

Parameter estimation based on discrete observations of fractional Ornstein–Uhlenbeck process of the second kind

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Abstract Fractional Ornstein–Uhlenbeck process of the second kind (fOU₂) is a solution of the Langevin equation $dX_t = -\theta X_t dt + dY_t^{(1)}$, $\theta > 0$ with a Gaussian driving noise $Y_t^{(1)} := \int_0^t e^{-s} dB_{a_s}$, where $a_t = He^{\frac{t}{H}}$ and B is a fractional Brownian motion with Hurst parameter $H \in (0, 1)$. In this article we consider the case $H > \frac{1}{2}$, and by using the ergodicity of fOU₂ process we construct consistent estimators for the drift parameter θ based on discrete observations in two possible cases: (i) the Hurst parameter H is known and (ii) the Hurst parameter H is unknown. Moreover, using Malliavin calculus techniques we prove central limit theorems for our estimators which are valid for the whole range $H \in (\frac{1}{2}, 1)$.

Keywords Fractional Ornstein–Uhlenbeck processes · Malliavin calculus · Multiple Wiener integrals · Central limit theorem (CLT) · Parameter estimation

Mathematics Subject Classification 60G22 · 60H07 · 62F99

1 Introduction

1.1 Motivation and overview

Assume $B = \{B_t\}_{t \geq 0}$ is a fractional Brownian motion with Hurst parameter $H \in (0, 1)$, i.e. a continuous, centered Gaussian process with covariance function

$$R_H(s, t) = \frac{1}{2} \left\{ s^{2H} + t^{2H} - |t - s|^{2H} \right\}, \quad s, t \geq 0.$$

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Consider the following Langevin equation with a drift parameter $\theta > 0$ and a driving noise N

$$dX_t = -\theta X_t dt + dN_t. \tag{1.1}$$

When the driving noise $N = B$ is a fractional Brownian motion, a solution of the Langevin equation (1.1) is called the fractional Ornstein–Uhlenbeck process of the first kind, (fOU₁) in short. The fractional Ornstein–Uhlenbeck process of the second kind is a solution of the Langevin Eq. (1.1) with a driving noise $N_t = Y_t^{(1)} := \int_0^t e^{-s} dB_{a_s}$, where $a_t = He^{\frac{t}{H}}$. Terms “of the first kind” and “of the second kind” are taken from Kaarakka and Salminen (2011). It is well known that the classical Ornstein–Uhlenbeck process, i.e. when the driving noise $N = W$ is a standard Brownian motion, has the same finite dimensional distributions as the Lamperti transformation (see 2.6 for definition) of Brownian motion. However, when one replaces Brownian motion with a fractional Brownian motion the solution of the Langevin equation (1.1) is different from the one that is obtained by the Lamperti transformation of a fractional Brownian motion, see Cheridito et al. (2003), Kaarakka and Salminen (2011). The motivation behind introducing the noise process $N = Y^{(1)}$ is related to the Lamperti transformation of fractional Brownian motion. We refer to Subsection 2.2.2 or (Kaarakka and Salminen 2011, Sect. 3) for more details.

Usually statistical models with fractional processes exhibit short (long) memory property for $H < \frac{1}{2}$ ($H > \frac{1}{2}$, respectively) and this is true for (fOU₁) processes. In contrast, the fOU₂ process always exhibits short range dependence regardless of the Hurst parameter H . This phenomenon makes fOU₂ an interesting process for modelling in many different disciplines. For example, for applications of short memory processes in econometric or in modelling the extremes of time series see Mynbaev (2011), Chavez-Demoulin and Davison (2012) respectively.

In this article we use ergodicity of the fOU₂ process to construct consistent estimator of the drift parameter θ based on observations of the process at discrete times. More precisely, assume that the process is observed at discrete times $0, \Delta_N, 2\Delta_N, \dots, N\Delta_N$ and let $T_N = N\Delta_N$ denote the length of the observation window. We show that:

- (i) **when H is known** one can construct a strongly consistent estimator $\hat{\theta}$, introduced in Theorem 3.2, with asymptotic normality property under the mesh conditions

$$T_N \rightarrow \infty, \quad \text{and} \quad N\Delta_N^2 \rightarrow 0$$

with arbitrary mesh Δ_N such that $\Delta_N \rightarrow 0$ as N tends to infinity.

- (ii) **when H is unknown** one can construct another strongly consistent estimator $\tilde{\theta}$, introduced in Theorem 5.1, with asymptotic normality property under the restricted mesh condition

$$\Delta_N = N^{-\alpha}, \quad \text{with} \quad \alpha \in \left(\frac{1}{2}, \frac{1}{4H-2} \wedge 1 \right).$$

1.2 History and further motivations

Statistical inference of the drift parameter θ based on a data recorded from continuous (discrete) trajectories of X is an interesting problem in the realm of mathematical statistics. In the case of diffusion processes with Brownian motion as a driving noise the problem is well studied (e.g. see Kutoyants (2004) and references therein among many others). However, the estimation of the drift parameter becomes very challenging with fractional processes as

a driving noise. This is mainly because of the fact that fractional Brownian motion B with Hurst parameter $H \neq \frac{1}{2}$ is neither a semimartingale nor a Markov process (we refer to the recent book [Prakasa Rao \(2010\)](#) for more details). In the case of the fractional Ornstein–Uhlenbeck process of the first kind, maximum likelihood (MLE) and least squares (LSE) estimators based on continuous observations of the process are considered in [Kleptsyna and Breton \(2002\)](#) and [Hu and Nualart \(2010\)](#) respectively. In this case it turns out that both MLE and LSE provide strongly consistent estimators. Moreover, the asymptotic normality of MLE is shown in [Bercu et al. \(2011\)](#) for values $H > \frac{1}{2}$ and for LSE in [Hu and Nualart \(2010\)](#) for values $H \in [\frac{1}{2}, \frac{3}{4}]$. In the case of the fractional Ornstein–Uhlenbeck process of the second kind, [Azmoodeh and Morlanes \(2013\)](#) showed that LSE is a consistent estimator using continuous observations. Moreover, they showed that a central limit theorem for LSE holds for the whole range $H > \frac{1}{2}$.

The main feature of this paper is to provide strongly consistent estimators for the drift parameter θ based on discrete observations of the process X together with CLTs using the modern approach of Malliavin calculus for normal approximations [Nourdin and Peccati \(2012\)](#). From practical point of view, it is very important to assume that we have a data collected from process X observed at discrete times. In addition to its applicability, such a demand makes the problem more delicate. Therefore, such a problem could not remain open for the fractional Ornstein–Uhlenbeck process of the first kind. In fact, estimation of the drift parameter θ for the fOU₁ process with discretization procedure of integral transform is considered in [Xiao et al. \(2011\)](#) assuming that the Hurst parameter H is known. In the same setup, [Brouste and Iacus \(2012\)](#) introduced an estimation procedure that can be used to estimate both the drift parameter θ and the Hurst parameter H based on discrete observations. In this paper, we also display a new estimation method that can be used to estimate the drift parameter θ of the fOU₂ process based on discrete observations when the Hurst parameter H is unknown (Theorem 5.1).

1.3 Plan

The paper is organized as follows. In Sect. 2 we give auxiliary facts on Malliavin calculus and fractional Ornstein–Uhlenbeck processes. Section 3 is devoted to estimation of the drift parameter when H is known. In Sect. 4 we give a short explanation how the Hurst parameter H can be estimated by using discrete observations. Section 5 deals with estimation of the drift parameter when H is unknown. Finally, some technical lemmas are collected to Appendix A.

2 Auxiliary facts

2.1 A brief review on Malliavin calculus

In this subsection we briefly introduce some basic facts on Malliavin calculus with respect to Gaussian processes needed in this paper. We also recall some results how Malliavin calculus can be used to obtain a central limit theorem for a sequence of multiple Wiener integrals. For more details on the topic, we refer to [Alos et al. \(2001\)](#), [Nualart \(2006\)](#), [Nourdin and Peccati \(2012\)](#). Let W be a Brownian motion and let $G = \{G_t\}_{t \in [0, T]}$ be a continuous centered Gaussian process of the form

$$G_t = \int_0^t K(t, s) dW_s,$$

where the *Volterra* kernel K , i.e. $K(t, s) = 0$ for all $s > t$, satisfies $\sup_{t \in [0, T]} \int_0^t K(t, s)^2 ds < \infty$. Moreover, we assume that for any s the function $K(\cdot, s)$ is of bounded variation on any interval $(u, T]$ for all $u > s$. A typical example of this type of Gaussian processes is a fractional Brownian motion. It is known that for $H > \frac{1}{2}$ the kernel takes the form

$$K_H(t, s) = c_H s^{\frac{1}{2}-H} \int_s^t (u-s)^{H-\frac{3}{2}} u^{H-\frac{1}{2}} du.$$

Moreover, we have the following inverse relation

$$W_t = B((K_H^*)^{-1}(\mathbf{1}_{[0,t]})), \tag{2.1}$$

where the operator K_H^* is defined as

$$(K_H^* \varphi)(s) = \int_s^T \varphi(t) \frac{\partial K_H}{\partial t}(t, s) dt.$$

Consider the set \mathcal{E} of all step functions on $[0, T]$. The Hilbert space \mathcal{H} associated to the process G is the closure of \mathcal{E} with respect to inner product

$$\langle \mathbf{1}_{[0,t]}, \mathbf{1}_{[0,s]} \rangle_{\mathcal{H}} = R_G(t, s),$$

where $R_G(t, s)$ denotes the covariance function of G . The mapping $\mathbf{1}_{[0,t]} \mapsto G_t$ can be extended to an isometry between the Hilbert space \mathcal{H} and the Gaussian space \mathcal{H}_1 associated with the process G . Consider next the space \mathcal{S} of all smooth random variables of a form

$$F = f(G(\varphi_1), \dots, G(\varphi_n)), \quad \varphi_1, \dots, \varphi_n \in \mathcal{H}, \tag{2.2}$$

where $f \in C_b^\infty(\mathbb{R}^n)$. For any smooth random variable F of the form (2.2), we define its Malliavin derivative $D^{(G)} = D$ as an element of $L^2(\Omega; \mathcal{H})$ by

$$DF = \sum_{i=1}^n \partial_i f(G(\varphi_1), \dots, G(\varphi_n)) \varphi_i.$$

In particular, $DG_t = \mathbf{1}_{[0,t]}$. We denote by $\mathbb{D}_G^{1,2} = \mathbb{D}^{1,2}$ the space of all square integrable Malliavin derivative random variables as the closure of the set \mathcal{S} of smooth random variables with respect to the norm

$$\|F\|_{1,2}^2 = \mathbb{E}|F|^2 + \mathbb{E}(\|DF\|_{\mathcal{H}}^2).$$

Consider next a linear operator K^* from \mathcal{E} to $L^2[0, T]$ defined by

$$(K^* \varphi)(s) = \varphi(s)K(T, s) + \int_s^T [\varphi(t) - \varphi(s)] K(dt, s),$$

where $K(dt, s)$ stands for the measure associated to the bounded variation function $K(\cdot, s)$. The Hilbert space \mathcal{H} generated by covariance function of the Gaussian process G can be represented as $\mathcal{H} = (K^*)^{-1}(L^2[0, T])$ and $\mathbb{D}_G^{1,2}(\mathcal{H}) = (K^*)^{-1}(\mathbb{D}_W^{1,2}(L^2[0, T]))$. Furthermore, for any $n \geq 1$ let \mathcal{H}_n be the n th Wiener chaos of G , i.e. the closed linear subspace of $L^2(\Omega)$ generated by the random variables $\{H_n(G(\varphi)), \varphi \in \mathcal{H}, \|\varphi\|_{\mathcal{H}} = 1\}$ where H_n is the n th Hermite polynomial. It is well known that the mapping $I_n^G(\varphi^{\otimes n}) = n!H_n(G(\varphi))$ provides a linear isometry between the symmetric tensor product $\mathcal{H}^{\otimes n}$ and the space \mathcal{H}_n . The random variables $I_n^G(\varphi^{\otimes n})$ are called *multiple Wiener* integrals of order n with respect to the Gaussian process G .

Let $\mathcal{N}(0, \sigma^2)$ denote the Gaussian distribution with zero mean and variance σ^2 . We use notation $\xrightarrow{\text{law}}$ for convergence in distribution. The next proposition provides a central limit theorem for a sequence of multiple Wiener integrals of fixed order.

Proposition 2.1 *Nualart and Ortiz-Latorre (2008)* Let $\{F_n\}_{n \geq 1}$ be a sequence of random variables in the q th Wiener chaos \mathcal{H}_q with $q \geq 2$ such that $\lim_{n \rightarrow \infty} \mathbb{E}(F_n^2) = \sigma^2$. Then the following statements are equivalent:

- (i) $F_n \xrightarrow{\text{law}} \mathcal{N}(0, \sigma^2)$ as n tends to infinity.
- (ii) $\|DF_n\|_{\mathcal{H}}^2$ converges in $L^2(\Omega)$ to $q\sigma^2$ as n tends to infinity.

2.2 Fractional Ornstein–Uhlenbeck processes

In this subsection we briefly introduce fractional Ornstein–Uhlenbeck processes although we mostly focus on fractional Ornstein–Uhlenbeck process of the second kind for which we also provide some new results. Our main references are [Cheridito et al. \(2003\)](#), [Kaarakka and Salminen \(2011\)](#).

2.2.1 Fractional Ornstein–Uhlenbeck processes of the first kind

Let $B = \{B_t\}_{t \geq 0}$ be a fractional Brownian motion with Hurst parameter $H \in (0, 1)$. To obtain a fractional Ornstein–Uhlenbeck process, consider the following Langevin equation

$$dU_t^{(H, \xi_0)} = -\theta U_t^{(H, \xi_0)} dt + dB_t, \quad U_0^{(H, \xi_0)} = \xi_0. \tag{2.3}$$

The solution of the SDE (2.3) can be expressed as

$$U_t^{(H, \xi_0)} = e^{-\theta t} \left(\xi_0 + \int_0^t e^{\theta s} dB_s \right). \tag{2.4}$$

Notice that the stochastic integral can be understood as a pathwise Riemann–Stieltjes integral or, equivalently, as a Wiener integral. Let \hat{B} denote a two sided fractional Brownian motion. The special selection

$$\xi_0 := \int_{-\infty}^0 e^{\theta s} d\hat{B}_s$$

leads to a unique (in the sense of finite dimensional distributions) stationary Gaussian process $U^{(H)}$ of the form

$$U_t^{(H)} = \int_{-\infty}^t e^{-\theta(t-s)} d\hat{B}_s. \tag{2.5}$$

Definition 2.1 [Kaarakka and Salminen \(2011\)](#) The process $U^{(H, \xi_0)}$ given by (2.4) is called a fractional Ornstein–Uhlenbeck process of the first kind with initial value ξ_0 . The process $U^{(H)}$ defined in (2.5) is called a stationary fractional Ornstein–Uhlenbeck process of the first kind.

Remark 2.1 It is shown in [Cheridito et al. \(2003\)](#) that the covariance function of the stationary process $U^{(H)}$ decays like a power function, and hence $U^{(H)}$ is ergodic. Furthermore, for $H \in (\frac{1}{2}, 1)$ the process $U^{(H)}$ exhibits long range dependence.

2.2.2 Fractional Ornstein–Uhlenbeck processes of the second kind

Now we define a new stationary Gaussian process $X^{(\alpha)}$ by means of Lamperti transformation of the fractional Brownian motion B . More precisely, we set

$$X_t^{(\alpha)} := e^{-\alpha t} B_{a_t}, \quad t \in \mathbb{R}, \tag{2.6}$$

where $\alpha > 0$ and $a_t = \frac{H}{\alpha} e^{\frac{\alpha t}{H}}$. We aim to represent the process $X^{(\alpha)}$ as a solution to the Langevin equation. To this end, we consider the process $Y_t^{(\alpha)}$ defined via

$$Y_t^{(\alpha)} := \int_0^t e^{-\alpha s} dB_{a_s}, \quad t \geq 0,$$

where again the stochastic integral can be understood as a pathwise Riemann–Stieltjes integral as well as a Wiener integral. Using the self-similarity property of fractional Brownian motion one can see that (Karakka and Salminen 2011, Proposition 6) the process $Y^{(\alpha)}$ satisfies a scaling property

$$\left\{ Y_{t/\alpha}^{(\alpha)} \right\}_{t \geq 0} \stackrel{\text{f.d.d.}}{=} \left\{ \alpha^{-H} Y_t^{(1)} \right\}_{t \geq 0}, \tag{2.7}$$

where $\stackrel{\text{f.d.d.}}{=}$ stands for equality in finite dimensional distributions. Using $Y^{(\alpha)}$, the process $X^{(\alpha)}$ can be viewed as a solution of the Langevin equation

$$dX_t^{(\alpha)} = -\alpha X_t^{(\alpha)} dt + dY_t^{(\alpha)}$$

with random initial value $X_0^{(\alpha)} = B_{a_0} = B_{H/\alpha} \sim \mathcal{N}(0, (\frac{H}{\alpha})^{2H})$. Taking into account the scaling property (2.7), we consider the following Langevin equation

$$dX_t = -\theta X_t dt + dY_t^{(1)}, \quad \theta > 0 \tag{2.8}$$

with $Y^{(1)}$ as the driving noise. The solution of the Eq. (2.8) is given by

$$X_t = e^{-\theta t} \left(X_0 + \int_0^t e^{\theta s} dY_s^{(1)} \right) = e^{-\theta t} \left(X_0 + \int_0^t e^{(\theta-1)s} dB_{a_s} \right) \tag{2.9}$$

with $\alpha = 1$ in a_t . Moreover, special selection $X_0 = \int_{-\infty}^0 e^{(\theta-1)s} dB_{a_s}$ for the initial value X_0 leads to the following unique stationary Gaussian process

$$U_t = e^{-\theta t} \int_{-\infty}^t e^{(\theta-1)s} dB_{a_s}. \tag{2.10}$$

Definition 2.2 Karakka and Salminen (2011) The process X given by (2.9) is called the fractional Ornstein–Uhlenbeck process of the second kind with initial value X_0 . The process U defined in (2.10) is called the stationary fractional Ornstein–Uhlenbeck process of the second kind.

For the rest of the paper we assume $H > \frac{1}{2}$ and we take $X_0 = 0$ in the general solution (2.9). Then the corresponding fractional Ornstein–Uhlenbeck process of the second kind takes the form

$$X_t = e^{-\theta t} \int_0^t e^{(\theta-1)s} dB_{a_s}, \tag{2.11}$$

and we have a useful relation

$$U_t = X_t + e^{-\theta t} \xi, \quad \xi = \int_{-\infty}^0 e^{(\theta-1)s} dB_{d_s}. \tag{2.12}$$

We start with a series of known results on fractional Ornstein–Uhlenbeck processes of the second kind required for our purposes.

Proposition 2.2 *Azmoodeh and Morlanes (2013)* Denote $\tilde{B}_t = B_{t+H} - B_H$ the shifted fractional Brownian motion and let X be the fractional Ornstein–Uhlenbeck process of the second kind given by (2.11). Then there exists a regular (see (Alos et al., 2001, page 767) for definition) Volterra kernel \tilde{L} such that

$$\{X_t\}_{t \in [0, T]} \stackrel{f.d.d}{=} \left\{ \int_0^t e^{-\theta(t-s)} d\tilde{G}_s \right\}_{t \in [0, T]} \tag{2.13}$$

where the Gaussian process \tilde{G} is given by

$$\tilde{G}_t = \int_0^t \left(K_H(t, s) + \tilde{L}(t, s) \right) d\tilde{W}_s$$

and \tilde{W} is a standard Brownian motion.

Remark 2.2 Notice that by a direct computation and applying Lemma 4.3 of Azmoodeh and Morlanes (2013), the inner product of the Hilbert space $\tilde{\mathcal{H}}$ generated by the covariance function of the Gaussian process \tilde{G} is given by

$$\langle \varphi, \psi \rangle_{\tilde{\mathcal{H}}} = \alpha_H H^{2H-2} \int_0^T \int_0^T \varphi(u) \psi(v) e^{(u+v)\left(\frac{1}{H}-1\right)} \left| e^{\frac{u}{H}} - e^{\frac{v}{H}} \right|^{2H-2} dv du$$

where $\varphi, \psi \in \tilde{\mathcal{H}}$ and $\alpha_H = H(2H - 1)$.

The following lemma plays an essential role in the paper. More precisely, we use this lemma to construct our estimators for drift parameter. In what follows, $B(x, y)$ denotes the complete Beta function with parameters x and y .

Proposition 2.3 *Azmoodeh and Morlanes (2013)* Let X be the fractional Ornstein–Uhlenbeck process of the second kind given by (2.11). Then

$$\frac{1}{T} \int_0^T X_t^2 dt \rightarrow \Psi(\theta), \quad T \rightarrow \infty$$

almost surely and in $L^2(\Omega)$, where

$$\Psi(\theta) = \frac{(2H - 1)H^{2H}}{\theta} B((\theta - 1)H + 1, 2H - 1). \tag{2.14}$$

Proposition 2.4 *Kaarakka and Salminen (2011)* The covariance function c of the stationary process U decays exponentially and hence U exhibits short range dependence. More precisely, we have

$$c(t) := \mathbb{E}(U_t U_0) = O \left(\exp \left(- \min \left\{ \theta, \frac{1-H}{H} \right\} t \right) \right), \quad \text{as } t \rightarrow \infty.$$

Let v_U be the variogram of the stationary process U , i.e.

$$v_U(t) := \frac{1}{2} \mathbb{E} (U_{t+s} - U_s)^2 = c(0) - c(t).$$

The following lemma tells us the behavior of the variogram function v_U near zero. For functions f and g , the notation $f(t) \sim g(t)$ as $t \rightarrow 0$ means that $f(t) = g(t) + r(t)$, where $r(t) = o(g(t))$ as $t \rightarrow 0$.

Lemma 2.1 *The variogram function v_U satisfies*

$$v_U(t) \sim Ht^{2H} \text{ as } t \rightarrow 0^+.$$

Proof Due to (Kaarakka and Salminen, 2011, Proposition 3.11) there exists a constant $C(H, \theta) = H(2H - 1)H^{2H(1-\theta)}$ such that

$$c(t) = C(H, \theta)e^{-\theta t} \left(\int_0^{a_t} \int_0^{a_0} (xy)^{(\theta-1)H} |x - y|^{2H-2} dx dy \right).$$

Denote the term inside parentheses by $\Phi(t)$. Then with some direct computations, one can see that

$$\Phi(t) = \frac{a_0^{2\theta H}}{\theta H} B((\theta - 1)H + 1, 2H - 1) + \frac{1}{2\theta H} \left(a_t^{2\theta H} - a_0^{2\theta H} \right) \int_0^{\frac{a_0}{a_t}} z^{(\theta-1)H} (1 - z)^{2H-2} dz.$$

Therefore,

$$\begin{aligned} c(t) &= \frac{(2H - 1)H^{2H}}{\theta} B((\theta - 1)H + 1, 2H - 1)e^{-\theta t} \\ &\quad + \frac{(2H - 1)H^{2H}}{2\theta} (e^{\theta t} - e^{-\theta t}) \int_0^{\frac{a_0}{a_t}} z^{(\theta-1)H} (1 - z)^{2H-2} dz \\ &= c(0) - (2H - 1)H^{2H} \times t \times \int_{\frac{a_0}{a_t}}^1 z^{(\theta-1)H} (1 - z)^{2H-2} dz + r(t) \end{aligned} \tag{2.15}$$

where $r(t) = o(t^{2H})$ as $t \rightarrow 0^+$. Hence, by use of the mean value theorem, we infer that as $t \rightarrow 0^+$ we have

$$\int_{\frac{a_0}{a_t}}^1 z^{(\theta-1)H} (1 - z)^{2H-2} dz \sim \frac{H H^{-2H}}{2H - 1} t^{2H-1}. \tag{2.16}$$

Substituting (2.16) into (2.15) we obtain the claim. □

The next lemma studies the regularity of sample paths of the fractional Ornstein–Uhlenbeck process of the second kind X . Usually Hölder constants are almost surely finite random variables and depend on bounded time intervals where the process is considered. The next lemma gives more probabilistic information on Hölder constants.

Lemma 2.2 *Let X be the fractional Ornstein–Uhlenbeck process of the second kind given by (2.11). Then for every interval $[S, T]$ and every $0 < \epsilon < H$, there exist random variables $Y_1 = Y_1(H, \theta)$, $Y_2 = Y_2(H, \theta, [S, T])$, $Y_3 = Y_3(H, \theta, [S, T])$, and $Y_4 = Y_4(H, \epsilon, [S, T])$ such that for all $s, t \in [S, T]$*

$$|X_t - X_s| \leq (Y_1 + Y_2 + Y_3) |t - s| + Y_4 |t - s|^{H-\epsilon}$$

almost surely. Moreover,

- (i) $Y_1 < \infty$ almost surely,
- (ii) $Y_k(H, \theta, [S, T]) \stackrel{\text{law}}{=} Y_k(H, \theta, [0, T - S]), \quad k = 2, 3,$
- (iii) $Y_4(H, \epsilon, [S, T]) \stackrel{\text{law}}{=} Y_4(H, \epsilon, [0, T - S]).$

Furthermore, all moments of random variables Y_2, Y_3 and Y_4 are finite, and $Y_2(H, \theta, [0, T]), Y_3(H, \theta, [0, T])$ and $Y_4(H, \epsilon, [0, T])$ are increasing in T .

Proof Assume $s < t$. By change of variables formula we obtain

$$X_t = e^{-t} B_{a_t} - e^{-\theta t} B_{a_0} - Z_t,$$

where

$$Z_t = e^{-\theta t} \int_0^t B_{a_u} e^{(\theta-1)u} du.$$

Therefore

$$\begin{aligned} |X_t - X_s| &\leq |B_{a_0}| |e^{-\theta t} - e^{-\theta s}| + e^{-t} |B_{a_t} - B_{a_s}| + |B_{a_s}| |e^{-t} - e^{-s}| \\ &\quad + \left| e^{-\theta t} \int_0^t B_{a_u} e^{(\theta-1)u} du - e^{-\theta s} \int_0^s B_{a_u} e^{(\theta-1)u} du \right| \\ &= I_1 + I_2 + I_3 + I_4. \end{aligned}$$

For the term I_1 , we obtain

$$I_1 \leq \theta |B_{a_0}| |t - s|$$

where $\theta |B_{a_0}|$ is almost surely finite random variable. Similarly for the term I_3 we get

$$I_3 \leq \sup_{u \in [S, T]} e^{-u} |B_{a_u}| |t - s|.$$

Note next that Z is a differentiable process. Hence for the term I_4 we get

$$I_4 \leq \left[\theta \sup_{u \in [S, T]} |Z_u| + \sup_{u \in [S, T]} e^{-u} |B_{a_u}| \right] |t - s|.$$

Moreover, by using (2.12), we have

$$|X_t| \leq |U_t| + |\xi|.$$

As a result we obtain

$$|Z_u| \leq |U_u| + |\xi| + |B_{a_0}| + |e^{-u} B_{a_u}|$$

which implies

$$I_4 \leq \left[\theta \sup_{u \in [S, T]} |U_u| + \theta |\xi| + \theta |B_{a_0}| + (\theta + 1) \sup_{u \in [S, T]} e^{-u} |B_{a_u}| \right] |t - s|.$$

Collecting the estimates for I_1, I_3 and I_4 we obtain

$$\begin{aligned} I_1 + I_3 + I_4 &\leq \left[2\theta |B_{a_0}| + \theta |\xi| \right] |t - s| \\ &\quad + \left[\theta \sup_{u \in [S, T]} |U_u| + (\theta + 2) \sup_{u \in [S, T]} e^{-u} |B_{a_u}| \right] |t - s|. \end{aligned}$$

Put

$$Y_1 = 2\theta|B_{a_0}| + \theta|\xi|, \quad Y_2(H, \theta, [S, T]) := \theta \sup_{u \in [S, T]} |U_u|$$

and finally

$$Y_3(H, \theta, [S, T]) := (\theta + 2) \sup_{u \in [S, T]} e^{-u}|B_{a_u}|.$$

Obviously the random variable Y_1 fulfils property (i). Notice also that U_t and $e^{-u}B_{a_t}$ are continuous, stationary Gaussian processes from which property (ii) follows. Moreover, all moments of supremum of a continuous Gaussian process on a compact interval are finite (see Lifshits (1995) for details on supremum of continuous Gaussian process). Hence it remains to consider the term I_2 . By Hölder continuity of the sample paths of fractional Brownian motion we obtain

$$\begin{aligned} I_2 &\leq e^{-t}C(\omega, H, \epsilon, [S, T])|a_t - a_s|^{H-\epsilon} \\ &\leq C(\omega, H, \epsilon, [S, T])|t - s|^{H-\epsilon}. \end{aligned}$$

To conclude, we obtain (see Nualart and Răşcanu (2002) and remark below) that the random variable $C(\omega, H, \epsilon, [S, T])$ has all the moments and $C(\omega, H, \epsilon, [S, T]) \stackrel{\text{law}}{=} C(\omega, H, \epsilon, [0, T - S])$. Now it is enough to take $Y_4 = C(\omega, H, \epsilon, [S, T])$. □

Remark 2.3 The exact form of the random variable $C(\omega, H, \epsilon, [0, T])$ is given by

$$C(\omega, H, \epsilon, [0, T]) = C_{H,\epsilon}T^{H-\epsilon} \left(\int_0^T \int_0^T \frac{|B_t - B_s|^{\frac{2}{\epsilon}}}{|t - s|^{\frac{2H}{\epsilon}}} dt ds \right)^{\frac{\epsilon}{2}},$$

where $C_{H,\epsilon}$ is a constant. Moreover, for all $p \geq 1$ there exists a constant $c_{\epsilon,p}$ such that $\mathbb{E}C(\omega, H, \epsilon, [0, T])^p \leq c_{\epsilon,p}T^{\epsilon p}$.

3 Estimation of the drift parameter when H is known

We start with the fact that the function Ψ is invertible. This fact allows us to construct an estimator for the drift parameter θ .

Lemma 3.1 *The function $\Psi : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ given by (2.14) is bijective, and hence invertible.*

Proof It is straightforward to see that Ψ is surjective. Hence the claim follows because for any fixed parameter $y > 0$, the complete Beta function $B(x, y)$ is decreasing in the variable x . □

We continue with the following central limit theorem.

Theorem 3.1 *Let X be the fractional Ornstein–Uhlenbeck process of the second kind given by (2.11). Then as T tends to infinity, we have*

$$\sqrt{T} \left(\frac{1}{T} \int_0^T X_t^2 dt - \Psi(\theta) \right) \xrightarrow{\text{law}} \mathcal{N}(0, \sigma^2)$$

where the variance σ^2 is given by

$$\sigma^2 = \frac{2\alpha_H^2 H^{4H-4}}{\theta^2} \int_{[0, \infty)^3} \left[e^{-\theta x - \theta|y-z|} e^{\left(1 - \frac{1}{H}\right)(x+y+z)} \times \left(1 - e^{-\frac{y}{H}}\right)^{2H-2} \left|e^{-\frac{x}{H}} - e^{-\frac{z}{H}}\right|^{2H-2} \right] dz dx dy. \tag{3.1}$$

The proof relies on two lemmas proved in the Appendix where we also show that $\sigma^2 < \infty$. The variance σ^2 is given as iterated integral over $[0, \infty)^3$ and the given equation is probably the most compact form.

Proof of Theorem 3.1 For further use put

$$F_T = \frac{1}{\sqrt{T}} I_2^{\tilde{G}}(\tilde{g}), \tag{3.2}$$

where the symmetric function \tilde{g} of two variables is given by

$$\tilde{g}(x, y) = \frac{1}{2\theta} \left[e^{-\theta|x-y|} - e^{-\theta(2T-x-y)} \right].$$

The notation $I_2^{\tilde{G}}$ refers to multiple Wiener integral with respect to \tilde{G} introduced in Subsection 2.1. By Proposition 2.2 we have

$$X_t \stackrel{\text{law}}{=} I_1^{\tilde{G}}(h(t, \cdot)), \quad h(t, s) = e^{-\theta(t-s)} \mathbf{1}_{s \leq t}.$$

Using product formula for multiple Wiener integrals and Fubini’s theorem we infer that

$$\begin{aligned} \frac{1}{T} \int_0^T X_t^2 dt &\stackrel{\text{law}}{=} \frac{1}{T} \int_0^T \|h(t, \cdot)\|_{\mathcal{H}}^2 dt + \frac{1}{T} I_2^{\tilde{G}} \left(\int_0^T (h(t, \cdot) \otimes h(t, \cdot)) dt \right) \\ &= \frac{1}{T} \int_0^T \mathbb{E} X_t^2 dt + \frac{1}{T} I_2^{\tilde{G}}(\tilde{g}). \end{aligned}$$

We get

$$\sqrt{T} \left(\frac{1}{T} \int_0^T X_t^2 dt - \Psi(\theta) \right) \stackrel{\text{law}}{=} \sqrt{T} \left(\frac{1}{T} \int_0^T \mathbb{E} X_t^2 dt - \Psi(\theta) \right) + F_T. \tag{3.3}$$

Next we note that (see [Azmoodeh and Morlanes 2013](#), Lemma 3.4)

$$\Psi(\theta) = \mathbb{E} U_0^2 = \frac{1}{T} \int_0^T \mathbb{E} U_t^2 dt.$$

Hence

$$\begin{aligned} \frac{1}{T} \int_0^T \mathbb{E} X_t^2 dt - \Psi(\theta) &= \frac{1}{T} \int_0^T (\mathbb{E} X_t^2 - \mathbb{E} U_t^2) dt \\ &= \mathbb{E} U_0^2 \frac{1}{T} \int_0^T e^{-2\theta t} dt - \frac{2}{T} \int_0^T e^{-\theta t} \mathbb{E}(U_t U_0) dt, \end{aligned}$$

and thus we obtain

$$\sqrt{T} \left(\frac{1}{T} \int_0^T \mathbb{E} X_t^2 dt - \Psi(\theta) \right) \rightarrow 0. \tag{3.4}$$

Therefore it suffices to show that

$$F_T \stackrel{\text{law}}{\rightarrow} \mathcal{N}(0, \sigma^2)$$

as T tends to infinity. Now by Lemmas 5.1 and 5.2 presented in the Appendix A we have

$$\|D_s F_T\|_{\mathcal{H}}^2 \xrightarrow{L^2(\Omega)} 2\sigma^2 \quad \text{and} \quad \mathbb{E}(F_T^2) = \frac{2}{T} \|\tilde{g}\|_{\mathcal{H}^{\otimes 2}}^2 \longrightarrow \sigma^2.$$

Hence the result follows by applying Proposition 2.1. □

Now we are ready to state the main result of this section.

Theorem 3.2 *Assume we observe the fractional Ornstein–Uhlenbeck process of the second kind X given by (2.11) at discrete time points $\{t_k = k\Delta_N, k = 0, 1, \dots, N\}$ and put $T_N = N\Delta_N$. Assume that $\Delta_N \rightarrow 0$, $T_N \rightarrow \infty$ and $N\Delta_N^2 \rightarrow 0$ as N tends to infinity. Define*

$$\widehat{\mu}_{2,N} = \frac{1}{T_N} \sum_{k=1}^N X_{t_k}^2 \Delta t_k \quad \text{and} \quad \widehat{\theta}_N := \Psi^{-1}(\widehat{\mu}_{2,N}), \tag{3.5}$$

where Ψ^{-1} is the inverse of the function Ψ given by (2.14). Then $\widehat{\theta}$ is a strongly consistent estimator of the drift parameter θ in the sense that as N tends to infinity, we have

$$\widehat{\theta}_N \longrightarrow \theta \tag{3.6}$$

almost surely. Moreover, we have

$$\sqrt{T_N}(\widehat{\theta}_N - \theta) \xrightarrow{law} \mathcal{N}(0, \sigma_\theta^2), \quad N \rightarrow \infty, \tag{3.7}$$

where

$$\sigma_\theta^2 = \frac{\sigma^2}{[\Psi'(\theta)]^2} \tag{3.8}$$

and σ^2 is given by (3.1).

Proof Applying Lemma 2.2 we obtain for any $\epsilon \in (0, H)$ that

$$\begin{aligned} \sqrt{T_N} \left| \widehat{\mu}_{2,N} - \frac{1}{T_N} \int_0^{T_N} X_t^2 dt \right| &= \frac{1}{\sqrt{T_N}} \left| \sum_{k=1}^N \int_{t_{k-1}}^{t_k} (X_{t_k}^2 - X_t^2) dt \right| \\ &\leq \frac{2}{\sqrt{T_N}} \left(\sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| \int_{t_{k-1}}^{t_k} |X_{t_k} - X_t| dt \right) \\ &\leq \frac{2Y_1(H, \theta)}{\sqrt{T_N}} \left(\sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| \int_{t_{k-1}}^{t_k} (t_k - t) dt \right) \\ &\quad + \frac{2}{\sqrt{T_N}} \left(\sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| Y_2(H, \theta, [t_{k-1}, t_k]) \int_{t_{k-1}}^{t_k} (t_k - t) dt \right) \\ &\quad + \frac{2}{\sqrt{T_N}} \left(\sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| Y_3(H, \theta, [t_{k-1}, t_k]) \int_{t_{k-1}}^{t_k} (t_k - t) dt \right) \\ &\quad + \frac{2}{\sqrt{T_N}} \left(\sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| Y_4(H, \epsilon, [t_{k-1}, t_k]) \int_{t_{k-1}}^{t_k} (t_k - t)^{H-\epsilon} dt \right) \\ &=: I_1 + I_2 + I_3 + I_4. \end{aligned}$$

We begin with last term I_4 . Clearly we have

$$\sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| Y_4(H, \epsilon, [t_{k-1}, t_k]) \leq N \sup_{u \in [0, T_N]} |X_u| Y_4(H, \epsilon, [0, T_N]).$$

By Remark 2.3, we have $\mathbb{E}Y_4(H, \epsilon, [0, T_N])^p \leq CT_N^{\epsilon p}$ for any $p \geq 1$. Hence, thanks to Markov’s inequality, we obtain for every $\delta > 0$ that

$$\mathbb{P} (N^{-\gamma} Y_4(H, \epsilon, [0, T_N]) > \delta) \leq \frac{C^p T_N^{\epsilon p}}{N^{\gamma p} \delta^p}.$$

Now by choosing $\epsilon < \gamma$ and p large enough we obtain

$$\sum_{N=1}^{\infty} \mathbb{P} (N^{-\gamma} Y_4(H, \epsilon, [0, T_N]) > \delta) < \infty.$$

Consequently, Borel-Cantelli Lemma implies that

$$N^{-\gamma} Y_4(H, \epsilon, [0, T_N]) \rightarrow 0$$

almost surely for any $\gamma > \epsilon$. Similarly, we obtain

$$N^{-\gamma} \sup_{u \in [0, T_N]} |X_u| \rightarrow 0$$

almost surely for any $\gamma > 0$. Consequently, we get

$$\frac{1}{N^{1+2\gamma}} \sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| Y_4(H, \epsilon, [t_{k-1}, t_k]) \rightarrow 0$$

almost surely for any $\gamma > \epsilon$. Note also that by choosing $\epsilon > 0$ small enough we can choose γ in such way that $1 + 2\epsilon < 1 + 2\gamma < \frac{3}{4} + \frac{H-\epsilon}{2}$. In particular, this is possible if $\epsilon < \min \{H - \frac{1}{2}, \frac{H}{5}\}$. With this choice we have

$$\begin{aligned} I_4 &\leq \frac{2}{H - \epsilon + 1} \sqrt{T_N} \Delta_N^{H-\epsilon} \frac{1}{N} \sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| Y_4(H, \epsilon, [t_{k-1}, t_k]) \\ &= \frac{2}{H - \epsilon + 1} \sqrt{T_N} \Delta_N^{H-\epsilon} N^{2\gamma} \frac{1}{N^{1+2\gamma}} \sum_{k=1}^N \sup_{u \in [t_{k-1}, t_k]} |X_u| Y_4(H, \epsilon, [t_{k-1}, t_k]) \\ &\rightarrow 0 \end{aligned}$$

almost surely, because the condition $N \Delta_N^2 \rightarrow 0$ and our choice of γ implies that

$$\sqrt{T_N} \Delta_N^{H-\epsilon} N^{2\gamma} = \left(N \Delta_N^{\frac{2H+1-2\epsilon}{1+4\gamma}} \right)^{2\gamma+\frac{1}{2}} \leq (N \Delta_N^2)^{2\gamma+\frac{1}{2}} \rightarrow 0.$$

Treating I_1, I_2 , and I_3 in a similar way, we deduce that

$$\sqrt{T_N} \left| \widehat{\mu}_{2,N} - \frac{1}{T_N} \int_0^{T_N} X_t^2 dt \right| \rightarrow 0 \tag{3.9}$$

almost surely. Moreover, we have convergence (3.6) by Lemma 2.3. To conclude the proof, we set $\mu = \Psi(\theta)$ and use Taylor’s theorem to obtain

$$\begin{aligned} \sqrt{T_N} (\hat{\theta}_N - \theta) &= \frac{d}{d\mu} \Psi^{-1}(\mu) \sqrt{T_N} (\hat{\mu}_{2,N} - \Psi(\theta)) \\ &\quad + R_1(\hat{\mu}_{2,N}) \sqrt{T_N} (\hat{\mu}_{2,N} - \Psi(\theta)) \\ &= \frac{d}{d\mu} \Psi^{-1}(\mu) \sqrt{T_N} \left(\frac{1}{T_N} \int_0^{T_N} X_t^2 dt - \Psi(\theta) \right) \\ &\quad + \frac{d}{d\mu} \Psi^{-1}(\mu) \sqrt{T_N} \left(\hat{\mu}_{2,N} - \frac{1}{T_N} \int_0^{T_N} X_t^2 dt \right) \\ &\quad + R_1(\hat{\mu}_{2,N}) \sqrt{T_N} (\hat{\mu}_{2,N} - \Psi(\theta)) \end{aligned}$$

for some reminder function $R_1(x)$ such that $R_1(x) \rightarrow 0$ when $x \rightarrow \Psi(\theta)$. Now continuity of $\frac{d}{d\mu} \Psi^{-1}$ and Ψ^{-1} implies that R_1 is also continuous. Hence the result follows by using (3.9), Theorem 3.1, Slutsky’s theorem and the fact that

$$\frac{d}{d\mu} \Psi^{-1}(\mu) = \frac{1}{\Psi'(\theta)}.$$

□

Remark 3.1 We remark that it is straightforward to construct strongly consistent estimator without the mesh restriction $\Delta_N \rightarrow 0$. However, in order to obtain central limit theorem using Theorem 3.1, one need to pose the condition $\Delta_N \rightarrow 0$ to get the convergence

$$\sqrt{T_N} \left| \hat{\mu}_{2,N} - \frac{1}{T_N} \int_0^{T_N} X_t^2 dt \right| \rightarrow 0.$$

Remark 3.2 Note that we obtained a consistent estimator which depends on the inverse of the function Ψ . However, to the best of our knowledge there exists no explicit formula for the inverse and hence the inverse has to be computed numerically.

Remark 3.3 Theorem 3.2 imposes different conditions on the mesh Δ_N . One possible choice for the mesh satisfying such conditions is $\Delta_N = \frac{\log N}{N}$.

Remark 3.4 Notice that we obtained strong consistency of the estimator $\hat{\theta}$ without assuming uniform discretization of the partitions. The uniform discretization will play a role in estimating the Hurst parameter H .

4 Estimation of the Hurst parameter H

There are different approaches to estimate the Hurst parameter H of fractional processes. Here we consider an approach which is based on filtering. For more details we refer to [Istas and Lang \(1997\)](#), [Coeurjolly \(2001\)](#).

Let $\mathbf{a} = (a_0, a_1, \dots, a_L) \in \mathbb{R}^{L+1}$ be a filter of length $L + 1$, $L \in \mathbb{N}$, and of order $p \geq 1$, i.e. for all indices $0 \leq q < p$,

$$\sum_{j=0}^L a_j j^q = 0 \quad \text{and} \quad \sum_{j=0}^L a_j j^p \neq 0.$$

We define the dilated filter \mathbf{a}^2 associated to the filter \mathbf{a} by

$$a_k^2 = \begin{cases} a_{k'}, & k = 2k' \\ 0, & \text{otherwise} \end{cases}$$

for $0 \leq k \leq 2L$. Assume that we observe the process X given by (2.11) at discrete time points $\{t_k = k\Delta_N, k = 1, \dots, N\}$ such that $\Delta_N \rightarrow 0$ as N tends to infinity. We denote the generalized quadratic variation associated to filter \mathbf{a} by

$$V_{N,\mathbf{a}} = \frac{1}{N} \sum_{i=0}^{N-L} \left(\sum_{j=0}^L a_j X_{(i+j)\Delta_N} \right)^2$$

and we consider the estimator \widehat{H}_N given by

$$\widehat{H}_N = \frac{1}{2} \log_2 \frac{V_{N,\mathbf{a}^2}}{V_{N,\mathbf{a}}}. \tag{4.1}$$

Assumption (A):

We say the filter \mathbf{a} of the length $L + 1$ and order p satisfies assumption (A) if for any real number r such that $0 < r < 2p$ and r is not an even integer, the following property holds:

$$\sum_{i=0}^L \sum_{j=0}^L a_i a_j |i - j|^r \neq 0.$$

Example 1 A typical example of a filter with finite order satisfying assumption (A) is $\mathbf{a} = (1, -2, 1)$ with order $p = 2$.

Theorem 4.1 *Let \mathbf{a} be a filter of the order $p \geq 2$ satisfying assumption (A) and put $\Delta_N = N^{-\alpha}$ for some $\alpha \in (\frac{1}{2}, \frac{1}{4H-2})$. Then*

$$\widehat{H}_N \longrightarrow H$$

almost surely as N tends to infinity. Moreover, we have

$$\sqrt{N}(\widehat{H}_N - H) \xrightarrow{\text{law}} \mathcal{N}(0, \Gamma(H, \theta, \mathbf{a}))$$

where the variance Γ depends on H, θ and the filter \mathbf{a} and is explicitly computed in Coeurjolly (2001) and also given in Brouste and Iacus (2012).

Remark 4.1 It is worth to mention that when $H < \frac{3}{4}$, it is not necessary to assume that the observation window $T_N = N\Delta_N$ tends to infinity whereas for $H \geq \frac{3}{4}$ condition $T_N \rightarrow \infty$ is necessary (see Istas and Lang 1997). Notice also that $H \geq \frac{3}{4}$ if and only if $\frac{1}{4H-2} \leq 1$.

Proof of Theorem 4.1 Let v_U denote the variogram of the process U . By Lemma 2.1 we have

$$v_U(t) = Ht^{2H} + r(t)$$

as $t \rightarrow 0^+$, where $r(t) = o(t^{2H})$. Moreover, $r(t)$ is differentiable and direct calculations show that for $\epsilon \in (0, 1)$

$$r^{(4)}(t) \leq G|t|^{2H+1-\epsilon-4}.$$

Hence the claim follows by following the proof in [Brouste and Iacus \(2012\)](#) for the fractional Ornstein–Uhlenbeck process of the first kind and applying results of [Istas and Lang \(1997, Theorem3\)](#). To conclude, we note that the given variance is also computed in [Coeurjolly \(2001, p. 223\)](#). □

5 Estimation of the drift parameter when H is unknown

In this section we consider $\Psi(\theta, H)$ instead of $\Psi(\theta)$ to take account the dependence on Hurst parameter H . Let $\mu = \Psi(\theta, H)$. Now implicit function theorem implies that there exists a continuously differentiable function $g(\mu, H)$ such that

$$g(\mu, H) = \theta$$

where θ is the unique solution to equation $\mu = \Psi(\theta, H)$. Hence for every fixed H , we have

$$\frac{\partial g}{\partial \mu}(\mu, H) = \frac{1}{\frac{\partial \Psi}{\partial \theta}(\theta, H)}.$$

Moreover, by chain rule we obtain

$$0 = \frac{d}{dH}g(\Psi(\theta, H), H) = \frac{\partial g}{\partial H} + \frac{\partial g}{\partial \mu} \frac{\partial \mu}{\partial H},$$

and note that here $\frac{\partial g}{\partial \mu}$ and $\frac{\partial \mu}{\partial H}$ are known from which we can compute $\frac{\partial g}{\partial H}$. Let $\widehat{\mu}_{2,N}$ be given by (3.5) and let \widehat{H}_N be given by (4.1) for some filter \mathbf{a} of order $p \geq 2$ satisfying assumption (A). We consider the estimator

$$\widetilde{\theta}_N = g(\widehat{\mu}_{2,N}, \widehat{H}_N) \tag{5.1}$$

for which we have the following result.

Theorem 5.1 *Assume $\Delta_N = N^{-\alpha}$ for some number $\alpha \in (\frac{1}{2}, \frac{1}{4H-2} \wedge 1)$. Then the estimator $\widetilde{\theta}_N$ given by (5.1) is strongly consistent, i.e. as N tends to infinity, we have*

$$\widetilde{\theta}_N \longrightarrow \theta \tag{5.2}$$

almost surely. Moreover, we have

$$\sqrt{T_N} (\widetilde{\theta}_N - \theta) \xrightarrow{law} \mathcal{N}(0, \sigma_\theta^2), \tag{5.3}$$

where the variance σ_θ^2 is given by (3.8).

Proof First note that

$$\begin{aligned} \sqrt{T_N} (\widetilde{\theta}_N - \theta) &= \sqrt{T_N} (g(\widehat{\mu}_{2,N}, \widehat{H}_N) - g(\widehat{\mu}_{2,N}, H)) \\ &\quad + \sqrt{T_N} (g(\widehat{\mu}_{2,N}, H) - g(\mu, H)). \end{aligned} \tag{5.4}$$

Now convergence

$$\sqrt{T_N} (g(\widehat{\mu}_{2,N}, H) - g(\mu, H)) \xrightarrow{law} \mathcal{N}(0, \sigma_\theta^2)$$

is in fact Theorem 3.2. Moreover, by Taylor’s theorem we get

$$\begin{aligned} \sqrt{T_N} \left(g(\hat{\mu}_{2,N}, \hat{H}_N) - g(\hat{\mu}_{2,N}, H) \right) &= \frac{\partial g}{\partial H}(\hat{\mu}_{2,N}, H) \sqrt{T_N} (\hat{H}_N - H) \\ &\quad + \frac{\partial g}{\partial H}(\hat{\mu}_{2,N}, H) R_2(\hat{\mu}_{2,N}, \hat{H}_N) \sqrt{T_N} (\hat{H}_N - H) \end{aligned}$$

for some reminder function R_2 which converges to zero as $(\hat{\mu}_{2,N}, \hat{H}_N) \rightarrow (\mu, H)$. Therefore, by continuity and Theorem 4.1 we obtain

$$\sqrt{T_N} \left(g(\hat{\mu}_{2,N}, \hat{H}_N) - g(\hat{\mu}_{2,N}, H) \right) \rightarrow 0$$

in probability. Hence, we also have

$$\sqrt{T_N} (\hat{\theta}_N - \theta) \xrightarrow{\text{law}} \mathcal{N}(0, \sigma_\theta^2)$$

by Slutsky’s theorem. To conclude the proof, we obtain (5.2) from Eq. (5.4) by continuous mapping theorem. □

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Appendix

Computations used in the paper

Lemma 5.1 For F_T given by (3.2) and the variance σ^2 given by (3.1) we have

$$\|D_s F_T\|_{\mathcal{H}}^2 \xrightarrow{L^2(\Omega)} 2\sigma^2 \tag{5.5}$$

as T tends to infinity.

Proof It is sufficient to show that as T tends to infinity, we have

$$\mathbb{E} \left[\|D_s F_T\|_{\mathcal{H}}^2 - \mathbb{E} \|D_s F_T\|_{\mathcal{H}}^2 \right]^2 \rightarrow 0. \tag{5.6}$$

Indeed, since

$$\lim_{T \rightarrow \infty} \mathbb{E} \|D_s F_T\|_{\mathcal{H}}^2 = 2 \lim_{T \rightarrow \infty} \mathbb{E}(F_T^2),$$

we obtain that (5.6) implies (5.5). Now we have

$$D_s F_T = \frac{2}{\sqrt{T}} I_1^{\tilde{G}}(\tilde{g}(s, \cdot)).$$

Hence, using Remark 2.2, we can write

$$\begin{aligned} \|D_s F_T\|_{\mathcal{H}}^2 &= \frac{4\alpha_H H^{2H-2}}{T} \int_0^T \int_0^T I_1^{\tilde{G}}(\tilde{g}(u, \cdot)) I_1^{\tilde{G}}(\tilde{g}(v, \cdot)) \\ &\quad \times e^{(u+v)\left(\frac{1}{H}-1\right)} \left| e^{\frac{u}{H}} - e^{\frac{v}{H}} \right|^{2H-2} \text{d}v \text{d}u. \end{aligned}$$

Let now $K(u, v)$ denote the kernel associated to the space $\tilde{\mathcal{H}}$ i.e.

$$K(u, v) = e^{(u+v)\left(\frac{1}{H}-1\right)} \left| e^{\frac{u}{H}} - e^{\frac{v}{H}} \right|^{2H-2}. \tag{5.7}$$

Using multiplicative formula for multiple Wiener integrals we see that

$$\begin{aligned} I_1^{\tilde{G}}(\tilde{g}(u, \cdot)) I_1^{\tilde{G}}(\tilde{g}(v, \cdot)) &= \langle \tilde{g}(u, \cdot), \tilde{g}(v, \cdot) \rangle_{\tilde{\mathcal{H}}} + I_2^{\tilde{G}}(\tilde{g}(u, \cdot) \tilde{\otimes} \tilde{g}(v, \cdot)) \\ &=: A_1(u, v) + A_2(u, v). \end{aligned}$$

Here A_1 is deterministic and A_2 has expectation zero. Hence, in order to have (5.6), we need to show that

$$\mathbb{E} \left[\frac{1}{T} \int_0^T \int_0^T A_2(u, v) K(u, v) dv du \right]^2 \rightarrow 0. \tag{5.8}$$

Therefore, by applying Fubini’s Theorem, it suffices to show that

$$\begin{aligned} \frac{1}{T^2} \int_{[0, T]^4} \mathbb{E} [A_2(u_1, v_1) A_2(u_2, v_2)] \\ \times K(u_1, v_1) K(u_2, v_2) du_1 dv_1 du_2 dv_2 \rightarrow 0 \end{aligned} \tag{5.9}$$

as T tends to infinity. First we get

$$\begin{aligned} \mathbb{E} [A_2(u_1, v_1) A_2(u_2, v_2)] &= 2 \int_{[0, T]^4} (\tilde{g}(u_1, \cdot) \tilde{\otimes} \tilde{g}(v_1, \cdot))(x_1, y_1) (\tilde{g}(u_2, \cdot) \tilde{\otimes} \tilde{g}(v_2, \cdot))(x_2, y_2) \\ &\times K(x_1, x_2) K(y_1, y_2) dx_1 dy_1 dx_2 dy_2. \end{aligned}$$

By plugging into (5.9) we obtain that it suffices to have

$$\begin{aligned} \frac{1}{T^2} \int_{[0, T]^8} (\tilde{g}(u_1, \cdot) \tilde{\otimes} \tilde{g}(v_1, \cdot))(x_1, y_1) (\tilde{g}(u_2, \cdot) \tilde{\otimes} \tilde{g}(v_2, \cdot))(x_2, y_2) \\ \times K(x_1, x_2) K(y_1, y_2) K(u_1, v_1) K(u_2, v_2) \\ dv_1 du_2 dv_2 du_1 dx_1 dy_1 dx_2 dy_2 \rightarrow 0 \end{aligned} \tag{5.10}$$

as T tends to infinity. Here we have

$$(\tilde{g}(u, \cdot) \tilde{\otimes} \tilde{g}(v, \cdot))(x, y) = \frac{1}{2} [\tilde{g}(u, x) \tilde{g}(v, y) + \tilde{g}(u, y) \tilde{g}(v, x)].$$

Note first that for every $0 \leq x, y \leq T$, we have that

$$e^{-\theta(2T-x-y)} \leq e^{-\theta|x-y|}.$$

As a consequence, we can omit the term $e^{-\theta(2T-x-y)}$ on function $\tilde{g}(x, y)$. This implies that instead of

$$(\tilde{g}(u_1, \cdot) \tilde{\otimes} \tilde{g}(v_1, \cdot))(x_1, y_1) (\tilde{g}(u_2, \cdot) \tilde{\otimes} \tilde{g}(v_2, \cdot))(x_2, y_2)$$

it is sufficient to consider the following integrand:

$$\begin{aligned}
 & e^{-\theta|u_1-x_1|} e^{-\theta|v_1-y_1|} e^{-\theta|u_2-x_2|} e^{-\theta|v_2-y_2|} \\
 & + e^{-\theta|u_1-x_1|} e^{-\theta|v_1-y_1|} e^{-\theta|u_2-y_2|} e^{-\theta|v_2-x_2|} \\
 & + e^{-\theta|u_1-y_1|} e^{-\theta|v_1-x_1|} e^{-\theta|u_2-x_2|} e^{-\theta|v_2-y_2|} \\
 & + e^{-\theta|u_1-y_1|} e^{-\theta|v_1-x_1|} e^{-\theta|u_2-y_2|} e^{-\theta|v_2-x_2|}.
 \end{aligned}
 \tag{5.11}$$

Next we consider the first term and show that

$$\begin{aligned}
 & \frac{1}{T^2} \int_{[0,T]^8} e^{-\theta|u_1-x_1|} e^{-\theta|v_1-y_1|} e^{-\theta|u_2-x_2|} e^{-\theta|v_2-y_2|} \\
 & \times K(x_1, x_2) K(y_1, y_2) K(u_1, v_1) K(u_2, v_2) \\
 & du_1 dv_1 du_2 dv_2 dx_1 dy_1 dx_2 dy_2 \rightarrow 0.
 \end{aligned}
 \tag{5.12}$$

In what follows C is a non-important constant which may vary from line to line. First it is easy to prove that

$$\int_0^T e^{-\theta|x-y|} dx \leq C,
 \tag{5.13}$$

where constant does not depend on y or T . Moreover, by change of variable we obtain

$$\int_0^T K(x, y) dx \leq 2HB(1 - H, 2H - 1)
 \tag{5.14}$$

for every y and T . Consider now the iterated integral in (5.12). The value of the integral depends on the order of the variables, and eight variables can be ordered in $8! = 40320$ ways. However, it is clear that without loss of generality we can choose the smallest variable, say y_2 , and integrate over region $\{0 < y_2 < u_1, u_2, v_1, v_2, x_1, x_2, y_1 < T\}$. Other cases can be treated similarly with obvious changes. Assume now that the smallest variable is y_2 and denote the second smallest variable by r_7 , i.e.

$$r_7 = \min(u_1, u_2, v_1, v_2, x_1, x_2, y_1).$$

Integrating first with respect to y_2 and applying upper bound $e^{\theta y_2} \leq e^{\theta r_7}$ together with (5.14), we obtain that

$$\begin{aligned}
 & \int_{[0,T]^7} \int_0^{r_7} e^{-\theta|u_1-x_1|} e^{-\theta|v_1-y_1|} e^{-\theta|u_2-x_2|} e^{-\theta v_2 + \theta y_2} \\
 & \times K(x_1, x_2) K(y_1, y_2) K(u_1, v_1) K(u_2, v_2) dy_2 du_1 dv_1 du_2 dv_2 dx_1 dy_1 dx_2 \\
 & \leq C \int_{[0,T]^7} e^{-\theta|u_1-x_1|} e^{-\theta|v_1-y_1|} e^{-\theta|u_2-x_2|} e^{-\theta v_2 + \theta r_7} \\
 & \times K(x_1, x_2) K(u_1, v_1) K(u_2, v_2) du_1 dv_1 du_2 dv_2 dx_1 dy_1 dx_2.
 \end{aligned}$$

Next we integrate with respect to y_1 . In the case when $r_7 = y_1$, we have

$$\int_0^{r_6} e^{-\theta(v_1+v_2-2y_1)} dy_1 \leq C e^{-\theta(v_1+v_2-2r_6)} \leq C,$$

where r_6 is the third smallest variable, and in the case when $r_7 \neq y_1$, we obtain by (5.13)

$$\int_0^T e^{-\theta|v_1-y_1|} e^{-\theta v_2 + \theta r_7} dy_1 \leq C.$$

Hence we obtain upper bound

$$\begin{aligned} & \int_{[0,T]^7} e^{-\theta|u_1-x_1|} e^{-\theta|v_1-y_1|} e^{-\theta|u_2-x_2|} e^{-\theta v_2+\theta r_7} \\ & \quad \times K(x_1, x_2)K(u_1, v_1)K(u_2, v_2) du_1 dv_1 du_2 dv_2 dx_1 dy_1 dx_2 \\ & \leq C \int_{[0,T]^6} e^{-\theta|u_1-x_1|} e^{-\theta|u_2-x_2|} \\ & \quad \times K(x_1, x_2)K(u_1, v_1)K(u_2, v_2) du_1 dv_1 dv_2 dx_1 du_2 dx_2. \end{aligned}$$

Next we integrate first with respect to variables v_1 and v_2 and then with respect to variables u_1 and u_2 . Together with estimates (5.13) and (5.14) this yields

$$\begin{aligned} & \int_{[0,T]^6} e^{-\theta|u_1-x_1|} e^{-\theta|u_2-x_2|} \\ & \quad \times K(x_1, x_2)K(u_1, v_1)K(u_2, v_2) dv_1 dv_2 du_1 du_2 dx_1 dx_2 \\ & \leq C \int_{[0,T]^4} e^{-\theta|u_1-x_1|} e^{-\theta|u_2-x_2|} K(x_1, x_2) du_1 du_2 dx_1 dx_2 \\ & \leq C \int_{[0,T]^2} K(x_1, x_2) dx_1 dx_2 \\ & \leq CT \end{aligned}$$

which gives (5.12). It remains to note that other three terms in (5.11) can be treated with the same arguments since only the "pairing" of variables in terms of form $e^{-\theta|x-y|}$ changes. Thus we have (5.10) and implications (5.10)⇒(5.8)⇒(5.6) ⇒(5.5) complete the proof. □

Lemma 5.2 For F_T given by (3.2) and σ^2 given by (3.1) we have

$$\mathbb{E}[F_T^2] \longrightarrow \sigma^2 \tag{5.15}$$

as T tends to infinity.

Proof Using isometry we obtain

$$\mathbb{E}[F_T^2] = \frac{2}{T} \|\tilde{g}\|_{\mathcal{H}^{\otimes 2}}^2 =: \frac{2I_T}{T}$$

where

$$\begin{aligned} I_T &= \alpha_H^2 H^{4H-4} \int_{[0,T]^4} \tilde{g}(u_1, v_1)\tilde{g}(u_2, v_2)e^{\left(\frac{1}{H}-1\right)(u_1+v_1+u_2+v_2)} \\ & \quad \times \left| e^{\frac{u_2}{H}} - e^{\frac{u_1}{H}} \right|^{2H-2} \left| e^{\frac{v_2}{H}} - e^{\frac{v_1}{H}} \right|^{2H-2} du_1 du_2 dv_1 dv_2. \end{aligned}$$

Recall that

$$\tilde{g}(x, y) = \frac{1}{2\theta} e^{-\theta|x-y|} - \frac{1}{2\theta} e^{-\theta(2T-x-y)}.$$

We first show that we can omit the second term $\frac{1}{2\theta}e^{-\theta(2T-x-y)}$ in the function \tilde{g} . To see this, we have

$$\begin{aligned} & \int_{[0,T]^4} e^{-\theta(2T-u_1-v_1)} \tilde{g}(u_2, v_2) e^{\left(\frac{1}{H}-1\right)(u_1+v_1+u_2+v_2)} \\ & \quad \times \left| e^{\frac{u_2}{H}} - e^{\frac{u_1}{H}} \right|^{2H-2} \left| e^{\frac{v_2}{H}} - e^{\frac{v_1}{H}} \right|^{2H-2} du_1 du_2 dv_1 dv_2 \\ & \leq C(\theta) \int_{[0,T]^4} e^{-\theta(2T-u_1-v_1)} e^{\left(\frac{1}{H}-1\right)(u_1+v_1+u_2+v_2)} \\ & \quad \times \left| e^{\frac{u_2}{H}} - e^{\frac{u_1}{H}} \right|^{2H-2} \left| e^{\frac{v_2}{H}} - e^{\frac{v_1}{H}} \right|^{2H-2} du_1 du_2 dv_1 dv_2 \\ & = C(\theta) \left[\int_0^T \int_0^T e^{-\theta(T-v)+\left(\frac{1}{H}-1\right)(v+u)} \left| e^{\frac{u}{H}} - e^{\frac{v}{H}} \right|^{2H-2} dv du \right]^2. \end{aligned}$$

By change of variables $\tilde{v} = T - v, \tilde{u} = T - u$, and then $x = e^{-\frac{\tilde{v}}{H}}, y = e^{-\frac{\tilde{u}}{H}}$ we infer that this is the same as

$$\left[\int_{e^{-\frac{T}{H}}}^1 \int_{e^{-\frac{T}{H}}}^1 x^{(\theta-1)H} y^{-H} |y-x|^{2H-2} dx dy \right]^2.$$

Let now $x < y$. By change of variable $z = \frac{x}{y}$ we obtain

$$\begin{aligned} & \int_{e^{-\frac{T}{H}}}^1 \int_{e^{-\frac{T}{H}}}^y x^{(\theta-1)H} y^{-H} |y-x|^{2H-2} dx dy \\ & \leq \int_0^1 \int_0^1 y^{\theta H-1} z^{(\theta-1)H} (1-z)^{2H-2} dz dy \\ & \leq \frac{1}{\theta H} B((\theta-1)H+1, 2H-1) \end{aligned}$$

which converges to zero when divided with T tending to infinity. The case $x > y$ can be treated in a similar way, and hence it is sufficient to consider the function

$$\frac{1}{2\theta} e^{-\theta|x-y|}$$

instead of $\tilde{g}(x, y)$. We shall use L'Hopital's rule to compute the limit. Taking derivative with respect to T , we obtain

$$\begin{aligned} \frac{dI_T}{dT} &= \frac{\alpha_H^2 H^{4H-4}}{\theta^2} \int_{[0,T]^3} e^{-\theta|T-u_1|} e^{-\theta|u_2-v_2|} e^{\left(\frac{1}{H}-1\right)(T+u_1+u_2+v_2)} \\ & \quad \times \left| e^{\frac{u_2}{H}} - e^{\frac{u_1}{H}} \right|^{2H-2} \left| e^{\frac{T}{H}} - e^{\frac{v_1}{H}} \right|^{2H-2} du_1 du_2 dv_1 dv_2. \end{aligned}$$

By change of variables $x = T - u_1, y = T - u_2$ and $z = T - v_1$, this reduces to

$$\begin{aligned} \frac{dI_T}{dT} &= \frac{\alpha_H^2 H^{4H-4}}{\theta^2} \int_{[0,T]^3} e^{-\theta x} e^{-\theta|y-z|} e^{\left(1-\frac{1}{H}\right)(x+y+z)} \\ & \quad \times \left(1 - e^{-\frac{y}{H}}\right)^{2H-2} \left| e^{-\frac{x}{H}} - e^{-\frac{z}{H}} \right|^{2H-2} dz dx dy. \end{aligned}$$

Therefore, we have

$$\lim_{T \rightarrow \infty} \frac{dI_T}{dT} = \frac{\alpha_H^2 H^{4H-4}}{\theta^2} \int_{[0, \infty)^3} e^{-\theta x} e^{-\theta|y-z|} e^{\left(1-\frac{1}{H}\right)(x+y+z)} \times \left(1 - e^{-\frac{y}{H}}\right)^{2H-2} \left|e^{-\frac{x}{H}} - e^{-\frac{z}{H}}\right|^{2H-2} dz dx dy.$$

We end the proof by showing that this triple integral, denoted by I , is finite. By use of the obvious bound $e^{-\theta|z-y|} \leq 1$ we infer that

$$\begin{aligned} I &\leq \int_{[0, \infty)^3} e^{-\theta x} e^{\left(1-\frac{1}{H}\right)(x+y+z)} \\ &\quad \times \left(1 - e^{-\frac{y}{H}}\right)^{2H-2} \left|e^{-\frac{x}{H}} - e^{-\frac{z}{H}}\right|^{2H-2} dz dx dy \\ &= \left[\int_0^\infty e^{\left(1-\frac{1}{H}\right)y} \left(1 - e^{-\frac{y}{H}}\right)^{2H-2} dy \right] \\ &\quad \times \left[\int_0^\infty \int_0^\infty e^{-\theta x} e^{\left(1-\frac{1}{H}\right)(x+z)} \left|e^{-\frac{x}{H}} - e^{-\frac{z}{H}}\right|^{2H-2} dz dx \right] \\ &= I_1 \times I_2. \end{aligned}$$

For the term I_1 , we obtain by change of variable $u = e^{-\frac{y}{H}}$ that

$$I_1 = C \int_0^1 u^{-H} (1-u)^{2H-2} du < \infty.$$

For the term I_2 , we obtain by change of variables $u = e^{-\frac{x}{H}}$ and $v = e^{-\frac{z}{H}}$ that

$$\begin{aligned} I_2 &= C \int_0^1 \int_0^1 u^{(\theta-1)H} v^{-H} |u-v|^{2H-2} du dv \\ &= \left[\int_0^1 \int_0^u + \int_0^1 \int_u^1 \right] u^{(\theta-1)H} v^{-H} |u-v|^{2H-2} dv du \\ &= I_{2,1} + I_{2,2}. \end{aligned}$$

For the term $I_{2,1}$, we obtain by change of variable $z = \frac{u}{v}$ that

$$I_{2,1} = C \int_0^1 u^{\theta H-1} \int_0^1 z^{-H} (1-z)^{2H-2} dz du = \frac{1}{\theta H} B(1-H, 2H-1).$$

Similarly for the term $I_{2,2}$, we get by change of variable $z = \frac{u}{v}$ that

$$I_{2,2} = \frac{1}{\theta H} B((\theta-1)H+1, 2H-1).$$

□

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