

Automated Detection of Coronal Loops Using a Wavelet Transform Modulus Maxima Method

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Abstract We propose and test a wavelet transform modulus maxima method for the automated detection and extraction of coronal loops in extreme ultraviolet images of the solar corona. This method decomposes an image into a number of size scales and tracks enhanced power along each ridge corresponding to a coronal loop at each scale. We compare the results across scales and suggest the optimum set of parameters to maximize completeness, while minimizing detection of noise. For a test coronal image, we compare the global statistics (*e.g.* number of loops at each length) to previous automated coronal-loop detection algorithms.

Keywords Corona, structures · Active regions, structure

1. Introduction

Historically, images of the Sun have always presented a myriad of features far in advance of our physical understanding of their existence and evolution. New data with increased sensitivity, spatial and temporal resolution are continually challenging our theoretical models of the solar atmosphere. This general statement is especially true in the case of coronal

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loops as observed by the fleet of extreme ultraviolet (EUV) imagers launched since the mid 1990s. Beginning with the *Extreme ultraviolet Imaging Telescope* (EIT: Delaboudinière *et al.*, 1995), this list includes the *Transition Region and Coronal Explorer* (TRACE: Handy *et al.*, 1999) and recently the *Extreme Ultraviolet Imager* (EUVI: Howard *et al.*, 2008) on-board the twin *Solar Terrestrial Earth Relations Observatory* (STEREO: Kaiser *et al.*, 2008) spacecraft. From these data, our current models suggest these coronal loops trace out hot coronal plasma (≈ 1 MK) up to heights of around 50 Mm above the surface of the Sun, but questions remain over their temperature and density profile, temporal evolution, and 3D structure.

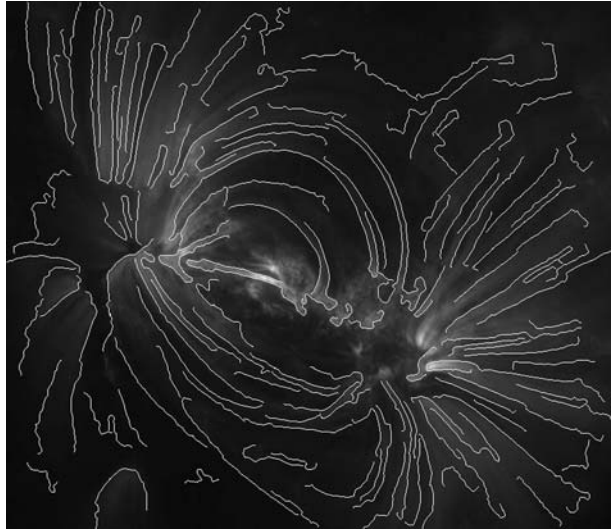
Studies of these open questions are limited by the ability to extract the loop system of interest from the hundreds of others typically present in a solar EUV image. Furthermore each individual loop as viewed by any current instrument is probably a collection of multiple strands, and there will be any number of these strands overlapping along the instrument line of sight. With the launch of STEREO, solar physicists can now view the corona from multiple angles, enabling the 3D reconstruction of coronal loops to try to overcome this problem. The largest problem in this reconstruction is the identification of the same feature from each spacecraft. In the ideal world, the same feature would be tracked in image sequences as observed from each spacecraft and the 3D reconstruction would become a mathematically simple problem. This would enable the scientist to proceed with studying the physical parameter of interest. Ideally this extraction should be automated in order to make the process instantly repeatable. From these issues, there arises a natural requirement for the automated detection of loops based on a statistical approach. Additionally, and perhaps most pressing, the expected data load from the *Solar Dynamics Observatory* (SDO) will be overwhelming without automated feature recognition (McAteer *et al.*, 2004, 2005).

Aschwanden *et al.* (2008) describe five such feature-recognition algorithms for coronal-loop identification and apply each algorithm to an example image from TRACE. This technique of comparing algorithms to each other, as well as to some ground truth, is the best method for rigorous testing of any feature-recognition algorithm. In this paper we apply a sixth such technique to the same TRACE image as described in Section 2. The algorithm is described in detail in Section 3. The technique is called the 2D Wavelet Transform Modulus Maxima (WTMM) method and was originally developed by Arneodo and colleagues as a multifractal analysis formalism (Arneodo, Decoster, and Roux, 2000; Arneodo *et al.*, 2003; Khalil *et al.*, 2006). It has been expanded to characterize the anisotropic nature of complex structures (Khalil *et al.*, 2006, 2009; Snow *et al.*, 2008) and also to perform an automated and objective segmentation of image features of interest from a noisy background (Khalil *et al.*, 2007; Caddle *et al.*, 2007). In solar physics, the technique has been used recently to study the complexity of solar active-region magnetic fields (McAteer, Gallagher, and Ireland, 2005; Conlon *et al.*, 2008, 2010; Kestener *et al.*, 2010; McAteer, Gallagher, and Conlon, 2009), X-ray solar-flare emission (McAteer *et al.*, 2007), and in tracking coronal mass ejections (Byrne *et al.*, 2009). In Section 4 we explain the natural advantages of applying this to coronal-loop identification and compare the global statistics against those identified in Aschwanden *et al.* (2008). Finally, in Section 5 we discuss the benefits and drawbacks of the WTMM algorithm and suggest future extensions.

2. Observations

In order to directly compare our results to those previously published, we have tested our algorithm on the same sub-image as in Aschwanden *et al.* (2008). This image was observed

Figure 1 Edge detection of coronal loops for the underlying test image. The image is 800×600 pixels, corresponding to 296×222 Mm. Maxima chains were obtained from the wavelet transform modulus maxima, using the analyzing wavelets defined in Equation (2) at a size scale of ≈ 11 pixels.



on 19 May 1998, 22:21:43 UT, with an exposure time of 23.172 seconds with a passband centered on 171 \AA . The data were de-biased, flat-fielded, and de-spiked to remove cosmic-ray spikes from the image. The resulting image is shown with our identified loops in Figure 1 and presents three of the main problems with detecting coronal loops: Firstly, they are curvilinear structures: point sources, straight lines, and areas are well behaved in the $[x, y]$ CCD plane; curves are much more difficult. Secondly, many of the loops rise almost vertically through the solar atmosphere. The rapid density drop off in the corona (with a scale height of about 50 Mm) results in a distinct lack of loop tops along many of these features. This makes it difficult to track from footpoint to footpoint. Thirdly there is a collection of small-scale features near the bunch of footpoints, which are typically not of interest in studies of coronal-loop systems.

3. Method

Our method of coronal-loop identification is based on the 2D Wavelet Transform Modulus Maxima method. The continuous nature of the wavelet transform allows us to scan all size scales continuously in order to take full advantage of the space–scale information available. This allows us to perform the segmentation of objects of interest in total objectivity, without any prior knowledge on the size or morphology of the objects.

Image segmentation with continuous wavelets is based on the derivative of a 2D smoothing function (filter) acting as an “edge detector”. Let us consider two wavelets that are, respectively, the partial derivatives with respect to x and y of a 2D smoothing (Gaussian) function,

$$\phi_{\text{Gau}}(x, y) = e^{-(x^2+y^2)/2} = e^{-|x|^2/2}, \quad (1)$$

namely

$$\begin{aligned} \psi_1(x, y) &= \partial \phi_{\text{Gau}}(x, y) / \partial x \quad \text{and} \\ \psi_2(x, y) &= \partial \phi_{\text{Gau}}(x, y) / \partial y. \end{aligned} \quad (2)$$

For any function $f(x, y) \in L^2(\mathbf{R})$ (where $L^2(\mathbf{R})$ consists of all square-integrable functions), the continuous wavelet transform of f with respect to ψ_1 and ψ_2 is expressed as a vector (Mallat and Zhong, 1992; Mallat and Hwang, 1992):

$$\begin{aligned} \mathbf{T}_\psi[f](\mathbf{b}, a) &= \begin{pmatrix} T_{\psi_1}[f] = a^{-2} \int d^2\mathbf{x} \psi_1(a^{-1}(\mathbf{x} - \mathbf{b}))f(\mathbf{x}) \\ T_{\psi_2}[f] = a^{-2} \int d^2\mathbf{x} \psi_2(a^{-1}(\mathbf{x} - \mathbf{b}))f(\mathbf{x}) \end{pmatrix} \\ &= \nabla\{T_{\phi_{\text{Gau}}}[f](\mathbf{b}, a)\} = \nabla\{\phi_{\text{Gau},\mathbf{b},a} * f\}. \end{aligned} \tag{3}$$

Thus, Equation (3) amounts to defining the 2D wavelet transform as the gradient vector of $f(\mathbf{x})$ smoothed by dilated versions $\phi_{\text{Gau}}(a^{-1}\mathbf{x})$ of the Gaussian filter. The wavelet transform can be written in terms of its modulus $\mathcal{M}_\psi[f](\mathbf{b}, a)$ and argument $\mathcal{A}_\psi[f](\mathbf{b}, a)$

$$\mathcal{M}_\psi[f](\mathbf{b}, a) = \sqrt{(T_{\psi_1}[f](\mathbf{b}, a))^2 + (T_{\psi_2}[f](\mathbf{b}, a))^2}, \tag{4}$$

$$\mathcal{A}_\psi[f](\mathbf{b}, a) = \text{Arg}(T_{\psi_1}[f](\mathbf{b}, a) + iT_{\psi_2}[f](\mathbf{b}, a)). \tag{5}$$

The modulus maxima of the wavelet transform, or intensity gradient maxima, are defined by the positions where the modulus of the wavelet transform [$\mathcal{M}_\psi[f](\mathbf{b}, a)$] *i.e.* the gradient, is locally maximal. These WTMM are automatically organized as maxima chains, which act as contour lines of the smoothed image at the considered scales. This is a non-trivial process, which has been discussed and solved empirically by Arneodo, Decoster, and Roux (2000). At a given scale, the algorithm scans all of the boundary lines that correspond to the highest values of the gradient, *i.e.* the maxima chains. For each size scale considered, the algorithm outputs the $[x, y]$ pixel location chain of each edge detected. Each chain then corresponds to a single extracted edge in the image. We identify each edge as a single coronal loop, as shown in Figure 1. This process is repeated continuously at all scales. Further detail on the actual code is given in the [appendix](#).

4. Results

In this section we describe the nature of the WTMM method to trace out and connect edges as one of the major advantages of this method. We compare automatically detected loops against those manually identified in Aschwanden *et al.* (2008). In order to quantitatively compare this technique against those previously published, we also collate a number of global image statistics: the total number of loops detected is an indication of the completeness of the algorithm, the maximum loop length is an indication of the upper limit on the detection, the scaling index of the cumulative distribution shows how the algorithm performs across size scales.

4.1. Loop Tracing

Figure 2a shows an individual loop corresponding to one of the large complete coronal loops near the center of the image. The modulus and angle information used to detect this loop is shown in Figure 2b with the same color scale as Figure 2a. As the loop is traced from one footpoint to the other, both the modulus and angle information change smoothly. The modulus is strongest at the footpoints (where the intensity in the original image is strongest) and drops off by a factor of five at the loop tops. The angle starts off at 0° at the right-most footpoint (*i.e.* pointing to the right of the image, along the x -axis), increases to $\pi/2$ at the

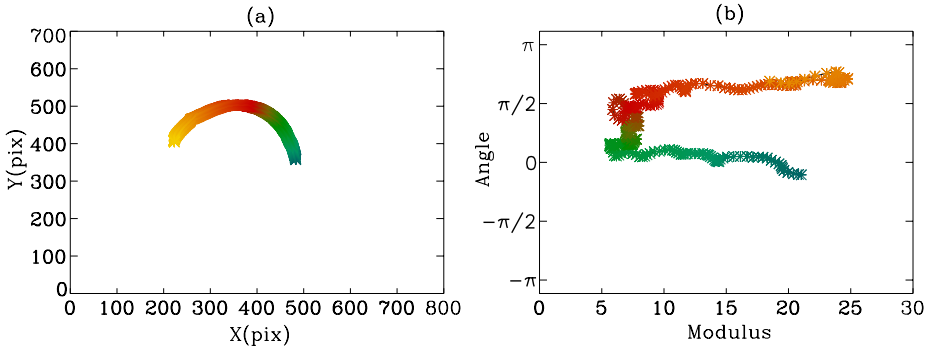


Figure 2 A complete coronal loop. (a) The $[x, y]$ location of each pixel along the loop from one footpoint. The color scheme traces the loop from the first footpoint (green), through the loop top (red) to the second footpoint (yellow). The axes are in pixel units, where one pixel ≈ 370 km. (b) The Modulus [Equation (4)] and Angle [Equation (5)] information at each point along the loop with the same color scheme as in (a), where a positive (negative) angle corresponds to a counter-clockwise (clockwise) rotation from 0 (along the positive x -direction).

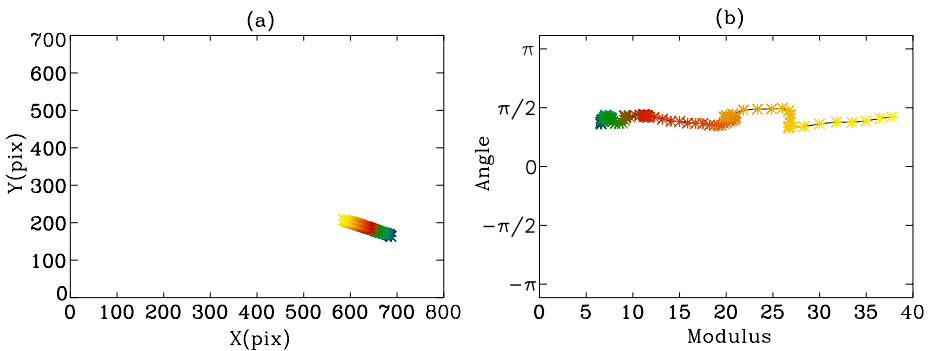


Figure 3 An incomplete coronal loop. (a) The $[x, y]$ location of each pixel along the loop from one footpoint. The color scheme traces the loop from the first footpoint (yellow), up through the atmosphere (red, green) where it presumably connects to another footpoint, possibly elsewhere in the image. The axes are in pixel units, where one pixel ≈ 370 km. (b) The Modulus [Equation (4)] and Angle [Equation (5)] information at each point along the loop with the same color scheme as in (a), where a positive (negative) angle corresponds to a counter-clockwise (clockwise) rotation from 0 (along the positive x -direction).

loop top (*i.e.*, pointing straight up along the y -axis), and increases to π at the second footprint (*i.e.*, pointing to the left, along the x -axis). It is not a completely smooth transition; this is most evident at the loop top, where the angle varies more rapidly. This variation of angle and modulus hence provides two parameters for loop identification. Most importantly, although the modulus may drop off dramatically at the loop top (a manifestation of the decreased signal), the variation in angle is much smoother. A second example, for tracing a partial loop, is displayed in Figure 3 for a loop leg in the lower right of the image. In this case, the angle stays constant ($\approx \pi/2$) along the leg, while the modulus drops off with increasing distance from the footprint.

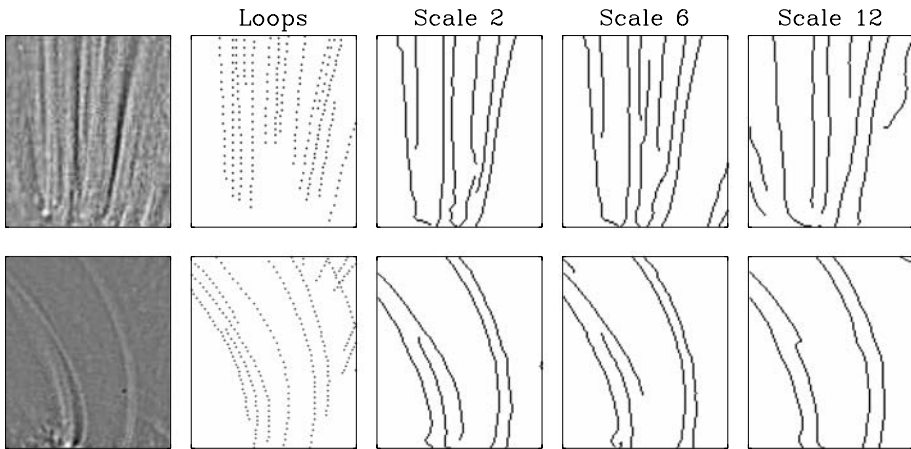


Figure 4 Comparison of (left to right) the smooth-subtracted data, manually identified loops, edges from scale 2 of the WTMM, edges from scale 6 of WTMM, edges from scale 12 of WTMM. Top row is for a sub-image extracted at the left-most footprints of the loop system consisting of mostly vertical loops, bottom row for a sub-image extracted near the right-most footprints of the system and consists of more curved loops. Each sub-image is 100×150 pixels ($37\,000 \times 55\,500$ km).

4.2. Comparison of Loop Locations

A more detailed analysis of extracted loops can be carried out by comparing the edges against those loop manually identified in Aschwanden *et al.* (2008). Figure 4 shows such a comparison for two sub-images of the data. These regions are picked as they illustrate two important spatial regimes of loop identification and demonstrate some of the main positive aspects and drawbacks of the WTMM method. Figure 4 (top) contains a bunch of loop legs. These are mostly linear and display a relatively constant intensity along each loop inside this window. These two aspects should make these features easier to detect however none of the WTMM scales completely extract the manual coordinates. Each of the scales displayed pick out seven–eight of thirteen manually detected loops. The largest scale displayed (scale 12; right column) is unable to pick out any of the weaker, shorter loops however it does produce smooth edges. The smallest scale (scale 2; middle column) identifies more of the smaller loops, but at the expense of producing a more ragged edge. Scale 6 still contains many of small loops, but also produces a smooth edge. Scale 6 is also the only scale that captures the loops in the bottom right of this sub-image. Figure 4 (bottom) shows a selection of the curvi-linear features. As these are curves in $[x, y]$ pixel space and show a drop in intensity, they are normally more difficult to extract. However the ability of the WTMM method to chain in angle space assists in overcoming this problem. Of the seven curved, manually detected features, each scale reproduces four–five. The large, but weak, loop near the center of the sub-image is not detected at any scale. The wavelet transform power lies below the threshold for identification. Lowering our threshold to extract this loop results in pulling out an unmanageable number of weak non-loop features. Scale 2 and scale 6 are similar; the main difference is the ability of scale 6 to begin to detect the loop in the top left. At scale 12, the smoothing window is so large that the tracking algorithm jumps from one large loop to a second large loop. Clearly there is a trade-off to be made between completeness and quality of extracted edges.

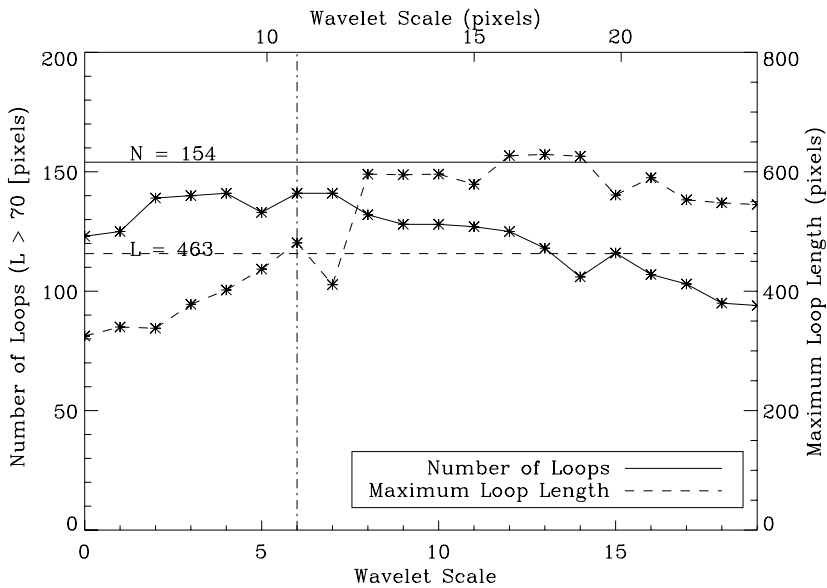


Figure 5 The completeness and upper limits across each wavelet scale. The bottom axis corresponds to the wavelet scale order; the top axis is the same value converted to pixel units. The total number of loops with size greater than 70 pixels [N] is plotted as a solid line and asterisks according to the left axis. The solid horizontal line is plotted at the desired complete value of $N = 154$. The maximum loop length [L] is plotted as a dashed line with asterisks according to the right axis. The dashed horizontal line is plotted at the desired upper limit of $L = 463$. Both desired values are from the manually identified loops in Aschwanden *et al.* (2008).

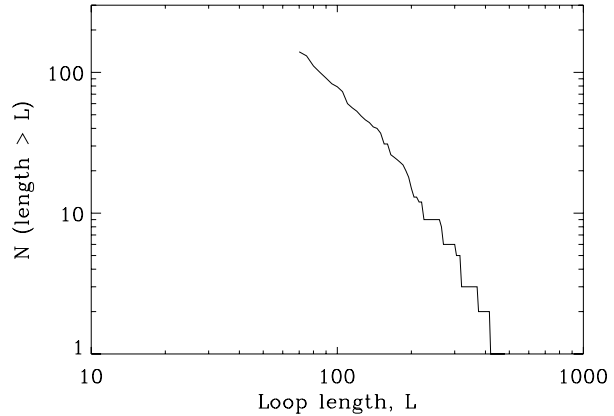
4.3. Number of Loops and Maximum Loop Length

A full quantitative comparison of the positive and negatives of each scale is best achieved by comparing the global statistics. The total number of detected loops and maximum loop length at each wavelet scale size is displayed in Figure 5. The bottom axis shows the wavelet scale, with the corresponding scale in pixels displayed on the top axis. The solid line shows how the total number of loops greater than 70 pixels (left axis) varies with each wavelet scale size. The horizontal solid line shows the ground truth (*i.e.* manually extracted) of $N = 154$. It is clear that the wavelet algorithm consistently underperforms, but it is noticeable that the plot peaks at scale size 6. At this scale we can reach a completeness ratio of 0.91. The dashed line shows how the maximum loop length (right axis) varies with each scale size. The horizontal dashed line shows the ground truth (*i.e.* manually extracted) value of $L = 463$. As expected, the ability of the algorithm to detect the largest loops increases with increased wavelet scale size. It underperforms at small scale size, and over-detects larger structures at large scales sizes. The best performance is also at wavelet scale 6, where we reach a detection ratio of 1.05. It is comforting to note that we achieve maximum accuracy and completeness in these two parameters at the same scale size.

4.4. Scaling Index of Cumulative Frequency Distribution

The cumulative distribution of loops with length greater than 70 pixels is displayed in Figure 6 for scale 6 (≈ 11 pixels). The scaling index [β] of this plot is indicative of the distribution of loop lengths in the image and reflects the tendency of smaller loops of naturally

Figure 6 The cumulative distribution of loops with length greater than 70 pixels.



outnumber larger loops. Aschwanden *et al.* (2008) report a ground truth of $\beta = -2.8$ for the manually traced loops, with the five algorithms producing values between -2.0 and -3.2 . For the WTMM algorithm in Figure 6, $\beta = -2.78 \pm 0.1$.

5. Conclusions and Future Work

EUV images of the corona provide a multitude of information regarding the plasma properties of coronal-loop systems. A number of algorithms currently exist that attempt to extract the locations of these loops automatically. These algorithms are discussed in detail in Aschwanden *et al.* (2008) but here we compare a brief outline of each one against our WTMM method.

The Oriented Connectivity Method (OCM: Lee, Newman, and Gary, 2006a) consists of a preprocessing step to remove non-loop candidates, a linkage step based on magnetic-field extrapolation guide, and post-processing to spline fit and link loop segments. The Dynamic Aperture-Based Loop Segmentation Method (DAM: Lee, Newman, and Gary, 2006b) replaces the linkage step of the OCM with a search for loops by connecting pixels which have a similar Gaussian cross-sectional profile (Cardeo *et al.*, 2003) and orientation. The Oriented-Directivity Loop Tracing Method (ODM: Aschwanden *et al.*, 2008) is essentially a local-directivity version of the OCM, tracing loops locally in two directions from their tops to their footpoints. The Ridge Detection by Automated Scaling (RAS: Inhester, Feng, and Wiegmann, 2008) is a multiscale “ridgel” extension of the OCM. Finally, the Unbiased Detection of Curvi-Linear Structures (UDM: Steger, 1996) is a more generic method consisting the determination of the centroid of the loop structures from a second derivative perpendicular to the loop, and extended (Raghupathy, 2004) to connect structures using a generalized Radon transform. Aschwanden *et al.* (2008) show that the OCM and DAM both successfully extract the large loop features. The ODM, RAS, and UDM codes contain many more free parameters and extracted more segmented, small loop-like features.

The WTMM method presented here offers many natural advantages over other techniques: Firstly, it has naturally directional linkage. The algorithm works by tracking along edges (perpendicular to the gradient) hence the resulting edges require minimum post-processing linkage. Secondly it is naturally multiscale, enabling a scientist to extract features in the size-scale range of interest. Essentially we negate the need for preprocessing by simply studying the scale at which the feature of interest occurs. These two features combine to provide a

good localization and a single response to each edge in the image. These advantages are evident in the good comparison of global statistics against existing algorithms from Aschwanden *et al.* (2008). We reach a completeness of 0.91 and detection ratio of 1.05, with a scaling index for the cumulative distribution of $\beta = -2.78 \pm 0.1$. These three statistics compare favorably to those in Aschwanden *et al.* (2008). The WTMM contains a few free parameters, which we have attempted to optimize for loop detection. The most important free parameter is the choice of the form of the mother wavelet. We choose the Derivative of Gaussian (DOG) as it is well studied mathematically and has previously been successful in tracking edges. We note that other, more naturally curvi-linear, multiscale algorithms exist (*e.g.* curvelets, ridgelets) which may assist in providing better identification of individual coronal loops. The other main free parameters are the choice of thresholds in tracking the edges (see the [appendix](#) for details). Finally we propose that our sixth wavelet scale (corresponding to ≈ 11 pixels) seems to optimize our ability to pick out the small-scale features, retain smoothness in the extracted coordinates, and best agrees with the expected global statistics

An obvious extension of this work is to use the multiscale aspect in a soft-thresholding sense: currently we decide on an optimum scale, there is the possibility to instead decide on a best scaling index. This would consist of attempting to track each feature across scale. A noisy feature is expected to contain a lot of power at small scales, with very little power at larger scales. For a real feature, particularly the large, complete coronal loops, there is expected to be much less change across scales. This may assist in removing non-loop features near the footpoints of the loops system. We also note that our extracted loops, even at scale 6, are probably too ragged for stereoscopy and a degree of post-processing smoothing may still be necessary. When applied to EUVI images from both STEREO Ahead and Behind data, this may allow for a 3D reconstruction of these loop systems with less manual labor. We expect algorithms such as the one studied here will be vital in the SDO era for the tracking of loops across a sequence of images, and hence in the expanding field of coronal seismology.

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Appendix

A.1. Light Weight Notations for WTMM Formulae

Let $f(x, y)$ be the input image to be analyzed and T the wavelet transform vector of f at scale a . The wavelet transform components are

$$f_x = \partial_x(f \star \phi_a), \quad (6)$$

$$f_y = \partial_y(f \star \phi_a). \quad (7)$$

The square modulus of the wavelet transform vector is $M^2 = f_x^2 + f_y^2$. Let us note higher-order derivative of f this way using multiple x - or y -indices: $\partial_x f_x = f_{xx}$, $\partial_y f_x = f_{xy}$, *etc.*

A.2. WTMM Edge Definition

The WTMM are defined as the locations of the points where there is a maximum of the wavelet transform modulus along the direction of the wavelet transform vector, *i.e.*, WTMM

are the location of the greatest slope in the $f \star \phi_a$ landscape. The steepest slope line is not exactly orthogonal to the WTMM edge (only in particular cases). To find those points, we need to evaluate the scalar quantity (N) defined as the dot product

$$N = \nabla(M^2) \bullet \mathbf{T}. \tag{8}$$

At each WTMM location, N is zero and N changes its sign when moving along the direction \mathbf{T} and crossing the WTMM. This is not enough to clearly identify a maximum, we also require that the second derivative along the direction \mathbf{T} must be strictly negative, *i.e.*

$$N' = \partial_x N f_x + \partial_y N f_y < 0. \tag{9}$$

The quantities N and N' are computed exactly:

$$N = \nabla(M^2) \bullet \mathbf{T} = 2f_x^2 f_{xx} + 4f_x f_y f_{xy} + 2f_y^2 f_{yy}. \tag{10}$$

N' is given by

$$\begin{aligned} N' = \partial_x N f_x + \partial_y N f_y = & 4f_x^2 f_{xx}^2 + 4f_y^2 f_{yy}^2 + 2f_x^3 f_{xxx} + 2f_y^3 f_{yyy} \\ & + 4f_x^2 f_{xy}^2 + 4f_y^2 f_{xy}^2 \\ & + 8f_x f_y f_{xy} f_{yy} + 8f_x f_y f_{xy} f_{xx} \\ & + 6f_x^2 f_y f_{xxy} + 6f_x f_y^2 f_{xyy}. \end{aligned}$$

A.3. WTMM Computation Algorithm

Finally, the full algorithm is summed up. It is essentially a FOR loop over pixel location. Note that the final edge image is made of the pixels containing a WTMM. The modulus at the “exact” WTMM location is adjusted using a polynomial fit using points along direction \mathbf{T} .

Algorithm 1 WTMM edges computation algorithm

Require: $f(i, j)$ the input image

Require: a scale parameter

compute WT of f at scale a : $\mathbf{T} = \nabla(f \star G_\sigma)$

compute N and N'

for pixel $(i, j) \in$ image range **do**

if $N' < 0$ and N changes sign in 3×3 -neighborhood **then**

 pixel (i, j) is labelled as a WTMM

 perform a third-order polynomial interpolation along direction \mathbf{T} to get an accurate wavelet transform modulus value at maximum

end if

end for

WTMM edge image

Several free parameters have been statistically and/or empirically adjusted over the years. The following have been used for this study.

- *Spatial resolution: Finite-size and edge effects.* The minimum scale at which the wavelet transform is still resolved is seven pixels. On the other hand, in order to avoid artificial effects from the edges of the images, only the central 72% of the original wavelet transformed image should be kept for analysis. A methodical calculation of this parameter is carried out in Arneodo, Decoster, and Roux (2000).
- *Scale resolution.* For an image of 1024×1024 pixels, usually 50 wavelet scales are considered, from seven pixels to ≈ 200 pixels, in \log_2 steps (*i.e.* very high resolution for small scales and low resolution at high scales). The seven-pixel lower limit restricts our resolution and is adopted to reduce the computational overhead.

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