a key to a planet's past because their abundances and isotope ratios represent the earliest record of atmospheres of terrestrial planets. Their evolved isotope ratios contain records of cataclysmic events that the evolving atmospheres experienced. These could be collisions with planetary embryos, comets, mass-fractionating atmospheric escape processes, outgassing from the planetary interior and geologic upheavals (Becker et al. 2003; Baines et al. 2013; Lammer et al. 2018).

Table 1 summarizes and compares the observed atmospheric isotope ratios in the atmospheres of Venus, Earth and Mars. The large Martian atmospheric ¹⁴N/¹⁵N disequilibrium (Fox and Hać 1997; Füri and Marty 2015) between the atmosphere and mantle and the similar but smaller disequilibrium on Earth (Cartigny and Marty 2013) combined with the fractionation of the significantly heavier Xe isotopes (e.g., Pepin 1991, 2000; Cassata 2017; Avice et al. 2018a, 2018b; Zahnle et al. 2019; Avice and Marty 2020; this issue) are a strong indication for different mass- and/or ionization-related atmospheric escape processes and atmosphere-magnetosphere environments of early Earth and Mars. These processes were shaped by the young Sun's radiation and plasma conditions (Lammer et al. 2018; and references therein). A detailed discussion on the nitrogen and Xe disequilibria in Earth's and Mars' atmosphere can be found in Avice and Marty (2020; this issue).

Additionally to Earth, icy moons such as Triton or the dwarf planet Pluto also possess N₂dominated atmospheres. For example, Titan has a 1.45 bar atmosphere with a N₂-dominated atmosphere (Strobel and Shemansky 1982; Coustenis and Taylor 2008). However, the origin of such atmospheres is most likely related to early photolysis of accreted and outgassed NH₃ from subsurface H₂O-NH₃ oceans with contributions from refractory organics (Miller et al. 2019), leading to a very different evolution compared to Earth's N₂ atmosphere (Coustenis and Taylor 2008; Mandt et al. 2009, 2014; Miller et al. 2019). The origin and evolution of atmospheres that are constrained by ¹⁴N/¹⁵N isotope ratios at frozen worlds like Titan, Triton, and Pluto are discussed in Scherf et al. 2020; this issue). Here we focus mainly on noble gases and moderate volatile rock-forming element evolution on planetary embryos, early Venus, Earth and Mars that originated in an inner planetary system.

The present atmospheric ³⁶Ar/³⁸Ar and ²⁰Ne/²²Ne noble gas isotope ratios can provide historical records on the growth rate of proto-Venus and Earth in the presence of the nebula (Lammer et al. 2020b) and hence their early atmospheric evolution, which is also the case for

Atmospheric isotope ratios	Venus	Earth	Mars
D/H	$(160 \pm 20) \times 10^{-4a}$ 0.064 - 0.08 (above 70 km altitude ^b)	$(1.49 \pm 0.03) \times 10^{-4}$ 1.56×10^{-4} (sea water ^d)	$(7.58-10.9) \times 10^{-4e}$ < 1.99×10^{-4} (interior ^f)
³ He/ ⁴ He	Not measured; expected: $< 3 \times 10^{-4g}$	1.37×10^{-6h}	Not measured
¹² C/ ¹³ C	88^{+2i}_{-1}	89 (inorganic ^h)	$85.1 \pm 0.3^{e,j}$
¹⁴ N/ ¹⁵ N	273_{-46}^{+70k}	272 ^h	173 ± 11^{1}
¹⁶ O/ ¹⁷ O	Not measured	2520 ^h	$2577 \pm 12^{e,j}$
¹⁶ O/ ¹⁸ O	$500 \pm 25^{\mathbf{k}}$	489 ^h	$462 \pm 2.5^{e,j}$
²⁰ Ne/ ²² Ne	11.8 ± 0.7^{a}	9.78 ^m	10.1 ± 0.7^{n}
²¹ Ne/ ²² Ne	Not measured; expected: < 0.067 ^a	0.029 ^m	Not measured ^h
³³ S/ ³² S	Not measured; expected: $\approx 8 \times 10^{-3h}$	8.01×10^{-3h}	Not measured ^h
³⁴ S/ ³² S	Not measured; expected: $\approx 0.04^{h}$	0.045 ^h	Not measured ^h
³⁶ Ar/ ³⁸ Ar	$5.6 \pm 0.6^{\circ}$	$5.32 \pm 0.33^{\mathrm{p}}$	4.2 ± 0.1^{n}
⁴⁰ Ar/ ³⁶ Ar	$\begin{array}{c} 1.03 \pm 0.04^k \\ 1.19 \pm 0.04^i \end{array}$	$298.56 \pm 0.31^{\text{q}}$	$(1.714 \pm 0.17) \times 10^{3e,j}$
¹³⁶ Xe/ ¹³² Xe	Not measured; expected: $\approx 1^{h}$	$0.3294 \pm 0.0004^{\mathrm{p}}$	$\begin{array}{c} 0.3451 \pm 0.0023^{p} \\ \text{ALH84001:} \ 0.3450 \pm 0.002^{r} \end{array}$
¹²⁹ Xe/ ¹³² Xe	Not measured; expected: $\approx 3^{h}$	$0.983 \pm 0.001^{\mathrm{p}}$	$\begin{array}{c} 2.522 \pm 0.006^{p} \\ \text{ALH84001:} \ 2.151 \pm 0.032^{r} \end{array}$

 Table 1
 Atmospheric isotope ratios of the terrestrial planets in the Solar System

^aDonahue et al. (1982), ^bBertaux et al. (2007), ^cLécuyer et al. (1998), ^dMichael (1988), ^eWebster et al. (2013a, 2013b), ^fHallis et al. (2012), ^gvon Zahn et al. (1983), ^hBaines et al. (2013), ⁱIstomin et al. (1980a), ^jAvice et al. (2018a, 2018b), ^kHoffman et al. (1980), ^lWong et al. (2013), ^mMarty (2012), ⁿPepin (1991), Atreya et al. (2013), ^oIstomin et al. (1980b), ^pOzima and Podosek (2002), ^qLee et al. (2006), ^rCassata (2017)

Mars (e.g., Jakosky et al. 1994, 2017; Kurokawa et al. 2018). However, there are yet some relevant uncertainties on some of the noble gas isotope ratios at these planets. For Venus, one should also note that due to large uncertainties, the measured value of 20 Ne/ 22 Ne-ratio is not precise enough for accurate scientific statements related to Venus' atmospheric evolution (von Zahn et al. 1983; Johnson and Oliveira 2019; Lammer et al. 2020b). New precise measurements by future spacecraft are necessary (Dandouras et al. 2020; this issue). It is also interesting to note that the 40 Ar/ 36 Ar ratio measured in the 4.1 Ga old Martian meteorite ALH84001 is \leq 1400 (Murty and Mohapatra 1997; Grady et al. 1998) and therefore lower compared to the modern atmospheric value that has been measured both in meteorites and in-situ to be \approx 1700 (Webster et al. 2013a, 2013b; Avice et al. 2018a, 2018b). One can expect that those ratios rise over time because of atmospheric escape and/or outgassing (Swindle 2002).

A reproduction of the primordial and steam atmosphere related D/H ratios yields information that may lead to a better understanding of the oxidation states and initial H_2O inventories of terrestrial planets. The terrestrial water (oceans) and meteorites are enriched by a factor of about 6 in D relative to the protosolar nebula value (Robert et al. 2000; Marty 2012; Pahlevan et al. 2019). Besides heterogeneities in the disk, mass fractionation during atmospheric escape processes from early terrestrial planets can enrich water in oceans in D, since the lighter isotopes are preferentially concentrated in the escaping upper atmosphere relative to ocean water (Genda and Ikoma 2008).

The preservation of a carbonaceous chondritic D/H signature with a mean statistical D/H ratio of $(150 \pm 10) \times 10^{-4}$ (e.g., Robert et al. 2000) in Earth's seawater (Marty 2012) provides evidence that early Earth's ocean underwent minimal D-enrichment via equilibration with H₂ followed by hydrodynamic escape. According to the recent results of Pahlevan et al. (2019) the close match of ≈ 10 - 20% between Earth's seawater and carbonaceous chondrites further constrains the prior existence of an atmospheric H₂-inventory of any origin on post-giant-impact Earth to < 20 bars. Moreover, the observed oxidation of silicate Earth occurred before or during the crystallization of the final magma ocean (Pahlevan et al. 2019). Unless D escape is very efficient on Venus, the present H escape flux (averaged over a solar cycle) cannot be larger than about 10^7 cm⁻² s⁻¹, if the present water vapour in Venus' atmosphere is a remnant of water deposited billion years ago (Donahue 1999). On the other hand, these authors argued that, if the escape flux remains constant and is as large as about 3×10^7 cm⁻² s⁻¹, water would be the remnant of H₂O outgassed only about 500 million years ago. If the hydrogen escape fluxes were $> 3 \times 10^7$ cm⁻² s⁻¹ during the past 100 Myrs then the present day D/H ratio in Venus' atmosphere could have been created since a massive resurfacing event a few 100 Myrs ago (Grinspoon 1993).

Hydrogen escape rates based on the loss of suprathermal H atoms and ions on present Venus were studied by Lammer et al. (2006). These authors modelled photochemically produced suprathermal H atom escape rates in the order of $\approx 3.8 \times 10^{25} \text{ s}^{-1}$ which is in agreement with the estimates of Donahue and Hartle (1992).

Ion escape rates were uncertain before the arrival of Venus Express, and model-based H⁺ escape rates in the order of $\approx 7.0 \times 10^{25} \text{ s}^{-1}$ largely overestimated the escape from Venus over the planets nightside by assuming ion acceleration by an outward electric polarization force related to ionospheric holes (Hartle and Grebowsky 1993). A recent study that investigated the H⁺ escape from Venus during 8 years of data from the Venus Express ion mass analyser instrument obtained much lower H⁺ ion loss rates in the order of $\approx 7.6 \times 10^{24} \text{ s}^{-1}$ at solar minimum and $\approx 2.1 \times 10^{24} \text{ s}^{-1}$ during solar maximum (Persson et al. 2018). These authors found that the decrease of the H⁺ ion escape from solar minimum to maximum is caused mainly by the change in the flow direction in the magnetotail. During solar maximum a significant Venus-ward ion flow close to the planet and its magnetotail was observed (Persson et al. 2018). Thus, the present day Venus total hydrogen escape is in the order of $\approx 4.5 \times 10^{25} \text{ s}^{-1}$ ($\leq 10^7 \text{ cm}^{-2} \text{ s}^{-1}$) with suprathermal H atom escape as the dominant process.

The findings of lower H^+ ion escape rates by Persson et al. (2018) compared to previous estimates indicates that Venus' present D/H ratio is most likely a remnant of the ancient evaporated water inventory and subsequently outgassed water vapour that were modified by various non-thermal atmospheric escape processes during time periods that lasted longer than the last resurfacing event which occurred about 200 - 1000 Myr ago (e.g., Strom et al. 1994; Basilevsky et al. 1997; Kreslavsky et al. 2015).

Other isotopes such as ${}^{16}O/{}^{17}O$, ${}^{16}O/{}^{18}O$ may reveal information related to a common kinship of the planets (e.g., Clayton 1993; Yurimoto et al. 2006), ${}^{33}S/{}^{32}S$ and ${}^{34}S/{}^{32}S$ isotopes contain evidence on the past and current planetary volcanic activity and magmatic composition (e.g., Mahaffy et al. 2012), outgassing from the interior is recorded by 4 He, 40 Ar and ${}^{129,131-136}$ Xe (e.g., Allegre et al. 1987; Marty and Allé 1994; Watson et al. 2007), and ${}^{12}C/{}^{13}C$ can be related to a possible biological activity (e.g., Galimov 1985, 2003). As one can see in Table 1, several isotope ratios have large error bars when their measurements were analysed in laboratory settings, or were not measured/detected by spacecraft in

the past. Therefore, new future in-situ measurements of these isotopes as discussed in Dandouras et al. (2020; this issue) in planetary atmospheres are necessary for constraining and understanding the state and evolution of terrestrial planets in the Solar System and beyond.

3 The Disk Phase: H₂-Dominated Envelopes on Low Mass Planets

In the Solar System, there are two well-separated classes of planets: rocky planets (Mercury, Venus, Earth, and Mars with average densities between 3.7 and 5.5 g cm⁻³) and gas/ice giants (Jupiter, Saturn, Uranus, and Neptune with average densities between 0.7 and 1.6 g cm⁻³). The rocky planets have a mass of about one Earth-mass (M_{Earth}), or less, while gas giants have masses of tens to hundreds of M_{Earth} . However, the largest population of exoplanets discovered to date has a mass/radius falling between those of the Earth and Neptune (e.g., Mullally et al. 2015), demonstrating that the variety of planets existing in nature is much larger than previously thought.

3.1 Low Mass Exoplanets with Primordial Atmospheres

The mass-radius, or equivalently mass-density, diagram is the first diagnostic enabling the basic classification of a planet, e.g. identify whether a planet hosts an H_2 -dominated atmosphere or not. The rapidly increasing sample of Neptune-like planets and smaller for which we have a measurement of both mass and radius indicates that these planets present a surprisingly large variety of densities (see Fig. 1).

This can be clearly seen in the left panel of Fig. 1 showing the position of the known exoplanets with a mass smaller than Neptune in the mass vs. radius diagram, color-coded by the planetary equilibrium temperature (T_{eq}). This plot clearly indicates the presence of a large spread in planetary radius, and thus density, at a given mass, particularly for planets with masses larger than 1.5-2.0 M_{Earth} . The same can be seen also when considering only systems for which masses and radii are known to better than 25% (Fig. 1, right panel), indicating that this large spread in bulk density is not the result of large uncertainties on the planetary parameters.

Although planetary cores with various densities/compositions certainly exist, this spread is most likely tightly related to the existence of a large variety of planetary atmospheres that can range from being H₂-dominated primordial atmospheres to CO_2/H_2O -dominated steam atmospheres, or secondary atmospheres, thus more compact. Furthermore, this spread does not seem to be directly correlated with the amount of irradiation they receive from the host star. The origin of this phenomenon is unclear and several factors are likely to play a role, such as protoplanetary nebula lifetime (e.g., Stökl et al. 2015, 2016), mass accretion time (e.g., Schiller et al. 2018), migration history, collisions (e.g., Bonomo et al. 2019), received amount of high-energy irradiation (e.g., Fossati et al. 2017; Owen and Wu 2016; Zahnle and Catling 2017; Jin and Mordasini 2018; Kubyshkina et al. 2018b, 2018c), possible biases in the measurement of planetary radii due to the presence of high-altitude aerosols (e.g., Lammer et al. 2016; Cubillos et al. 2017a, 2017b).

It has been shown that these planets are likely to have Earth-like density cores (Owen and Wu 2016) and the observations conducted to date clearly show that a number of them still host an H_2 -dominated atmosphere, most likely accreted from the protoplanetary nebula, despite their close distance to the host star. It is therefore possible that also the inner planets of the Solar System began their evolution with a rocky core surrounded by an H_2 -dominated envelope.

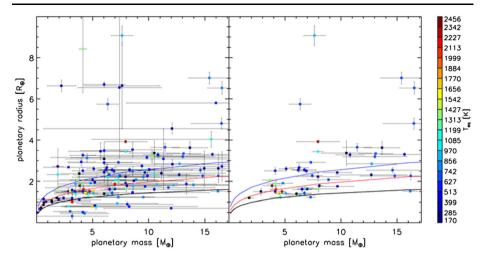


Fig. 1 Left: known exoplanets with masses smaller than Neptune in the mass vs radius plane, in Earth units. The color indicates the planetary equilibrium temperature T_{eq} . The black, red, and blue lines indicate compositions corresponding to 100% Fe, 100% rock, and 100% water, respectively (from Zeng et al. 2016). Right: same as left panel, but considering only planets with a mass and radius known to better than 25%. Data from the NASA exoplanet archive (https://exoplanetarchive.ipac.caltech.edu/)

3.2 Evidence of Captured Nebula Gas During Early Earth's and Venus' Accretion

Besides disk lifetime (e.g. Montmerle et al. 2006; Bollard et al. 2017; Wang et al. 2017) and the young Sun's X-ray, EUV and FUV fluxes and plasma winds (Tu et al. 2015; see also Güdel 2020; this issue), the accretion time is the relevant factor on whether a protoplanet captured enough mass to be able to accumulate and sustain an H₂-dominated envelope after the disk disappeared (e.g., Ikoma and Genda 2006; Lammer et al. 2016, 2020b; Stökl et al. 2015, 2016). In the case of the early terrestrial planets in the Solar System, scientists of various research areas debate since decades whether it was possible for proto-Venus and Earth to accrete enough mass during the disk lifetime for accumulating nebula gas before the disk disappeared. If a protoplanet reached a mass $> 0.5 M_{Earth}$ before the disk evaporated it could have captured a small H₂-dominated envelope that remains after the nebula dissipated (e.g., Ikoma and Genda 2006; Stökl et al. 2015, 2016; see also Lammer et al. 2020a; this issue).

Captured nebula gas was used in past attempts to reproduce the present atmospheric ${}^{36}\text{Ar}/{}^{38}\text{Ar}$ and ${}^{20}\text{Ne}/{}^{22}\text{Ne}$ noble gas ratios of Venus and Earth by various researchers (e.g., Sekiya et al. 1980a, 1980b; Sasaki and Nakazawa 1989; Harper and Jacobsen 1996; Pepin 1991, 2000, Porcelli and Pepin 2000; Porcelli et al. 2001; Gillmann et al. 2009; Odert et al. 2018; Lammer et al. 2020b). In the case of the Earth, there is data from ${}^{20}\text{Ne}/{}^{22}\text{Ne}$ isotope ratios in mantle material indicating that a fraction of noble gases has been trapped in Earth's interior in solar composition from the protoplanetary disk (Mizuno et al. 1980; Porcelli et al. 2001; Yokochi and Marty 2004; Mukhopadhyay 2012). Some researchers argued that these ${}^{20}\text{Ne}/{}^{22}\text{Ne}$ isotope ratios between 12.52-12.75 could have also been reproduced through implantation of solar wind onto accreted material (Trieloff et al. 2000; Podosek 2003; Feigelson et al. 2002; Raquin and Moreira 2009; Moreira and Charnoz 2016; Péron et al. 2016, 2017). Later on, such solar wind related noble gas isotopes could then be

outgassed during the magma ocean phase. A recent study, however, measured 20 Ne/ 22 Ne values up to 13.03 ± 0.04 from deep mantle plumes and determined a ratio for the primordial plume mantle of 13.23 ± 0.22 (Williams and Mukhopadhyay 2019). These values are only slightly fractionated but comparable with the nebular (solar) ratio. Even if the debate carries on, this can be seen as a robust evidence for a reservoir of nebular gas that is preserved in the deep mantle. If the accreting proto-Venus and Earth captured little H₂-dominated envelopes from the disk then these primordial atmospheres were lost during some time after disk dispersal.

Below we discuss various atmospheric escape processes, thier efficiencies and capabilities for the fractionation of various elements (i.e., isotopes and volatile elements, including rock-forming ones).

4 Atmospheric Escape Processes and Their Relevance for Elemental Fractionation

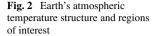
Losses of atmospheric gases to space are caused by a large number of physical processes taking place in the thermosphere, exosphere, and in the case of the Earth in the magnetosphere. These processes are mostly driven by inputs from the Sun in the forms of its high-energy radiative spectrum and the solar wind. For this reason, loss processes were more important during earlier phases of the Solar System when the Sun was more active.

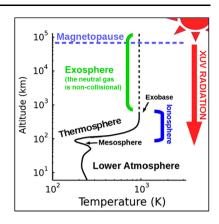
4.1 Upper Atmospheres

The upper atmospheres of planets are of primarily importance for determining how rapidly atmospheric gases are lost to space and for determining the relative losses of different elements and isotopes. Understanding how Earth's upper atmosphere has evolved is important for understanding the effects of losses on isotopic fractionation. Various atmospheric layers of interest are outlined in Fig. 2. The main layers in a planetary atmosphere can be defined by the temperature structure. On Earth, above the troposphere and stratosphere, starting at approximately 50 km altitude, is a region called the mesosphere; in this region, the temperature decreases with altitude due to cooling by CO_2 molecules emitting infrared radiation (IR) to space. Above this is the thermosphere, starting at approximately 100 km altitude, which is heated by the Sun's X-ray, extreme ultraviolet and ultraviolet radiation, and therefore has a positive temperature gradient. On Earth, only negligible amounts of CO_2 reach the thermosphere, hence no efficient IR-coolant as on the CO_2 -dominated atmospheres of Venus and Mars that more or less prevents the formation of an Earth-like thermosphere on these planets.

The top of the thermosphere is the exobase which is ≈ 500 km for the modern Earth but only about 210 km at present Venus and Mars (Lammer et al. 2018; and references therein). The exobase is the point above which the gas density is so low that the mean free paths of particles are long enough compared to the scale height of the atmosphere making the gas non-collisional. This non-collisional region, for the neutral gas above the exobase, is the exosphere. An additional important layer is the thermosphere-ionosphere region, which starts in the upper mesosphere and extends to the exobase. This is the domain in which most of the Sun's X-ray and extreme ultraviolet radiation is absorbed, which causes photoionization of the atmospheric constituents.

In the thermosphere, the composition of the gas not only changes due to molecular diffusion, but also due to photodissociation and photoionization, driven by the absorption of





solar radiation at wavelengths below approximately 200 nm. For the modern Earth, the dissociation of O_2 and N_2 molecules, combined with molecular diffusion, mean that the gas in the exobase is primarily composed of atomic O and N.

The absorption of this radiation has several additional effects on the gas, with the most important being heating. When the atmosphere is in a steady state, this heating of the upper atmosphere is balanced by cooling from infrared emission to space, particularly by CO_2 . Locally within the gas however, heating and cooling do not balance, and downward thermal conduction is in fact the main process responsible for cooling the upper thermosphere.

For Venus and Mars, the same basic set of physical processes are important in the upper atmospheres, with the only additional important process being heating by the absorption of solar infrared radiation by CO_2 (Fox and Bougher 1991). The high CO_2 concentration in these atmospheres also means that infrared cooling is much stronger and the thermospheres are therefore much cooler (e.g., de Pater and Lissauer 2015; Bauer and Lammer 2004). The higher concentrations of CO_2 expected in the Earth's atmosphere at earlier times (Walker 1985; Tian et al. 2005) would also likely have this effect, cooling the atmosphere and reducing the effects of atmospheric loss (Kulikov et al. 2007; Johnstone et al. 2018).

As illustrated in Fig. 3, it is common to break down the vertical structure of an atmosphere into the homosphere and heterosphere. In the homosphere, which for the Earth extends from the surface to the homopause at approximately 120 km altitude, turbulent mixing of the gas causes the mixing ratios of all long lived species (those that are not rapidly created and destroyed by chemical processes) to be uniform with altitude regardless of the particle masses. This process is often known as eddy diffusion. In the homosphere, diffusion is slower than mixing, allowing heavier particles such as CO_2 molecules and noble gases such as Ne, Ar, and even Xe atoms to reach the lower thermosphere. Without eddy diffusion, the abundances of heavier particles would rapidly drop with increasing altitude as molecular diffusion separates the particles by mass. This is indeed what happens above the homopause, with the abundances of heavier particles becoming extremely low in the mid and upper thermospheres. Here, diffusion is faster than mixing, and heavy species (e.g., N₂, Xe, etc.) decrease in density with altitude more rapidly than lighter species. This process contributes to causing lighter particles such as H atoms to be lost from the atmosphere more efficiently than heavier particles.

An important concept related to the escape of H atoms is diffusion-limited escape. This is itself not a loss process but an upper limit on how rapidly H atoms can be lost from the upper atmosphere of a planet that is dominated by an other species (N, O, etc.). This limit

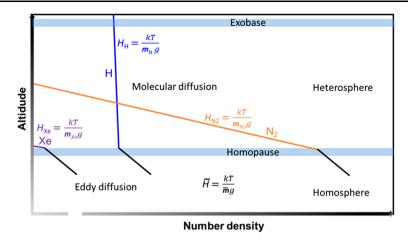


Fig. 3 Illustration of density-altitude structures of H atoms, N_2 molecules and Xe isotopes in the homosphere-heterosphere transition region. In the homosphere the atmosphere is mixed and all species are distributed by the same average scale height (homosphere), while above the homopause all species are distributed by their own scale heights as a function of mass. The density of heavy elements decreases very fast compared to light elements

exists because a minor species (i.e. H atoms) cannot be lost from the upper atmosphere more rapidly than they can be transported upwards by molecular diffusion at the homopause.

For instance, H atom escape on present Earth, Venus and Mars is limited by upward diffusion through the upper atmosphere (e.g. Hunten 1973; Kasting and Pollack 1983). On Earth, H₂O, H₂, and CH₄ move upward in the atmosphere. At \approx 17 km, the cool temperatures at the tropopause freeze out most of the water vapour, preventing the upward transport of some hydrogen. Near the homopause hydrogen bearing molecules (i.e. CH₄) are dissociated by UV photons so that atomic and molecular hydrogen particles are produced. These H and H₂ particles diffuse upward through the upper atmosphere to the exobase level from where they can escape by thermal Jeans escape. On present Earth, the "bottleneck" for hydrogen-mixing ratio at the homopause, which is \approx 5.0 × 10⁻⁷ at present Venus (Kasting and Pollack 1983) and \approx 1.2 × 10⁻⁵ and \approx 3.0 × 10⁻⁵ at present Earth and Mars, respectively (Krasnopolsky and Feldman 2001). For such low water vapour mixing ratios at the homopause the exobase level at Mars and Venus, but also at Earth are thermal Jeans escape.

Higher homopause water mixing ratios lead to higher escape rates that are not controlled by diffusion through the upper atmosphere anymore. Kasting and Pollack (1983) found that for H₂O mixing ratios at the homopause that are $\geq 5.0 \times 10^{-4}$ the stellar/solar EUV flux that is absorbed above the homopause level in a hydrostatic atmosphere becomes larger than the energy that is carried away by the thermally escaping particles. At such high H₂O mixing ratios at the homopause level the hydrogen-dominated upper atmosphere becomes unstable against expansion and Jeans escape transitions to hydrodynamic escape. According to Kasting and Pollack (1983), for H₂O mixing ratios at the homopause that are $\geq 10^{-3}$ the hydrogen escape flux reaches values that are comparable to energy limited escape. High H₂O-mixing ratios at the homopause level in the order of 0.1 and higher can originate in steam atmospheres (e.g., Lebrun et al. 2013; Massol et al. 2016; Salvador et al. 2017; Benedikt et al. 2020). Under such conditions thermal escape of H atoms from the upper atmosphere and the corresponding fractionation of isotopes and other heavier minor and/or trace species is then controlled by thermal hydrodynamic escape (e.g., Hunten 1973; Zahnle and Kasting 1986; Hunten et al. 1987; Zahnle et al. 1990; Odert et al. 2018; Lammer et al. 2020a, 2020b).

4.2 Atmospheric Loss Processes and Their Influence on Isotope Fractionation

Loss processes are often broken down into thermal and non-thermal processes. Thermal processes are related to the thermal energy of the gas and can be primarily broken down into Jeans escape and hydrodynamic escape. Non-thermal processes are essentially all others. Most escape processes preferentially remove lighter isotopes, though they differ in how efficiently they fractionate by mass. The ability of an escape process to fractionate isotopes is often described by a fractionation factor f, with f > 1 meaning that the heavier isotope is preferentially removed, f < 1 meaning that the escape process preferentially removes the lighter isotope, and f = 1 describing the case of no fractionation. The Rayleigh distillation relationship is commonly used to describe fractionation of an isotope ratio due to a separate process. This relationship connects the initial, n_0 , and current, n, inventory of the constituent and a fractionation factor f of a particular escape process, "x" (e.g., Lunine et al. 1999; Donahue 1999; Mandt et al. 2009, 2014, 2015):

$$\frac{n_0}{n} = \left(\frac{R}{R_0}\right)^{\frac{1}{1-f_X}},\tag{1}$$

where *R* is the current ratio of the heavy to light isotope abundance, and R_0 is the initial ratio. The ratio R/R_0 defines the degree of enrichment in the heavier isotope in comparison to the primordial value. The upper limit of the enrichment of the heavier isotope can be found as

$$\frac{R}{R_0} \le \left(\frac{\varphi t}{n+1}\right)^{1-f_{\rm X}},\tag{2}$$

where φ is the maximum escape flux and *t* is time. Different escape processes have various possible maximum escape fluxes and different fractionation coefficients that can depend on the structure and the composition of the upper atmosphere. Below, we briefly summarize the main atmospheric escape processes and their influence on the atmospheric fractionation.

4.2.1 Fractionation by Thermal Atmospheric Escape

Jeans escape: when an upper atmosphere is hydrostatic and is not able to flow away from the planet hydrodynamically, it is still possible for particles in the high-speed tail of the thermal distribution of particles at the exobase to have speeds exceeding the escape velocity. The loss of such particles is known as Jeans escape. Jeans escape is highly mass fractionating only if the escape is not diffusion-limited. Only if Jeans escape is the bottleneck, this process leads to highly mass-dependent fractionation. For the present solar system planets, Jeans escape rates are quite low, so that it takes a long time to fractionate a large atmospheric reservoir assuming present escape rates (Mandt et al. 2014) though the process would have been much more rapid at earlier times when the Sun was much more active. The fractionation factor f_{Jeans} for Jeans escape can be written as (e.g., Mandt et al. 2009)

$$f_{\text{Jeans}} = \sqrt{\frac{m_1}{m_2}} \left[e^{(\lambda_1 - \lambda_2)} \frac{(1 + \lambda_2)}{(1 + \lambda_1)} \right],$$
(3)

where m_1 and m_2 are the masses of the respective isotopes and $\lambda_{1,2} = (GM_{\rm pl}m_{1,2})/(k_{\rm B}T_{\rm ex}r_{\rm ex})$ are the Jeans escape parameters of the lighter and heavier isotope at the exobase level, respectively. Here G is the Newton gravitational constant, $M_{\rm pl}$ is the planetary mass, $k_{\rm B}$ is the Bolzmann constant, $T_{\rm ex}$ the exobase temperature and $r_{\rm ex}$ the exobase distance.

Hydrodynamic escape: if the upper atmosphere is heated by the EUV flux of the planet's host star to very high temperatures, it is possible that the gas in the upper atmosphere flows away from the planet with a speed at the exobase exceeding the escape speed. Such a flow has the form of a transonic Parker wind (e.g. Parker 1958; Tian et al. 2005; Johnstone et al. 2019). Whether an atmosphere reaches this state depends on the gas temperature in the upper atmosphere and the mass of the planet, with hydrodynamic escape taking place significantly more readily for lower mass planets. The source of the energy that heats the gas is typically the X-ray and ultraviolet field of the star. Very early in a planet's history, immediately after the dissipation of the primordial nebula, it is also possible for hydrodynamic escape to be powered by the initial accretion energy of the envelope gas (Stökl et al. 2015) and the bolometric radiative emission of the star (Owen and Wu 2016), though the latter is unlikely to have been significant for the young solar system planets due to their large distances from the Sun. Previous works have found that hydrodynamic escape removes all isotopes very efficiently, which leads to small fractionations (e.g., Sekiya et al. 1980a; Mandt et al. 2009). For hydrodynamic escape, the Jeans escape parameter λ decreases to very small values (Chamberlain 1963; Volkov et al. 2011; Erkaev et al. 2015). In such a case, f_{Jeans} converges towards f_{hy} , i.e.

$$f_{\rm hy} = \sqrt{\frac{m_1}{m_2}},\tag{4}$$

due to the second part in Eq. (3) converging towards one for very small lambdas. When hydrogen atoms escape hydrodynamically, however, it will drag heavier elements with it. This effect can subsequently lead to a fractionation of lighter isotopes from heavier isotopes (e.g., Zahnle and Kasting 1986; Hunten et al. 1987; Zahnle et al. 1990; Odert et al. 2018). If, besides hydrogen, one assumes a second major species, e.g. oxygen, in the hydrodynamically escaping atmosphere, which is also dragged by hydrogen, then the fractionation factors f_{drag}^i for the second major species and f_{drag}^j for an arbitrary amount of additional minor species can be written as (e.g., Odert et al. 2018)

$$f_{\rm drag}^{\rm i} = 1 - \frac{g \left(m_{\rm i} - m_{\rm H}\right) b_{\rm H,i}}{F_{\rm H} k_{\rm B} T \left(1 + \xi_{\rm j}\right)},\tag{5}$$

and

$$f_{\rm drag}^{j} = \frac{1 - \frac{g(m_{\rm j} - m_{\rm H})b_{\rm H,j}}{F_{\rm H}k_{\rm B}T} + \frac{b_{\rm H,j}}{b_{\rm H,i}}\xi_{\rm i}\left(1 + \xi_{\rm i}\right) + \frac{b_{\rm H,j}}{b_{\rm i,j}}\xi_{\rm i}f_{\rm drag}^{\rm i}}{1 + \frac{b_{\rm H,j}}{b_{\rm i,i}}\xi_{\rm i}},\tag{6}$$

where g is the gravitational acceleration of the planet, T the upper atmosphere temperature, m_i , m_H , and m_j the masses of the masses of the respective atoms, $b_{H,i}$, $b_{O,j}$, $b_{H,j}$ the binary diffusion parameters, F_H the hydrodynamic escape flux of hydrogen atoms, and ξ_i the mixing ratio of the second major species to hydrogen.

To summarize thermal escape and fractionation processes, one finds that hydrodynamic outflow is relatively inefficient at fractionating elements and isotopes with a small mass difference (i.e., D/H...), while it can be efficient if a light species such as atomic hydrogen drags heavier ones, whereas Jeans escape fractionates efficiently but total atmospheric

escape is much less efficient. As a rule of thumb, high-rate loss processes, such as hydrodynamic escape, are very efficient in removing the atmospheric gas, but are inefficient fractionators in case of light elements where it takes a long time to change the isotopic ratio. On the other hand, a low-rate loss process such as Jeans escape can be an efficient fractionator, but take a very long time to influence a large inventory.

4.2.2 Fractionation by Non-thermal Atmospheric Escape

Ion pick-up: above the exobase, neutral atmospheric particles follow ballistic trajectories. If these particles reach the region above the magnetopause or ionopause, they can directly interact with the solar wind, where they can be ionized by solar EUV radiation and interactions with the solar wind. These ions are picked up by the solar wind and are lost from the planet. If a planet does not have an intrinsic magnetic field, as is the case for modern Mars and Venus, then the solar wind can directly interact with the ionosphere of the planet, thus making the ion pick-up escape more efficient. For a magnetized planet, atmospheric particles have first to reach the location of the magnetopause to be picked up by the solar wind, which can reduce the ion pick-up escape rates in comparison to the non-magnetized case. Due to diffusive separation of atmospheric particles by mass above the homopause (see also Fig. 2), isotopes with lighter masses reach higher altitudes and will therefore be preferentially removed from the top of the atmosphere. For non-thermal escape such as ion pick-up, the diffusive separation fractionation factor f_{ds} can be defined as follows (Lunine et al. 1999; Mandt et al. 2014), i.e.

$$f_{\rm ds} = \exp\left[-\frac{g(\Delta z)(m_2 - m_1)}{k_{\rm B}T}\right],\tag{7}$$

where $\Delta z = r_{ex} - r_{h}$ is the distance between the homopause r_{h} and the exobase r_{ex} levels (see also Fig. 2) if the planet has no intrinsic magnetic field, otherwise the exobase can be substituted by the magnetopause.

Sputtering: in this process, atmospheric atoms or molecules are ejected due to collisions between solar wind protons and backscattered pick-up ions with the planetary atmosphere (e.g., Johnson 1990). This process efficiently fractionates the isotopes through diffusion as a function of the temperature in the upper atmosphere and the distance between the exobase and the homopause. Like with other loss processes, the efficiency of the fractionation by sputtering depends on the exact structure of the upper atmosphere, with expanded atmospheres fractionating more easily. Sputtering preferentially removes the lighter atom or particle. On low mass planetary bodies such as Mars, ion pick-up and sputtering are likely the two fractionation processes that led to the observed ³⁸Ar/³⁶Ar fractionation (Jakosky et al. 2017). For atmospheric sputtering the diffusive separation, fractionation is relevant and can be calculated in the same way as for ion pick-up escape in Eq. (7).

Plasma instabilities: plasma instabilities, such as Kelvin-Helmholz instability, can drive an escape of plasma vortexes from the ionospheric boundary from planets which have no intrinsic magnetic dynamo (i.e., Venus, Mars, ...). A wave-like structure of initially small amplitudes can grow and form vortexes. These vortexes detach from the ionized upper atmosphere and escape from the planet in a plasma cloud (Terada et al. 2002; Amerstorfer et al. 2007). Since plasma instabilities remove a large number of ions at once, they are not very efficient isotope fractionators.

Cold ion-outflows: ions and free electrons produced by the absorption of solar extreme ultraviolet and X-ray photons are able to escape by a separate process. The expansion of the lighter electrons produces an ambipolar electric field that accelerates the ions away from the

planet causing an outflow of ions and electrons (Ganguli 1996). For magnetised planets like the Earth, this flow takes place at the magnetic poles. This process can accelerate the planetary ions and, if they reach a sufficiently high velocity, lead to them escaping the planet's gravity (e.g., Hartle and Grebowsky 1993; Lundin 2011). Heavier ions can also be dragged by other atmospheric constituents through this process, as was probably the case for the easily ionized Xe⁺ isotopes in the Earth's and Mars' atmosphere (e.g., Zahnle et al. 2019; Avice and Marty 2020; this issue).

Photochemical escape: this type of escape is related to chemical reactions that produce hot atoms with high energies, so called suprathermal or "hot" atoms. If the energy of an atom is high enough, it can exceed the energy of the gravitational binding to the planet, and then the atom is lost (e.g., Amerstorfer et al. 2017). The reactions can include various chemical reactions, including dissociative recombination or charge exchange. Photochemically produced suprathermal atom escape can preferentially remove lighter isotopes. On Mars, photochemical processes lead to removal of nitrogen and oxygen atoms, and similar to the fractionation by ion pick-up, the fractionation is driven by the diffusive separation of different isotopes by mass above the homopause (e.g., Jakosky et al. 1994).

On Earth, He escapes due to exothermic charge exchange reactions between He⁺ ions and the major atmospheric constituents N₂, O₂, and O (e.g., Lie-Svendsen and Rees 1996) and on Venus photochemical processes are relevant for the D/H fractionation (Lammer et al. 2006).

Impact erosion: early in the planet's history, atmospheric gas can be removed by impacts due to a large mass of gas removed simultaneously (e.g., Schlichting et al. 2015; Schlichting and Mukhopadhyay 2018); this process is not an efficient fractionator of isotopes related elements in atmospheres and is discussed in more detail in Sect. 5.3.

5 Impact Erosion of Primordial Atmospheres, Early Crust, and Mantle

Understanding impact and collision processes and their consequences during accreting planetary embryos and protoplanets is a vital component in the overall picture, since those events happened during all stages of terrestrial planet formation. The currently established picture distinguishes several main stages, where theoretical knowledge about the early phases from dust to planetary embryos is still limited and various competing theories are discussed (e.g. Raymond et al. 2014; Johansen and Lambrechts 2017; Izidoro and Raymond 2018; Lammer et al. 2020a; this issue). However, it is relatively well established that the final phase of accretion, lasting for millions to several 10s of Myr is marked by giant collisions among planetary embryos and remaining planetesimals, which shape the bulk properties of eventually formed planets.

5.1 Impact Related Erosion of the Primordial Atmosphere

As discussed above, growing bodies accrete nebular gas once they become sufficiently massive, where cores larger than one Earth-mass accrete an H₂-dominated envelope that most likely remain during their whole lifetime (Johnstone et al. 2015; Owen and Wu 2016), while protoplanets below 0.5 M_{Earth} probably loose most of their primordial atmosphere, which is largely boiled off after the surrounding gas dissipates and potentially replaced by an outgassed one later (Stökl et al. 2016). Protoplanets with higher masses in orbits inside the habitable zone keep primordial atmospheres that can be lost afterwards by EUVdriven hydrodynamic escape and erosion by large impactors (e.g., Schlichting et al. 2015; Schlichting and Mukhopadhyay 2018; Lammer et al. 2020b). It is important to mention that existing work on the interplay of internal and external (impact) heating, magma ocean and (proto-)atmosphere evolution of growing terrestrial planets has focused on roughly Earthmass objects, even though smaller bodies – the size of planetary embryos and below – are the stepping stones on the way to fully grown planets.

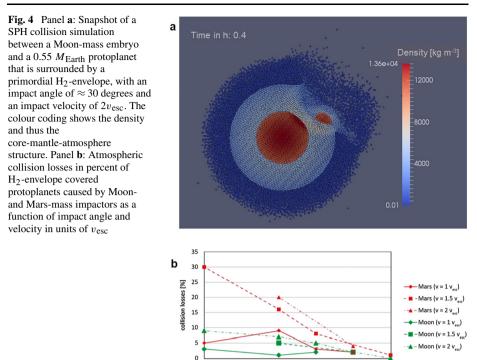
Studies on impact efficiencies are controversial since all the work on that subject is theoretical. The first impact studies onto atmosphere and ocean-covered protoplanets were often based on 1-D hydrodynamic codes (Genda and Abe 2003, 2005; Inamdar and Schlichting 2016), analytical models in combination with such codes (Schlichting et al. 2015), and also 3-D hydrocode simulations of impacts of small bodies (Shuvalov 2009) up to head-on collisions in the Earth- to super-Earth mass range (Liu et al. 2015). An important result is that smaller impactors are - per unit mass - much more efficient in eroding atmosphere than larger ones (e.g. Schlichting et al. 2015; Schlichting and Mukhopadhyay 2018, for an overview). The mass loss efficiency (the ratio of lost atmospheric to impactor mass) is largest for small bodies, e.g. in the kilometer range for current Earth, and decreases with mass by orders of magnitude up to the giant impact regime (\sim 1000 km bodies) This is mainly because for small impactors the majority of the energy is available for disturbing the atmosphere locally; above the tangent plane of the impact location, the gas can be entirely lost. It should also be noted that smaller impactors can also deliver material to the planet.

For impactors with energy deposition larger than what is required for ejecting all material above the tangent plane, considerable parts of their kinetic energy have to be used to drive a strong shock through the target body to eject atmosphere globally, once the shock-induced ground velocity is sufficient to accelerate atmosphere components to escape velocity. The presence of oceans on the surface has been found to aid the escape of overlying atmosphere, due to their lower shock impedance compared to a solid surface (Genda and Abe 2005). In recent 3-D smoothed-particle hydrodynamics (SPH) giant impact simulations, atmosphere losses of $\approx 5\%$ per collision for Moon-mass/size, and $\approx 10-20\%$ for Mars-mass/size impactors in collisions with a $\approx 0.5-0.75 M_{Earth}$ protoplanet with primordial atmospheric mass fractions of ≥ 0.001 have been obtained (Lammer et al. 2020b).

The impact scenarios in this study cover the full range of impact angles and common impact velocities from 1 to $2 \times v_{esc}$. An SPH collision snapshot of a Moon-mass impactor that crashes into an H₂-covered protoplanet with 0.5 M_{Earth} is shown in panel (a) in Fig. 4 (Burger et al. 2018; Lammer et al. 2020b). Panel (b) in Fig. 4 shows eroded primordial H₂-dominated atmospheres in percent from protoplanets with masses $\approx 0.5 - 0.75 M_{Earth}$ caused by Moon-mars and Mars-mass impactors as a function of impact angle and escape velocity (Lammer et al. 2020b). The different results shown in Lammer et al. (2020b) indicate that studying atmospheric losses caused by giant impacts only by assuming head-on scenarios (e.g. Liu et al. 2015) may be insufficient to understand the full diversity and efficiency of atmospheric losses by such events. These results confirm that - like for small bodies - even in the giant impact regime lower-mass bodies are more efficient (normalized to impacting mass) than larger ones in eroding atmosphere (Schlichting et al. 2015; Schlichting and Mukhopadhyay 2018).

5.2 Collisional Erosion of Crust/Mantle Material by Large Impactors

During its accretion phase, in particular during the so-called giant impact phase, a growing protoplanet is expected to experience energetic collisions with other planetary embryos (e.g., Quintana et al. 2016). Such a collision can then remove part of the newly formed crust



from the planet which will either re-accrete or be lost to space (Carter et al. 2018). This process will therefore preferentially remove incompatible lithophile elements from the growing planet leading to a fractionation between siderophile and lithophile elements (O'Neill and Palme 2008; O'Neill et al. 2020; this issue). While Fe/Mg or Fe/Si, of which Fe is siderophile and Mg and Si are lithophile, will be fractionated by collisional erosion, other ratios such as K/U, which are both lithophile elements, will not. This process could therefore account for the high terrestrial Fe/Mg ratio of 2.1 ± 0.1 compared to the solar ratio of 1.9 ± 0.1 (O'Neill and Palme 2008). Another ratio that was suggested to be fractionated by collisional erosion is Earth's ¹⁴²Nd/¹⁴⁴Nd ratio (Jellinek and Jackson 2015). Here, however, a more recent study found that ¹⁴²Nd/¹⁴⁴Nd might not have been fractionated at all (Bouvier and Boyet 2016).

15

30

45

impact angle α [°]

60

75

n

Whether collisional erosion can account for the high terrestrial Fe/Mg ratio was numerically studied by Carter et al. (2015) and Bonsor et al. (2015). Even though they both conclude that collisional erosion can fractionate Earth's Fe/Mg from a chondritic to the present-day terrestrial value, they also found that this fractionation can only be reached if Earth was an extreme case. The growing planetary embryos that exhibit a different range of core mass fractions through giant impacts must have had a significant dynamical excitation to finally create the present-day Fe/Mg ratio of Earth (Bonsor et al. 2015; Carter et al. 2015).

However, the escape of Mg from a planetary embryo through hydrodynamic drag of hydrogen could additionally alter the Fe/Mg ratio, which might make the Earth a less extreme case than expected in these studies (Benedikt et al. 2020). A more recent study by Carter et al. (2018) found that a growing protoplanet on average can lose about one third of its crust through collisional erosion which in turn would lead to the loss of up to 20% of its heat producing elements if the crust does not re-accrete onto the planet. Such an alteration of the planetary heat budget might further also have implications on the habitability of a planet since this could strongly affect the evolution of tectonic activity (Jellinek and Jackson 2015; O'Neill et al. 2020; this issue).

Collisional erosion could further also have been important in view of Mercury's high iron abundance, as was recently simulated by Chau et al. (2018). These authors found that an energetic impactor could alter the iron distribution within a planet by up to 25%. According to their study, Mercury's abnormally high abundance of iron could therefore be the result of multiple collisions. Another recent study by Clement et al. (2019), however, concludes that, even though planets with elevated core fractions were produced in 90% of their simulations through the fragmentation of collided bodies, there is only a probability of $\leq 1\%$ to simultaneously reproduce Mercury and Venus after such an event.

5.3 Elemental Fractionation Caused by Large Impact Events

To summarize the impact erosion processes, one finds two scenarios where large impactors will change the composition of a growing proto-planet. Both cases are illustrated in Fig. 4. If a proto-planet accreted $> 0.5 M_{Earth}$ inside the disk, elements in solar abundance and magma ocean-related outgassed moderately volatile rock-forming elements and noble gases will be embedded as trace elements in a captured H₂-dominated primordial atmosphere.

As illustrated in Fig. 5a, large impactors that mainly affect the atmosphere can then contribute to the fractionation of moderately volatile elements such as K and/or Mg compared to U or Fe, which remain in the crust/mantle. If a large impactor crashes into the body, crustal /mantle material including refractory elements (e.g., U, Th) might also be lost to space if not re-accreted by the protoplanet (see Fig. 5b), (e.g. Palme and O'Neill 2003; see also O'Neill et al. 2020; this issue). Depending on the impactor and its parameters, impact erosion can be a strong fractionating process.

Impacts, however, will not fractionate isotopes of atmospheric noble gases (e.g. ³⁶Ar, ³⁸Ar, ²⁰Ne, ²²Ne, etc.) if they only affect the atmosphere (Fig. 5a) but contribute to a depletion of volatile elements from a planetary body. If impactors with different isotopic compositions and ratios are involved in the accretion of a protoplanet (Fig. 5b), noble gas isotope ratios, however, can be modified by their mixing with compositions, abundances and ratios of the impacting bodies as discussed in Sect. 6. Below we discuss atmospheric escape and fractionation of such elements from planetary embryos, which experience magma oceans and the therewith related outgassing of volatile elements.

6 Escape of Rock-Forming Elements and Volatiles from Planetary Embryos

Terrestrial planets accrete mass by numerous collisions between small objects that accumulate to larger ones, including Moon and Mars-mass planetary embryos that accrete onto full planets (e.g., Johansen and Lambrechts 2017; Izidoro and Raymond 2018; Lammer et al. 2020a, this issue). Growing planetesimals, which are embedded in the circumstellar disk, with a size between hundred to thousand kilometers develop global magma oceans, caused by the heating of short-lived radioactive elements, like ²⁶Al (e.g., Lichtenberg et al. 2016; 2017; Young et al. 2019), and mutual collisions (e.g., Safronov and Zvjagina 1969; Wetherill 1980; Tonks and Melosh 1993). Schaefer and Fegley (2007) and Fegley et al. (2016) performed chemical kinetic and equilibrium calculations for modeling the chemistry of elements (e.g., K, Si, Mg, Fe, Ca, Al, Na, S, P, Cl, F) released by heating of different types of carbonaceous, ordinary and enstatite chondrites as a function of temperature and pressure

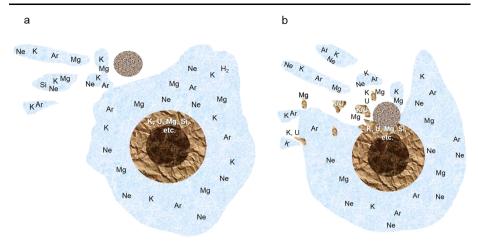


Fig. 5 Illustration of the two impact related erosion scenarios. (**a**) Impact erosion of atmospheres, and (**b**) direct impacts and related erosion of crustal and mantle material. Impacts that mainly affect a primordial atmosphere can contribute to the modification of ratios between moderately volatile rock-forming elements that can also be in a primordial atmosphere and refractory elements that remain in the surface, such as U or Th (panel **a**), while direct impacts can also erode crustal and mantle material (panel **b**)

of early atmospheres. These studies yield significant amounts of major rock-forming and other volatile elements, such as Rb and Zn, and noble gases (Ar, Ne, etc.) in primordial and/or catastrophically outgassed steam atmospheres at high temperatures (1000-2500 K).

When the disk dissipates, after a short but efficient boil-off phase (Owen and Wu 2016; Lammer et al. 2018) the magma oceans of the former disk-embedded bodies solidify and steam atmospheres are catastrophically outgassed (see Fig. 6). Depending on the redox state these catastrophically outgassed atmospheres contain H_2 , CO (reduced interior) and/or H_2O and CO₂ (oxidized interior; Nikolaou et al. 2019). The outgassed amount per element depends on the bulk composition of a particular planetary embryo and the solidification process of the mentioned magma ocean (e.g., Elkins-Tanton 2008, 2012; Salvador et al. 2017; Nikolaou et al. 2019).

Recently, Benedikt et al. (2020) studied the losses of magma ocean-related catastrophically outgassed steam atmospheres and outgassed rock-forming elements as well as Ar and Ne noble gases from planetary embryos with masses between 1.0 M_{Moon} and 1.5 M_{Mars} along the EUV-flux evolution tracks of slow and moderate/fast rotating young G-type stars (Tu et al. 2015; Güdel 2020; this issue) at orbital distances between 0.5-1.5 AU.

Based on studies from Wakita and Schmitt (1970), Righter and Chabot (2011), and Elkins-Tanton (2012), one can assume magma ocean depths for planetary embryos within the mass ranges mentioned above between 500 - 1000 km for surface temperatures of 1500 - 3000 K. Because of the low mass of planetary embryos and these high surface temperatures, according to 1-D radiative-convective atmosphere models (Marcq 2012; Marcq et al. 2017; Pluriel et al. 2019), it was found that the steam atmosphere heights are increasing up to a few thousands of kilometers (Benedikt et al. 2020). The main process, which is dragging heavier atoms into space, is related to the dissociation of H₂O molecules around the 1 µbar (Guo 2019) level, so that atomic hydrogen dominates the upper atmosphere and is lost hydrodynamically, thereby potentially dragging heavier elements.

For planetary embryos with masses $\leq 1 M_{Moon}$, no dense silicate atmosphere or steam atmosphere can build up (Hin et al. 2017; Benedikt et al. 2020; Young et al. 2019), the embryo will continuously lose outgassed elements very fast to space. Therefore, one can

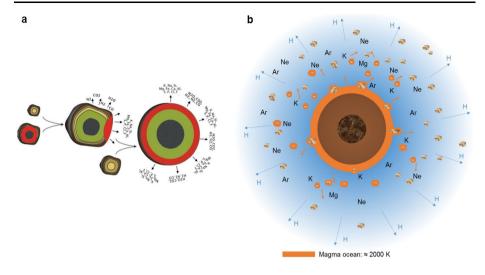


Fig. 6 (a) Illustration of planetary embryo growth after Elkins-Tanton (2012) due to collisions between differentiated planetesimals/planetary embryos. The red areas correspond to magma from which various elements evaporate and can be lost directly at bodies with masses $\leq M_{\text{Moon}}$ (Benedikt et al. 2020; Young et al. 2019) and by dragging of hydrodynamically escaping H atoms dissociated either from H₂O molecules originated from catastrophically outgassed steam atmospheres during magma ocean solidification (Benedikt et al. 2020) or (**b**) from remnants of captured H₂-dominated nebula-based envelopes (Lammer et al. 2020b)

expect that these bodies will be strongly depleted in volatiles, noble gases and outgassed rock-forming elements.

One should note that in the recent study by Young et al. (2019) it is assumed that hydrostatic conditions for an outgassed rock vapour atmosphere above a magma ocean with a surface temperature of 1830 K on a planetary embryo with half of the mass of Pluto build up. As discussed in Benedikt et al. (2020), and shown in Fig. 7, under such extreme conditions where the ratio between the gravitational energy and the thermal energy, i.e. the Jeans parameter λ is < 6 (in the case of Young et al. 2019 the λ at the surface is 1.38) it is impossible that a stable atmosphere can build up at such low mass and hot bodies. In such cases, direct escape of outgassed rock-forming elements occurs unless there is a steady supply from the interior to compensate for these fast losses. For silicate atmospheres, this means that the entire mantle will eventually depleted in the outgassed elements. The assumption of a hydrostatic equilibrium as assumed in the study of Young et al. (2019) is thus only justified for cases in which the Jeans parameter of the outgassed elements at the surface of an embryo exceed \approx 6 (Erkaev et al. 2015; Volkov et al. 2011).

Figure 7 shows the evolution of the atmospheric temperature (red curves) and Jeans parameter (green curves) by using the same embryo and silicate atmosphere parameters with 34 amu as given in Young et al. (2019) for an assumed hydrostatic (dotted lines) and realistic non-hydrostatic (solid lines) atmosphere.

In the realistic non-hydrostatic case, the atmospheric temperature is decreasing much faster due to adiabatic cooling of the outward flowing bulk gas. At about 9 embryo radii the atmosphere is heated again because of absorption of the EUV radiation that is scaled in Benedikt et al. (2020) to the value expected for the early Solar System at the time when the disk dissipated (Tu et al. 2015; Güdel 2020; this issue). Because of the low λ value of ≈ 1.4 , the hydrodynamic escaping bulk atmosphere is supersonic just form the surface.

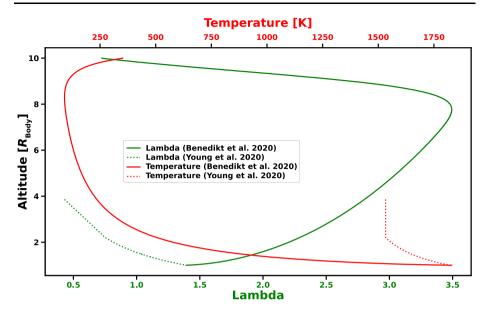


Fig. 7 Change of temperature (red lines) and Jeans parameter lambda (green lines) for the assumed but not realistic hydrostatic (dotted lines) outgassed silicate atmosphere (Young et al. 2019) and a non-hydrostatic (solid lines) silicate atmosphere (Benedikt et al. 2020) around a planetary embryo with 0.5 M_{Pluto} (after Benedikt et al. 2020)

According to Benedikt et al. (2020), more massive embryos with masses > 0.5 M_{Mars} also lose outgassed heavy trace elements when they are incorporated into solidified magma ocean-related steam atmospheres. As long as steam atmospheres do not condense and outgassed rock-forming elements reach the upper atmosphere, they can also be lost by the dragging of escaping H atoms.

Figure 8 shows an example of escape times of magma ocean-related catastrophically outgassed steam atmospheres and embedded trace elements for a planetary embryo with a mass of 0.5 M_{Mars} at various orbital distances, partial surface pressures of 50/5 (green), 100/10 (red), 150/30 (blue) bar H₂O/CO₂ and surface temperatures of 1500 K. The bold blue line corresponds to the condensation temperature of a steam atmosphere based on the results of Lebrun et al. (2013). As one can see in Fig. 7, depending on the partial pressure and surface temperature, embryos with a mass of 0.5 M_{Mars} can lose steam atmospheres at an orbital distance of 1 AU within the condensation time. For heavier planetary embryos of up to $\approx 1.5 M_{\text{Mars}}$ almost all outgassed trace elements can be lost together with the escaping steam atmosphere within ≈ 10 Myr (Benedikt et al. 2020), which lies within the time frame of the formation of the first protocrust on Mars (Bouvier et al. 2018; Kruijer et al. 2020).

This indicates that frequent smaller impactors (Maindl et al. 2015) keep shallow magma oceans molten most likely over longer periods so that the steam atmosphere can longer remain in vapour form compared to the modelled time scales given in Lebrun et al. (2013). A more active young Sun-like star leads to a shorter escape time for the whole steam atmosphere as for the slow rotator in Fig. 8.

Figure 9 shows an example of the remaining amount of various elements on a planetary embryo with 0.5 M_{Mars} at 1 AU, with a magma ocean-related steam atmosphere partial surface pressure of 50 bar H₂O and 5 bar CO₂. If all outgassed elements escape, the particular volatiles are depleted by $\approx 75\%$ within the planetary embryo. The remaining amount of Ne

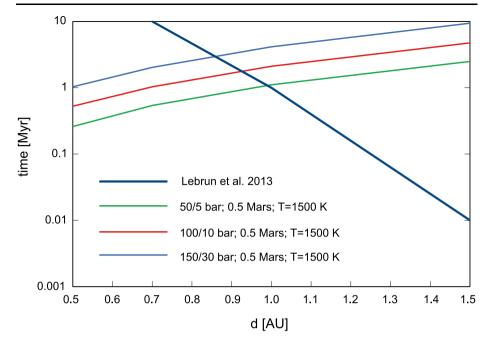


Fig. 8 Escape time *t* of steam atmospheres for a planetary embryo with a mass of 0.5 M_{Mars} , various partial surface pressures of H₂O/CO₂, for a surface temperature of 1500 K and for a weakly active slowly rotating young G-type star in dependence of the orbital distance *d*. The solid dark blue line in each panel shows the steam atmosphere condensation time based on the study of Lebrun et al. (2013) (after Benedikt et al. 2020)

and Ar isotopes in mol/g and ²⁰Ne/²²Ne and ³⁶Ar/³⁸Ar ratios are also shown. One can see that on such low mass bodies the Ne and Ar isotopes only escape to space but do not fractionate due to the high loss rates. Isotope fractionation, however, between the vapour and melt of small planetary embryos could at least fractionate the isotopic ratios of Mg and Si at near-equilibrium (Young et al. 2019).

It should also be noted, that the depletions and ratios of the volatile rock-forming elements of the larger planetary embryos with masses $\geq 0.5 M_{\text{Mars}}$ most likely represent upper values. First, it is unclear whether 100% of the outgassed elements will indeed be incorporated into the steam atmosphere or if a fraction remains inside the magma ocean. Second, the outgassed rock-forming elements have to overcome a bottleneck from the lower atmosphere to the upper atmosphere, the so-called cold-trap in the mesosphere; otherwise, they will reach the condensation temperature. For noble gases such as Ne or Ar, however, the atmospheric cold-trap of larger bodies does not exist due to their low condensation temperatures and rain back to the surface.

One should also note that for reduced atmospheres with high contents of H_2 and CO molecules, the atmosphere would only marginally bound on such a low mass body. For surface temperatures of ≥ 2000 K, reduced H_2 and CO dominated atmospheres would be unbound. Therefore, outgassed atmospheres that are highly reduced will escape from such low mass bodies even faster than the oxidized steam atmospheres shown in Fig. 9.

Because of this, it seems reasonable that large planetary embryos should be more depleted in noble gases (i.e., Ne and Ar) than in rock-forming elements such as K, Si, Mg, Fe, Ca, Al, Na, S, P, Cl, F due to their significantly higher condensation temperatures.

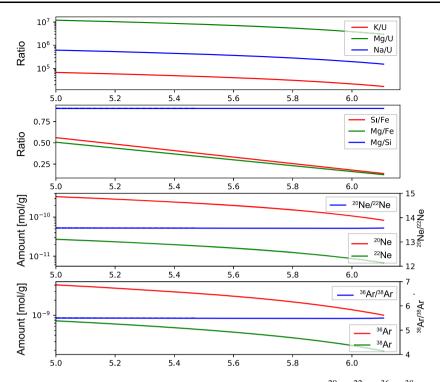


Fig. 9 Escape related evolution of the K/U, Mg/U, Na/U, Si/Fe, Mg/Fe, Mg/Si, 20 Ne/ 22 Ne, 36 Ar/ 38 Ar ratios and losses of Ar and Ne as a function of time for a 50 bar H₂O and 5 bar CO₂ outgassed steam atmosphere from a planetary embryo with 0.5 M_{Mars} at 1 AU. The embryo's steam atmosphere is exposed to the EUV flux of a slowly rotating young G-type star after the magma ocean solidification (assumed to be at ≈ 5 Myr). About 75% of the total initial amounts of the elements are lost more or less within 1 Myr with surface temperatures corresponding to 1500 K. The lines end when the steam atmospheres are lost from the embryo. The ratios were calculated for mol/g and U and Fe assumed not to escape but remain at the body with their initial amounts (after Benedikt et al. 2020)

In general, these studies also agree with the recent experimental and theoretical based results of Sossi and Fegley (2018) and Sossi et al. (2019), that atmospheric escape caused by one large evaporation event at Earth is ineffective in fractionating moderately volatile element abundances to an extend that match the observations. Instead, the present moderately volatile element abundance in terrestrial planets most likely reflects accretion from several lower-mass embryos where each of them experienced different degrees of volatilization. One can conclude from these recent studies that a late accretion of chondritic material via planetary embryos that underwent melting processes might therefore not necessarily lead to a volatile rich planet.

7 Evolution of Earth's and Venus' Ar and Ne Noble Gas Isotope and K/U Bulk Elemental Ratios

Chondritic meteorites have long been linked to the bulk composition of the silicate part of Solar System terrestrial planets (e.g. Taylor 1964; Anders and Grevesse 1989; Javoy et al. 2010; Dauphas 2017; Schiller et al. 2018). Dauphas (2017) analyzed isotopes

of lithophile (O, Ca, Ti, Nd), highly siderophile (Ru) and moderately siderophile (Cr, Ni and Mo) elements from Earth's mantle and found that they record different accretion stages. According to analysed isotope compositions by Dauphas (2017), the first 60% of proto-Earth consisted of about 50% enstatite-like and 5% chondrite-like material, whereas all of the remaining accreted material consisted of enstatite-like impactors. According to Budde et al. (2019), most isotopic compositions can also be explained if mixtures of carbonaceous chondritic and non-carbonaceous chondritic material composed the building blocks of the accreting Earth. On the other hand, the oxygen isotopic composition of the Earth is nearly identical to that of enstatite chondrites. Oxygen is the most abundant element in rocks and one of the most abundant element in Earth depending on the estimated FeO content in the mantle and the Fe content in the core. Unless the Earth has a large enstatite chondrite component, the oxygen-isotopic composition and the iron content of the Earth cannot be matched (Dauphas 2017). The Dauphas (2017) O-isotope model and composition hypothesis agrees also with the two-component models of Ringwood (1979), Wänke and Dreibus (1988) and Lodders (2000), which assume almost the same percentage of O-isotope related enstatite chondritic material.

Recently, Schiller et al. (2018) investigated the ⁴⁸Ca/⁴⁴Ca isotope evolution of the composition of different Solar System objects such as Lunar and Martian meteorites, ordinary and carbonaceous chondrites, angrites, eucrites, ureilites, as well as very young chondrules from ordinary and carbonaceous chondrites and its role as a precursor element to rocky planets. The results of this study indicate that the ⁴⁸Ca/⁴⁴Ca composition of the earliest material in the inner disk originated from a ureilite-like composition and evolved to a terrestrial within the first million years after the formation of the Sun. Calculated accretion rates of the inner protoplanetary disk inferred from the Ca isotope composition of the involved objects reveal that the bulk disk composition at \approx 5-6 Myr after the formation of the Solar System consisted of a mixture of about 60% ureilite-like and 40% outer Solar System CI chondritic material (Schiller et al. 2018).

In contrast to this, 36 Ar/ 38 Ar and 20 Ne/ 22 Ne noble gases of the present Earth atmosphere and isotopes of water (D/H), and nitrogen (14 N/ 15 N) can also be reproduced by assuming a volatile content equivalent to a mixture of about 98% non-chondritic dry material with a contribution of about 2% undepleted carbonaceous CI-CM-like material (Marty 2012). Although, the results for the proto-Earth obtained by Dauphas (2017) and Schiller et al. (2018) might seem to contradict the findings of Marty (2012) representing present Earth, it is not necessarily the case. As shown and discussed in Odert et al. (2018), Benedikt et al. (2020) and in Sect. 6, the building blocks of terrestrial planets were most likely depleted in noble gases and other moderately volatile rock-forming elements. Additionally, one should not forget that proto-Earth and Venus underwent different evolutionary processes since the evaporation of the disk, including the possible escape of a primordial atmosphere, magma oceans and impactors with a depletion-modified abundance of their composition. These processes certainly played a crucial role in shaping the noble gas isotope, elemental and volatile evolution of these planets over their history.

Since decades, various researchers tried to model the compositional evolution of early Venus and Earth by reproducing the observed isotope ratios of the non-radiogenic noble gases Ar and Ne in their atmospheres (e.g., Cameron 1983; Pepin 1991, 1997; Becker et al. 2003; Gillmann et al. 2009). For instance, Gillmann et al. (2009) studied the escape and fractionation of ³⁶Ar/³⁸Ar, and ²⁰Ne/²²Ne noble gas isotope ratios of Venus' early atmosphere with a simple hydrodynamic escape model that is based on approaches and inputs provided by earlier studies of Hunten (1987), Kasting and Pollack (1983), Zahnle and Kasting (1986) and Chassefière (1996a, 1996b, 1997).

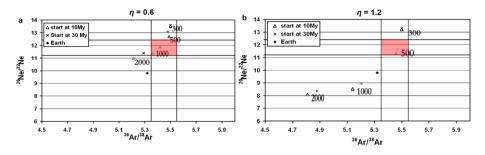


Fig. 10 Reproduction attempts of 36 Ar/ 38 Ar and 20 Ne/ 22 Ne isotope ratios according to different evolution scenarios related to atmospheric escape of early Venus with assumed average thermospheric temperatures between 300 - 2000 K. The two horizontal and vertical lines limit the ranges of the measured values of the isotopic ratios of Ne and Ar. They are delineated with the shaded area, which marks the scenarios that are in agreement with the measured isotopic ratios. a) Assumed initial atmospheric hydrogen content based on H₂O-amounts of 4 - 14 terrestrial oceans (TOs) for an η value of 0.6. b) Assumed initial atmospheric hydrogen does not escape in these scenarios (Gillmann et al. 2009)

In this approach Gillmann et al. (2009) introduced a factor $\eta = \varepsilon \zeta + \Phi_{sw}/\Phi_{EUV}$, that depends on the heating efficiency ε of the solar EUV flux Φ_{EUV} , generally assumed to be $\approx 15\%$ (Chassefière 1996b; Shematovich et al. 2014; Lammer et al. 2014), a geometrical amplification factor ζ related to the EUV heated and extended upper atmosphere and the relative contribution Φ_{sw} of the solar wind energy deposition to the total energy deposition (solar wind + EUV) near the exobase level. Since at the time of this study only very rough estimates existed on the young Sun's activity evolution, Gillmann et al. (2009) assumed η values between 0.3 - 20. The introduced geometrical amplification factor ζ was also based on arbitrary assumptions within a range of 4 - 8, since no upper atmosphere structures have been modelled in detail in these pioneering but simplified models.

Figure 10 shows resulting ³⁶Ar/³⁸Ar and ²⁰Ne/²²Ne isotope fractionation patterns in Venus atmosphere, obtained by the simple model approaches presented in Gillmann et al. (2009). One can see that their obtained evolutionary scenarios are consistent within the present error bar ranges. As mentioned in Sect. 2, one should note that the currently available ²⁰Ne/²²Ne ratio in Venus atmosphere and especially the quoted errors should be considered as very unprecise.

Since no upper atmosphere structure and temperature profiles were modelled in Gillmann et al. (2009), these authors assumed various average thermospheric temperatures between 300 and 2000 K. For high thermospheric temperatures > 1000 K, they obtained too large isotopic fractionations, because the densities of Ne and Ar and the escape fluxes are too large in their model approach. One can also see that for average thermosphere temperatures < 400 K, Ne is not fractionated enough to fit the present fractionation value due to its low abundance at the escape level. One should also note that the amount of hydrogen necessary to obtain the fractionation results for assumed η values of 0.6 (Fig. 10a) and 1.2 (Fig. 10b) contained 6 - 14 terrestrial oceans (TO) and 30 - 40 TOs, respectively (Gillmann et al. 2009).

More recently, Lichtenberg et al. (2016) studied the expected relative contribution of the solar wind energy deposition Φ_{sw} on early Venus, by assuming that the young Sun was either a weak or a moderately active young G star. These authors assumed a hydrogen-dominated upper atmosphere and modelled the precipitation of energetic hydrogen atoms (ENAs) that originate due to charge exchange between the planet's extended exosphere and solar wind protons (Chassefière 1996b, 1997). Chassefière (1997) and Gillmann et al. (2009) expected that these ENAs deposit a huge amount of energy into the upper atmosphere and

as a result enhance the H escape rate that influenced the noble gas fractionation. According to Lichtenberg et al. (2016), however, ENAs deposit their energy and heat mainly above the main EUV energy deposition layer. Although these particles modify the thermal structure of the upper atmosphere, the enhancement of the thermal escape rates caused by ENAs remains negligible (Lichtenberg et al. 2016). Since several pioneering studies including that of Cameron (1983), Pepin (1991) and Gillmann et al. (2009) lacked knowledge on the X-ray and EUV evolution of the young Sun/stars, magma ocean physics or impact studies, assumed arbitrarily reservoirs of H₂-molecules from primordial atmospheres and/or catastrophically outgassed H₂O inventories, one can consider that these arbitrarily chosen assumptions are insufficiently accurate and hence outdated.

Recently, Odert et al. (2018) tried to reproduce the 36 Ar/ 38 Ar and 20 Ne/ 22 Ne isotope ratios of Venus' atmosphere within a more realistic parameter space and found that EUV-driven hydrodynamic escape of an H₂-envelope that remained around proto-Venus after the disk evaporated can reproduce the observed atmospheric Ne and Ar isotope ratios within their large error bars. In the study of these authors, the EUV activity evolution of the young Sun corresponded to a slowly to moderately rotating young Sun (Tu et al. 2015; Lammer et al. 2020b). These findings stimulated Lammer et al. (2020b) to apply more complex models to figure out whether one can reproduce the measured atmospheric 36 Ar/ 38 Ar, 20 Ne/ 22 Ne, and, in case of the Earth, also 36 Ar/ 22 Ne noble gas isotope ratios and bulk planet K/U elemental ratios of Earth and Venus to constrain their early evolution.

While the atmospheric 36 Ar/ 22 Ne ratio of ≈ 18.8 in Earth's present day atmosphere cannot be explained by a pure nebula origin (Marty 2012; Marty and Allé 1994), the ³⁶Ar/³⁸Ar, ²⁰Ne/²²Ne ratios have been proposed to originate from the nebula only (e.g., Sasaki et al. 1989; Pepin 1991). Lammer et al. (2020b) assumed a reasonable set of initial protoplanet and impactor compositions based on the hypotheses of Dauphas (2017) and Schiller et al. (2018) and figured out under which protoplanetary maximum masses, at the time when the disk dissipated, Earth's present day atmospheric ³⁶Ar/²²Ne ratio can be reproduced. Additionally to these noble gas ratios, these authors also investigated whether the measured K/U ratios of both planets can be reproduced through the same process. As discussed in Sect. 6, elements like K are released at high temperatures (≈ 1500 - 2500 K) from magma oceans that form below primordial atmospheres (Schaefer and Fegley 2010; Fegley et al. 2016; see also O'Neill et al. 2020; this issue). That a fraction of K compared to its initial composition was lost during accretion can be seen in the K/U ratio of each planet (e.g. ≈ 13800 for the bulk silicate Earth (BSE) (Arevalo et al. 2009; Lodders et al. 2009), and \approx 7000 on Venus' surface (e.g., Surkov et al. 1976; Abdrakhimov and Basilevsky 2002), compared to a K/U abundance ratio of ≈ 64650 in the solar photosphere (Lodders et al. 2009) and ≈ 68930 in unprocessed chondrites (Lodders et al. 2009).

The model efforts carried out in Lammer et al. (2020b) are illustrated in Fig. 11 and summarized below:

Protoplanetary masses between 0.5 - 0.8 M_{Earth} are placed within the disk at Venus and Earth orbit locations. In case a primordial H₂-envelope survives hydrodynamic boil-off during disk dispersal, the atmosphere mass and upper atmosphere structure including its photospheric radius and effective EUV absorption radius are modelled by a hydrodynamic upper atmosphere model and embedded trace elements (³⁶Ar, ³⁸Ar, ²⁰Ne, ²²Ne, K) in solar composition, and outgassed from an assumed 1000 km deep magma ocean in abundances according to different assumed compositions - i.e., ureilite (UR), enstatite (EN), carbonaceous chondrites (CC); as illustrated in Fig. 10, are included into the remaining H₂-dominated primordial atmosphere.

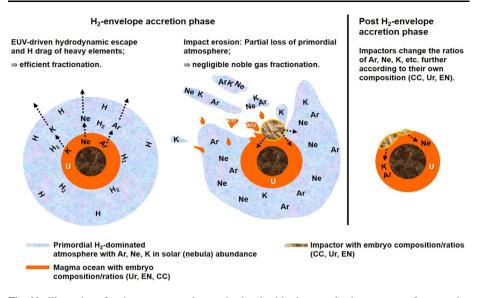


Fig. 11 Illustration of main processes and scenarios involved in the reproduction attempts of present day atmospheric ³⁶Ar, ³⁸Ar, ²⁰Ne, ²²Ne, noble gas isotope ratios and bulk K/U ratios on Venus and Earth by Lammer et al. (2020b)

- Hydrodynamic upper atmosphere models (Lammer et al. 2014; Kubyshkina et al. 2018b, 2018c) that can calculate the EUV-driven escape of H atoms of the primordial atmosphere and an evolution code that drag the heavier trace elements (³⁶Ar, ³⁸Ar, ²⁰Ne, ²²Ne, etc.) (Zahnle and Kasting 1986; Odert et al. 2018) out of the gravity field of accreting protoplanets are applied for various expected EUV-flux evolution scenarios from a slowly to a moderately and a fast rotating young Sun.
- Atmospheric mass loss of remaining H₂-envelopes caused by Moon- to Mars-mass impactors (Sect. 5) is considered additionally to the EUV-driven thermal escape and modelled with a multi-material SPH code (Schäfer et al. 2016), allowing simulation of gases, non-viscous fluids (Burger et al. 2018), granular media (Schäfer et al. 2017), solid bodies (Maindl et al. 2015; Burger and Schäfer 2014), and porous continua (Haghighipour et al. 2018). The masses of these impactors are assumed in Lammer et al. (2020b) to be admixtures of CC, UR, EN, and are added to the protoplanetary mass where they are also altering the ³⁶Ar, ³⁸Ar, ²⁰Ne, ²²Ne, and K content of proto-Earth and Venus.
- Evolution of the photospheric radii within the escaping H₂-envelope according to mass additions of accreting bodies, gravity change and atmospheric escape as well as its effects on the isotope/element ratio evolution is also included in the study of Lammer et al. (2020b).

The illustrations in Fig. 11 show the main phases and processes that are included in the recent model simulation attempts of Lammer et al. (2020b) for reproducing proto-Earth's and Venus' atmospheric Ar and Ne noble gas isotope and bulk K/U ratios. By reproducing present day measurements within their error bars these authors were able to constrain the growth rate of both planets in the presence of the solar nebular as well as a certain range for the EUV flux evolution of the young Sun.

Lammer et al. (2020b) present several model runs that cannot reproduce the abovementioned elemental ratios and some that could. Here we present three cases for which

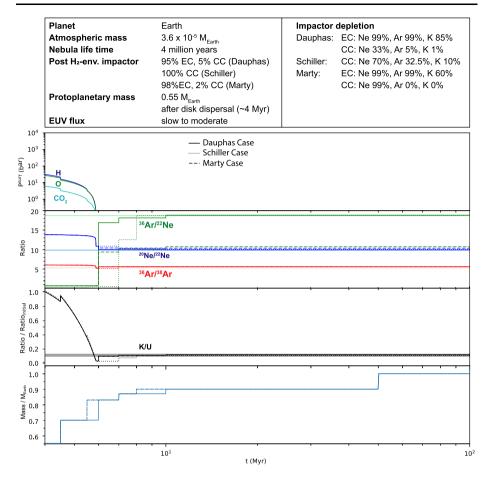


Fig. 12 Reproduction attempts of Earth's present-day atmospheric ${}^{36}\text{Ar}/{}^{38}\text{Ar}$, ${}^{20}\text{Ne}/{}^{22}\text{Ne}$, ${}^{36}\text{Ar}/{}^{22}\text{Ne}$ and bulk K/U ratios by assuming different post-H₂-envelope impactor compositions (after Lammer et al. 2020b). The rise of Earth's final mass at ≈ 50 Myr is related to the Moon-forming impact event for which it is assumed that the final impactor does not significantly alter the composition of proto-Earth

the atmospheric Ar and Ne isotope ratios and Earth's K/U bulk ratio can be reproduced if one assumes that proto-Earth accreted enough mass within the disk to accrete a small H_2 -dominated primordial atmosphere which was subsequently lost by EUV-driven hydrodynamic escape.

Figure 12 shows evolution scenarios for early Earth with an initial protoplanetary mass of 0.55 M_{Earth} at the end of the disk at ≈ 4 Myr (Bollard et al. 2017; Wang et al. 2017) orbiting a young Sun that evolves along an EUV activity that corresponds to an intermediate rotating young G-type star (Tu et al. 2015; Güdel 2020; this issue). In such a case proto-Earth would have a post-H₂-envelope impactor composition of $\approx 95\%$ EN and $\approx 5\%$ CC (Dauphas 2017; solid lines), 100% CC (Schiller et al. 2018; dotted lines) and 98% EC and 2% CC (Marty 2012; dashed lines). Earth's atmospheric ³⁶Ar/²²Ne ratio of ≈ 18.8 could be reproduced by the model methods applied in Lammer et al. (2020b) for the Dauphas and Schiller scenarios but with different assumed growth rates and depletions of the post-H₂-envelope impactors. A significant difference in these production attempts is the amount of volatile rich CCs. The

more CCs are delivered to the growing protoplanet the faster it has to grow so that the volatile elements can escape together with the captured primordial H₂-dominated atmosphere. If one uses Ar and Ne isotope ratios corresponding to CI material instead of CCs the outcomes only change marginally yielding more or less similar results. If proto-Earth captured such a tiny primordial atmosphere, the initial composition suggested by Marty (2012) cannot reproduce Earth's present atmospheric ³⁶Ar/²²Ne ratio. However, if proto-Earth accreted less mass so that no primordial atmosphere remained after disk dispersal the composition derived in Marty (2012) cannot be excluded. Due to this, Lammer et al. (2020b) concluded that the reproduction of the present day atmospheric Ar and Ne isotope ratios, by including a nebula component that escaped during a few million years after the disk dissipated, yields an upper limit for the mass of the proto-Earth of ≤ 0.58 M_{Earth} at the time when the disk evaporated.

As mentioned before, an important constraint for the composition of the post-nebula accretion is concealed in the atmospheric ratio of 36 Ar/ 22 Ne. Lammer et al. (2020b) found that different admixtures of UR, EC and CC-like material, for the impactors that were delivered after the primordial atmospheres escaped, could reproduce the atmospheric Ar and Ne ratios only if the impactors are depleted in the studied elements. As discussed in Sect. 6 and shown in Benedikt et al. (2020), one can expect that lower mass planetary embryos that developed their own magma ocean phases should be highly depleted in noble gases and outgassed volatiles.

By assuming an initial proto-Earth composition (95% EC, 5% CC) as suggested by Dauphas (2017), one can see in Fig. 12 that Earth's atmospheric Ar and Ne isotope ratios can only be reproduced if the accreting EC material is strongly depleted in noble gases (\approx 99%) and volatile elements, while the remaining CCs are less depleted (Ne: \approx 33%; Ar \approx 5%). In such a case, most of the delivered CCs would be unfractionated material instead of highly depleted larger planetary embryos. For reproducing, the atmospheric Ar and Ne isotope ratios, if one assumes that $\approx 100\%$ of CCs were involved in early Earth's post-nebula accretion as suggested by Schiller et al. (2018), one needs impactors that are more efficiently depleted in these volatile elements (Ne: $\approx 70\%$; Ar $\approx 32.5\%$). In such a case, a larger fraction of the CC-related building blocks should have been bigger bodies that underwent magma ocean related outgassing events thereby losing much of its volatiles. If one assumes an initial proto-Earth composition of $\approx 98\%$ volatile poor material and $\approx 2\%$ CCs, as suggested by Marty (2012), then Earth's atmospheric 36 Ar/ 22 Ne ratio can only be reproduced if the Earth accreted slower with most of its mass being accreted after the disk dissipated. In such a case, no primordial atmosphere was captured by proto-Earth. One can see from Fig. 12 that such a scenario cannot reproduce the ³⁶Ar/²²Ne ratio if only even a tiny primordial atmosphere was present for a few million years after the disk evaporated. One should also note that the assumed depletions in the model runs of Lammer et al. (2020b) for the post- H_2 -envelope impactor material are not chosen completely arbitrarily. For instance, planetary building blocks where only one noble gas element is strongly depleted such as Ne, but Ar and/or moderately volatile elements such as K are not depleted at all, are physically not realistic (Benedikt et al. 2020).

Figure 13 shows a best-case evolution scenario obtained by Lammer et al. (2020b) that can reproduce the present known atmospheric Ar and Ne isotope and measured surface K/U ratios on Venus by assuming a similar initial composition and the same young Sun EUV flux evolution track as assumed for the successful reproduction of the elements on Earth in Fig. 13.

In this particular scenario, proto-Venus accreted its whole mass during the disk lifetime and was released from the disk at ≈ 4 Myr with an H₂-dominated primordial atmosphere with a hydrogen partial pressure of ≈ 1000 bar. The escaping H atoms dragged

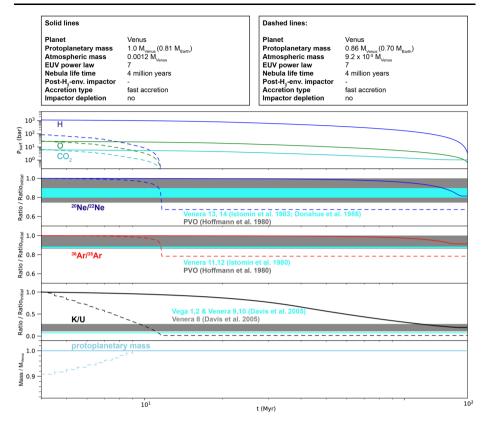


Fig. 13 Successful (solid lines) and failed (dashed-lines) reproduction attempts of the atmospheric 36 Ar/ 38 Ar and 20 Ne/ 22 Ne isotope and K/U surface ratios at Venus. A totally accreted Venus at the time when the disk evaporated can reproduce the present observations if the captured primordial H₂-dominated atmosphere is exposed to a slowly to less than moderately rotating young Sun (solid lines). All ratios are normalized to the initial ratio. A lower mass proto-Venus with a mass of about 0.86 M_{Venus} over-fractionates these noble gas isotopes and yields lower K/U ratios as measured by Venera 8, 9, 10, and Vega 1 and 2 (after Lammer et al. 2020b)

away the heavier trace species and fractionated ³⁶Ar/³⁸Ar, ²⁰Ne/²²Ne from the initially solardominated ratios to the ones in today's Venus' atmosphere (Lammer et al. 2020b). Besides the noble gas isotope ratios, the corresponding K/U ratio lies slightly above Vega 1, 2 and Venera 9, 10 data but within the range of the Venera 8 landing site (Basilevsky et al. 1997; Abdrakhimov and Basilevsky 2002).

The dashed lines in Fig. 13 show an example modelled by Lammer et al. (2020b) that failed to reproduce the present-day ratios for Venus, even within the large uncertainties of the analysed measurements. In this failed attempt a less massive proto-Venus with about 0.86 M_{Venus} (0.69 M_{Earth}) leads to an under fractionation of the Ar, Ne isotope and K/U ratios. It should also be noted that for Venus' case Lammer et al. (2020b) did not model the ³⁶Ar/²²Ne ratio due to the large error bars. New precise measurements in future Venus missions are necessary so that one can test whether the planets building blocks can be linked to material related to EC or to a mixture of EC and CC with remnants from a primordial atmosphere.

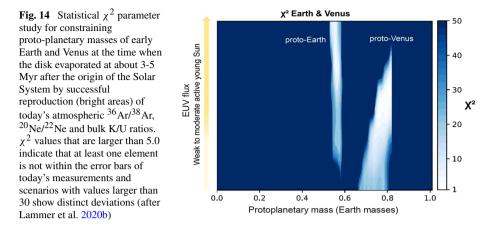


Figure 14 shows the statistical outcome of many modelled proto-Earth and proto-Venus masses at the end of the disk lifetime (3 - 5 Myr) by Lammer et al. (2020b). The white areas illustrate the parameter space of EUV flux and protoplanetary mass in which the ³⁶Ar/³⁸Ar and ²⁰Ne/²²Ne noble gas isotope ratios and bulk planet K/U elemental ratios were successfully reproduced if the protoplanets accreted enough mass during the disk life time so that small primordial atmospheres were captured and lost afterwards, for particular compositions on both planets within their error bars.

One can see that the present day ratios can be reproduced if proto-Earth and Venus captured primordial atmospheres before the disk evaporated between 3 - 5 Myr with masses of about 0.53 - 0.58 M_{Earth} and 0.68 - 0.81 M_{Earth} , respectively, if the young Sun followed an EUV flux evolution path between a slowly and moderately rotating young G-type star (Tu et al. 2015; Lammer et al. 2020b; Güdel 2020, this issue). One should also note that in Earth's case the atmospheric ³⁶Ar/²²Ne ratio can also be reproduced as in Fig. 14 for the initial compositions as suggested by Dauphas (2017) and/or Schiller et al. (2018) if one adjusts the assumed depletion of the studied elements in the post-H₂-envelope impactors. For proto-Earth-masses larger than $\approx 60\%$ of its present mass at the time when the disk evaporated, Lammer et al. (2020b) could not reproduce the atmospheric Ar and Ne isotope ratios. Furthermore, for proto-Earth masses $\geq 80\%$ the accreting Earth would not have been able to lose its captured primordial atmosphere. These findings are also in agreement with the study of Pahlevan et al. (2019) that the accordance within ≈ 10 –20 % between Earth's seawater and carbonaceous chondrites in D/H indicate that the prior existence of early Earth's atmospheric H₂ inventory on post-giant-impact Earth was < 20 bars (see also Sect. 2).

According to the resent work of Burkhardt et al. (2019) an analysis of measured Ti and Sr isotopic compositions of Ca, Al-rich inclusions (CAIs) from the Allende CV3 chondrite indicate that CAI-like material was the carrier of the nucleosynthetic isotopic anomalies in most minor and trace elements. These authors found that a variable admixing of CAI-like refractory material to an average inner solar nebula component could explain the planetaryscale Ti and Sr isotope anomalies and the elemental and isotopic difference between noncarbonaceous (NC) and carbonaceous (CC) nebular reservoirs for these elements. If this is the case, then the differences between meteorites and Earth may simply reflect the timing or spatial location for adding this CAI-like material to the solar nebula. The findings by Lammer et al. (2020a) that the post-H₂_envelope impactors should have also be at least partially related to the CC-reservoir agrees also with the dating of terrestrial planet building blocks from CAIs, from which it is inferred that CCs finished their accretion at times around $\approx 3 - 4$ Myr after the formation of the Solar System (Kruijer et al. 2017).

Compared to both large terrestrial planets in the Solar System, the evolution for the lower mass planet Mars was different and one cannot constrain it in a similar way as in the case of Venus and Earth. Mars' mass is too low to capture a relevant H₂-envelope during the phase the small planet grew within the gas disk. Noble gases could have also been lost and dragged away by the escaping H atoms as long as hydrogen, which originated from solidified magma ocean-related outgassed steam atmospheres, dominated the upper atmosphere (Pepin 1994; Pepin and Porcelli 2002). Afterwards ionization related fractionation through thermal and non-thermal escape (i.e., ion pick up, sputtering, dissociative recombination) of the planet's CO_2 atmosphere resulted in today's observed Ar, Ne and N isotopic compositions (e.g., Jakosky et al. 1994; Hutchins and Jakosky 1996). In the next section, the isotopic fractionation and escape of the Martian atmosphere and initial water inventory are discussed.

8 Isotopic Fractionation as a Constraint on the Atmospheric Evolution and Water Inventory on Mars

The small gravity and absence of an intrinsic magnetic field allow efficient atmospheric escape from Mars induced by the solar EUV radiation and the solar wind. Its isotopic compositions show heavy-isotope enrichment in hydrogen (D/H), carbon (¹³C/¹²C), nitrogen (¹⁵N/¹⁴N), and noble gases (²²Ne/²⁰Ne, ³⁶Ar/³⁸Ar, and non-radiogenic Xe isotopes) except Kr compared to their possible origins (nebula gas, chondrites, and comets), as an outcome of atmospheric escape. As mentioned above, noble gases are particularly useful to constrain the atmospheric evolution since they are mostly partitioned into the atmosphere and do not form molecules so that the escape processes are limited to hydrodynamic escape and to atmospheric sputtering for lower mass bodies without magnetospheres (e.g., Pepin 1991, 1994; Jakosky et al. 1994). Nitrogen behaves similarly to noble gases as it is chiefly partitioned into the atmosphere but ion-molecule photochemical reactions also induce atmospheric escape (e.g., Fox and Dalgarno 1983; Fox 1993a, 1993b; Fox and Hać 1999; see also Avice and Marty 2020; this issue). Carbon is partitioned into multiple reservoirs - the atmosphere, carbonate, organic compounds, CO₂/CH₄-clathrate, and CO_2 ice (e.g., Kurahashi-Nakamura and Tajika 2006; Longhi 2006; Hu et al. 2015; Wray et al. 2016) so that the isotopic fractionation is more complicated.

As briefly discussed in Sect. 4, fractionation mechanisms include diffusive separation by mass between the homopause and exobase (e.g., McElroy and Yung 1976; Jakosky et al. 1994, 2017). The number densities of photochemically long-lived atmospheric species follow their own scale height above the homopause. As the escape processes remove atoms, molecules, and ions from the vicinity of the exobase, the diffusive separation causes mass-dependent fractionation. Since atmospheric sputtering is so energetic that mass-dependent fractionation in the removal process itself is limited, the isotopic fractionation through sputtering is chiefly determined by the diffusive separation (Jakosky et al. 1994).

The diffusive separation between the homopause and exobase was recently probed with the help of the MAVEN spacecraft (Jakosky et al. 2017; Slipski et al. 2018). Measurements of the N₂/⁴⁰Ar ratio indicated substantial variations at the homopause ($\approx 60 - 140$ km) and exobase ($\approx 140 - 200$ km) altitudes. This represents an unresolved combination of real geophysical variations with time and location but confirms the concept of elemental and isotopic fractionation. Photochemical processes fractionate isotopic ratios due to the difference in reaction rates between isotopic species. Both photodissociation of CO and

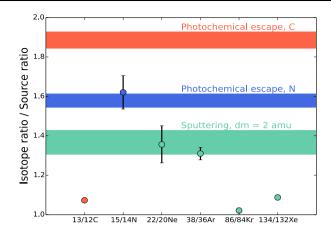


Fig. 15 Isotope ratios of volatile elements in the present-day Martian atmosphere with respect to their source isotope ratio (data points). We assumed the source values of Martian mantle informed from Martian meteorites for C and N (Wright et al. 1992; Mathew and Marti 2001), and solar abundances for noble gases (Pepin 1991; Pepin et al. 2012). Data for the Martian atmosphere are from Mahaffy et al. (2013a), Pepin (1991), Wong et al. (2013), Atreya et al. (2013), Conrad et al. (2016). The expected steady state values are shown as horizontal ranges. We assumed the upper atmospheric temperature to be 200 K and the homopause-exobase separation 60-80 km to calculate the fractionation factor for the diffusive separation and adopted additional fractionation for photochemical escape of suprathermal C and N atoms that originate mainly from dissociative recombination of N_2^+ and CO^+ molecules, f = 0.62 and 0.74, respectively (Manning et al. 2008; Hu et al. 2015)

dissociative recombination of CO^+ and CO_2^+ could induce preferential loss of ¹²C rather than ¹³C (Fox and Hać 1999; Hu et al. 2015). Dissociative recombination of N_2^+ causes isotopic fractionation between ¹⁴N and ¹⁵N (Wallis 1989; Fox and Hać 1997; Manning et al. 2008).

8.1 Isotopic Compositions of the Present-Day Martian Atmosphere

The observed isotopic compositions of N₂, Ne, and of the non-radiogenic Ar ($\delta^{15}N = 572 \pm 82\%$, ${}^{20}Ne/{}^{22}Ne = 10.1 \pm 0.7$, and ${}^{36}Ar/{}^{38}Ar = 4.2 \pm 0.1$, Pepin 1991; Wong et al. 2013; Atreya et al. 2013) are well described by the steady state where the loss via atmospheric escape balances the supply via, for example, volcanic outgassing (Fig. 15), and is given by $I/I_0 = 1/f$, where I, I_0 , and f are the isotope ratio in the atmosphere, the isotope ratio of the source, and the net fractionation factor.

The estimated lifetimes of N (via sputtering and photochemical escape) and Ne (via sputtering) are shorter than one billion years, supporting the hypothesis of steady state (Jakosky et al. 1994; Manning et al. 2008; Kurokawa et al. 2018). The lifetime of Ar is comparable to the age of Mars (Jakosky et al. 1994), but its isotopic ratio being close to the value of steady state suggests that the lifetime was shorter on earlier time (Kurokawa et al. 2018). Being close to the steady state suggests

- the model for their isotopic fractionation is valid, and
- the present-day isotopic compositions do not reflect the earlier period (i.e. before ≈ 4 billion years ago (Ga).

The observed ¹³C/¹²C ratio being smaller than the ratio expected from the steady state with sputtering and/or photochemical escape (Fig. 14) suggests that atmospheric

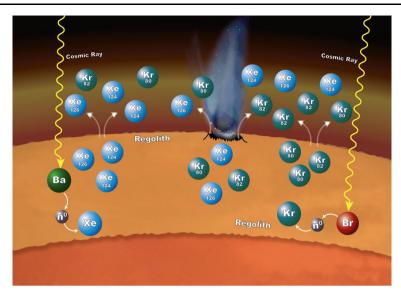


Fig. 16 Illustration of the cosmic ray-related surface chemistry that takes place in the Martian regolith, as an explanation why particular Xe and Kr isotopes are more abundant in the atmosphere than expected (Conrad et al. 2016). (NASA/GSFC/JPL-Caltech)

escape is not the only sink of an ancient dense atmosphere. Other sinks include carbonate precipitation (Hu et al. 2015; Wray et al. 2016), organic compounds deposition, and CO₂-clathrate formation following the basal melting of CO₂-ice caps (Longhi 2006; Kurahashi-Nakamura and Tajika 2006). Hu et al. (2015) considered atmospheric escape, carbonate precipitation, and concluded that Mars at \approx 3.8 Ga had an atmospheric pressure \leq 1 bar. Based on an experimental study (Eiler et al. 2000), part of the atmospheric CO₂ may convert into CO₂ ice without significant isotopic fractionation ($\delta^{13}C_{ice-vapor} < -0.4\%$ at 135 K). Considering additional sinks may allow a denser atmosphere.

The isotopic composition of Xe is also less fractionated compared to the value expected from the steady state with atmospheric sputtering. It suggests that the present-day Xe is a remnant of early hydrodynamic loss (Pepin 1994) or a result of ionization-induced escape during the first 500 Myr (Zahnle et al. 2019; see also Avice and Marty 2020; this issue). Ionization-induced escape can lead to high Xe⁺ escape rates, while the escape of other heavy noble gases such as the hardly ionizable Kr is not significant. Due to the ionization of Xe, its cross-section becomes significantly larger, which makes it possible to escape and fractionate, whereas Kr stays neutral and, therefore, cannot escape due to its lower cross-section (Avice and Marty 2020; this issue). Furthermore, as illustrated in Fig. 16, atmosphere-surface interaction processes in combination with cosmic ray exposure and its related chemistry in the Martian surface can also contribute to an enhancement of some of the Xe and Kr isotopes in the Martian atmosphere (Conrad et al. 2016).

Because of the present thin Martian atmosphere, galactic cosmic rays penetrate into the surface material. According to these authors, if cosmic rays hit a barium (Ba) atom, Ba can lose neutrons. Xe atoms can catch some of those neutrons to form the isotopes ¹²⁴Xe and ¹²⁶Xe. In a similar way, Br can lose neutrons to Kr, leading to the formation of ⁸⁰Kr and ⁸²Kr isotopes. Impacts and abrasion allow that these isotopes enter the atmosphere, where they underwent escape by impact erosion.

8.2 Hydrodynamic Escape from Early Mars

Models of the upper atmosphere have shown that hydrodynamic escape prohibits early Mars from developing a stable atmosphere (Tian et al. 2009; Erkaev et al. 2014; Odert et al. 2018). The primordial H₂/He atmosphere that originated from the solar nebula (Pepin 1991, 1994; Stökl et al. 2015, 2016; Saito and Kuramoto 2018) was lost due to a short and efficient boiloff phase during disk dispersal (Lammer et al. 2018). This was followed by a catastrophically outgassed steam atmosphere during the solidification of a magma ocean (Elkins-Tanton 2008; Lebrun et al. 2013; Erkaev et al. 2014) that has been lost early on, and a secondary atmosphere later from volcanic outgassing (Lammer et al. 2013).

Hydrodynamic escape involves mass-dependent elemental and isotopic fractionation: heavier species are bound by the gravity more strongly and enrich in the remaining atmosphere (e.g., Hunten et al. 1987; Pepin 1991, 1994; Odert et al. 2018). The isotopic compositions of trapped non-radiogenic Xe in Martian meteorites Allan Hills 84001 (ALH84001) formed 4.1 Ga and North West Africa 7034 (NWA7034) formed 4.4 Ga are indistinguishable from current atmospheric Xe (Cassata 2017; see also Avice and Marty 2020; this issue), suggesting that the fractionation has completed before 4.4 Ga. If atmospheric Xe fractionated before 4.4 Ga and was later influenced by impact erosion and replenishment through time (de Niem et al. 2012; Kurokawa et al. 2018; Sakuraba et al. 2019), atmospheric escape of Xe in the later period via a photo-ionized hydrogen wind (Zahnle et al. 2019) mentioned before might compensate the replenishment.

8.3 Atmospheric Evolution After the Build-up of a Secondary Atmosphere

Since except Xe the present-day isotopic compositions do not probe the early atmospheric history, Kurokawa et al. (2018) utilized the isotopic compositions in the atmosphere 4.1 Ga recorded in ALH84001 (Mathew and Marti 2001). While a dense atmosphere preserves primitive isotopic compositions in their model, a thin atmosphere on early Mars is severely influenced by stochastic impact events and following escape-induced fractionation (Fig. 17).

Though the timing when Mars started to build up a stable atmosphere depends on the initial spin rate and related activity of the Sun (Tu et al. 2015; Güdel 2020; this issue) as well as on the volcanic outgassing activity on early Mars (Amerstorfer et al. 2017; Lammer et al. 2018), the less fractionated N (δ^{15} N = 7‰) and Ar (36 Ar/ 38 Ar > 5.0) recorded in the ALH84001 meteorite (Mathew and Marti 2001) which formed at 4.1 Ga (Lapen et al. 2010) suggest that a denser atmosphere whose pressure was > 0.5 bar has already existed at the time (Kurokawa et al. 2018).

Other geochemical and geological estimates of paleo-pressures also point to a moderately dense atmosphere whose pressure is comparable to or slightly ≤ 1 bar (Kite 2019 and references therein). Elevated radiogenic Ar (40 Ar/ 36 Ar) and Xe (129 Xe/ 132 Xe) ratios recorded in ALH84001 (Cassata et al. 2010; Cassata 2017) are consistent with the Martian atmosphere originating at ≈ 4.1 Ga from late outgassing. The atmospheric 40 Ar on Mars is mainly the product of the past cumulative outgassing of 40 K from Mars' mantle and crust. This Ar isotope can therefore be used as a tracer of the past evolution of volatiles in Mars' crust, mantle and atmosphere. Similar as 36 Ar and 38 Ar, 40 Ar can also escape from the atmospheric escape, which is more efficient in the early evolution, does not significantly affect the present abundance of the atmospheric 40 Ar, which indicates that the outgassing of this isotope from the mantle occurred relatively late in the Martian evolution, making 40 Ar to a tracer of outgassing but not for atmospheric escape.

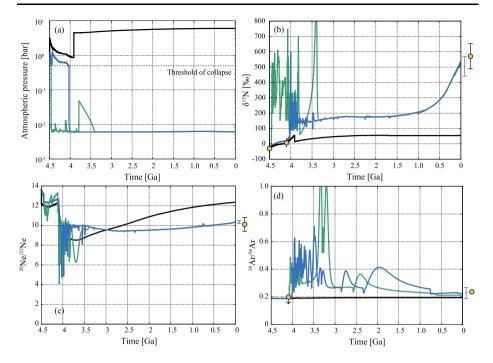


Fig. 17 Results of Monte Carlo simulations for the evolution of (a) atmospheric pressure, (b) N, (c) Ne, and (d) Ar isotope ratios (modified after Kurokawa et al. 2018). Three cases are shown: the atmosphere has been dense for 4.5 Gas (black), the atmosphere becomes thin after 4.1 Ga (blue), and the atmosphere becomes thin before 4.1 Ga (green). Data points in blue show the average value in the cases where the atmosphere collapsed at some point. Data points in yellow are the value of Martian atmosphere (Wong et al. 2013; Pepin 1991; Atreya et al. 2013; Mathew and Marti 2001)

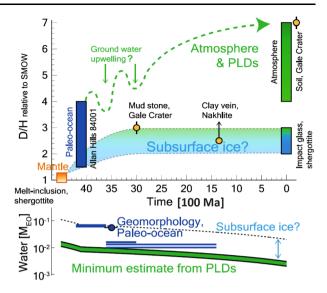
8.4 Loss of Water on Mars

The D/H ratio recorded in ALH84001 (Boctor et al. 2003; Greenwood et al. 2008; Usui et al. 2017) is in the range of $\delta D = 500 - 3000\%$, which means that more than 1/3 of the surface water could have been lost at 4.1 Ga (Fig. 18). The surface water has experienced further loss to reach $\delta D = 5000 - 6000\%$ (Usui et al. 2012; Kurokawa et al. 2014; Mahaffy et al. 2013b; Villanueva et al. 2015). An intermediate D/H ratio ($\delta D = 1000 - 2000\%$) recorded in impact glasses in Martian meteorites suggests that the majority of near-surface water might have been stored as subsurface ice isolated from the atmosphere (Usui et al. 2015; Kurokawa et al. 2016; Grimm et al. 2017).

On the other hand, if one assume for a Noachean ocean volume of about 500 m global equivalent layer (GEL) or 0.05 times the terrestrial ocean mass (M_{TO}) (Di Ahille and Hynek 2010), the water loss before 4.1 Ga is > 160 m GEL, which might exceed the integrated oxygen escape (Jakosky et al. 2018). This imbalance between H and O escape drives the oxidation of the Martian surface environment (Lammer et al. 2003b; Zahnle et al. 2008).

That Mars lost large amounts of water was also suggested by Lammer et al. (2003a), who estimated the present and past H₂O-ice reservoirs, which are and were in exchange with the atmosphere, by using the observed D/H ratio in the atmospheric water vapour, measured D/H ratios in Martian SNC meteorites and D/H isotope ratios based on a study regarding asteroid and cometary water delivery to early Mars by Lunine et al. (2003).

Fig. 18 Evolution of D/H (upper panel) and water volume (lower panel) on Mars (modified after Kurokawa et al. 2014, 2016; Usui et al. 2017). Data for D/H are from Usui et al. (2012, 2015, 2017), Boctor et al. (2003), Greenwood et al. (2008), Webster et al. (2013a, 2013b), Mahaffy et al. (2015), Villanueva et al. (2015), Krasnopolsky (2015), and Aoki et al. (2015). Water volume was computed by assuming Rayleigh fractionation with f = 0



These authors estimated the present H₂O-ice reservoir to be equivalent to about 11 - 27 m GEL. Here, the upper value more or less agrees with a more recent study by Carr and Head (2015) who found that a 34 m GEL resides in the Martian polar-layered deposits and shallow ground ice. By considering such a GEL value, a reservoir exchangeable with the atmosphere on Mars 3.5 Ga with a thickness of between ≈ 60 m is estimated (Lammer et al. 2003a; Carr and Head 2015).

By using these reservoirs from 3.5 Ga, Lammer et al. (2003a) found a global ocean at ≈ 4.5 Ga with a depth of ≈ 115 m GEL, if the initial isotopic ratio of Mars' water was 1.2 times the terrestrial sea water ratio (Lunine et al. 2003). A GEL of ≈ 115 m at ≈ 4.5 Ga is also in agreement with the initial water inventory in Mars mantle at the end of its accretion as estimated by Dreibus and Wänke (1987) by comparing the solubility of H₂O and HCl in molten silicates. In this way, Dreibus and Wänke (1987) obtained an H₂O content in the Martian mantle of ≈ 36 ppm corresponding to a global Martian ocean with a depth of about 130 m. The difference with respect to the initial value (di Achille and Hynek 2010) may have escaped unfractionated to space between 3.5 - 4.5 Ga and/or partly stored in the subsurface (Orosei et al. 2018) (see also Fig. 18).

8.5 Redox State of the Martian Atmosphere in Relation to the Atmospheric Escape

The preferential loss of hydrogen into space causes the oxidation of the Martian surface (e.g. Lammer et al. 2003b). Additionally, carbon escape may have dominated against oxygen escape on early Mars (Tian et al. 2009; Amerstorfer et al. 2017) which may have also oxidized Mars' surface. If this is the case, the redox state of the early Martian atmosphere could be more reducing than modern Mars. Thus, it is worth considering the possibility of an early reducing Martian atmosphere.

First, the outgassing composition of presumably early Martian volcanos could have been more reducing than Earth's since the oxygen fugacity of the Martian mantle could be up to 3 log-units lower than of the terrestrial mantle and possibly close to the IW-buffer especially during its early period (Wadhwa 2001; Herd et al. 2002; Righter et al. 2008;

McSween et al. 2009). When assuming an oxygen fugacity of log $f_{O2} = IW+2$ and a temperature of ≈ 1500 K, with ≈ 100 m GEL and ≈ 1 bar CO₂, the outgassing composition should be roughly H₂: H₂O: CO: CO₂ ≈ 10 : 1: 20: 1, which is much more reducing than Earth's present volcanic flux (e.g. Zolotov 2003). Although the oxygen fugacity of the Martian crust may vary and could be higher than in the mantle (Wadhwa 2001), the Martian basaltic volatiles could be dominated by H₂ and CO relative to CO₂ based on a thermodynamic calculation of Gaillard and Scaillet (2014). Given the H₂-rich volcanic flux, the early Martian atmosphere should have been very reducing and possibly rich in CO (Sholes et al. 2017). This is due to CO₂ continuously being photolyzed into CO, whereas the oxidation of CO back into CO₂ is not efficient if the volcanic flux provides sufficient reducing agents into the atmosphere. This might then lead to a CO-runaway (Yung and DeMore 1998; Zahnle et al. 2008).

The reducing early Martian atmosphere may provide a solution to explain why the atmospheric CO₂ of Mars ($\delta^{13}C_{PDB} = +46 \pm 4\%$; Webster et al. 2013a, 2013b) is so enriched in ¹³C compared to that of the Martian mantle ($\delta^{13}C_{PDB} = -25 \pm 5\%$; Wright et al. 1992). Galimov (2000) originally proposed that CO₂ could have been a minor component in the early Martian atmosphere as compared with CH₄ or CO from a thermodynamic point of view. CH₄ and CO are depleted in ¹³C and preferentially escaped into space, whereas the ¹³C-enriched CO₂ is ultimately deposited as carbonate. Also, as mentioned above, Hu et al. (2015) argue that photodissociation of CO may significantly induce preferential loss of ¹³C. These authors estimate that the current atmospheric ¹³C/¹²C ratio and rock and soil carbonate measurements indicate an early Martian atmosphere with a surface pressure < 1 bar at ≈ 3.8 Ga ago. Hu et al. (2015) points out that only scenarios with large amounts of carbonate formation in open lakes would permit higher values of up to 1.8 bar.

These isotopic evidence and studies on atmospheric escape (Lammer et al. 2018) as well as a recent investigation on the formation of Martian geologic features, including degraded impact craters, lakes and valley networks by Palumbo et al. (2020) indicate that Mars most likely did not need a warm and wet Noachian climate, with surface temperatures above-freezing that allowed rainfall.

In these scenarios, the atmospheric ¹³C-enrichment is caused by preferential escape of ¹³C into space though time. Based on the isotopic studies of carbonate globules in ALH84001 (Eiler et al. 2002; Halevy et al. 2011), the enrichment in ¹³C may have possibly occurred before 3.9 Ga, when the geomagnetic field may still have existed Lillis et al. (2013) and prohibited carbon escape into space via pick-up-ion sputtering induced by the solar wind. On the other hand, C atoms most likely escaped efficiently during this period by thermal escape (Tian et al. 2009; Amerstorfer et al. 2017). Moreover, due to the higher EUV flux of the young Sun, early Mars upper atmosphere expanded to larger distances so that the intrinsic magnetic field may not have shielded exospheric particles against escape. A reducing atmosphere (a smaller CO₂ mixing ratio) leads to a higher temperature in the upper atmosphere due to inefficient cooling and therefore allows efficient escape (Kulikov et al. 2007; Tian et al. 2009; Johnstone et al. 2018). Further experimental and model studies are required to test the reducing early Martian atmosphere scenario.

An additional alternative mechanism for early ¹³C-enrichment recorded in ALH84001 than atmospheric escape is isotopic fractionation by solar UV photodissociation of CO₂, which may cause significantly large carbon isotopic fractionation of about 200‰ based on theoretical calculations of absorption spectra of the CO₂ isotopologues (Schmidt et al. 2013). Such a process is also relevant if Mars' early atmosphere would have been lost during younger ages compared to that of ALH84001. Although the effect has not yet been tested by laboratory experiment, solar UV could dissociate ¹²CO₂ preferentially,

and thus provide ¹³C-enriched carbonate precipitating from the remaining CO₂. In this case, highly ¹³C-depleted CO is produced in the atmosphere. Under this reducing condition and the presence of water vapor, CO could be converted mainly into formalde-hyde and removed from the atmosphere (Bar-Nun and Hartman 1978; Pinto et al. 1980; Bar-Nun and Chang 1983). If this scenario is correct, then ¹²C preferentially deposited as organics into sediment, but did not escape into space.

9 Conclusion

We discussed the relevance of various loss processes that led most likely to today's atmospheric Ar, and Ne noble gas isotope anomalies and the depletion of moderately volatile rock-forming elements on the terrestrial planets. A complex interplay between the young Sun's EUV flux, the accreted protoplanetary mass within the disk lifetime, a possible accumulated primordial H₂-dominated atmosphere, the impacting bodies that were involved in the accretion before and after disk dispersal, as well as the depletion of volatile elements from planetary embryos and the final protoplanets led eventually to today's measured ratios. The discovery of many low mass exoplanets that have captured a primordial H₂-dominated atmosphere that never escaped to space during billions of years indicates that the accretion of smaller planets is also a fast process. One finds from early atmosphere evolution models that measured atmospheric Ar and Ne isotope ratios can be reproduced if proto-Venus and Earth were released from the solar nebula at ≈ 4 Myr with a small captured H₂-dominated envelope and maximum protoplanetary masses at that time of $\approx 0.85 - 1.0 M_{Venus}$ and 0.53 $- 0.58 M_{Earth}$, respectively.

If proto-Earth accreted such a mass from enstatite-like material during the disk lifetime, a tiny primordial atmosphere would have been lost soon after the disk dissipated and the presentday atmospheric Ar and Ne isotope ratios can be reproduced if the post-nebula impactors contained $\leq 5 \approx 95$ % strongly depleted enstatite chondrites. If higher amounts of carbonaceous chondrites were involved in early Earth's accretion then the present atmospheric Ar and Ne ratios can only be reproduced if the delivered carbonaceous chondritic material taht was delivered after the disk dissipated was also highly depleted in these noble gases and/or had to be mainly delivered as long as the primordial atmosphere was yet escaping. If proto-Earth accreted less mass during the disk lifetime atmospheric escape cannot explain the present-day ratios, then these ratios have to be explained differently (Marty, 2012).

Impact erosion and outgassing of volatiles from magma pools, oceans, and subsequent escape to space should have depleted the planetary embryos that acted in the final accretion of the early terrestrial planets. Latest reproduction attempts of atmospheric Ar and Ne isotope ratios and bulk K/U ratios on Venus agree with the findings related to early Earth's evolution that the young Sun was less active than a moderate rotator and EUV fluxes corresponding to a faster rotating young Sun were very unlikely.

Thus, new precise measurements by future Venus' space missions of atmospheric Ar, Ne, Kr, and Xe noble gas isotope ratios and more data of moderately volatile elements distributed over a wide surface area would give a precise evolution scenario of Venus and would further help to constrain the Sun's history. For terrestrial exoplanets one can expect that different accretion scenarios and different radiation fields of their host stars will result in different heat producing elements and hence in different tectonic regimes (see also O'Neill et al. 2020; this issue), which can have tremendous implication for the planet's evolution and habitability.

From D/H analysis, one finds that Mars lost a huge fraction of its initial water inventory during the first 500 Myr after its origin. The low gravity and high EUV fluxes of the young Sun led also to high escape rates of its bulk atmosphere. After the EUV flux decreased at pprox4-4.1 Ga, volcanic outgassing and impacts could have resulted in the build-up of a denser atmosphere of < 1 bar, which is also constrained by analysis of 36 Ar/ 38 Ar and 13 C/ 12 C data. Analysis of the atmosphere enriched in ¹³C isotopes compared to Mars' mantle may be an indication that early Mars was more reduced than today. CO_2 could have been a minor component of the early Martian atmosphere as compared with CO or CH₄ from a thermodynamic point of view. Both molecules are depleted in ¹³C and preferentially escaped into space, whereas 13 C-enriched CO₂ is ultimately deposited as carbonate. CO₂ could have been released from surface reservoirs during periods of climate change or impacts and may have modified the atmospheric surface pressure several times during the planet's history. Since \approx 3.5 Ga, when the young Sun's EUV flux decreased to lower values, Mars' atmospheric evolution until today was determined by a complex interplay of ion escape, atmospheric sputtering and suprathermal hot atom escape, carbonate precipitation, and serpentinization that led to today's low surface pressure.

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